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# The Partonic Nature of Instantons

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David Tong



Work in progress with Ben Collie

Shifmania, May 2009



Happy Birthday  
Misha!!



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# What are Instantons Made of?

- Instantons have a number of collective coordinates:
- SU(N) Yang-Mills
  - 4 translations, 1 scale size,  $4N-5$  orientation
  - $4N$  in total
- $CP^N$  Sigma-Model
  - 2 translations, 1 scale size,  $2N-3$  orientations
  - $2N$  in total
- An old idea: Instantons are composed of  $N$  partons

Belavin et al '79

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# The Partonic Nature of Instantons

- The conjecture of partons usually framed for
    - $d=3+1$  dimensions (for Yang-Mills)
    - $d=1+1$  dimensions (for sigma-models)
  - Here we revisit this idea in context of
    - $d=4+1$  dimensions (for Yang-Mills)
    - $d=2+1$  dimensions (for sigma-models)
  - Theories are non-renormalizable (effective theories)
  - Instantons are particle-like excitations
  - Is the single instanton a multi-particle state?
    - ...and how can we tell?
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# 5d Yang-Mills

- There exists a UV completion of gauge theories in  $d=4+1$  dimensions Seiberg '96
    - (at least with supersymmetry)
    - ...but we don't know much about it
  - With  $N=2$  supersymmetry, the UV completion is the  $(2,0)$  theory in  $d=5+1$  dimensions
  - The UV completion has  $N^3$  degrees of freedom. Klebanov and Tseytlin '96
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# The Instanton

- The instanton is the KK mode from the sixth dimension

$$M_{\text{inst}} = \frac{8\pi^2}{e^2} = \frac{1}{R}$$

Rozali, '97

- Proposal: The instanton is an N-particle state
  - The N partons inside the instanton are the remnant of the UV degrees of freedom which comprise the (2,0) theory.
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# Circumstantial Evidence

- Turn up the heat and look at cross-over in free-energy

Itzhaki et al '98

$$F \sim N^2 T^5 \longrightarrow RN^3 T^6$$

- The transition happens at the temperature

$$T \sim \frac{1}{NR} \sim \frac{1}{e^2 N}$$

- This had to be the case: this is where the 5d theory is strongly coupled, so this is where we need new UV degrees of freedom.

# More Circumstantial Evidence

- Anomaly coefficient for  $G = ADE$  is conjectured to be Intriligator, '00

$$c_2(G) \times |G|$$

dual coxeter number  
= number of instanton moduli

dimension of group

- Are the partons in the adjoint of  $G$ ?
  - What is the mechanism that confines them?
- Quantizing the scale size of the instanton gives a continuous spectrum...natural for a multi-particle state.
  - But the moduli space is *not* that of  $N$  free objects...does this contain clues about confining mechanism?

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# A Toy Model

- These questions are difficult to answer in case of Yang-Mills
  - The  $CP^N$  sigma-model provides (as always!) a nice toy-model where we can see how some of these issues are resolved.
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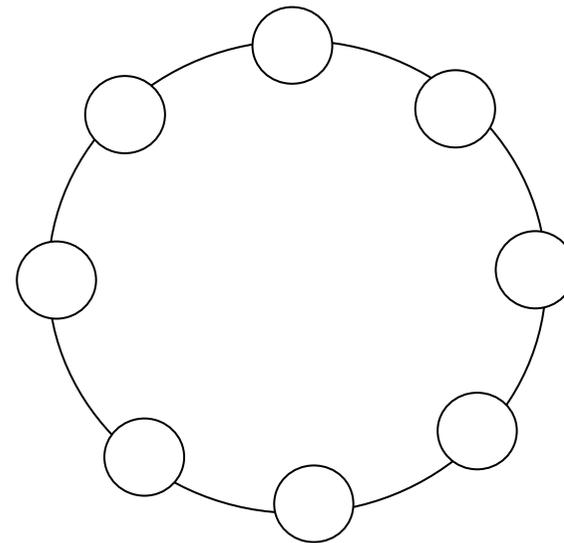
# A Toy Model

- Our toy will be in  $d=2+1$  dimensions
- It is a gauge theory with  $N=4$  supersymmetry
  - Vector multiplet:  $V = (A_\mu, \phi_i, \text{fermions}) \quad i = 1, 2, 3$
  - Hypermultiplet:  $Q = (q, \tilde{q}, \text{fermions})$

Intriligator and Seiberg, '96

- $U(1)^N + N$  hypermultiplets

- gauge coupling =  $g^2$
- mass of hypers =  $m$



# The Low-Energy Dynamics

- The hypermultiplets have physical mass

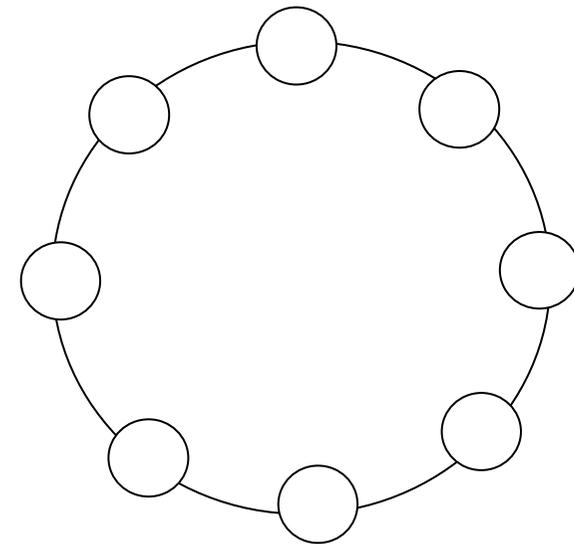
$$M^a = (\phi^a - \phi^{a+1} + m)$$

“bare” mass

vevs of vector multiplets

- At low energies, we want an effective action for the massless vector multiplets

$$\phi_i^a \quad \text{and} \quad F_{\mu\nu}^a = g^2 \epsilon_{\mu\nu\rho} \partial_\rho \sigma^a$$

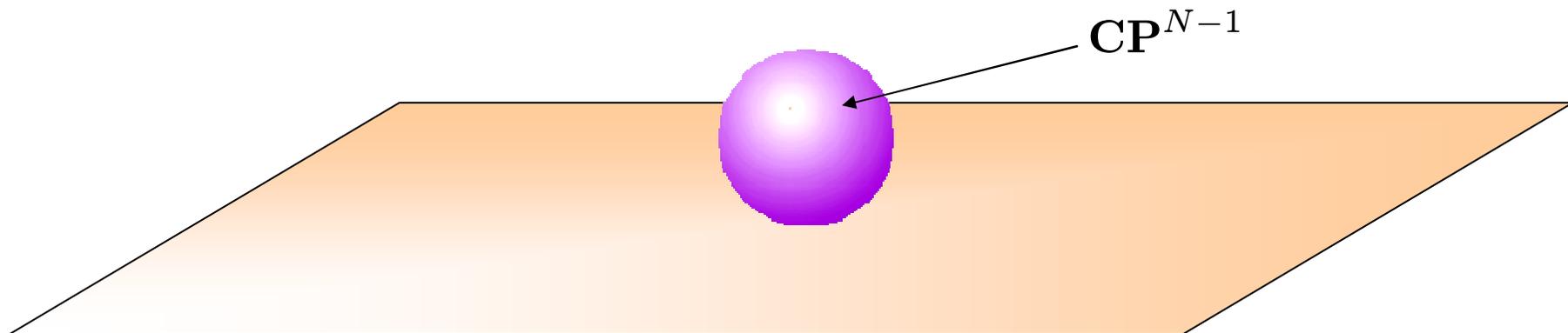


# Low-Energy Dynamics

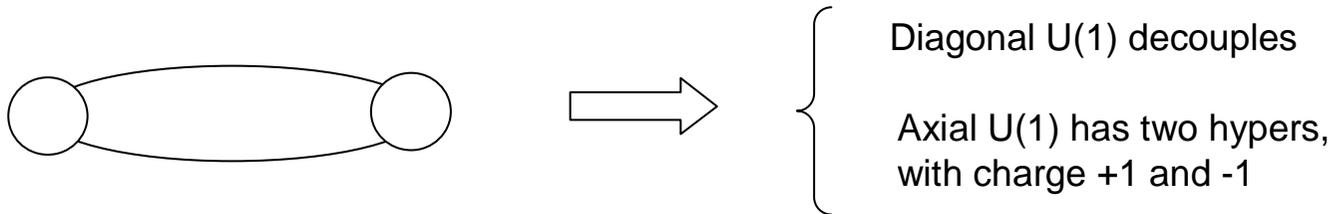
Intriligator and Seiberg, '96

- We integrate out the hypermultiplets
- This induces interactions for vector multiplets
- The low-energy dynamics is a sigma-model with target space

$$\mathbf{R}^4 \times T^* \mathbf{C}P^{N-1}$$



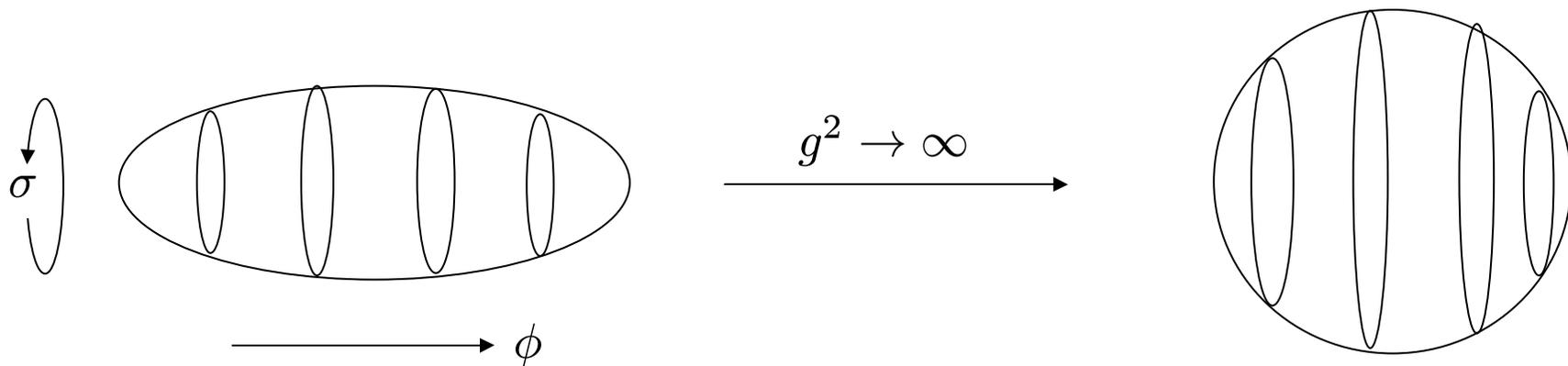
# An Example: $\mathbb{C}P^1$



Integrate out hypers:

$$\mathcal{L}_{\text{eff}} = \frac{1}{g_{\text{eff}}^2} (\partial\phi)^2 + g_{\text{eff}}^2 (\partial\sigma)^2$$

$$\frac{1}{g_{\text{eff}}^2} = \frac{1}{g^2} + \frac{1}{m-\phi} + \frac{1}{m+\phi}$$



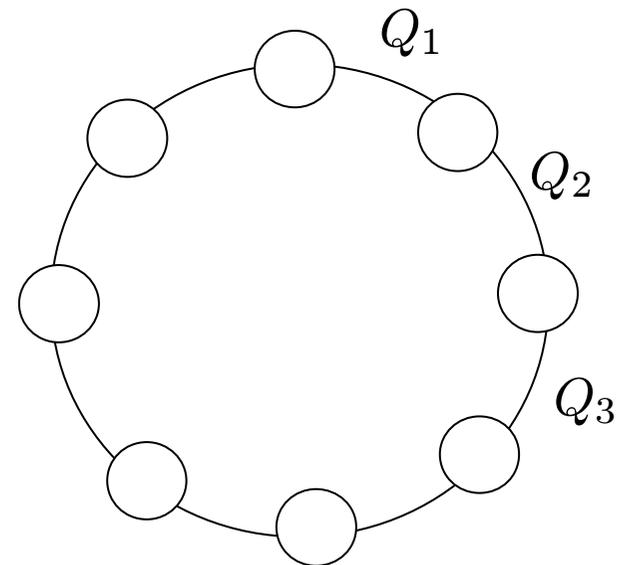
# The Soliton

- The low-energy theory has a soliton: it is a sigma-model lump
  - 2 translation modes
  - 1 scale size
  - 2N-3 orientation modes

$$\partial_\mu \phi = g_{\text{eff}}^2(\phi) \epsilon_{\mu\nu} \partial^\nu \sigma$$

- What is this soliton in the microscopic theory?
  - It is BPS
  - Mass = Nm
- It is an N-particle state.
  - The soliton is made of the objects that we thought we'd integrated out!

$$Q_1 Q_2 Q_3 \dots Q_N$$



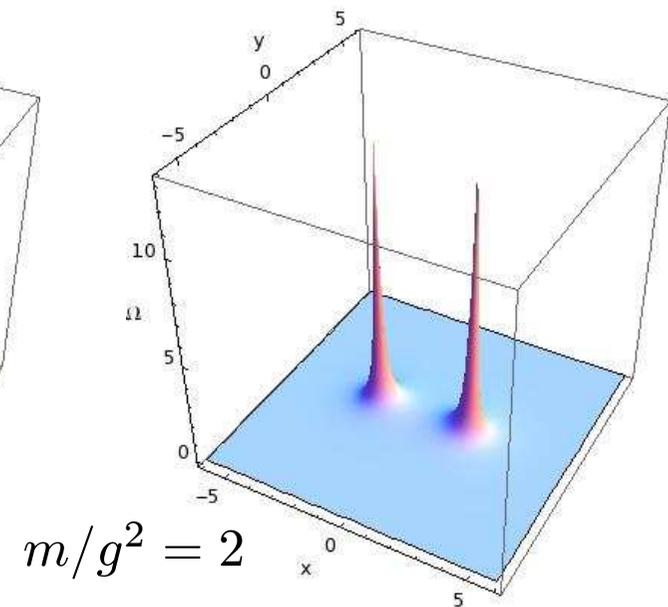
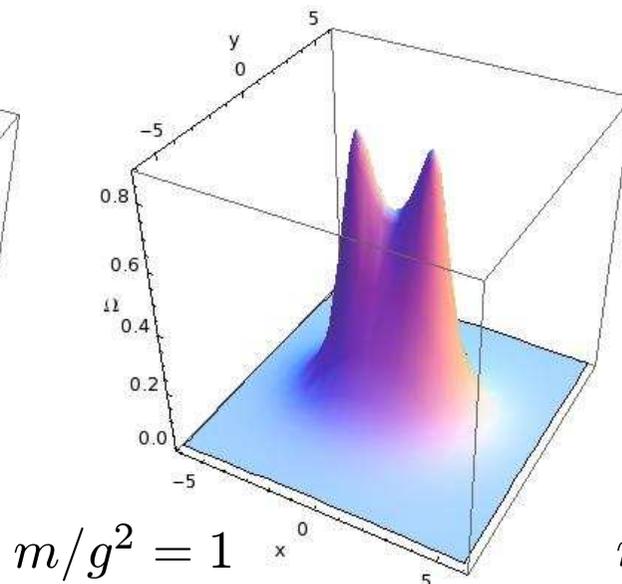
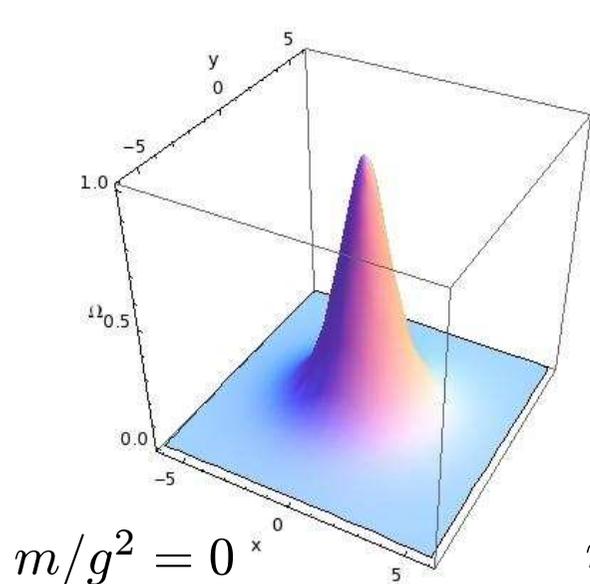
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# Reconstructing the UV Physics

- Question: Suppose we just have access to the IR physics on the Coulomb branch. What does the soliton tell us about the UV completion?
  - Answer: Pretty much everything!
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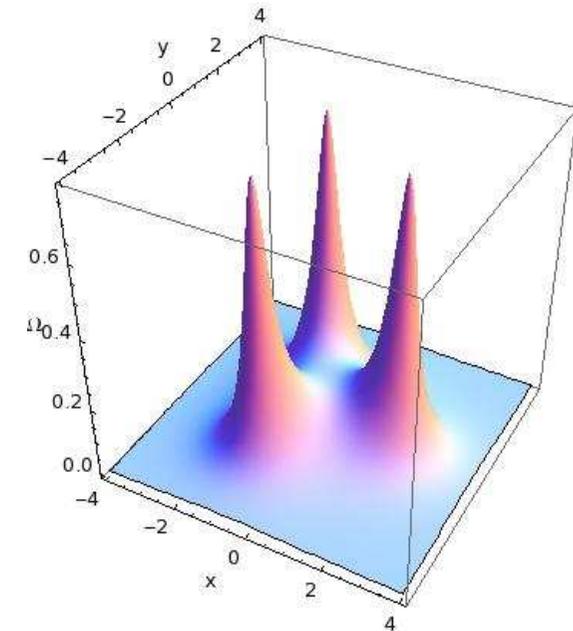
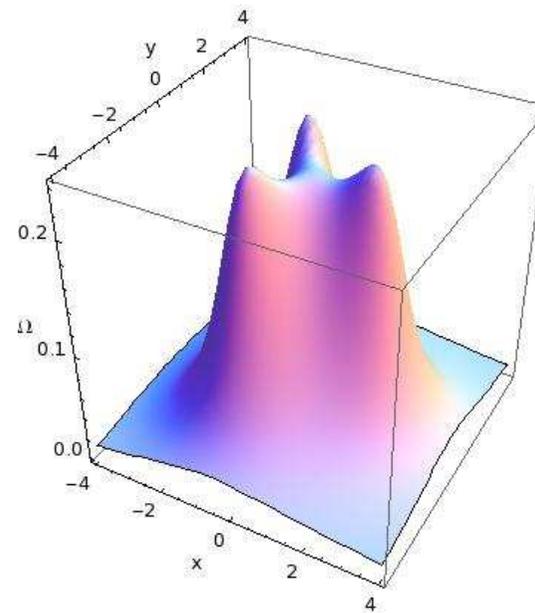
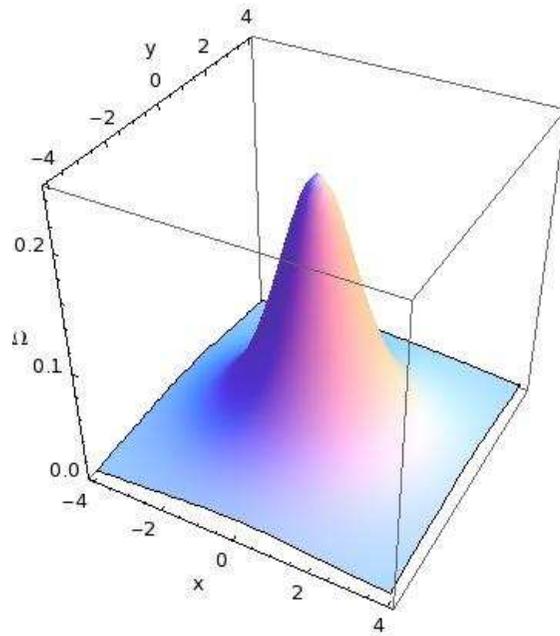
# Seeing the Partons

- We get the round metric on  $\mathbf{CP}^{N-1}$  only in the limit  $g^2 \rightarrow \infty$
- If we study solitons on the squashed target space, the solitons dramatically reveal themselves. Here's a soliton on  $\mathbf{CP}^1$



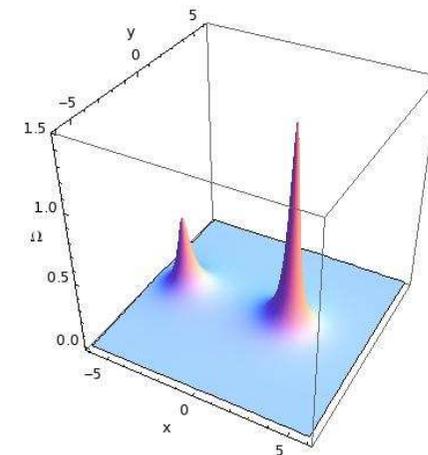
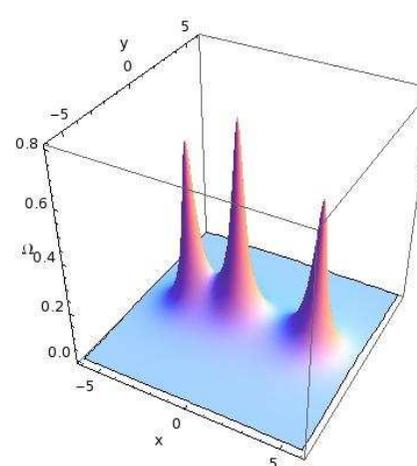
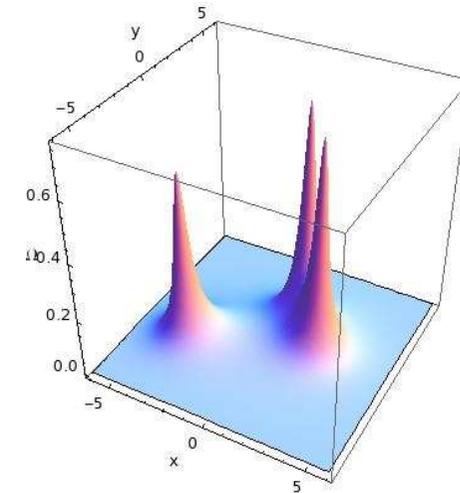
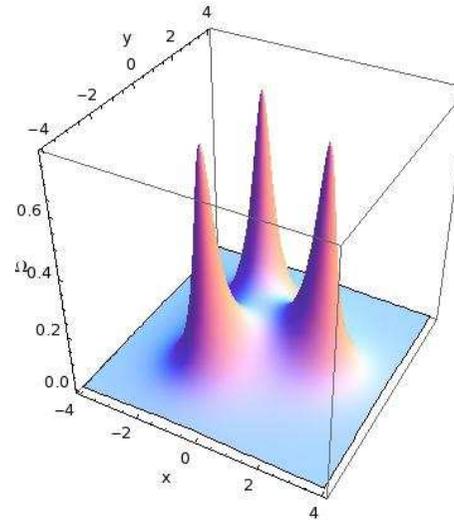
# Seeing the Partons

- Here's a single soliton on a squashed  $\mathbf{CP}^2$



# Changing the Orientation

- Keep the “scale size” of the lump fixed
- Change the “orientation modes”
- Watch the partons move



# Confinement of Partons

- Why does the low-energy theory only include the N-particle state?

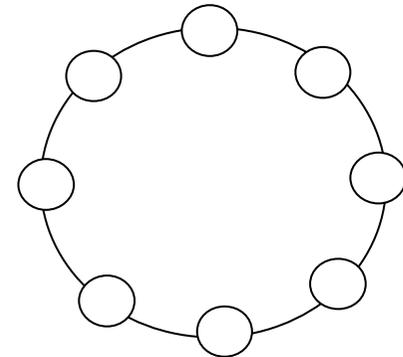
$$Q_1 Q_2 \dots Q_N$$

- Answer from microscopic theory: logarithmic confinement

- In  $d=2+1$  dimensions, electric charges have  $E \sim \frac{1}{r}$
- This gives log divergent mass
- Only gauge singlet states have finite mass

- This log divergence re-appears when partons move.

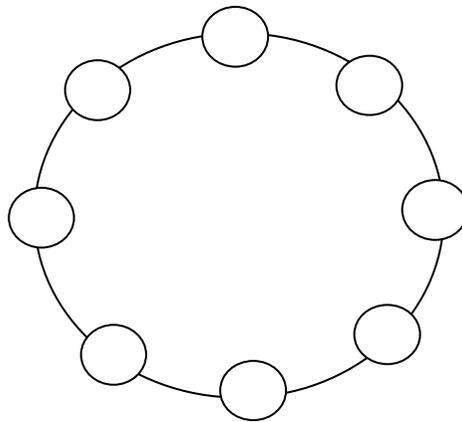
- Seen in IR theory as a log divergence in moduli space metric. Only modes which don't change dipole moments have finite norm



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# Parton Quantum Numbers

- Can we reconstruct the quantum numbers of the partons?
  - In other words, can we reconstruct the full UV theory?



# Dual Bogomolnyi Equations

- The Bogomolnyi equations for  $\mathbf{CP}^1$  are

$$\partial_\mu \phi = g_{\text{eff}}^2(\phi) \epsilon_{0\mu\nu} \partial^\nu \sigma$$

- They have a moduli space of solutions: dimension  $2kN$  for  $k$  soliton sector
- Can rewrite this equation in dual variables  $F_{\mu\nu} = g_{\text{eff}}^2(\phi) \epsilon_{\mu\nu\rho} \partial_\rho \sigma$

$$F_{0\mu} = \partial_\nu \phi$$

# Dual Bogomolnyi Equations

$$F_{0\mu} = \partial_\nu \phi$$

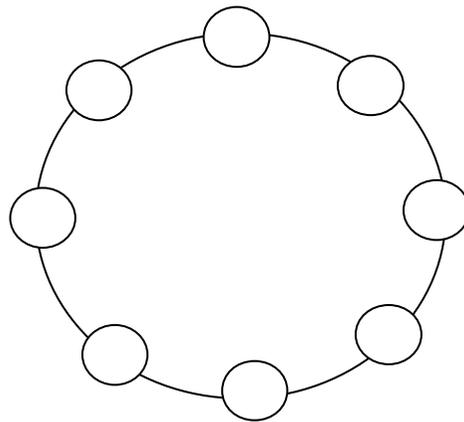
- The dual Bogomolnyi equations have *no* smooth solutions
- But we can reproduce the exact soliton solution if we introduce electric sources

$$\partial_\mu \left( \frac{1}{g_{\text{eff}}^2} F_{0\mu} \right) = \sum_{n=1}^k \delta(z - z_n^+) - \delta(z - z_n^-)$$

In soliton description, these  
are collective coordinates.  
Here, they are sources.

# Dual Bogomolnyi Equation

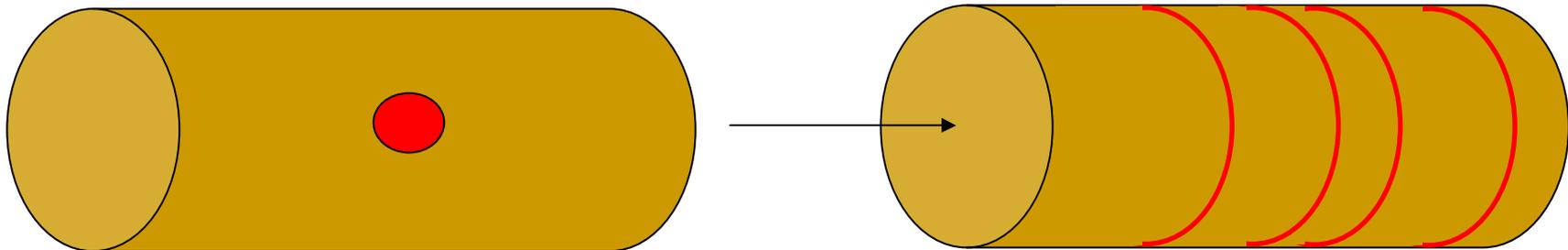
- The dual formulation of the Bogomolnyi equation provides an explicit map between a soliton and fundamental fields
- It also works for  $\mathbf{CP}^{N-1}$
- Can reconstruct quantum numbers of partons which determines the UV microscopic theory



# Calorons: A Red Herring?

- There is one way that instantons are known to split into N partons
  - Calorons
- Put  $d=4+1$  theory on circle. Add a Wilson line.
  - Instantons  $\rightarrow$  N monopole strings

Lee and Yi, Van Baal



- Doesn't seem possible that this can happen in non-compact space
- Also happens for lumps in the toy model
  - But the calorons have nothing to do with the true partons

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# Summary: Questions, not Answers

- Toy model in  $d=2+1$ 
    - Explicit demonstration that solitons can be thought of as multi-particle states
    - A study of the soliton allows us to reconstruct the UV behaviour
  - Real Interest:  $d=4+1$  Yang-Mills
    - Does the instanton solution hold clues about the constituents of the  $(2,0)$  theory?
    - What is the confinement mechanism?
      - No hint of log confinement...partons are probably not merons
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Happy Birthday  
Misha!!

