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By

E.G. Kalnins
and
W. Miller, Jr.

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HYPERGEOMETRIC EXPANSIONS OF HEUN POLYNOMIALS*

E.G. KALNINS† AND W. MILLER, JR.‡

Abstract. The product of two Heun polynomials is expanded in terms of products of two Jacobi polynomials. This is done by making crucial use of group theory and the knowledge of separable coordinate systems on the n-sphere. The expansion presented includes as a special case hypergeometric function expansions of Heun polynomials that have been derived previously by other methods.

Key words. multivariable orthogonal polynomials, the n-sphere, Heun functions

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1. Introduction. Any Fuchsian equation of second order with four singularities can be reduced to the form

\[
\frac{d^2w}{dx^2} + \left[\frac{\gamma}{x-e_1} + \frac{\delta}{x-e_2} + \frac{\epsilon}{x-e_3}\right] \frac{dw}{dx} + \frac{\alpha \beta x - q}{(x-e_1)(x-e_2)(x-e_3)} w = 0
\]

where \(\alpha + \beta - \gamma - \delta - \epsilon + 1 = 0\).

The singularities are located at \(x = e_1, e_2, e_3\) and \(\infty\) and have indices depending upon \(\alpha, \ldots, \epsilon\). The constant \(q\) is known as the accessory parameter. This is Heun's equation [1] and solutions may be characterised by the \(P\) symbol [2].

\[
P \left\{ \begin{array}{cccc}
e_1 & e_2 & e_3 & \infty \\
0 & 0 & 0 & \alpha \\
1-\gamma & 1-\delta & 1-\epsilon & \beta \end{array} \right\}
\]

Power series expansions for the solutions of Heun's equation have been studied by Heun for various arguments [1], [3]. There turn out to be 96 distinct types of power series. Alternatively, solutions of Heun's equations can be expanded in series of hypergeometric functions. Such expansions were studied by Svartholm [4] and Erdelyi [5]. Typically such expansions have the form

\[
P \left\{ \begin{array}{cccc}
e_1 & e_2 & e_3 & \infty \\
0 & 0 & 0 & \alpha \\
1-\gamma & 1-\delta & 1-\epsilon & \beta \end{array} \right\} = \sum_{m<0} A_m P \left\{ \begin{array}{cccc}
0 & 0 & \lambda + m & x \\
1-\gamma & 1-\delta & 1-\epsilon & \beta \end{array} \right\}
\]

where \(\lambda + \mu = \gamma + \delta - 1 = \alpha + \beta - \epsilon\). Two types of expansion were given;

(i) Series of type I for which \(\lambda = \alpha, \mu = \beta - \epsilon\). These series converge outside an ellipse with foci at \(e_1, e_2\) and which passes through \(e_3\). There are three distinct expansions of this type.

(ii) Series of type II for which \(\mu = 0, \gamma - 1, \delta - 1\) or \(\gamma + \delta - 2\).

*Supported in part by the National Science Foundation under grant DMS 88-23054
†Department of Mathematics and Statistics, University of Waikato Hamilton, New Zealand
‡School of Mathematics, 127 Vincent Hall, University of Minnesota, Minneapolis, MN 55455, U.S.A.
In all these expansions the coefficients \( A_m \) satisfy three term recurrence relations

\[
(1.4) \quad b_0 A_0 + c_1 A_1 = 0 \\
a_r A_{r-1} + b_r A_r + c_{r+1} A_{r+1} = 0, \quad r = 1, 2, \ldots
\]

where \( a_r, b_r, c_r \) are known expressions in \( r \) and \( c_r \neq 0 \). If \( q \) is chosen from a number of characteristic values then expansions of this type converge. In this article we derive some of these expansions for the case of Heun polynomials from considerations based on group theory and its connection with separation of variables solutions of the Laplace-Beltrami eigenvalue equation on the \( n \)-sphere. The method used makes a judicious choice of coordinates on the \( n \)-sphere. The expansions that are first derived are for products of Heun polynomials as sums of products of Jacobi polynomials. The coefficients in the expansions obey three term recurrence relations. The corresponding single variable expansions are then obtained by allowing one of the variables to take a fixed value. This paper is an extension of [8] in which the motivation and background can be found.

2. Derivation of the expansion formula. The graphical calculus of separable coordinates for the Laplace-Beltrami eigenvalue equation on the \( n \)-sphere has been completely worked out by Kalnins and Miller [6], [7]. To derive an expansion for Heun polynomials we consider coordinate systems corresponding to graphs of the type

\[
\begin{array}{ccc}
\varepsilon_1 & \varepsilon_2 & \varepsilon_3 \\
\vdots & \vdots & \vdots \\
S_{n_1} & S_{n_2} & S_{n_3}
\end{array}
\]

on the \( n \) sphere, \( n = n_1 + n_2 + n_3 + 2 \). A suitable choice of coordinates is

\[
(1.5) \quad s_i = u_1 w_i, \quad i = 1, \ldots, n_1 + 1 \\
s_{j+n_1+1} = u_2 t_j, \quad j = 1, \ldots, n_2 + 1 \\
s_{k+n_1+n_2+2} = u_3 z_k, \quad k = 1, \ldots, n_3 + 1
\]

where

\[
\sum_{i=1}^{n_1+1} w_i^2 = 1, \quad \sum_{j=1}^{n_2+1} t_j^2 = 1, \quad \sum_{k=1}^{n_3+1} z_k^2 = 1
\]

and

\[
(1.6) \quad u_i^2 = \frac{(x - e_i)(y - e_i)}{(e_j - e_i)(e_k - e_i)}, \quad i = 1, 2, 3, \quad i, j, k \text{ pairwise distinct.}
\]
The metric on the $n$ sphere is

\begin{align}
(1.7) \quad ds^2 &= -\frac{(x-y)}{4} \left[ \frac{dx^2}{(x-e_1)(x-e_2)(x-e_3)} - \frac{dy^2}{(y-e_1)(y-e_2)(y-e_3)} \right] \\
&\quad + \frac{(x-e_1)(y-e_1)}{(e_2-e_1)(e_3-e_1)} \sum_{i=1}^{n_1+1} dw_i^2 + \frac{(x-e_2)(y-e_2)}{(e_3-e_2)(e_1-e_2)} \sum_{j=1}^{n_2+1} dt_j^2 \\
&\quad + \frac{(x-e_3)(y-e_3)}{(e_2-e_3)(e_1-e_3)} \sum_{k=1}^{n_3+1} dz_k^2.
\end{align}

The coordinate systems chosen for $w_i, t_j, z_k$ can be taken to be, say, spherical coordinates in each case, corresponding to the graph [6].

\[\begin{array}{cccc}
0 & 1 & \cdots & 1 \\
0 & 1 & \cdots & 1 \\
\cdots & \cdots & \cdots & \cdots \\
0 & 1 & \cdots & 1
\end{array}\]

$n_i$ boxes

We then seek eigenfunctions $\psi$ of the Laplacian satisfying

\begin{align}
(1.8) \quad \Delta \psi &= -J(J + n_1 + n_2 + n_3 + 1)\psi,
\end{align}

where $J$ is a non-negative integer. In the coordinates we have chosen, this equation has the form

\begin{align}
(1.9) \quad \Delta \psi &= -\frac{4}{(x-y)}[(x-e_1)(x-e_2)(x-e_3)] \left[ \frac{\partial^2}{\partial x^2} + \frac{1}{2} \left[ \frac{n_1+1}{x-e_1} + \frac{n_2+1}{x-e_2} + \frac{n_3+1}{x-e_3} \right] \frac{\partial}{\partial x} \right] \psi \\
&\quad - (y-e_1)(y-e_2)(y-e_3) \left[ \frac{\partial^2}{\partial y^2} + \frac{1}{2} \left[ \frac{n_1+1}{y-e_1} + \frac{n_2+1}{y-e_2} + \frac{n_3+1}{y-e_3} \right] \frac{\partial}{\partial y} \right] \psi \\
&\quad + \left[ \frac{(e_1-e_2)(e_1-e_3)}{(x-e_1)(y-e_1)} \Delta_1 + \frac{(e_2-e_1)(e_2-e_3)}{(x-e_2)(y-e_2)} \Delta_2 + \frac{(e_3-e_1)(e_3-e_2)}{(x-e_3)(y-e_3)} \Delta_3 \right] \psi \\
&\quad = -J(J + n_1 + n_2 + n_3 + 1)\psi,
\end{align}

where $\Delta_k$ is the Laplacian on the sphere $S_{n_k}$.

If we seek eigenfunctions such that

\begin{align}
(1.10) \quad \Delta_i \psi &= -\ell_i(\ell_i + n_i - 1)\psi, \quad i = 1, 2, 3,
\end{align}

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where the $\ell_i$ are non-negative integers, then writing

$$ (1.11) \quad \psi = \prod_{i=1}^{3} [(x-e_i)(y-e_i)]^{\ell_i/2} \phi, $$

we find (1.9) has the form

$$ (1.12) \quad \begin{aligned} &- \frac{4}{(x-y)} \left\{ (x-e_1)(x-e_2)(x-e_3) \left( \frac{\partial^2}{\partial x^2} + \left[ \frac{\ell_1 + \frac{1}{2}(n_1+1)}{x-e_1} + \frac{\ell_2 + \frac{1}{2}(n_2+1)}{x-e_2} \right) \phi + A x \phi \\
&+ \frac{\ell_3 + \frac{1}{2}(n_3+1)}{x-e_3} \right\} \phi = -J(J+N+1)\phi \\
&- (y-e_1)(y-e_2)(y-e_3) \left( \frac{\partial^2}{\partial y^2} + \left[ \frac{\ell_1 + \frac{1}{2}(n_1+1)}{y-e_1} + \frac{\ell_2 + \frac{1}{2}(n_2+1)}{y-e_2} \right) \phi + A y \phi \right\} \phi = -J(J+N+1)\phi \\
&+ \frac{\ell_3 + \frac{1}{2}(n_3+1)}{y-e_3} \right\} \phi \end{aligned} $$

where

$$ A = \frac{1}{4} (L + N + 1) L $$

and

$$ L = \ell_1 + \ell_2 + \ell_3, \quad N = n_1 + n_2 + n_3. $$

The corresponding separable solutions have the form

$$ (1.13) \quad \psi = u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} \varphi_{1}^{\ell_1, \ell_2, \ell_3}(x) \varphi_{2}^{\ell_1, \ell_2, \ell_3}(y) \varphi_{3}^{\ell_1, \ell_2, \ell_3}(w, t, z) $$

where a complete set of functions $\varphi_{\ell_1, \ell_2, \ell_3}(w, t, z)$ can be taken as

$$ (1.14) \quad \varphi_{\ell_1, \ell_2, \ell_3}(w, t, z) = \varphi_{\ell_1}(w) \varphi_{\ell_2}(t) \varphi_{\ell_3}(z) $$

and typically,

$$ (1.15) \quad \varphi_{\ell_1}(w) = \prod_{j=0}^{n_1-2} C_{K_j - K_{j+1}}^{1/2(n_1-j-1)} \frac{1}{j!} \frac{1}{(\pi)^{1/2}} \frac{(\sin \theta_{n_1-j})^{K_j+1}}{K_j^{1/2(n_1-j-1)}} e^{\pm i K_{n_1-j-1} \theta}, $$
for $\ell_1 = K_0 \geq K_1 \geq \cdots \geq K_{n_1-1} \geq 0$, and

\begin{equation}
\Delta_{(k)} \Theta_{\ell_1}(w) = -K_k(K_k + n_1 - k - 1)\Theta_{\ell_1}(w)
\end{equation}

where $C^w_\mu(z)$ is a Gegenbauer polynomial. The coordinates on $S_{n_1}$ are

\begin{equation}
\begin{aligned}
  w_1 &= \sin\theta_{n_1} \cdots \sin\theta_2 \sin\theta_1 \\
  w_2 &= \sin\theta_{n_1} \cdots \sin\theta_2 \cos\theta_1 \\
  &\vdots \\
  w_{n_1} &= \sin\theta_{n_1} \cos\theta_{n_1-1} \\
  w_{n_1+1} &= \cos\theta_{n_1}
\end{aligned}
\end{equation}

and the operator $\Delta_{(k)}$ is given by

\begin{equation}
\Delta_{(k)} = \sum_{r < \ell \leq n_1+1-k} I_{r\ell}^2, \quad I_{r\ell} = w_r \frac{\partial}{\partial w_\ell} - w_\ell \frac{\partial}{\partial w_r}, \quad k = 0, \ldots, n_1 - 1.
\end{equation}

(The $\Delta_{(k)}$ are the second order symmetry operators for $\Delta_1$ whose eigenvalue equations (1.16) characterize the separable coordinates (1.17), see [6], [7].) The corresponding separation equations are

\begin{equation}
[-4(\lambda - e_1)(\lambda - e_2)(\lambda - e_3) \left[ \frac{d^2}{d\lambda^2} + \left( \frac{\ell_1 + \frac{1}{2}(n_1 + 1)}{\lambda - e_1} - \frac{\ell_2 + \frac{1}{2}(n_2 + 1)}{\lambda - e_2} + \frac{\ell_3 + \frac{1}{2}(n_3 + 1)}{\lambda - e_3} \right) \frac{d}{d\lambda} + (J - L)(J + L + N + 1)\lambda + 4q \right] \Phi^\epsilon_{J\ell_1\ell_2\ell_3q}(\lambda) = 0
\end{equation}

where $\lambda = x, y$ according as $\epsilon = 1, 2$, respectively. This is Heun's equation of the form (1.1) with $\gamma = \ell_1 + \frac{1}{2}(n_1 + 1)$, $\delta = \ell_2 + \frac{1}{2}(n_2 + 1)$, $\epsilon = \ell_3 + \frac{1}{2}(n_3 + 1)$, $\alpha = \frac{1}{2}(L - J)$, $\beta = \frac{1}{2}(L + J + N + 1)$. The solutions for the functions $\Phi^\epsilon_{J\ell_1\ell_2\ell_3q}(\lambda)$ are Heun polynomials which for fixed $J$ will form a complete set of basis functions once the eigenvalues $q$ have been calculated. To calculate the eigenvalues it is convenient to observe that in the coordinate system (1.5) the operator $\mathcal{M}$ whose eigenvalue $\chi$ is

\begin{equation}
\chi = (e_1 + e_2 + e_3)[\ell_1^2 + \ell_2^2 + \ell_3^2 + \ell_1 n_1 + \ell_2 n_2 + \ell_3 n_3 - J(J + N + 1)]
+ 2\ell_1 \ell_2 e_3 + 2\ell_1 \ell_3 e_2 + 2\ell_2 \ell_3 e_1 - \ell_1 e_1 - \ell_2 e_2 - \ell_3 e_3
+ \ell_1 n_2 e_3 + \ell_1 n_3 e_2 + \ell_2 n_1 e_3 + \ell_2 n_3 e_1 + \ell_3 n_1 e_2 + \ell_3 n_2 e_1 - 4q
\end{equation}
is given by [6], [7]

\[(1.21)\]
\[\mathcal{M} = (e_1 + e_2) \sum_{p \in P} \sum_{q \in Q} I_{pq}^2 + (e_2 + e_3) \sum_{q \in Q} \sum_{r \in R} I_{rq}^2 + (e_1 + e_3) \sum_{p \in P} \sum_{r \in R} I_{pr}^2\]

\[P = \{1, \ldots, n_1 + 1\}, Q = \{n_1 + 2, \ldots, n_1 + n_2 + 2\},\]
\[R = \{n_1 + n_2 + 3, \ldots, n_1 + n_2 + n_3 + 3\}\]

That is, \(\mathcal{M}\) is the second order symmetry operator for the Laplacian ([\(\mathcal{M}, \Delta\] = 0) which corresponds to the separable coordinates \(x, y\). Expression (1.20) gives the relationship between the eigenvalue \(\chi\) and \(q\). (The terms involving the \(\ell_j\) result from consideration of the factor \(u_{1}^{\ell_1} u_{2}^{\ell_2} u_{3}^{\ell_3}\).)

The basis functions on the sphere \(S_n\) corresponding to coordinates of the graph can also be expanded in terms of the basis functions of the coordinate system corresponding to the graph [6],

\[
\begin{array}{c}
0 \ 1 \\
\vdots \\
\vdots \\
\vdots \\
\vdots \\
S_{n_3} \\
0 \ 1 \\
\vdots \\
\vdots \\
S_{n_1} \ \ \ S_{n_2}
\end{array}
\]

i.e., the coordinates (1.5) with

\[(1.22)\]
\[u_1 = \sin \theta \cos \phi, \quad u_2 = \sin \theta \sin \phi, \quad u_3 = \cos \theta\]

and the infinitesimal distance

\[(1.23)\]
\[ds^2 = d\theta^2 + \sin^2 \theta d\phi^2 + \sin^2 \theta \cos^2 \phi \sum_{i=1}^{n_1+1} dw_i^2 + \sin^2 \theta \sin^2 \phi \sum_{j=1}^{n_2+1} dt_j^2 + \cos^2 \theta \sum_{k=1}^{n_3+1} dz_k^2.\]

Eigenfunction solutions of (1.8) in these coordinates are

\[(1.24)\]
\[\psi = (\sin \theta)^{M} (\cos \theta)^{\ell_3} (\sin \phi)^{\ell_2} (\cos \phi)^{\ell_1} \times P_{(J-M-\ell_3)/2}^{M+\frac{1}{2}(n_1+n_2), \ell_3+\frac{1}{2}(n_3-1)}(\cos 2\theta) \times P_{(M-\ell_1-\ell_2)/2}^{\ell_2+\frac{1}{2}(n_2-1), \ell_1+\frac{1}{2}(n_1-1)}(\cos 2\phi) \Theta_{\ell_1 \ell_2 \ell_3}(w, t, z) = \psi_{JM} \Theta_{\ell_1 \ell_2 \ell_3}\]
where \( P^{\alpha,\beta}(z) \) are Jacobi polynomials. Here \( J = L + 2j \) and \( M = L + 2m \) where \( j = 0,1,\ldots, m = 0,1,\ldots j-1,j \). The eigenfunctions satisfy

\[
\Delta' \psi = -M(M+n_1+n_2) \psi,
\]

where

\[
\Delta' = \sum_{i>j} I_{ij}^2.
\]

and \( i,j \) range from 1 to \( n_1+n_2+2 \).

Note that in terms of the Cartesian coordinates \( u_1, u_2, u_3 \) on the 2-sphere \( (u_1^2+u_2^2+u_3^2 = 1) \) these eigenfunctions take the form

\[
\psi_{JM} = u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} (u_1^2 + u_2^2)^{(M-\ell_1-\ell_2)/2} \times P^M_{(J-M-\ell_3)/2} (1 - 2u_1^2 - 2u_2^2) \times P_{(M-\ell_1-\ell_2)/2} \left( \frac{2u_1^2}{u_1^2 + u_2^2} - 1 \right) = u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} \Phi_{jm},
\]

i.e., the form \( u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} \Phi(u_1^2, u_2^2) \) where \( \Phi \) is a polynomial.

This remark leads to another way of viewing the Heun and Jacobi bases. In the equation \( \Delta \psi = -J(J+N+1) \psi \) with \( \Delta \psi \) given by (1.9) and \( \Delta_k \) replaced by the values \(-\ell_k(\ell_k+n_k-1), k = 1,2,3 \) we set \( \psi = u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} \Phi(x_1, x_2) \) and introduce the new coordinates \( x_1 = u_1^2, x_2 = u_2^2 \). The eigenvalue equation for \( \Phi \) reads

\[
H \Phi = -j(j+G-1) \Phi,
\]

where

\[
H = \sum_{i,j=1}^2 (x_i \delta_{ij} - x_i x_j) \frac{\partial^2}{\partial x_i \partial x_j} + \sum_{i=1}^2 (\gamma_i - G x_i) \frac{\partial}{\partial x_i}.
\]

Here \( G = \gamma_1 + \gamma_2 + \gamma_3 \) and in this particular case

\[
\gamma_i = \ell_i + \frac{1}{2}(n_i + 1), \quad i = 1,2,3
\]

\[
j = \frac{1}{2}(J-L) = 0,1,2,\ldots.
\]

This coincides with equation (1.4) in [8]. In particular \( H \) maps polynomials of maximum degree \( m_i \) in \( x_i \) to polynomials of the same type. Furthermore, it is easy to see that the
polynomial eigenfunctions of $H$ form a basis for the space of all polynomials $f(x_1, x_2)$ and that the spectrum of $H$ acting on this space is exactly \{-$j(j + G - 1)$ : $j = 0, 1, \ldots$\}. It is also shown in [8] that $H = \Delta_2 + \Lambda_2$ where $\Delta_2$ is the Laplace Beltrami operator on $S_2$ and

\begin{equation}
\Lambda_2 = \sum_{i=1}^{2} \left[ \gamma_i - \frac{1}{2} + \left( \frac{3}{2} - G \right) x_i \right] \frac{\partial}{\partial x_i}.
\end{equation}

Moreover, $H$ is self-adjoint with respect to the inner product

\begin{equation}
(f_1, f_2) = \iint_{x_1, x_2 > 0, 1 - x_1 - x_2 > 0} f_1(x) \overline{f_2(x)} d\omega
\end{equation}

where

\begin{equation}
d\omega = x_1^{\gamma_1 - 1} x_2^{\gamma_2 - 1} (1 - x_1 - x_2)^{\gamma_3 - 1} dx_1 dx_2 : (H f_1, f_2) = (f_1, H f_2).
\end{equation}

Here $f_1, f_2$ are polynomials in $x = (x_1, x_2)$. For fixed $j$ the polynomials

\begin{equation}
\Phi_{jm}(x_1, x_2) = (x_1 + x_2)^m P_{j - m}^{\gamma_1 + \gamma_2 + 2m - 1, \gamma_3 - 1} (2x_1 + 2x_2 - 1) \\
\times P_m^{\gamma_2 - 1, \gamma_3 - 1} \left( \frac{2x_1}{x_1 + x_2} - 1 \right), \quad m = 0, 1, \ldots, j
\end{equation}

form an orthogonal basis for the eigenspace corresponding to eigenvalue $-j(j + G - 1)$. (This is the orthogonal basis of Proriol [9] and of Karlin and McGregor [10]). Similarly the Heun polynomials $\Phi_{j\ell_1 \ell_2 \ell_3 q}(x) \Phi_{j\ell_1 \ell_2 \ell_3 q}(y)$ where $q$ runs over the possible eigenvalues, form an alternate orthogonal basis for this same space. Moreover as pointed out in [11] these bases correspond to spherical and ellipsoidal coordinates on the 2-sphere and are the only coordinates in which $\Delta_2$ separates.

With this point of view we are operating on $S_2$ rather than $S_n$ and our two distinguished orthogonal bases are the only ones possible rather than two out of a multiplicity of separable systems on $S_n$ for large $n$. The principal advantage of this new point of view is that the eigenfunctions are obviously polynomials in $x_1, x_2$ and that the only requirement on the constants $\gamma_1, \gamma_2, \gamma_3$ to ensure orthogonality is that they be strictly positive. Thus the $\ell_i$ and $n_i$ need not be integers; it is only required that $2\ell_i + n_i + 1 > 0$.

In the following our expansion formulas are valid for all real $\gamma_i > 0$. In the special case $\gamma_1 = \gamma_2 = \gamma_3 = \frac{1}{2}$ we have $H = \Delta_2$, the Laplace-Beltrami operator on $S_2$. In this case the eigenvalue equation $\Delta_2 \Phi = -j(j + \frac{1}{2}) \Phi$ admits the Lie algebra $so(3)$ as a symmetry algebra. A basis for $so(3)$ is \{u_1 \partial_{u_2} - u_2 \partial_{u_1}, u_3 \partial_{u_1}, u_3 \partial_{u_2}\} where $u_3 = \pm (1 - u_1^2 - u_2^2)^{\frac{1}{2}}$. This extra symmetry is associated with the fact that there are additional polynomial
solutions of the eigenvalue equation (see §3 of reference [8]). In particular the equation admits polynomial solutions of the form \( f(u_1, u_2) \) and the spectrum of \( \Delta_2 \) acting on the space of all such polynomials is \(-j(j + \frac{1}{2})\) where now \(2j = 0, 1, 2, \ldots\). Furthermore there exist solutions of the form \( u_3 g(u_1, u_2) \) with \( g \) a polynomial and with the same eigenvalues. The dimension of each eigenspace is \(2j + 1\) rather than \(j + 1\) for the general case. In this special case the eigenfunctions corresponding to spherical coordinates are just the spherical harmonics whereas those corresponding to ellipsoidal coordinates are products of Lamé polynomials. For the solution of the problem of expanding the Lamé basis in terms of a spherical harmonic basis see [11], [12], [13].

Returning to the case of general \( \ell_i, n_i \) we consider the problem of expanding the Heun basis (1.13) in terms of the Jacobi polynomial basis (1.24), (1.27), (1.34):

\[
\psi = u_1^{\ell_1} u_2^{\ell_2} u_3^{\ell_3} \Phi_1^{\ell_1 \ell_2 \ell_3} (x) \Phi_2^{\ell_1 \ell_2 \ell_3} (y) = \sum_{m=0}^{j} \xi_m \psi_{J\ell_1 \ell_2 \ell_3 M}(\theta, \phi).
\]

Three term recurrence relations for the expansion coefficients \( \xi_m \) (where \( M = \ell_1 + \ell_2 + 2m \)) can be deduced by requiring that

\[
\mathcal{M} \psi = \chi \psi \tag{1.36}
\]

Using the recurrence formulas for Jacobi polynomials this relation can be deduced. Indeed to do this we need the action of the various pieces of \( \mathcal{M} \) on the Jacobi bases \( \psi_{J\ell_1 \ell_2 \ell_3 M}(\theta, \phi) \).

We have

\[
\mathcal{M} \psi_{J\ell_1 \ell_2 \ell_3 M}(\theta, \phi) = \sum_{r=-1}^{+1} X_r \psi_{J\ell_1 \ell_2 \ell_3 M+2r}(\theta, \phi) \tag{1.37}
\]

where

\[
X_1(m, j) = \frac{4(e_1 - e_2)(\gamma_1 + \gamma_2 + \gamma_3 + m + j - 1)(\gamma_3 - m + j - 1)(m + 1)(\gamma_1 + \gamma_2 + m - 1)}{(\gamma_1 + \gamma_2 + 2m - 1)(\gamma_1 + \gamma_2 + 2m)}
\]

\[
X_{-1}(m, j) = \frac{4(e_1 - e_2)(\gamma_1 + \gamma_2 + m + j - 1)(-m + j + 1)(\gamma_2 - 1)(\gamma_1 - 1)}{(\gamma_1 + \gamma_2 + 2m - 1)(\gamma_1 + \gamma_2 + 2m - 2)}
\]

\[
X_0(m, j) = \frac{2(e_1 - e_2)[m^2 + m(\gamma_1 + \gamma_2 - 1) - j^2 - j(\gamma_1 + \gamma_2 + \gamma_3 - 1)](\gamma_1 + \gamma_2 - 2)(\gamma_1 - \gamma_2)}{(\gamma_1 + \gamma_2 + 2m - 2)(\gamma_1 + \gamma_2 + 2m)}
\]

\[
\quad + \frac{4(e_1 - e_2)m\gamma_3(\gamma_1 - \gamma_2)(m + \gamma_2)}{(\gamma_1 + \gamma_2 + 2m - 2)(\gamma_1 + \gamma_2 + 2m)}
\]

\[
\quad + 2(e_1 + e_2)[-m^2 - m(\gamma_1 + \gamma_2 - 1) + j^2 + j(\gamma_1 + \gamma_2 + \gamma_3 - 1)]
\]

\[
\quad + 4e_3[m^2 + m(\gamma_1 + \gamma_2 - 1)] + 4q.
\]

9
Keys to deriving this result are the following recurrence formulas for Jacobi polynomials $P_n^{\alpha,\beta}(x)$

(1.39) \[ xP_n^{\alpha,\beta} = AP_{n-1}^{\alpha,\beta} + BP_n^{\alpha,\beta} + CP_{n+1}^{\alpha,\beta} \]

\[ A = \frac{2(n + \alpha)(n + \beta)}{(2n + \alpha + \beta + 1)(2n + \alpha + \beta)}, \quad B = \frac{(\beta - \alpha)(\beta + \alpha)}{(2n + \alpha + \beta + 2)(2n + \alpha + \beta)} \]

\[ C = \frac{2(n + 1)(n + \alpha + \beta + 1)}{(2n + \alpha + \beta + 2)(2n + \alpha + \beta + 1)} \]

\[ (1 - x^2) \frac{d}{dx} P_n^{\alpha,\beta} = AP_{n-1}^{\alpha,\beta} + BP_n^{\alpha,\beta} + CP_{n+1}^{\alpha,\beta} \]

\[ A = \frac{2(n + \alpha)(n + \beta)(n + \alpha + \beta + 1)}{(2n + \alpha + \beta + 1)(2n + \alpha + \beta)}, \quad B = \frac{2n(\alpha - \beta)(n + \alpha + 1)}{(2n + \alpha + \beta + 2)(2n + \alpha + \beta)} \]

\[ C = -\frac{2n(n + 1)(n + \alpha + \beta + 1)}{(2n + \alpha + \beta + 2)(2n + \alpha + \beta + 1)} \]

Now substituting the expansion (1.35) into the eigenvalue equation $M\psi = \chi\psi$ and using (1.37) we find the three term recurrence relation

(1.40) \[ X_1(m - 1, j)\xi_{m-1} + (X_0(m, j) - \chi)\xi_m + X_{-1}(m + 1, j)\xi_{m+1} = 0 \]

where $m = 0, 1, \ldots, j$. Consequently the $j + 1$ independent eigenvalues $q$ are calculated from the determinant

(1.41) \[ \begin{vmatrix} X_0(j, j) - \chi & X_1(j - 1, j) & X_0(j - 1, j) - \chi & X_1(j - 2, j) \\ X_1(j, j) & X_0(j - 1, j) - \chi & X_1(j - 2, j) & \ddots \\ \vdots & \ddots & \ddots & \ddots \end{vmatrix} = 0 \]

To obtain the expansions in terms of one variable from (1.35) we proceed as follows. For the two choices of $u_i$, $i = 1, 2, 3$ given by (1.6) and (1.22) take $y = e_3$, $\theta = \frac{\pi}{2}$. Then the expansion has the form

(1.42) \[ \Phi^1_{J\ell_1,\ell_2,\ell_3}(x) = \sum_{M=\ell_1+\ell_2}^{J-\ell_3} \tilde{\gamma}_M P_{\frac{1}{2}(M-\ell_1-\ell_2)}^{\frac{1}{2}(n_2-1),\frac{1}{2}(n_1-1)}(\cos 2\phi) \]

where

\[ \cos 2\phi = 2 \frac{(x - e_1)}{(e_2 - e_1)} - 1. \]
This is an expansion of type 2 with \( \mu = 0 \). A different type of expansion can be obtained by taking \( \phi = \pi/2 \) and \( y = e_1 \). The resulting expression has the form

\[
(1.43) \quad \Phi^1_{J\ell_1 \ell_2 \ell_3 q}(x) = \sum_{M=\ell_1+\ell_2}^{J-\ell_3} \tilde{\gamma}_M (\sin \theta)^{M-\ell_1-\ell_2} \\
\times P^{M+\frac{i}{2}(n_1+n_2), \ell_3+\frac{i}{2}(n_3-1)}_{\frac{J-M-\ell_3}{2}}(\cos 2\theta)
\]

where

\[
\cos 2\theta = -2 \frac{(x-e_2)}{(e_2-e_3)} - 1.
\]

In both these examples the dependence of the \( \tilde{\gamma}_M \) and \( \tilde{\gamma}_M \) coefficients on the indices \( \ell_1, \ell_2, \ell_3, q \) has been suppressed.

This second type of expansion of a Heun polynomial appears to be new. Nothing that was done in the derivation of expansions (except the limits of summation on \( r \)) could not be extended to the representation of Heun functions when \( J, \ell_1, \ell_2, \ell_3 \) are complex. Consequently representations of such functions in terms of expansions whose coefficients obey three term recurrence relations can be derived. The convergence of series of this type will be discussed elsewhere.

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