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Hadrons in the Light-Front Field Theory

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# Introduction

1. Light-Front Formalism
2. Spin-0 Particles: Pion
3. Pion Pole Approximation
4. Covariance Restoration in the Light-Front
5. Conclusions

# Light-Front Formalism

- Light-Front Coordinates

$$x^+ = t + z \quad x^+ = x^0 + x^3 \quad \implies \text{Time}$$

$$x^- = t - z \quad x^- = x^0 - x^3 \quad \implies \text{Position}$$

- Four-Vector  $\implies x^\mu = (x^0, x^1, x^2, x^3) = (x^+, x^-, x_\perp)$

- Metric Tensor

$$g^{\mu\nu} = \begin{pmatrix} 0 & 2 & 0 & 0 \\ 2 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad \text{and} \quad g_{\mu\nu} = \begin{pmatrix} 0 & \frac{1}{2} & 0 & 0 \\ \frac{1}{2} & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

- Scalar product

$$\begin{aligned} x \cdot y &= x^\mu y_\mu = x^+ y_+ + x^- y_- + x^1 y_1 + x^2 y_2 \\ &= \frac{1}{2}(x^+ y^- + x^- y^+) - \vec{x}_\perp \vec{y}_\perp \end{aligned}$$

$$p^+ = p^0 + p^3$$

$$p^- = p^0 - p^3$$

$$p^\perp = (p^1, p^2)$$

- Dirac Matrix

$$\gamma^+ = \gamma^0 + \gamma^3 \implies \text{Electr. Current } J^+ = J^0 + J^3$$

$$\gamma^- = \gamma^0 - \gamma^3 \implies \text{Electr. Current } J^- = J^0 - J^3$$

$$\gamma^\perp = (\gamma^1, \gamma^2) \implies \text{Electr. Current } J^\perp = (J^1, J^2)$$

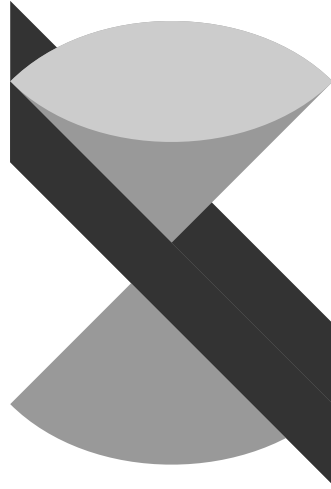


Fig. 1: Ligh-Front

- $p^\mu x_\mu = \frac{p^+ x^- + p^- x^+}{2} - \vec{p}_\perp \vec{x}_\perp$
- $x^+, x^-, x_\perp^\vec{\phantom{x}} \implies p^+, p^-, p_\perp^\vec{\phantom{p}}$
- $p^- \implies$  **Light-Front Energy**
- $p^2 = p^+ p^- - p_\perp^2 \implies p^- = \frac{p_\perp^2 + m^2}{p^+}$
- **Bosons:**  $\implies S_F(p) = \frac{1}{p^2 - m^2 + i\epsilon}$
- **Fermions:**  $\implies S_F(p) = \frac{\not{p} + m}{p^2 - m^2 + i\epsilon} + \frac{\gamma^+}{2p^+}$
- Ref: [Phys. Rept. 301, \(1998\) 299-486](#)  
S. J. Brodsky, H.C. Pauli and S.S. Pinsky

## Integration Light-Front

- **Example**

$$I_1^{cov} = \int d^4k \frac{1}{(k^2 - m^2 + i\epsilon)^3} = \frac{\pi^3}{2im^2} \neq 0$$

$$I_1^{FL} = \int d^2k_{\perp} dk^+ dk^- \frac{1}{(k^+k^- - k_{\perp}^2 - m^2 + i\epsilon)^3} = 0 !!!$$

- **Double Pole**  $\implies k^- = \frac{k_{\perp}^2 + m^2 - i\epsilon}{k^+}$

- **Pole Dislocation Method**  $\implies p'^+ = p^+ + \delta$

- **Boson Electromagnetic Current**

- **Breit Frame**  $\implies q^- = 0, q^+ \implies 0_+, \vec{q}_{\perp} \neq 0$

- $J^+ = J^- + TRC$

- $J_{\perp} \propto q^+ \Rightarrow 0$

**J.P.B.C. de Melo, J.H. Sales et T. Frederico**

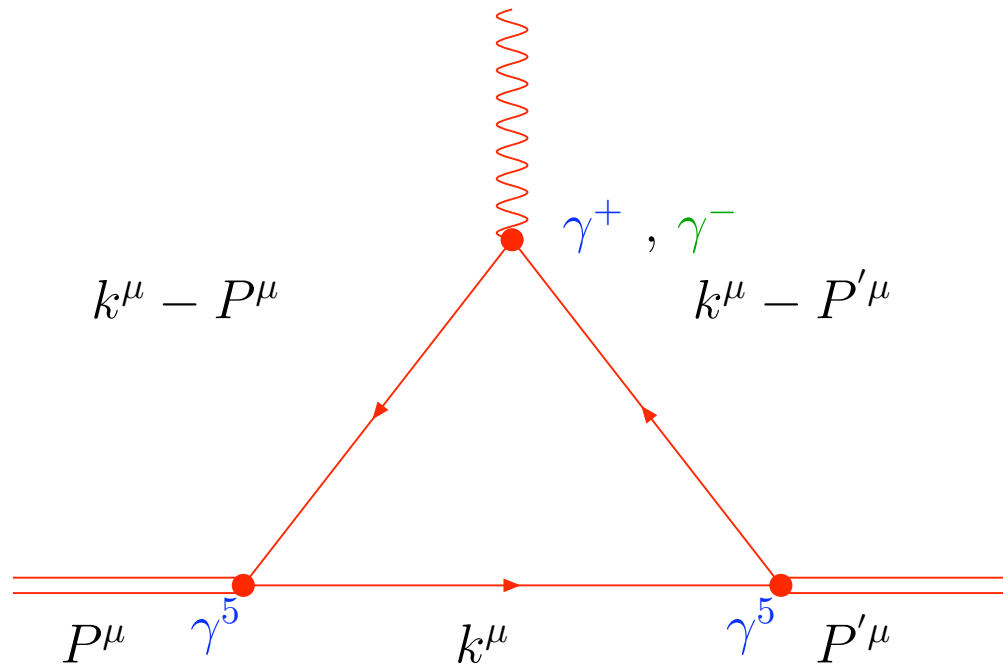
**Nucl. Phys. B631, (1998) 574c-579c.**

- **Ward-Takahashi Identity**  $\implies$  **Pair Contribution**

**H.W. Naus, J.P.B.C. de Melo et T. Frederico,**

**Few-Body Syst. 24, 1998, 99-107**

- Pion



- Frame

$$q^+ = -q^- = \sqrt{-q^2} \sin \alpha$$

$$q_x = \sqrt{-q^2} \cos \alpha, \quad q_y = 0$$

$$q^2 = q^+ q^- - (q_\perp)^2 .$$

- **Breit Frame** ( $\alpha = 0$ )  $\implies q^+ \rightarrow 0, q^- = 0; \vec{q} \neq 0$

- $J_\pi^+ = J^0 + J^3 \implies$  **No Pair Term Contribution**

- $J_\pi^- = J^0 - J^3 \implies$  **Pair Term Contribution**

- J.P.B.C. de Melo, T. Frederico and H.L. Naus,  
Phy.Rev. **C59** (1999) 2278

## Effective Lagrangian to Vertex $\pi \rightarrow q\bar{q}$

$$\mathcal{L}_I = -i \frac{m}{f_\pi} \vec{\pi} \cdot \bar{q} \gamma^5 \vec{\tau} q$$

- Electromagnetic Current

$$J^\mu = -i 2e \frac{m^2}{f_\pi^2} N_c \int \frac{d^4 k}{(2\pi)^4} \text{Tr} [S(k) \gamma^5 \times S(k - P') \gamma^\mu S(k - P) \gamma^5 \Lambda(k, P') \Lambda(k, P)]$$

$$S(p) = \frac{1}{\not{p} - m + i\epsilon}$$

- Symmetric Vertex Function

$$\Lambda(k, P) = \frac{N}{(k^2 - m_R^2 + i\epsilon)} + \frac{N}{((P - k)^2 - m_R^2 + i\epsilon)}$$

- Ref. Nucl. Phys. **A 707** (2002) 399-424

- Nonsymmetric Vertex Function

$$\Lambda(k, P) = \frac{N}{((P - k)^2 - m_R^2 + i\epsilon)}$$

- J.P.B.C. de Melo, T. Frederico and H.L. Naus, Phy.Rev. **C59** (1999) 2278

## Electromagnetic Form Factor

$$j^\mu = e(P^\mu + P'^\mu)F_\pi(q^2)$$

- ” $J^+$ ” Component of the Electromagnetic Current

- Integration Intervals in  $k^-$

I)  $0 < k^+ < P^+$

II)  $P^+ < k^+ < P'^+$

- Form Factor

$$F_\pi(q^2) = F_\pi^{(I)}(q^2, \alpha) + F_\pi^{(II)}(q^2, \alpha)$$

$$F_\pi^{(I)}(q^2, \alpha) = \frac{-im^2 N_c}{(2\pi)^4 (P^+ + P'^+) f_\pi^2} \int_0^{P^+} \frac{d^2 k_\perp dk^+ dk^- \Pi(k, P, P')}{k^+ (P^+ - k^+) (P'^+ - k^+)}$$

$$F_\pi^{(II)}(q^2, \alpha) = \frac{-im^2 N_c}{(2\pi)^4 (P^+ + P'^+) f_\pi^2} \int_{P^+}^{P'^+} \frac{d^2 k_\perp dk^+ dk^- \Pi(k, P, P')}{k^+ (P^+ - k^+) (P'^+ - k^+)}$$

$$\begin{aligned} \Pi(k, P, P', ) &= \frac{\text{Tr}[\mathcal{O}^+] \Lambda(k, P) \Lambda(k, P')}{(k^- - k_{on}^- + i\epsilon)(P^- - k^- - (P - k)_{on}^- + \frac{i\epsilon}{P^+ - k^+})} \\ &\quad \times \frac{1}{(P'^- - k^- - (P' - k)_{on}^- + i\epsilon)} \end{aligned}$$

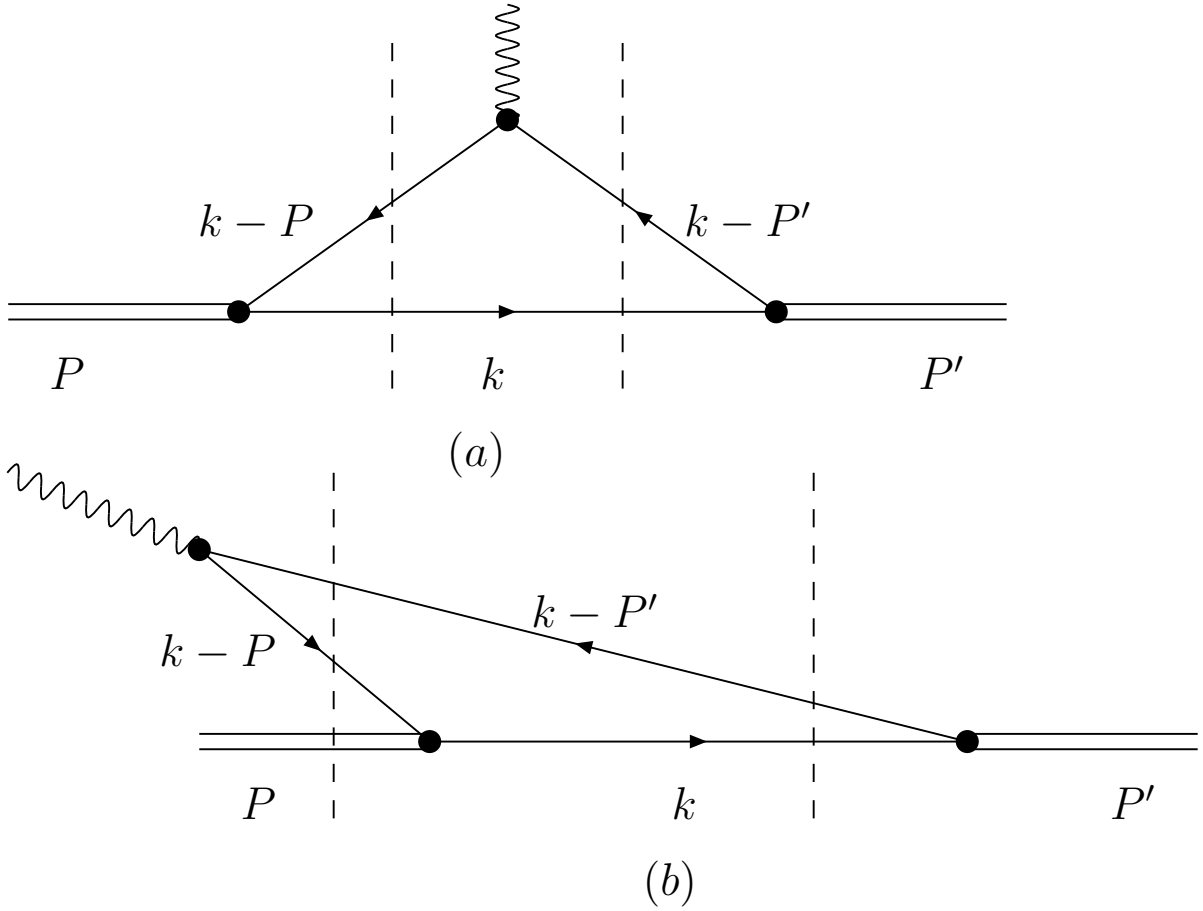


Fig. 2: Light-front time-ordered diagrams for the current: (a)  $F_\pi^{(I)}$  and (b)  $F_\pi^{(II)}$

- Dirac Propagator in the Light-Front

$$\frac{\not{k} + m}{k^2 - m^2 + i\epsilon} = \frac{\not{k}_{on} + m}{k^+(k^- - k_{on}^- + \frac{i\epsilon}{k^+})} + \frac{\gamma^+}{2k^+}$$

- Instantaneous Term Contribution
- Wave Function in the Light-Front
- Bethe-Salpeter Amplitude

$$\Psi(k, P) = \frac{m}{f_\pi} \frac{\not{k} + m}{k^2 - m^2 + i\epsilon} \gamma^5 \Lambda(k, P) \frac{\not{k} - \not{P} + m}{(k - P)^2 - m^2 + i\epsilon}$$

### The Wave Function Appear after Elimination

- i) Instantaneous Terms
- ii) Factors with Gamma Matrix in the Numerator
- iii) Factors  $k^+$  and  $(P^+ - k^+)$  in the Denominator

$$\begin{aligned} \Phi(k^+, \vec{k}_\perp; P^+, \vec{P}_\perp) &= i \mathcal{N} \int \frac{dk^-}{2\pi} \frac{1}{(k^- - k_{on}^- + \frac{i\epsilon}{k^+})} \\ &\times \frac{1}{(P^- - k^- - (P - k)_{on}^- + \frac{i\epsilon}{P^+ - k^+})} \\ &\times \left( \frac{1}{k^2 - m_R^2 + i\epsilon} + \frac{1}{(P - k)^2 - m_R^2 + i\epsilon} \right) \end{aligned}$$

$$\mathcal{N} = \sqrt{N_c} C \frac{m}{f_\pi}$$

- Integration  $k^-$

$$\Phi(k^+, \vec{k}_\perp; P^+, \vec{P}_\perp) = \frac{P^+}{m_\pi^2 - M_0^2} \left[ \frac{\mathcal{N}}{(1-x)(m_\pi^2 - \mathcal{M}^2(m^2, m_R^2))} + \frac{\mathcal{N}}{x(m_\pi^2 - \mathcal{M}^2(m_R^2, m^2))} \right]$$

- Where  $x = k^+/P^+$ , with  $0 \leq x \leq 1$

$$\mathcal{M}^2(m_a^2, m_b^2) = \frac{k_\perp^2 + m_a^2}{x} + \frac{(P - k)_\perp^2 + m_b^2}{1-x} - P_\perp^2$$

- Free Mass  $M_0^2 = \mathcal{M}^2(m^2, m^2)$
- Valence Electromagnetic Form Factor

$$\begin{aligned} F_\pi^{(WF)}(q^2, \alpha) &= \frac{1}{2\pi^3(P'^+ + P^+)} \int_0^{P^+} \frac{d^2 k_\perp dk^+ \Phi(k^+, \vec{k}_\perp; P'^+, \frac{\vec{q}_\perp}{2})}{k^+(P^+ - k^+)(P'^+ - k^+)} \\ &\times \left[ k_{on}^- P^+ P'^+ + \frac{1}{2} \vec{k}_\perp \cdot \vec{q}_\perp (P'^+ - P^+) - \frac{1}{4} k^+ q_\perp^2 \right] \\ &\times \Phi(k^+, \vec{k}_\perp; P^+, -\frac{\vec{q}_\perp}{2}) \end{aligned}$$

- $k_{on}^- = (k_\perp^2 + m^2)/k^+$
- Normalization Constant  $(C) \implies F_\pi(0) = 1$

- $F_\pi^{(WF)}$  for  $q^2 = 0 \implies$  Probability Density  $\eta$ .  
No Dependence in  $\alpha$  for the Components of the Valence  $q\bar{q}$
- If  $\eta < 1$ , in this case the Pair Terms is VIP
- Transverse Probability

$$f(k_\perp) = \frac{1}{4\pi^3 m_\pi} \int_0^{2\pi} d\phi \int_0^{m_\pi} \frac{dk^+ M_0^2}{k^+(m_\pi - k^+)} \Phi^2(k^+, \vec{k}_\perp; m_\pi, \vec{0})$$

- Integration of  $f(k_\perp)$  under  $\vec{k}_\perp$

$$\eta = \int_0^\infty dk_\perp k_\perp f(k_\perp)$$

- $\eta = 0.77$  Symmetric Vertex
- $\eta = 1.0$  NonSymmetric Vertex

## Pion Decay Constant

$$P_\mu \langle 0 | A_i^\mu | \pi_j \rangle = i m_\pi^2 f_\pi \delta_{ij}$$

$$A_i^\mu = \bar{q} \gamma^\mu \gamma^5 \frac{\tau_i}{2} q$$

- We Have

$$i P^2 f_\pi = \frac{m}{f_\pi} N_c \int \frac{d^4 k}{(2\pi)^4} \text{Tr}[P \gamma^5 S(k) \gamma^5 S(k-P)] \Lambda(k, P)$$

- $k^-$  Integration  $\implies f_\pi$  by the use of the Valence Component

$$f_\pi = \frac{m \sqrt{N_c}}{4\pi^3} \int \frac{d^2 k_\perp dk^+}{k^+ (m_\pi - k^+)} \Phi(k^+, \vec{k}_\perp; m_\pi, \vec{0})$$

## Model Parameters and Numerical Results

### NonSymmetric Vertex

- Quark Mass  $m_q$ :  $m = 0.220$  GeV
- Regulator Mass  $m_R$ :  $m = 0.946$  GeV
- Fit for  $f_\pi = 101.0$  MeV
- Pion Mass:  $0.140$  GeV
- Charge Radius

$$\langle r^2 \rangle = 6 \frac{\partial}{\partial q^2} F_\pi = 0.67 fm$$

### Symmetric Vertex

- Quark Mass  $m_q$ :  $m = 0.220$  GeV
- Regulator Mass  $m_R$ :  $m = 0.6$  GeV
- Fit for  $f_\pi^{exp} = 92.4$  MeV
- Pion Mass:  $0.140$  GeV
- Charge Radius

$$\langle r^2 \rangle = 6 \frac{\partial}{\partial q^2} F_\pi = 0.74 fm$$

10% Bigger than Experimental Value

$$(r_{exp} = 0.67 \pm 0.02 fm)$$

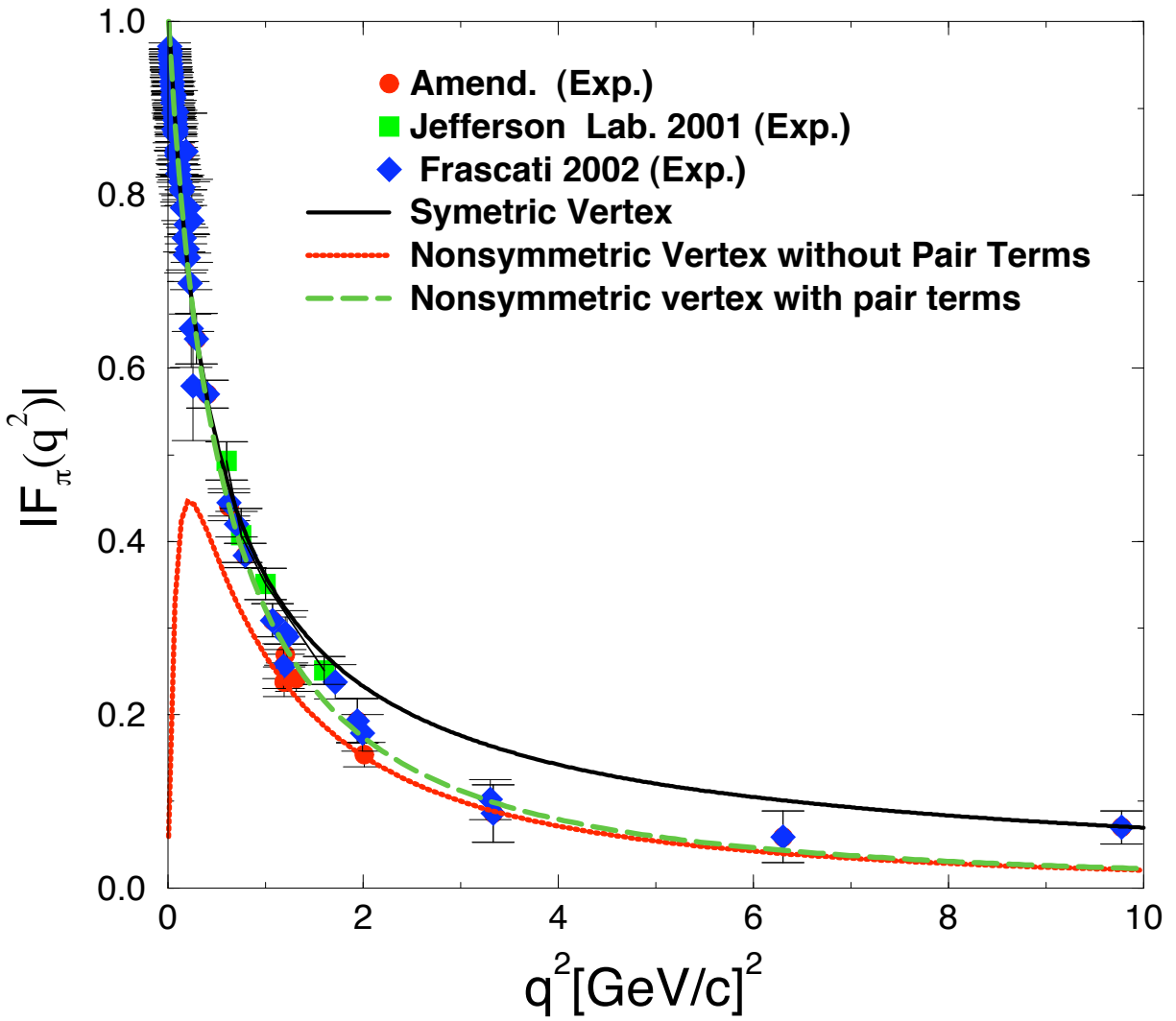


Fig. 3: Pion Form Factor

- $J_{\pi}^{-}$  Without Pair Terms
- $J_{\pi}^{-}$  With Pair Terms

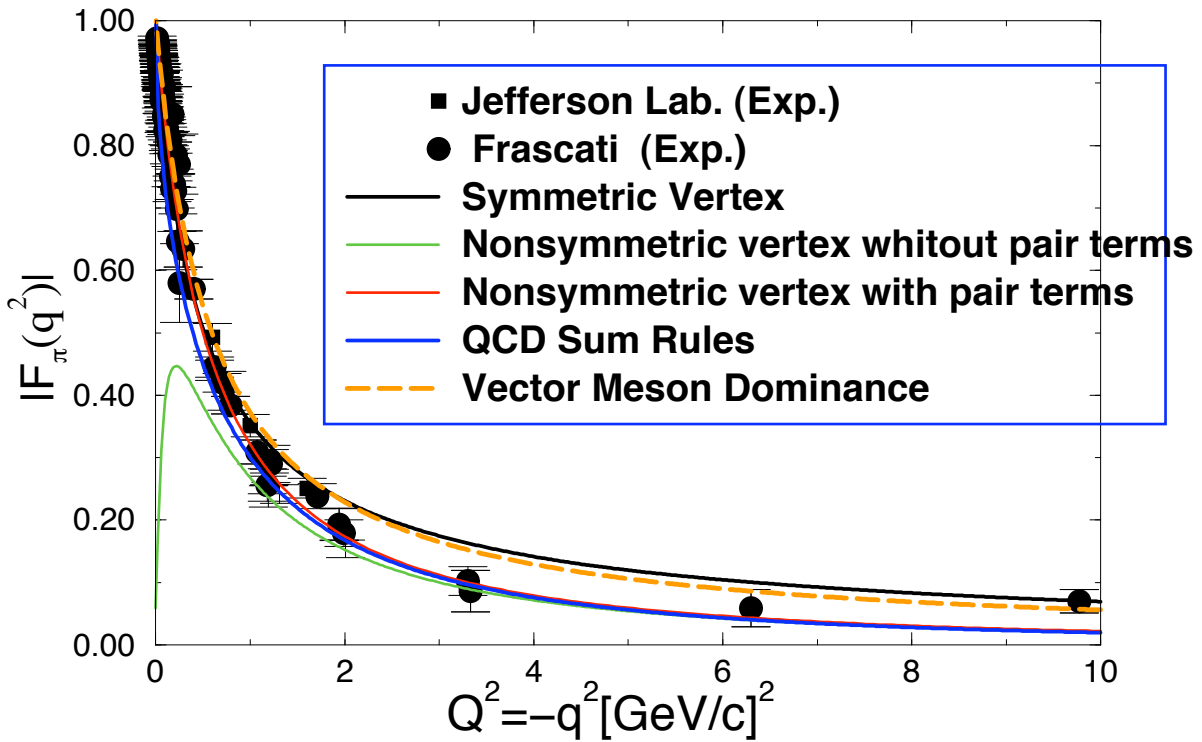


Fig. 4: Pion Form Factor

Ref. J. P. B. C. de Melo , [AIP Conf.Proc.739:553-556,2005](#)  
 and [hep-ph/0507265](#)

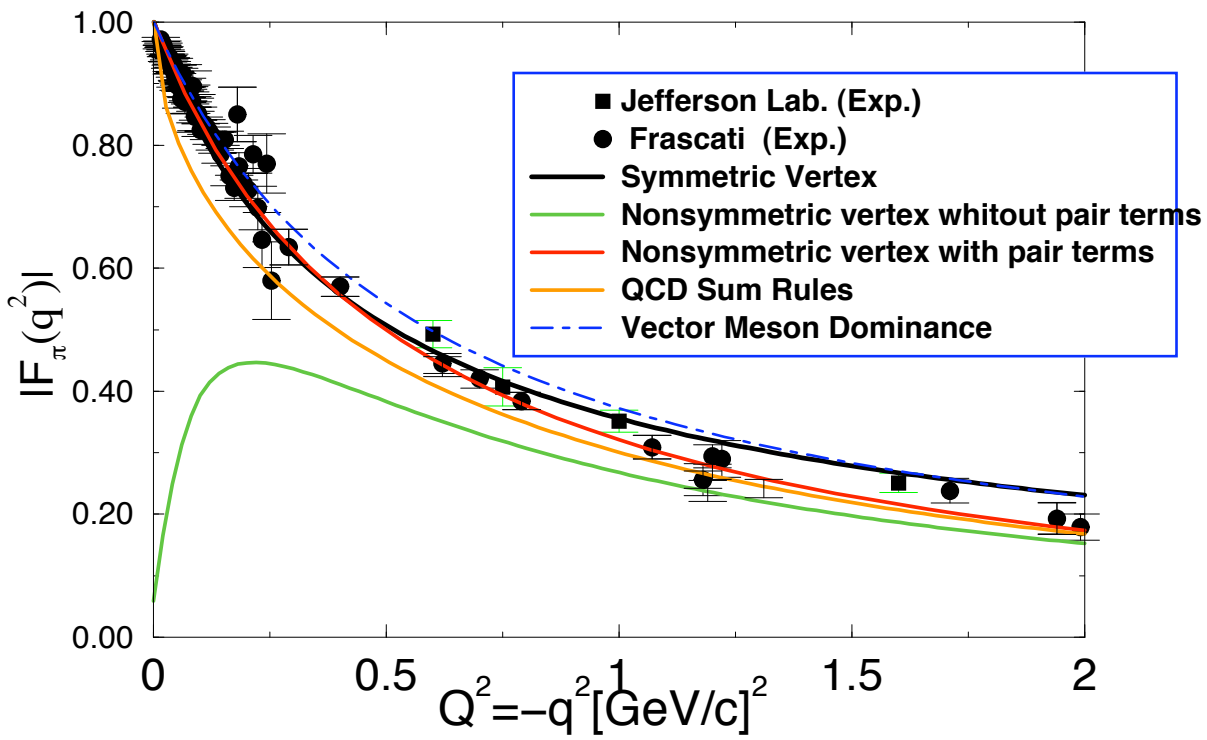


Fig. 5: Pion Form Factor at low energy

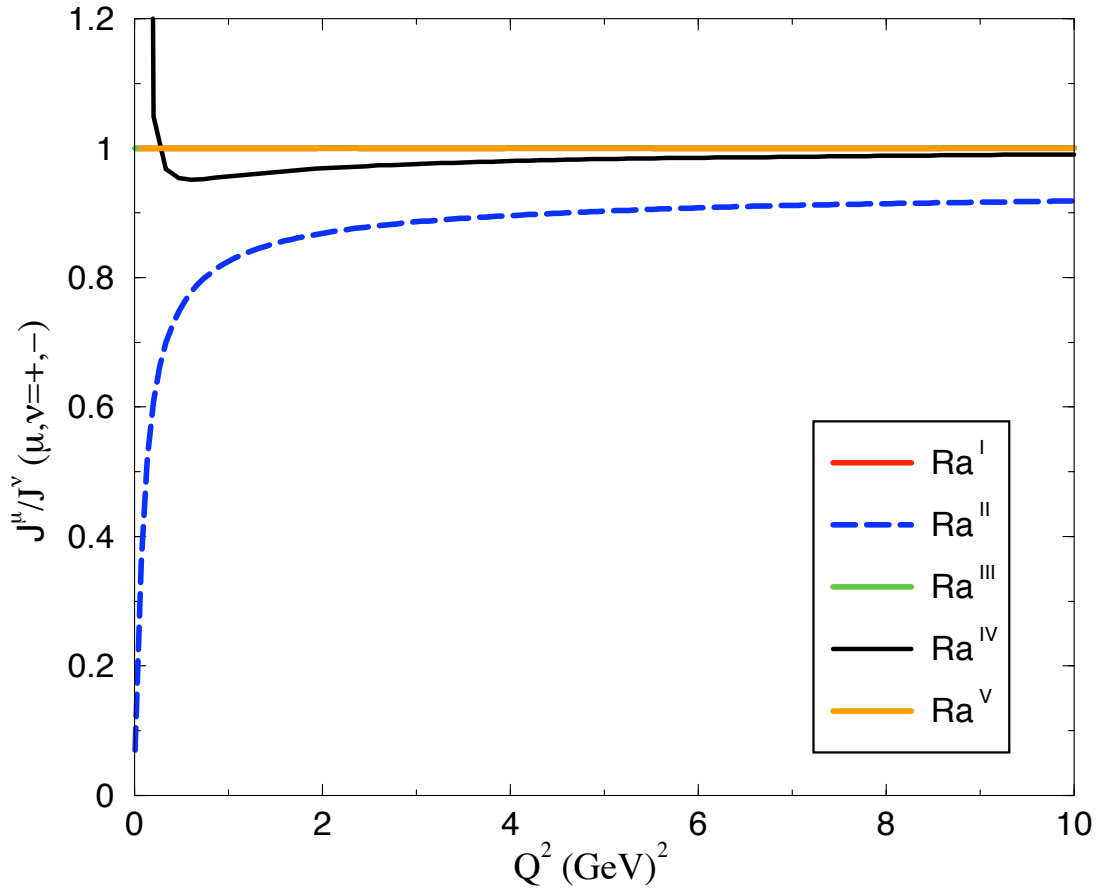


Fig. 6: Nonsymmetric Vertex.

$$Ra^I = \frac{J_{LF}^+}{J_{Cov}^+}$$

$$Ra^{II} = \frac{J_{LF}^-}{J_{Cov}^-}$$

$$Ra^{III} = \frac{J_{LF}^- + J_{LF}^{-(Pair)}}{J_{Cov}^-}$$

$$Ra^{IV} = \frac{J_{LF}^-}{J_{Cov}^+}$$

$$Ra^V = \frac{J_{LF}^- + J_{LF}^{-(Pair)}}{J_{Cov}^+}$$

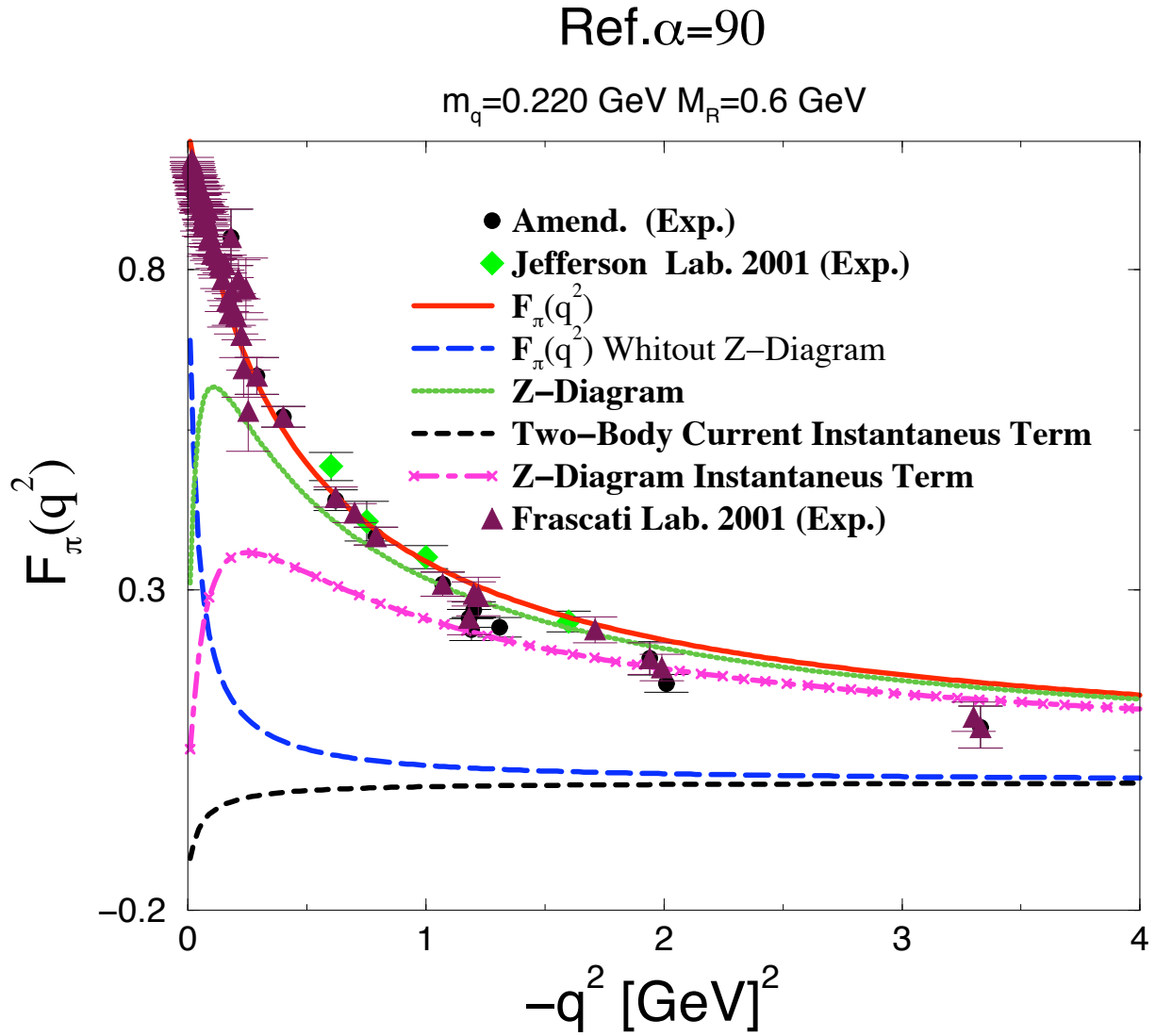


Fig. 7: Pion Form Factor

● Ref. de Melo, Frederico, Salmé and Pace  
 Nucl. Phys. A707 (2002) 399-424

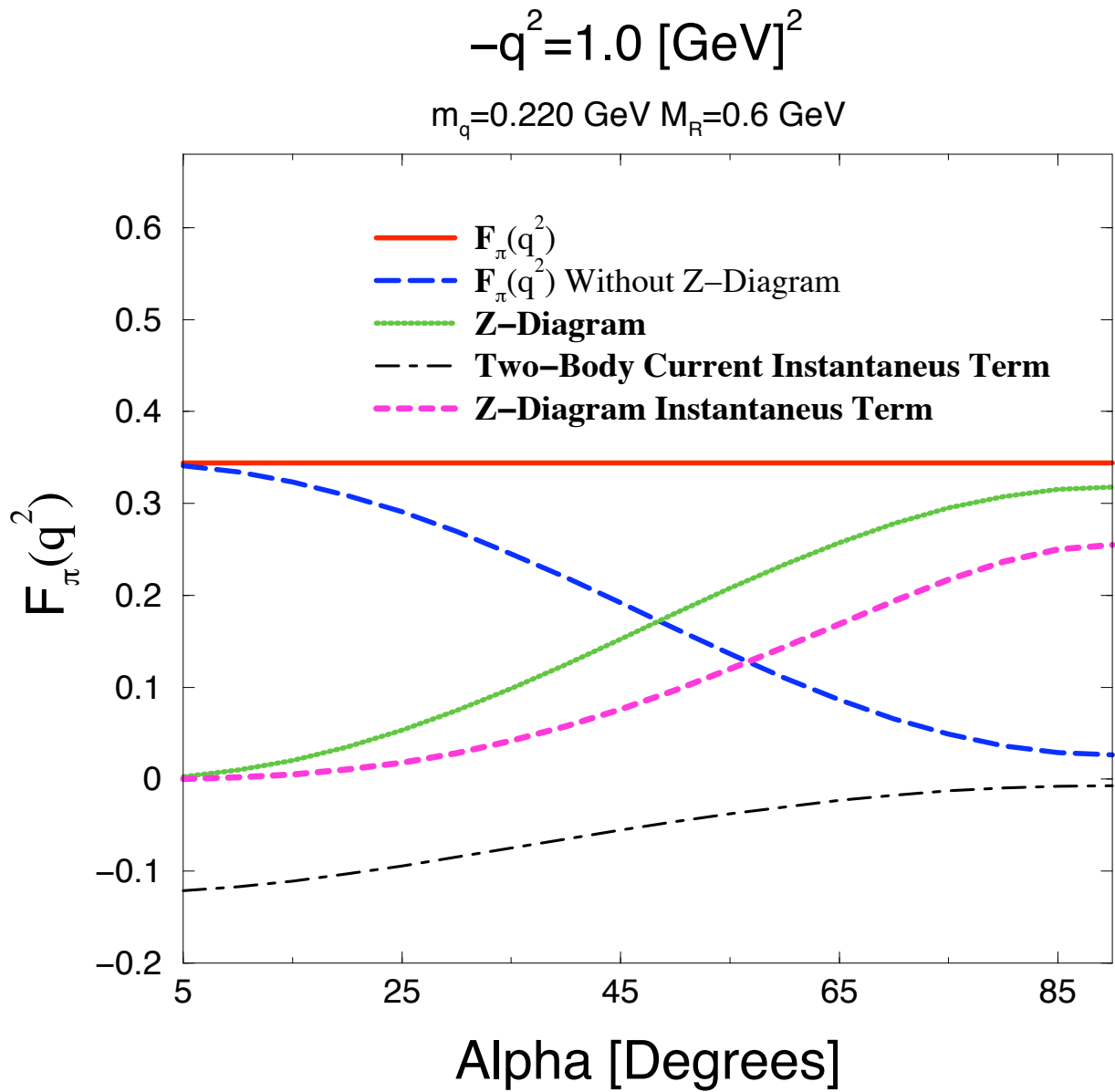


Fig. 8: Pion Form Factor vs  $\alpha$

- Ref. de Melo, Frederico, Salmé and Pace  
 Nucl. Phys. A707 (2002) 399-424

Ref.  $\alpha=45$

$m_q=0.220$  GeV  $M_R=0.6$  GeV

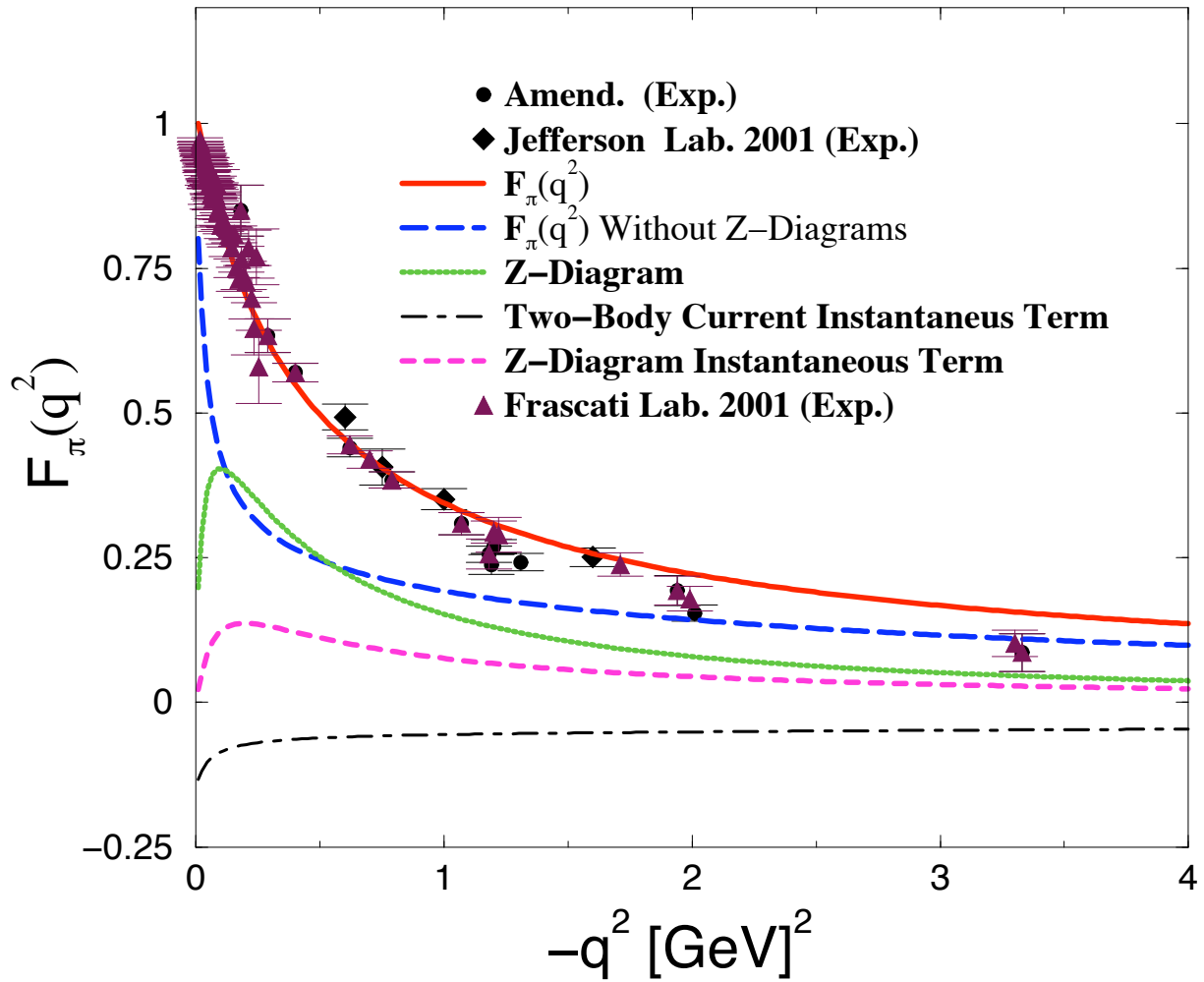


Fig. 9: Pion Form Factor with  $\alpha = 45$

• Ref. de Melo, Frederico, Salmé and Pace  
Nucl. Phys. A707 (2002) 399-424

## Pole Approximation

- Pole On-Shell Contribution

- $k_1^- = \frac{k_\perp^2 + m^2}{k^+}$

- Pole Z-Diagramm Contribution

- $k_3^- = \left( p'^- - \frac{(p-k)_\perp^2 + m^2}{(p'^+ - k^+)} \right)$

I)  $0 < k^+ < P^+$

II)  $P^+ < k^+ < P'^+$

- Form Factor

$$F_\pi(q^2) = F_\pi^{(1)}(q^2, \alpha) + F_\pi^{(2)}(q^2, \alpha)$$

$$F_\pi^{(1)}(q^2, \alpha) = \frac{-im^2 N_c}{(2\pi)^4 (P^+ + P'^+) f_\pi^2} \int_0^{P^+} \frac{d^2 k_\perp dk^+ dk^- \Pi_1(k, P, P')}{k^+ (P^+ - k^+) (P'^+ - k^+)}$$

$$F_\pi^{(2)}(q^2, \alpha) = \frac{-im^2 N_c}{(2\pi)^4 (P^+ + P'^+) f_\pi^2} \int_{P^+}^{P'^+} \frac{d^2 k_\perp dk^+ dk^- \Pi_2(k, P, P')}{k^+ (P^+ - k^+) (P'^+ - k^+)}$$

$$\begin{aligned} \Pi_1(k, P, P', ) &= \frac{\text{Tr}[\mathcal{O}^+] \Lambda(k, P) \Lambda(k, P')}{(P^- - k^- - (P - k)_{on}^- + \frac{i\epsilon}{P^+ - k^+})} \\ &\quad \times \frac{1}{(P'^- - k^- - (P' - k)_{on}^- + i\epsilon)} \end{aligned}$$

$$\Pi_2(k, P, P', ) = \frac{\text{Tr}[\mathcal{O}^+] \Lambda(k, P) \Lambda(k, P')}{(k^- - k_{on}^- + i\epsilon) (P^- - k^- - (P - k)_{on}^- + \frac{i\epsilon}{P^+ - k^+})}$$

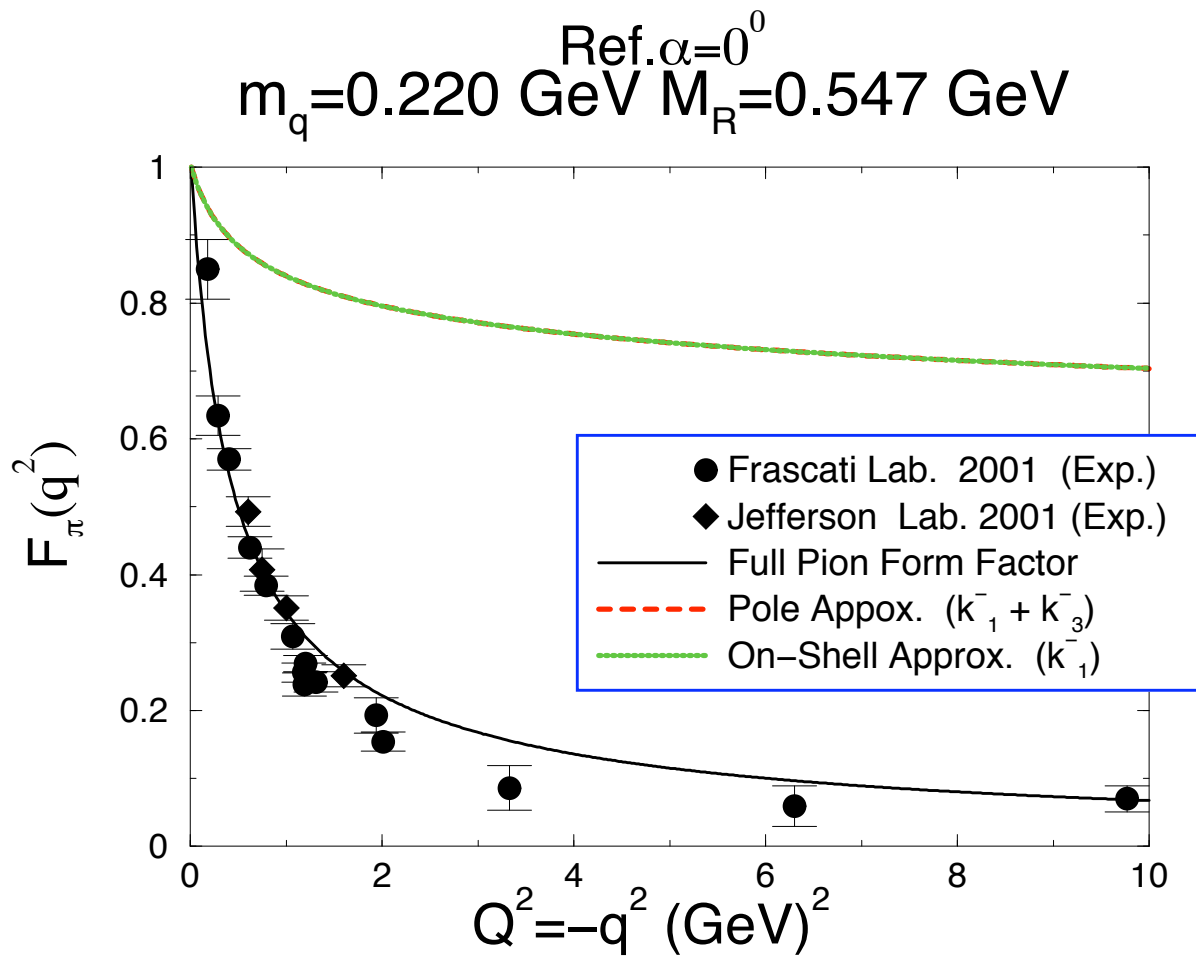


Fig. 10: Pion Pole Approximation

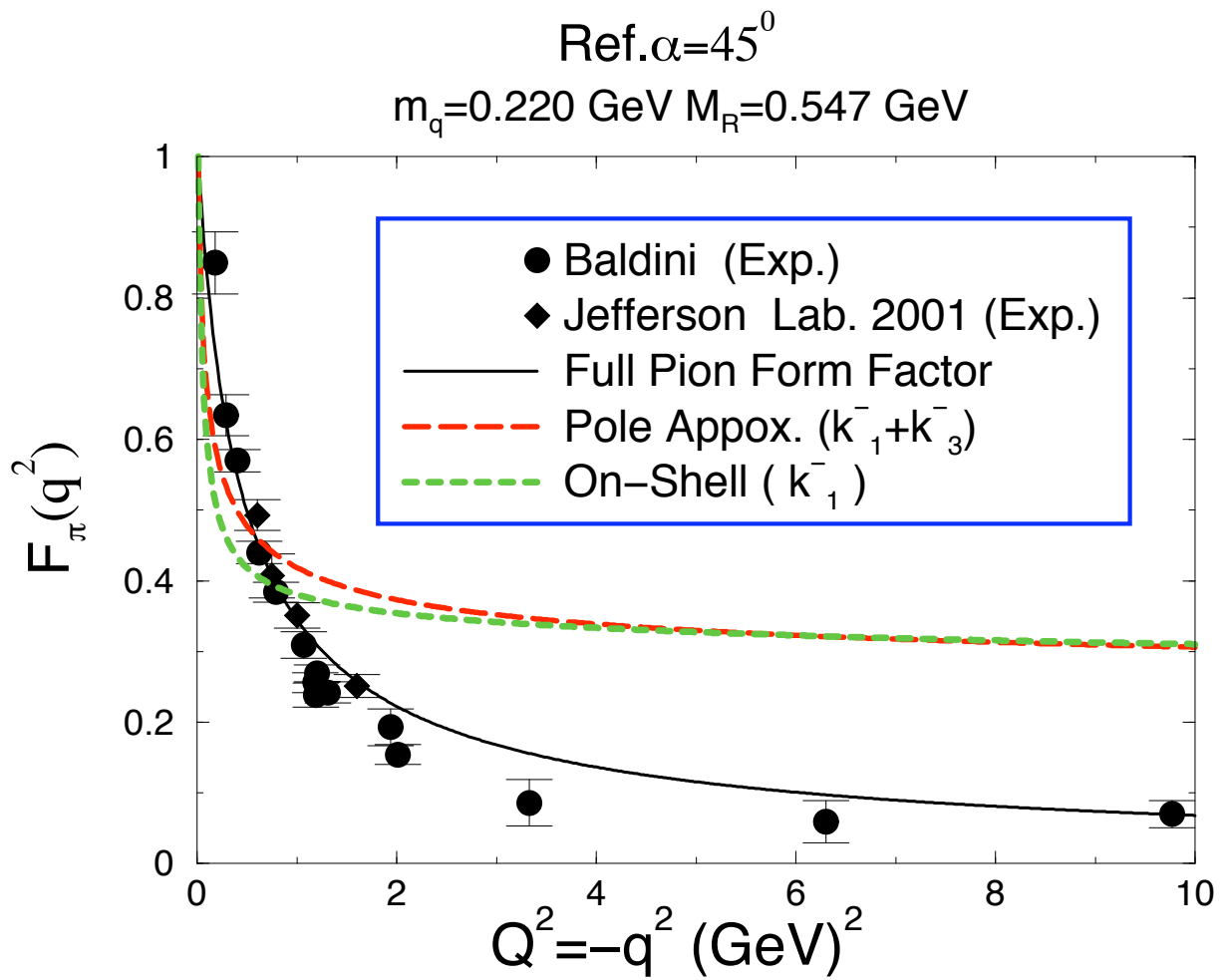


Fig. 11: Pion Pole Approximation

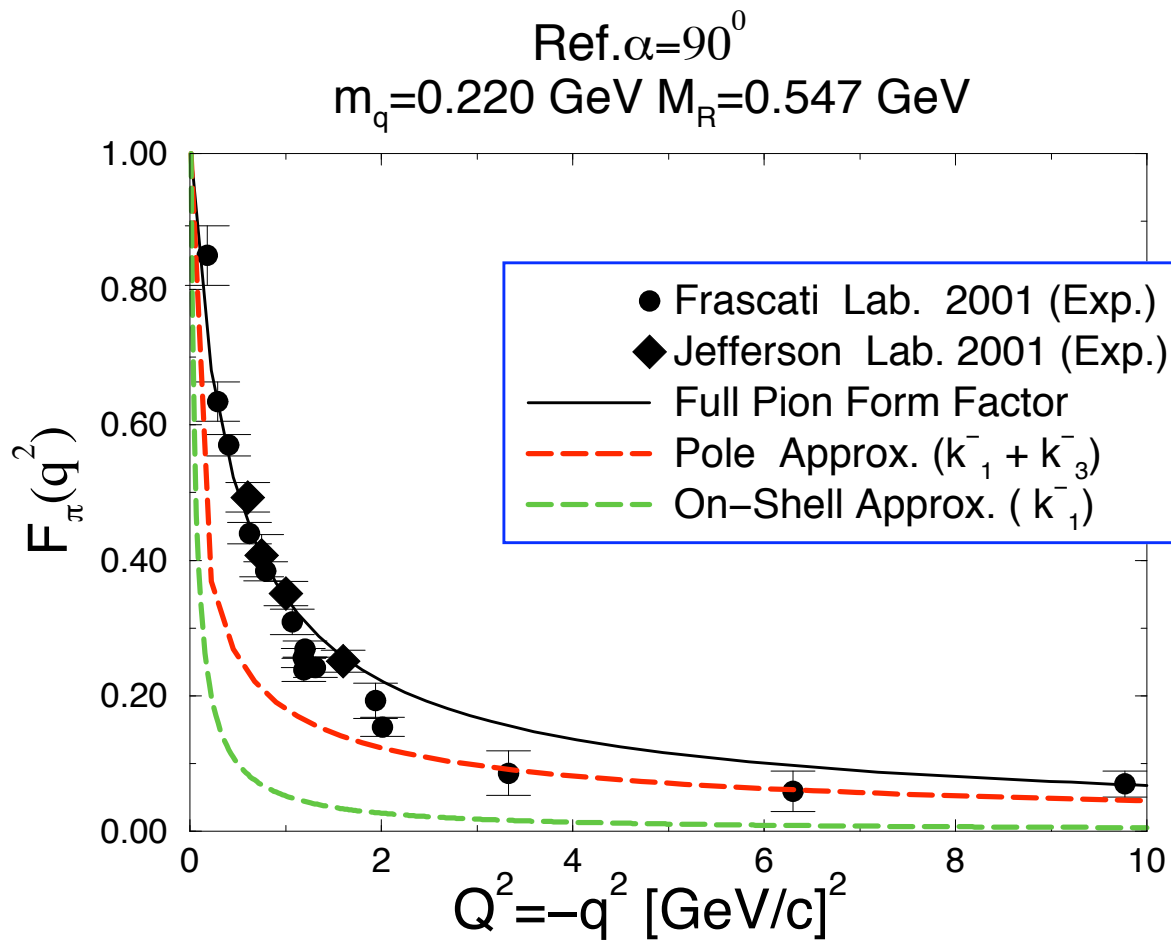


Fig. 12: Pion Pole Approximation

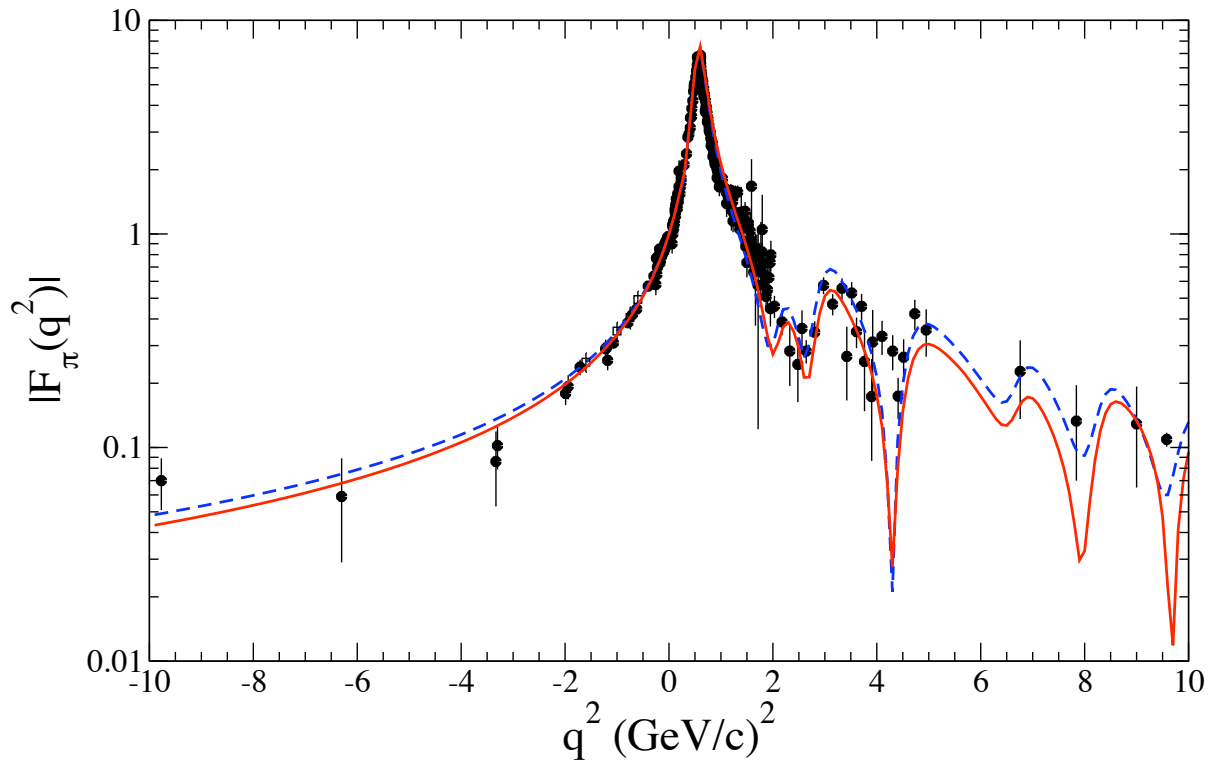


Fig. 13: Pion: Space and Time Like

• Ref. de Melo, Frederico, Salmé and Pace  
Phy. Rev. D707 (2006) 074013

## Conclusions

- Light-Front  $\implies$   $\left\{ \begin{array}{l} \textit{Bound States} \\ \textit{Covariance} \end{array} \right.$
- Rotational Invariance Broken  $\implies k^-$  Problematic
- Terms  $\left\{ \begin{array}{l} - \textit{Good} \\ - \textit{Bad} \end{array} \right.$
- Electromagnetic Current: “+”, “-” and “ $\perp$ ”
- Particles  $\left\{ \begin{array}{l} - \textit{Bosons} \\ - \textit{Pseudoscalar} \\ - \textit{Vector} \end{array} \right.$
- Pairs Terms Contribution  $\implies$  Full Covariance Restorate
- New Informations about Bound States  $q\bar{q}$