

**A PARABOLIC SYSTEM ARISING  
IN FILM DEVELOPMENT**

By

**Wenxiong Liu**

**IMA Preprint Series # 577**

August 1989

# A Parabolic System Arising In Film Development

Wenxiong Liu  
School of Mathematics  
University of Minnesota  
Minneapolis, MN 55455

*Key words and phrases:* Existence and uniqueness, edge enhancement, large time behavior

## §1 Introduction

In this paper we shall study the Cauchy problem:

$$P_t - P_{xx} = E(x)F(Q) - P + Q \quad \text{in } R^1 \times [0, \infty), \quad (1.1)$$

$$Q_t = P - Q \quad \text{in } R^1 \times [0, \infty), \quad (1.2)$$

$$P(x, 0) = Q(x, 0) = 0 \quad \text{for } x \in R^1, \quad (1.3)$$

$$\max_x |P(x, t)| < \infty, \quad \max_x |Q(x, t)| < \infty, \quad \text{for any } t < \infty. \quad (1.4)$$

where  $E(x)$  is a step function:

$$E(x) = \begin{cases} C_1 & x < 0 \\ C_2 & x > 0 \end{cases}, \quad \text{where } C_1 > C_2 > 0;$$

$F(u)$  is a  $C^2$  positive decreasing function,  $-\infty < u < \infty$ . System (1.1) arises in film development; for background, see [1].

Define dye  $D(x, T)$  at time  $T$  by

$$D(x, T) = E(x) \int_0^T F(Q(X, t)) dt \quad (1.5)$$

where  $(P, Q)$  is the solution of (1.1)-(1.4). Then one expects the profile of D-curve to be as indicated in Figure 1.

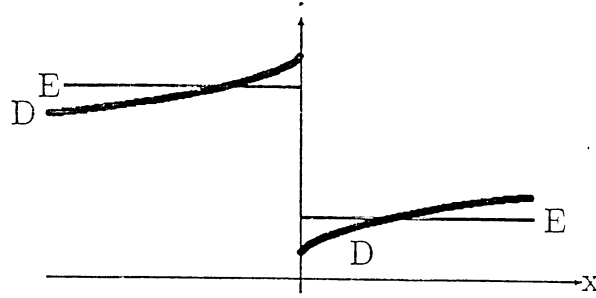


Figure 1

If it does, then we say that *edge enhancement* occurs (see [1]). In section 3, we shall prove that if  $T$  is sufficiently small, then this phenomenon does occur.

We shall be mainly interested however in the large time behavior of solutions. It turns out that this depends on the behavior of  $F(u)$  as  $u \rightarrow \infty$ . If  $F(u)$  reaches zero at a finite point  $u = M > 0$ , then both  $P$  and  $Q$  converge to  $M$  as  $t \rightarrow \infty$  under some additional assumption on  $F$  (Theorem 4.3). If  $F(u)$  approaches a positive limit as  $u \rightarrow \infty$ , then both  $P$  and  $Q$  increase linearly in  $t$  as  $t \rightarrow \infty$  (Theorem 4.10).

The most interesting case is when  $F(u)$  is strictly positive but converges to zero as  $u \rightarrow \infty$ ; then  $P$  and  $Q$  still converge to  $\infty$ ; more specifically, if  $F(u)/u^{-\gamma} \rightarrow \beta$ ,  $-F'(u)/u^{-\gamma-1} \rightarrow \gamma\beta$  as  $u \rightarrow \infty$  for some  $\beta > 0$  and  $\gamma > 0$ , then

$$\begin{aligned} P(x, t)/t^{\frac{1}{1+\gamma}} &\longrightarrow \frac{1}{2}g(0) \\ Q(x, t)/t^{\frac{1}{1+\gamma}} &\longrightarrow \frac{1}{2}g(0) \end{aligned} \quad \text{as } t \rightarrow \infty.$$

Here  $g(x)$  is the solution of

$$\begin{aligned} g''(x) + \frac{x}{2}g'(x) &= \frac{1}{1+\gamma}g - 2^\gamma\beta E(x)/g^\gamma, \\ g(+\infty) &= (2^\gamma C_2)^{\frac{1}{1+\gamma}}, \quad g(-\infty) = (2^\gamma C_1)^{\frac{1}{1+\gamma}}. \end{aligned} \quad (1.6)$$

## §2 Preliminary Results

Let  $W_p^{1,2}(K) = \{u : u, D_t u, D_x u \text{ and } D_{xx} u \text{ are all in } L^p(K)\}$ , where  $K \subset R^1 \times [0, \infty)$ . By a solution  $(P, Q)$  of (1.1)-(1.4), we mean that  $P \in W_p^{1,2}(K)$ ,  $Q \in W_p^{1,2}(K)$  for any  $p > 1$  and any compact subset  $K$ , and  $(P, Q)$  satisfies (1.1)-(1.4) ((1.1), (1.2) in a.e. sense).

In the sequel, various constants will be denoted by  $C$ . We shall first establish the existence and uniqueness of solutions of (1.1)-(1.4) for small time.

**Lemma 2.1** There exists a  $\delta > 0$ , such that the system (1.1)-(1.4) has a unique solution  $(P, Q)$  in  $R^1 \times [0, \delta]$ ; furthermore,  $P > 0$ ,  $Q > 0$ .

*Proof.* From (1.2), we get

$$Q(x, t) = \int_0^t e^{-(t-s)} P(x, s) ds, \quad (2.1)$$

so that (1.1) can be rewritten in the form

$$(P_t - P_{xx})(x, t) = E(x)F \left( \int_0^t e^{-(t-s)} P(x, s) ds \right) - P + \int_0^t e^{-(t-s)} P(x, s) ds. \quad (2.2)$$

We first solve (2.2) with  $P(x, 0) = 0$ . Let

$$B = C^0(R^1 \times [0, \delta]) \cap \{f : f \text{ is bounded in every strip } R^1 \times [0, t] \text{ for any } t < \delta\}.$$

For any  $\bar{P} \in B$ , let  $P$  be the unique solution of

$$\begin{aligned} (P_t - P_{xx})(x, t) &= E(x)F \left( \int_0^t e^{-(t-s)} \bar{P}(x, s) ds \right) - \bar{P} + \int_0^t e^{-(t-s)} \bar{P}(x, s) ds, \\ P(x, 0) &= 0. \end{aligned}$$

Define an operator  $L$  by  $L\bar{P} = P$ . It is easy to check that  $L$  maps  $B$  into  $B$ , and  $L$  is a contraction provided  $\delta$  is small. Therefore, there exists a unique fixed point  $P$  of  $L$ . Defining  $Q$  by (2.1), it is clear  $(P, Q)$  forms the unique solution of (1.1)-(1.4). By standard  $L^p$ -estimates (see [2]),  $P_t$  and  $P_{xx}$  are in  $L^p(K)$  for any  $p > 1$  and any compact subset  $K \subset R^1 \times [0, \delta]$ .

To prove that  $P(x, t) > 0$ , we represent  $P$  by means of the fundamental solution  $\Gamma(x, t; y, s) = \frac{1}{2\sqrt{\pi(t-s)}} e^{-\frac{(x-y)^2}{4(t-s)}}$ :

$$P(x, t) = \int_0^t \int_{R^1} \Gamma(x, t; y, s) (E(y)F(Q) - P(y, s) + Q(y, s)) dy ds. \quad (2.3)$$

Then

$$\begin{aligned} |P(x, t)| &\leq \int_0^t \int_{R^1} \Gamma(x, t; y, s) (E(y)F(Q) + C) dy ds \\ &\leq Ct. \end{aligned}$$

By (2.1), we also have the same estimate for  $Q(x, t)$ . Therefore,

$$\begin{aligned} P(x, t) &\geq \int_0^t \int_{R^1} \Gamma(x, t; y, s) (E(y)F(Q) - 2Cs) dy ds \\ &\geq Ct > 0, \end{aligned}$$

since  $F(Q) > \frac{1}{2} > 0$  for  $Q \leq Ct \leq C\delta$  if  $\delta$  is small. By (2.1), also  $Q(x, t) > 0$ .

We next prove some a priori estimates (for global solution).

**Lemma 2.2** If  $(P, Q)$  is a solution of (1.1)-(1.4) in  $R^1 \times [0, T)$ , then

$$P(x, t) > 0, \quad (2.4)$$

$$Q(x, t) > 0 \quad (2.5)$$

in  $R^1 \times (0, T)$ .

*Proof.* If the assertions are not true, then there exists a point  $(x_0, t_0)$  such that (2.4) and (2.5) hold for all  $t < t_0$  but one of the inequalities becomes an equality at  $(x_0, t_0)$ , i.e. either  $P(x_0, t_0) = 0$ , or  $Q(x_0, t_0) = 0$ . Note that  $Q(x_0, t_0) = 0$  is impossible by (2.1), so that  $P(x_0, t_0) = 0$ . Since

$$P_t - P_{xx} + P = E(x)F(Q) + Q > 0$$

in the strip  $\{0 \leq t \leq t_0\}$ , the maximum principle then implies that  $P \equiv 0$  in  $\{0 \leq t \leq t_0\}$ , a contradiction.

**Lemma 2.3** Let  $(P, Q)$  be a solution of (1.1)-(1.3) in  $R^1 \times [0, T)$ . Then

$$P(x, t) \leq CT e^T,$$

$$Q(x, t) \leq CT e^T$$

in  $R^1 \times [0, T)$ .

*Proof.* From (2.3) and the fact that  $P > 0$ , we have

$$\begin{aligned} P(x, t) &\leq \int_0^t \int_{R^1} \Gamma(x, t; y, s) (C_1 F(0) + Q(y, s)) dy ds \\ &\leq C_1 F(0)t + \int_0^t \max_x Q(x, s) ds \\ &= C_1 F(0)t + \int_0^t \int_0^s e^{-(t-s)} \max_x P(x, u) du ds \quad (\text{by (2.1)}) \\ &\leq C_1 F(0)t + \int_0^t \max_x P(x, s) ds. \end{aligned}$$

Gronwall inequality then gives

$$\max_x P(x, t) \leq C_1 F(0) t e^t \leq CT e^T.$$

By (2.1), the same inequality holds for  $Q$ .

The a priori estimates of Lemma 2.3 enable us to apply the local existence and uniqueness result of Lemma 2.1, step by step, in order to obtain a unique global solution :

**Theorem 2.4** There exists a unique solution of (1.1)-(1.4), for all  $t > 0$ ; furthermore,  $P > 0$ ,  $Q > 0$ .

We next establish additional regularities.

**Theorem 2.5** The solution  $(P, Q)$  has the following regularity properties:  $P_{tt}$  and  $P_{txx}$  are in  $L^p(K)$  for any  $p > 1$ ; furthermore,  $P_t$  and  $P_x$  are continuous up to  $\{t = 0\}$  except at  $(0, 0)$ . The same conclusions hold for  $Q$ .

*Proof.* Since  $(P, Q)$  is a solution of (1.1)-(1.3), for any  $\xi \in C_0^\infty(R^1 \times (0, \infty))$ , we have

$$-\int_0^\infty \int_{R^1} P(\xi_t + \xi_{xx}) = \int_0^\infty \int_{R^1} (E(x)F(Q) - P + Q)\xi.$$

Replacing  $\xi$  by  $\xi_t$ , and using integration by parts, we see that

$$\int_0^\infty \int_{R^1} P_t(\xi_t + \xi_{xx}) = -\int_0^\infty \int_{R^1} (E(x)F'(Q)Q_t - P_t + Q_t)\xi. \quad (2.6)$$

Equation (2.6) says that  $P_t$  is a weak solution of

$$P_{tt} - P_{txx} = (1 + E(x)F'(Q))Q_t - P_t. \quad (2.7)$$

Also

$$P_t(x, 0) = E(x)F(0).$$

Since  $P_t, Q_t$  are in  $L^p(K)$  for any compact subset  $K \subset R^1 \times [0, \infty)$ , standard  $L^p$ -estimates give us that  $P_{tt}$  and  $P_{txx}$  are in  $L^p(K)$  for any  $p > 1$ . By Sobolev imbedding theorems, we see that  $P_t$  and  $P_x$  are continuous in  $R^1 \times (0, \infty)$ . Noting that  $E(x)$  has a jump only at 0, we conclude that  $P_t$  and  $P_x$  are continuous up to  $\{t = 0\}$  except at  $(0, 0)$ . By (2.1), all assertions are also true for  $Q$ .

### §3 Edge Enhancement

In this section, we establish the edge enhancement phenomenon.

**Theorem 3.1** Let  $D(x, T)$  be the dye as defined in (1.6). Then, for any  $T$  sufficiently small, we have that

$$D(x, T) = E(x)g(x, T) + O(T^4)$$

where  $g(x, T) \in C^\infty(R^1)$  and  $\inf_{x, T} \frac{g(x, T)}{T} \geq c > 0$ ; furthermore,

$$g_x(x, T) \geq 0, \quad x \in R^1$$

$$g_{xx}(x, T) \geq 0 \quad \text{if } x < 0,$$

$$g_{xx} \leq 0 \quad \text{if } x > 0,$$

$$\lim_{x \rightarrow \infty} g(x, T) = F(0)T(1 + \frac{1}{6}C_2F(0)F'(0)T^2), \quad \lim_{x \rightarrow -\infty} g(x, T) = F(0)T(1 + \frac{1}{6}C_1F(0)F'(0)T^2).$$

*Proof.* Let  $f(x, t) = (1 + E(x)F'(Q))Q_t - P_t$ . Then from (2.10)

$$\begin{aligned} P_t(x, t) &= \int_{R^1} \Gamma(x, t; y, 0)E(y)F(0)dy + \int_0^t \int_{R^1} \Gamma(x, t; y, s)f(y, s)dy ds \\ &\triangleq I_1 + I_2. \end{aligned} \quad (3.1)$$

We compute

$$\begin{aligned}
I_1 &= F(0)[C_1 \int_{-\infty}^0 \Gamma(x, t; y, 0) dy + C_2 \int_0^{\infty} \Gamma(x, t; y, 0) dy] \\
&= F(0)[\frac{C_1 + C_2}{2} + (C_2 - C_1) \int_{-\frac{x}{\sqrt{2t}}}^0 \frac{1}{\sqrt{2\pi}} e^{-\frac{u^2}{2}} du], \tag{3.2}
\end{aligned}$$

and

$$I_2 \leq \int_0^t |f|_{L^\infty(\mathbb{R}^1 \times [0, \infty))} ds \leq Ct. \tag{3.3}$$

Now let us calculate  $Q(x, t)$  for small  $t$ :

$$\begin{aligned}
Q(x, t) &= \int_0^t e^{-(t-s)} P(x, s) ds \\
&= \int_0^t P(x, s) ds + \int_0^t P(x, s)(e^{-(t-s)} - 1) ds \\
&= \int_0^t P_t(x, s)(t-s) ds + O(t^3). \tag{3.4}
\end{aligned}$$

Hence

$$\begin{aligned}
\int_0^T Q(x, t) dt &= \frac{1}{2} \int_0^T P_t(x, s)(T-s)^2 ds + O(T^4) \\
&= \frac{1}{2} \int_0^T I_1(x, s)(T-s)^2 ds + \int_0^T I_2(x, s)(T-s)^2 ds + O(T^4) \quad (\text{by (3.1)}) \\
&= \frac{1}{2} \int_0^T I_1(x, t)(T-t)^2 dt + O(T^4) \quad (\text{by (3.3)}). \tag{3.5}
\end{aligned}$$

From (3.4), we know that  $Q(x, t) = O(t^2)$ . Therefore

$$F(Q(x, t)) = F(0) + F'(0)Q(x, t) + O(t^4).$$

Integrating in  $t$  and using (3.5) and (3.2), we get

$$\begin{aligned}
\int_0^T F(Q(x, t)) dt &= F(0)T + F'(0) \int_0^T Q(x, t) dt + O(T^4) \\
&\triangleq g(x, T) + O(T^4),
\end{aligned}$$

where

$$g(x, T) = F(0)T + \frac{1}{2} F'(0) F(0) \int_0^T [\frac{C_1 + C_2}{2} + (C_2 - C_1) \int_{-x/\sqrt{2t}}^0 \frac{1}{\sqrt{2\pi}} e^{-u^2/2} du] (T-t)^2 dt.$$

It is easy to check that  $g(x, T)$  has all the properties stated in Theorem 3.1. It is also clear that D-curve has the profile indicated in Figure 1. Hence, if  $T$  is sufficiently small, edge enhancement occurs.

## §4 Large Time Behavior

To study the large time behavior, we need to distinguish three cases:

$$\text{There exists an } M > 0, \text{ such that } F(M) = 0, \quad (4.1)$$

$$F(u) \rightarrow \beta > 0 \text{ as } u \rightarrow \infty, \quad (4.2)$$

$$F(u) \rightarrow 0 \text{ as } u \rightarrow \infty. \quad (4.3)$$

We shall begin with case 1.

**Lemma 4.1** In addition to (4.1), we also assume that  $F$  satisfies

$$1 + \max_u C_1 F'(u) \geq 0. \quad (4.4)$$

Then  $P_t > 0$ ,  $Q_t > 0$  in  $R^1 \times [0, \infty)$ .

*Proof.* By (1.2) and Lemma 2.3, we see that

$$|Q_t|_{L^\infty(R^1 \times [0,1])} \leq C < \infty, \quad (4.5)$$

Representing  $P_t$ , the solution of (2.7) with  $P_t(x, 0) = E(x)F(0)$ , in terms of the fundamental solution, we get

$$P_t(x, t) = \int_{R^1} \Gamma(x, t; y, 0) E(y) F(0) dy + \int_0^t \int_{R^1} \Gamma(x, t; y, s) [(1 + E(y)F'(Q))Q_t - P_t] dy ds \quad (4.6)$$

Hence

$$|P_t(x, t)| \leq C_1 F(0) + \tilde{C}t + \int_0^t \max_x |P_t(x, s)| ds, \quad (t \leq 1)$$

where  $\tilde{C} = \max_{x, 0 \leq t \leq 1} (1 + E(x)F'(Q(x, s))|Q_t(x, t)|) < \infty$  by (4.5). Gronwall's inequality then gives

$$\max_x |P_t(x, t)| \leq (C_1 + \tilde{C})e^t \leq C \quad (t \leq 1). \quad (4.7)$$

Using (4.5) and (4.7) in (4.6), we obtain

$$P_t(x, t) \geq C_2 F(0) - Ct > 0$$

in  $R^1 \times [0, \delta]$  for some small  $\delta > 0$ . Since  $Q_t = \int_0^t e^{-(t-s)} P_t(x, s) ds$ , we also have  $Q_t(x, t) > 0$ . On the other hand, by the assumption (4.4):

$$P_{tt} - P_{txx} + P_t = (1 + E(x)F'(Q))Q_t \geq 0$$

as long as  $Q_t \geq 0$ . The argument used in the proof of Lemma 2.2 now shows that  $P_t > 0$ , and  $Q_t > 0$  in  $R^1 \times [0, \infty)$ .

**Lemma 4.2** Under the same assumptions of Lemma 4.1, we have

$$P(x, t) < M, \quad Q(x, t) < M \text{ in } R^1 \times [0, \infty).$$

*Proof.* By Lemma 2.3, there is a  $\delta > 0$ , such that

$$P(x, t) < M \quad \text{in } R^1 \times [0, \delta].$$

Let  $\bar{P} = M - P(x, t)$ ,  $\bar{Q} = M - Q(x, t)$ , then

$$\begin{aligned} \bar{P} &> 0 \\ \bar{Q} &> 0 \end{aligned} \quad \text{in } R^1 \times [0, \delta].$$

From (1.1)

$$\bar{P}_t - \bar{P}_{xx} = -E(x)F(Q) - \bar{P} + \bar{Q}.$$

Since  $F(M) = 0$ ,

$$F(Q) = \int_0^1 F'(M + s(Q - M))ds(Q - M)$$

so that

$$\bar{P}_t - \bar{P}_{xx} + \bar{P} = [1 + E(x) \int_0^1 F'(M + s(Q - M))ds]\bar{Q} \geq 0$$

as long as  $\bar{Q} \geq 0$ . The same argument as before shows that  $\bar{P} \geq 0, \bar{Q} \geq 0$ .

Now we are in a position to state:

**Theorem 4.3** If  $F$  satisfies (4.1) and (4.4), then for any  $x \in R^1$

$$\begin{aligned} P(x, t) &\rightarrow M \\ Q(x, t) &\rightarrow M \end{aligned} \quad \text{as } t \rightarrow \infty.$$

*Proof.* By Lemma 4.1, the limits  $\lim_{t \rightarrow \infty} P(x, t) \triangleq \bar{P}(x)$ ,  $\lim_{t \rightarrow \infty} Q(x, t) \triangleq \bar{Q}(x)$  exist. By Lemma 4.2,  $\bar{P}(x) \leq M, \bar{Q}(x) \leq M$ .

Integrating  $Q_t = P - Q$  over  $(t - 1, t)$ , we get

$$Q(x, t) - Q(x, t-1) = P(x, t) - Q(x, t) + \int_{t-1}^t (P(x, s) - P(x, t))ds - \int_{t-1}^t (Q(x, s) - Q(x, t))ds.$$

Letting  $t \rightarrow \infty$ , we obtain

$$\bar{P}(x) - \bar{Q}(x) = 0. \tag{4.8}$$

Let  $\xi(x) \in C_0^\infty(R^1)$ . We multiply  $P_t - P_{xx} = E(x)F(Q) - P + Q$  by  $\xi$ , then integrate over  $R^1 \times (t - 1, t)$ , and finally let  $t \rightarrow \infty$  to get

$$- \int_{R^1} \bar{P}(x)\xi''(x) = \int E(x)F(\bar{P}(x))dx. \tag{4.9}$$

This means that  $\bar{P}(x)$  is a solution of

$$\bar{P}'' + E(x)F(\bar{P}) = 0. \tag{4.10}$$

and, in particular,  $\bar{P}'' \leq 0$ . But the function satisfying  $0 \leq \bar{P} \leq M, \bar{P}'' \leq 0$  is clearly a constant and since  $\bar{P}$  satisfies (4.10), this constant must be equal to  $M$ . This proves the assertions of the theorem.

**Remark 4.4** We conjecture that Theorem 4.3 is not true without assumption (4.4), as the following example suggests. Let

$$F(u) = \begin{cases} M - u & \text{if } u \leq M \\ 0 & \text{otherwise.} \end{cases}$$

Consider the system

$$\begin{aligned} p' &= KF(q) - p + q, \\ q' &= p - q, \\ p(0) &= q(0) = 0, \end{aligned} \tag{4.11}$$

where  $K > 1$ . This system corresponds to (1.1)-(1.4) with  $E(x) \equiv K$ , and certainly (4.4) is violated. Since  $q'' = p' - q' = KF(q) - 2q'$  and  $q'(0) = 0$ , we see that  $q'(t) = \int_0^t e^{-(t-s)} KF(q(s)) ds > 0$  for any  $t > 0$ . Noting that  $q'' = K(M - q) - 2q'$  for  $q \leq M$ , we can solve this linear ODE with initial conditions:  $q(0) = q'(0) = 0$ . From the resulting expression for  $q$ , we see that  $q(t_0) = M$  for some  $t_0 > 0$ . Since  $q' > 0$ , we conclude that  $\lim_{t \rightarrow \infty} q(t) > M$ . Hence the conclusions of Theorem 4.3 are not true in this case.

Now we turn to the remaining cases (4.2) and (4.3). We shall first consider the system

$$p' = KF(q) - p + q, \tag{4.12}$$

$$q' = p - q, \tag{4.13}$$

$$p(0) = q(0) = 0 \tag{4.14}$$

where  $K$  is a positive constant and  $F$  satisfies either (4.2) or (4.3). By Theorem 2.4, there exists a unique positive global solution of (4.11)-(4.13).

**Lemma 4.5** If  $F$  satisfies either (4.2) or (4.3), then there exist positive constants  $c, C$  and  $t_0$  such that

$$\begin{aligned} 0 < q' &\leq C \quad \text{for any } t \in \mathbb{R}^+, \\ |p'| &\leq C \quad \text{for any } t \in \mathbb{R}^+. \end{aligned} \tag{4.15}$$

Moreover, if  $F(u) \sim u^{-\gamma}$  at  $\infty$  for  $\gamma \geq 0$  (If  $\gamma = 0$ , we shall mean that  $F$  satisfies (4.2)), then

$$ct^{1/(1+\gamma)} \leq p(t) \leq Ct^{1/(1+\gamma)}, \quad ct^{1/(1+\gamma)} \leq q(t) \leq Ct^{1/(1+\gamma)} \tag{4.16}$$

for  $t \geq t_0$ .

*Proof.* The same argument of Remark 4.4 shows that  $q' > 0$  in  $\mathbb{R}^+$ . From this we deduce that

$$p' = KF(q) - q' \leq KF(q) \leq KF(0).$$

This, in turn, implies

$$q'(t) = \int_0^t e^{-(t-s)} p'(s) ds \leq KF(0),$$

and  $p' = KF(q) - q' \geq -q' \geq -KF(0)$ . We have thus proved (4.15).

Next we show that  $q(t) \rightarrow \infty$  as  $t \rightarrow \infty$ . Suppose this is not true, i.e.

$$0 < q(t) \leq M < \infty \quad \text{for any } t. \tag{4.17}$$

Integrating  $q'' = p' - q' = KF(q) - 2q'(t)$  over  $(0, t)$ , we get

$$q'(t) + 2q(t) = K \int_0^t F(q(s)) ds. \quad (4.18)$$

Since  $F(q(s)) \geq F(M)$ , we see that the RHS of (4.18) is larger than  $KF(M)t$ , whereas the LHS of (4.18) is less than  $KF(0) + 2M$  by (4.15) and (4.17). This is a contradiction if  $t$  is large.

Now integrating  $q''/F(q) = K - 2q'/F(q)$  over  $(0, t)$ , and using integration by parts, we obtain

$$\begin{aligned} \int_0^{q(t)} \frac{1}{F(s)} ds &= \frac{K}{2}t - \frac{q'(t)}{2F(q(t))} - \int_0^t \frac{F'(q)}{F^2(q)} q'^2 ds \\ &\geq \frac{K}{2}t - \frac{q'(t)}{2F(q(t))} \quad (\text{since } F' \leq 0) \\ &\geq \frac{K}{2}t - \frac{C}{F(q(t))} \quad (\text{since } 0 < q' \leq KF(0)). \end{aligned}$$

Since  $q(t) \rightarrow \infty$  as  $t \rightarrow \infty$ , we conclude from the above that

$$\begin{aligned} \frac{K}{2}t &\leq \frac{C}{F(q(t))} + \int_0^{q(t)} \frac{1}{F(s)} ds \\ &\leq Cq(t)^\gamma + Cq(t)^{1+\gamma} \\ &\leq Cq(t)^{1+\gamma}, \end{aligned}$$

provided  $t$  is large. This gives  $q(t) \geq ct^{1/(1+\gamma)}$ . Next integrating  $q'' = KF(q) - 2q'$  over  $(0, t)$ , and applying the result we just proved, we get

$$\begin{aligned} q(t) &= \frac{K}{2} \int_0^t F(q(s)) ds - \frac{q'(t)}{2} \\ &\leq C \int_0^t \frac{1}{q(s)^\gamma} ds + C \\ &\leq C \int_0^t s^{-\gamma/(1+\gamma)} ds + C \\ &\leq Ct^{\frac{1}{1+\gamma}} \end{aligned}$$

provided  $t$  is large. Since  $p(t)/q(t) = q'(t)/q(t) + 1$ , and  $q(t) \rightarrow \infty$  while  $q'$  remains bounded, we see that  $p(t)/q(t) \rightarrow 1$  as  $t \rightarrow \infty$ . Hence all the estimates for  $q$  hold for  $p$  as well. The proof is complete.

**Lemma 4.6** If  $F(u) \sim u^{-\gamma}$  at  $\infty$  ( $\gamma \geq 0$ ), then

$$p_1(t) \leq P(x, t) \leq p_2(t),$$

$$q_1(t) \leq Q(x, t) \leq q_2(t),$$

where  $(p_i, q_i)$  is a solution of (4.12)-(4.14) with  $K = K_i$ , and  $K_1$  is sufficiently small while  $K_2$  is sufficiently large.

Combining Lemma 4.5 and Lemma 4.6, we obtain

**Corollary 4.7** Under the same assumptions of Lemma 4.6, we can find positive constants  $c, C$  and  $t_0$ , such that

$$\begin{aligned} ct^{1/(1+\gamma)} &\leq P(x, t) \leq Ct^{1/(1+\gamma)}, \\ ct^{1/(1+\gamma)} &\leq Q(x, t) \leq Ct^{1/(1+\gamma)} \end{aligned}$$

for all  $x \in R^1$  and  $t \geq t_0$ .

*Proof of Lemma 4.6.* Let  $u_i = P - p_i$ ,  $v_i = Q - q_i$ . It is easy to see that

$$u_{it} - u_{ixx} + u_i = E(x)F(Q) - K_i F(q_i) + v_i \triangleq J_i(x, t) \quad (4.19)$$

$$v_{it} = u_i - v_i$$

$$u_i(x, 0) = v_i(x, 0) = 0.$$

Representing the solution  $u_i$  of (4.19) with  $u_i(x, 0) = 0$ , in terms of the fundamental solution, we get that for  $t$  small

$$\begin{aligned} u_1(x, t) &= \int_0^t \int_{R^1} \Gamma(x, t; y, s) [E(y)F(Q(y, s)) - K_1 F(q_1(s)) + v_1(y, s) - u_1(y, s)] dy ds \\ &\geq \int_0^t \int_{R^1} \Gamma(x, t; y, s) \left[ \frac{E(y) - 2K_1}{2} F(0) - Cs \right] dy ds \\ &\geq \int_0^t \left( \frac{C_2 - 2K_1}{2} - Cs \right) ds \\ &> 0 \quad (\text{if } K_1 < \frac{C_2}{2}). \end{aligned}$$

Hence also

$$v_1(x, t) = \int_0^t e^{-(t-s)} u_1(x, s) ds > 0 \quad (4.20)$$

for  $t$  small. Using mean value theorem, we see that

$$J_1 = (E(x) - K_1)F(Q) + [1 + K_1 F'(\lambda Q + (1 - \lambda)q)](Q - q) \geq 0 \quad (\text{for some } \lambda, 0 < \lambda < 1)$$

as long as  $Q - q = v_1 \geq 0$  and  $K_1$  is sufficiently small. Now we use the maximum principle to argue as same as in the proof of Lemma 2.2 to deduce that  $u_1(x, t) > 0$ ,  $v_1(x, t) > 0$  in  $R^1 \times (0, \infty)$ .

Next we consider  $u_2(x, t)$ ,  $v_2(x, t)$ . The same argument as in proving (4.20) shows that  $u_2(x, t) < 0$ ,  $v_2(x, t) < 0$  in  $R^1 \times (0, \delta]$  for some  $\delta > 0$ , provided  $K_2$  is sufficiently large. We shall next show that

$$J_2(x, t) \leq 0 \quad \text{in } R^1 \times [0, \infty) \quad (4.21)$$

as long as  $v_2(x, t) \leq 0$ . Hence the maximum principle gives that  $u_2(x, t) < 0$ ,  $v_2(x, t) < 0$  in  $R^1 \times (0, \infty)$ .

It remains to verify (4.21). By the assumptions on  $F$ , we have that  $|F'(u)| \rightarrow 0$  as  $u \rightarrow \infty$ . We choose  $N$  so large such that  $\frac{N}{2} \geq C_1 F(0)$  and  $1 + E(x)F'(u) \geq 0$  if  $u \geq \frac{N}{2}$ . If  $q_2 \geq N$ ,  $Q \leq \frac{N}{2}$ , then

$$J_2(x, t) \leq E(x)F(0) - \frac{N}{2} \leq 0.$$

If  $q_2 \geq N$ ,  $Q \geq \frac{N}{2}$ , then

$$\begin{aligned} J_2(x, t) &= (E(x) - K_2)F(q_2) + E(x)(F(Q) - F(q_2)) + Q - q_2 \\ &= (E(x) - K_2)F(q_2) + [1 + E(x)F'(\lambda q_2 + (1 - \lambda)Q)](Q - q_2) \\ &\leq 0 \end{aligned}$$

as long as  $v_2 = Q - q_2 \leq 0$ , because  $\lambda q_2 + (1 - \lambda)Q \geq \frac{N}{2}$  implies that  $1 + E(x)F'(\lambda q_2 + (1 - \lambda)Q) \geq 0$ .

Finally, if  $q \leq N$ , then

$$\begin{aligned} J_2(x, t) &\leq C_1 F(0) - K_2 F(N) + Q - q_2 \\ &\leq C_1 F(0) - K_2 F(N) \\ &\leq 0 \end{aligned}$$

as long as  $v_2 = Q - q_2 \leq 0$  and  $K_2$  is sufficiently large. The proof is complete.

**Lemma 4.8** Under the same assumptions of Lemma 4.6, we have that

$$\begin{aligned} |P_t| &\leq M \\ |Q_t| &\leq M \end{aligned} \quad \text{in } R^1 \times [0, \infty).$$

From (1.1) and (1.2), we then get

**Corollary 4.9**  $|P_{xx}| \leq C < \infty$ ,  $|Q_{xx}| \leq C < \infty$  in  $R^1 \times [0, \infty)$ .

*Proof of Lemma 4.8.* From (4.6), we obtain

$$\begin{aligned} |P_t(x, t)| &\leq C + C \int_0^t \int_{R^1} \Gamma(x, t; y, s) [|P - Q|(y, s) + |P_t|(y, s)] dy ds \\ &\leq C + C \int_0^t s e^{Cs} ds + \int_0^t \max_x |P_t(x, s)| ds \quad (\text{by Lemma 2.3}). \end{aligned}$$

Gronwall inequality then gives

$$\max_x |P_t(x, t)| \leq g(t) \tag{4.22}$$

where  $g(t)$  is a positive increasing function. Let  $u = P_t - M$ ,  $v = Q_t - M$ . It is easy to see that

$$u_t - u_{xx} + u = (1 + E(x)F'(Q))v + ME(x)F'(Q) \tag{4.23}$$

$$v_t = u - v. \tag{4.24}$$

By Corollary 4.7 and the fact that  $F'(u) \rightarrow 0$  as  $u \rightarrow \infty$ , we can find a  $\tilde{t}$ , such that

$$1 + E(x)F'(Q(x, t)) \geq 0 \quad \text{in } R^1 \times [\tilde{t}, \infty).$$

By (4.22), we can choose  $M$  so large such taht

$$\begin{aligned} u &< 0 \\ v &< 0 \end{aligned} \quad \text{in } R^1 \times [0, \tilde{t}].$$

We shall show that  $u < 0$ ,  $v < 0$  in  $R^1 \times [0, \infty)$ . If they are not true, then the argument as in the proof of Lemma 2.2, shows that there exists a point  $(x_0, t_0)$ , such that

$$\begin{aligned} u &< 0 \\ v &< 0 \end{aligned} \quad \text{in } R^1 \times [0, t_0)$$

and  $u(x_0, t_0) = 0$ . Obviously,  $\tilde{t} < t_0$ . Hence we see that the RHS of (4.23) is nonpositive in  $R^1 \times [\tilde{t}, t_0)$ . The maximum principle implies that  $u \equiv 0$  in the strip  $R^1 \times [\tilde{t}, t_0)$ , which is a contradiction. This proves our assertions. Similarly, we can prove  $P_t + M \geq 0$ ,  $Q_t + M \geq 0$  in  $R^1 \times [0, \infty)$ . The proof is complete.

**Theorem 4.10** If  $F$  satisfies (4.2), then

$$\lim_{t \rightarrow \infty} \frac{P(x, t)}{t} = \frac{C_1 + C_2}{4} \beta, \quad \lim_{t \rightarrow \infty} \frac{Q(x, t)}{t} = \frac{C_1 + C_2}{4} \beta \quad \text{uniformly in } x \text{ as } t \rightarrow \infty.$$

*Proof.* From (1.1) and (1.2)

$$(P + Q)_t - \frac{1}{2}(P + Q)_{xx} = E(x)F(Q) + \frac{1}{2}(P - Q)_{xx}.$$

Making a change of variables:  $y = \sqrt{2}x$ ,  $w = P + Q$ , we get

$$w_t - w_{yy} = E(y)F(Q) + (P - Q)_{yy}. \quad (4.25)$$

Also

$$w(y, 0) = 0.$$

Representing  $w$  in terms of the fundamental solution, we have

$$\begin{aligned} w(y, t) &= \int_0^t \int_{R^1} \Gamma(y, t; z, s) E(z) F(Q(z, s)) dz ds + \int_0^t \int_{R^1} \Gamma(y, t; z, s) (P - Q)_{zz}(z, s) dz ds \\ &\triangleq I_1 + I_2. \end{aligned} \quad (4.26)$$

Since by Corollary 4.7,  $F(Q(y, t)) \rightarrow \beta$  uniformly in  $y$  as  $t \rightarrow \infty$ , we can get

$$\begin{aligned} \frac{I_1}{t} &= \frac{1}{t} \int_0^t \int_{R^1} \frac{1}{\sqrt{2\pi}} e^{-u^2/2} E(y + \sqrt{2(t-s)}u) F(Q(y + \sqrt{2(t-s)}u, s)) ds du \\ &= \int_{R^1} \frac{1}{\sqrt{2\pi}} e^{-u^2/2} du \int_0^1 E(y + \sqrt{2t(1-v)}u) F(Q(y + \sqrt{2t(1-v)}u, tv)) dv \\ &\longrightarrow \int_0^\infty \frac{1}{\sqrt{2\pi}} e^{-u^2/2} C_1 \beta du + \int_{-\infty}^0 \frac{1}{\sqrt{2\pi}} e^{-u^2/2} C_2 \beta du \\ &= \frac{C_1 + C_2}{2} \beta \quad \text{uniformly in } y \text{ as } t \rightarrow \infty. \end{aligned} \quad (4.27)$$

We turn to the estimate of  $I_2$ . Write

$$\begin{aligned} I_2 &= \int_0^{t-1} \int_{R^1} \Gamma(y, t; z, s) (P - Q)_{zz}(z, s) dz ds + \int_{t-1}^t \int_{R^1} \Gamma(y, t; z, s) (P - Q)_{zz}(z, s) dz ds \\ &\triangleq J_1 + J_2. \end{aligned}$$

By Corollary 4.9,

$$\frac{|J_2|}{t} \leq \frac{C}{t} \int_{t-1}^t ds = \frac{C}{t}.$$

As to  $J_1$ , using integrations by parts, and noting that boundary terms disappear because  $(P - Q)_y$  and  $(P - Q)$  are bounded in  $y$  at  $\infty$ , we obtain

$$\begin{aligned} \frac{|J_1|}{t} &\leq \frac{C}{t} \int_0^{t-1} \int_{R^1} |\Gamma_t(y, t; z, s)| |P - Q| dz ds \\ &\leq \frac{C}{t} \int_0^{t-1} \int_{R^1} \left( \frac{1}{t-s} + \frac{(y-z)^2}{(t-s)^2} \right) \Gamma(y, t; z, s) |Q_t(z, s)| dz ds \\ &\leq \frac{C}{t} \int_0^{t-1} \frac{1}{t-s} ds \\ &= \frac{C \log t}{t} \longrightarrow 0, \quad \text{as } t \rightarrow \infty. \end{aligned}$$

From (4.26)-(4.27), and the estimates on the  $J_i$ , we conclude that

$$\frac{w(y, t)}{t} \rightarrow \frac{C_1 + C_2}{2} \beta \quad \text{uniformly in } y \text{ as } t \rightarrow \infty. \quad (4.28)$$

Noting that  $Q_t = P - Q$ , and  $Q_t$  is bounded whereas  $Q(x, t) \rightarrow \infty$  as  $t \rightarrow \infty$ , we get

$$\frac{P(x, t)}{Q(x, t)} = 1 + \frac{Q_t(x, t)}{Q(x, t)} \longrightarrow 1 \quad \text{uniformly in } x \text{ as } t \rightarrow \infty. \quad (4.29)$$

Combining (4.28) and (4.29), we obtain the assertions of the theorem.

## §5 Large Time Behavior(Continued)

In this section, we consider the case (4.3); more specifically, we shall assume that

$$\begin{cases} F(u)u^\gamma \rightarrow \beta > 0 \\ -F'(u)u^{\gamma+1} \rightarrow \gamma\beta > 0 \end{cases} \quad \text{as } u \rightarrow \infty. \quad (5.1)$$

We shall study the large time behavior of  $w \equiv P + Q$ . It will be convenient to rewrite (4.25) as

$$\begin{aligned} w_t - w_{yy} &= 2^\gamma E(y)F(w) + E(y)F(w)\left(\frac{F(Q)}{F(w)} - 2^\gamma\right) + (P - Q)_{yy} \\ &\triangleq 2^\gamma E(y)F(w) + f_1(y, t) + f_2(y, t). \end{aligned}$$

Consider the solution of

$$\begin{cases} u_t - u_{yy} = 2^\gamma E(y)F(w) \\ u(y, 0) = 0. \end{cases} \quad (5.2)$$

**Lemma 5.1** The function  $u$  satisfies

$$\frac{|u - w|_{L^\infty(R^1)}(t)}{t^{1/(1+\gamma)}} \longrightarrow 0 \quad \text{as } t \rightarrow \infty, \quad (5.3)$$

and consequently

$$ct^{\frac{1}{1+\gamma}} \leq u(y, t) \leq Ct^{\frac{1}{1+\gamma}} \quad (5.4)$$

for all  $y \in R^1$  and  $t \geq N$ ,  $N$  a large constant.

*Proof.* It is clear that  $u - w = I_1 + I_2$  where

$$I_i(y, t) = \int_0^t \int_{R^1} \Gamma(y, t; z, s) f_i(z, s) dz ds, \quad i = 1, 2.$$

$I_2$  is the same in the proof of Theorem 4.10 so that  $I_2(y, t)/t^{1/(1+\gamma)} \rightarrow 0$  uniformly in  $y$  as  $t \rightarrow \infty$ . To complete the proof of (5.3), it remains to show that  $I_1(y, t)/t^{1/(1+\gamma)} \rightarrow 0$  uniformly in  $y$  as  $t \rightarrow \infty$ . Noting that

$$h(t) \triangleq \max_y \left| \frac{F(Q)}{F(w)} - 2^\gamma \right| \rightarrow 0 \quad \text{as } t \rightarrow \infty$$

by (4.29) and (5.1), we get

$$\begin{aligned} |I_1(y, t)|/t^{1/(1+\gamma)} &\leq \frac{1}{t^{1/(1+\gamma)}} \int_0^t \int_{R^1} \Gamma(y, t; z, s) F(w) w^\gamma \frac{1}{w^\gamma} \left| \frac{F(Q)}{F(w)} - 2^\gamma \right| dz ds \\ &\leq \frac{C}{t^{1/(1+\gamma)}} \int_0^t \frac{h(s)}{s^{\gamma/(1+\gamma)}} ds \quad (\text{by Lemma 5.1}) \\ &= C \int_0^1 \frac{h(tv)}{v^{\gamma/(1+\gamma)}} dv \longrightarrow 0 \quad \text{as } t \rightarrow \infty. \end{aligned}$$

From (5.3) and Corollary 4.7, (5.4) follows.

Next we prove

**Lemma 5.3** There exist positive constants  $C$  and  $t_0$  such that

$$|u_y(y, t)| \leq Ct^{\frac{1}{1+\gamma} - \frac{1}{2}}$$

for all  $y \in R^1$  and  $t \geq t_0$ .

*Proof.* Representing  $u$  in terms of the fundamental solution, we have

$$u(y, t) = \int_0^t \int_{R^1} \Gamma(y, t; z, s) 2^\gamma E(z) F(w(z, s)) dz ds.$$

Differentiating the above equation in  $y$  and using Corollary 4.7, we see that

$$\begin{aligned}
|u_y(y, t)| &\leq C \int_0^t \int_{R^1} \Gamma(y, t; z, s) \frac{|x - y|}{(t - s)w^\gamma(z, s)} dz ds \\
&\leq C \int_0^t \frac{1}{(t - s)^{1/2} s^{\gamma/(1+\gamma)}} ds \\
&\leq C t^{\frac{1}{1+\gamma} - \frac{1}{2}}.
\end{aligned}$$

To study the large time behavior of  $u$ , we shall introduce the scaled function  $u_\alpha(y, t) = u(a + \alpha y, \alpha^2 t) / \alpha^{\frac{2}{1+\gamma}}$  for any  $a \in R^1$  and  $\alpha > 0$ . Then

$$\begin{aligned}
u_{\alpha t} - u_{\alpha y y} &= 2^\gamma \alpha^{\frac{2\gamma}{1+\gamma}} E(a + \alpha y) F(w) \\
&= \frac{2^\gamma E(a + \alpha y)}{u_\alpha^\gamma} F(u) u^\gamma + 2^\gamma E(a + \alpha y) \alpha^{\frac{2\gamma}{1+\gamma}} (F(w) - F(u)) \\
&\triangleq g_{1\alpha} + g_{2\alpha}.
\end{aligned} \tag{5.5}$$

In the above expressions, all functions are evaluated at  $(a + \alpha y, \alpha^2 t)$ . Using the assumption (5.1) and (5.3), we obtain

$$\begin{aligned}
|g_{2\alpha}(y, t)| &\leq C \alpha^{\frac{2\gamma}{1+\gamma}} |F'(\lambda u + (1 - \lambda)w)| |u - w| \quad (0 < \lambda < 1) \\
&\leq C \frac{1}{\alpha^{\frac{2}{1+\gamma}}} |u - w|_{L^\infty(R^1)}(\alpha^2 t) \longrightarrow 0 \quad \text{as } \alpha \rightarrow \infty, \text{ for any } t > 0;
\end{aligned} \tag{5.6}$$

here, as before, all functions are evaluated at  $(a + \alpha y, \alpha^2 t)$ . Note that, by (5.4),

$$ct^{\frac{1}{1+\gamma}} \leq u_\alpha(y, t) \leq Ct^{\frac{1}{1+\gamma}} \quad \text{for } \alpha \text{ sufficiently large.} \tag{5.7}$$

We next estimate  $g_{1\alpha}(y, t)$  as follows.

If  $\alpha^2 t \leq N$ ,  $N$  is the constant in (5.4), then

$$g_{1\alpha}(y, t) \leq C \alpha^{\frac{2\gamma}{1+\gamma}} \leq C \frac{1}{t^{\gamma/(1+\gamma)}}.$$

If  $\alpha^2 t \geq N$ , then by (5.7)

$$\begin{aligned}
g_{1\alpha}(y, t) &\leq \frac{C}{u_\alpha^\gamma(y, t)} \\
&\leq \frac{C}{t^{\gamma/(1+\gamma)}}.
\end{aligned}$$

We conclude that

$$g_{1\alpha}(y, t) \leq \frac{C}{t^{\gamma/(1+\gamma)}} \quad \text{for any } \alpha, t. \tag{5.8}$$

From (5.6) and (5.8), we see that the RHS of (5.5) is bounded in any compact subset  $K \subset R^1 \times (0, \infty)$ . Applying  $L^p$ -estimates, we get  $|u_\alpha|_{W_p^{1,2}(K)} \leq C < \infty$ . By Sobolev imbedding

theorem, we can find a sequence  $\alpha_n \rightarrow \infty$ , such that  $u_{\alpha_n}(y, t) \rightarrow \bar{u}(y, t)$  uniformly in  $K$  as  $\alpha_n \rightarrow \infty$ , for some function  $\bar{u}$ . By a diagonal argument, we may assume that

$$u_{\alpha_n}(y, t) \rightarrow \bar{u}(y, t) \quad \text{in } R^1 \times (0, \infty). \quad (5.9)$$

By (5.7), we have

$$ct^{1/(1+\gamma)} \leq \bar{u}(y, t) \leq Ct^{1/(1+\gamma)} \quad \text{for any } t > 0.$$

Since  $u_{\alpha_n}(y, 0) = 0$ , we can represent  $u_{\alpha_n}$  in terms of the fundamental solution to get

$$u_{\alpha_n}(y, t) = \int_0^t \int_{R^1} \Gamma(y, t; z, s) g_{1\alpha_n}(z, s) dz ds + \int_0^t \int_{R^1} \Gamma(y, t; z, s) g_{2\alpha_n}(z, s) dz ds.$$

Letting  $\alpha_n \rightarrow \infty$ , we see that (by (5.6) and (5.9))

$$\bar{u}(y, t) = \int_0^t \int_{R^1} \Gamma(y, t; z, s) 2^\gamma \beta \frac{E(z)}{\bar{u}^\gamma(z, s)} dz ds; \quad (5.10)$$

i.e.

$$\begin{cases} \bar{u}_t - \bar{u}_{yy} = 2^\gamma \beta E(y) / \bar{u}^\gamma & \text{in } R^1 \times (0, \infty) \\ \bar{u}(y, 0) = 0 \\ ct^{1/(1+\gamma)} \leq \bar{u}(y, t) \leq Ct^{1/(1+\gamma)}. \end{cases} \quad (5.11)$$

**Lemma 5.3** The solution of (5.11) is unique.

*Proof.* Suppose  $u_1, u_2$  are the two solutions of (5.11). Setting  $v = u_1 - u_2$ , we have

$$v_t - v_{yy} + g(y, t)v = 0 \quad \text{in } R^1 \times (0, \infty)$$

where  $g(y, t) = 2^\gamma \beta E(y) \gamma [\lambda u_1(y, t) + (1 - \lambda)u_2(y, t)]^{\gamma-1} / u_1^\gamma(y, t) u_2^\gamma(y, t) \geq 0$  (for some  $\lambda$ ,  $0 < \lambda < 1$ ). Set  $v_\epsilon = v + \epsilon$ . From the third equation of (5.13), we see that  $v_\epsilon > 0$  in  $R^1 \times [0, \delta]$  for some  $\delta > 0$ . Since

$$v_{\epsilon t} - v_{\epsilon yy} + g(y, t)v_\epsilon = \epsilon g(y, t) \geq 0,$$

the maximum principle gives us that  $v_\epsilon > 0$  in  $R^1 \times [0, \infty)$ . Letting  $\epsilon \rightarrow 0$ , we deduce that  $v \geq 0$ . Similarly, we can prove that  $v \leq 0$ . Hence  $v \equiv 0$ .

By Lemma 5.3, we conclude that

$$u_\alpha(y, t) \longrightarrow \bar{u}(y, t) \quad \text{as } \alpha \rightarrow \infty.$$

For  $\alpha > 0$ ,  $\bar{\alpha} > 0$ , we have that

$$u_{\alpha\bar{\alpha}}(y, t) \longrightarrow \bar{u}(y, t) \quad \text{as } \alpha \rightarrow \infty.$$

On the other hand

$$u_{\alpha\bar{\alpha}}(y, t) = \frac{1}{\bar{\alpha}^{1+\gamma}} u_\alpha(\bar{\alpha}y, \bar{\alpha}^2 t) \longrightarrow \frac{1}{\bar{\alpha}^{1+\gamma}} \bar{u}(\bar{\alpha}y, \bar{\alpha}^2 t) \quad \text{as } \alpha \rightarrow \infty.$$

Therefore,

$$\bar{u}(y, t) = \frac{1}{\bar{\alpha}^{1+\gamma}} \bar{u}(\bar{\alpha}y, \bar{\alpha}^2 t).$$

Let  $\bar{\alpha} = \frac{1}{\sqrt{t}}$ , we see that  $\bar{u}(y, t)$  must have the form:

$$\bar{u}(y, t) = t^{\frac{1}{1+\gamma}} g\left(\frac{y}{\sqrt{t}}\right). \quad (5.12)$$

It is easy to see that  $g$  must satisfy:

$$g''(x) + \frac{x}{2}g'(x) = \frac{1}{1+\gamma}g - 2^\gamma \beta \frac{E(x)}{g^\gamma}. \quad (5.13)$$

Set  $f(t) = \lim_{y \rightarrow \infty} \bar{u}(y, t)$ . Then rewrite (5.10) as

$$\bar{u}(y, t) = \int_0^t \int_{\mathbb{R}^1} \frac{1}{\sqrt{2\pi}} e^{-u^2/2} 2^\gamma \beta \frac{E(y + \sqrt{2(t-s)}u)}{\bar{u}(y + \sqrt{2(t-s)}u, s)} du ds.$$

Letting  $y \rightarrow \infty$ , we see that

$$f(t) = \int_0^t \frac{2^\gamma \beta C_2}{f^\gamma(s)} ds;$$

i.e.

$$f' = \frac{2^\gamma \beta C_2}{f^\gamma}, \quad f(0) = 0.$$

Hence  $f(t) = (2^\gamma \beta C_2)^{\frac{1}{1+\gamma}} t^{\frac{1}{1+\gamma}}$ . Comparing with (5.12), we see that

$$\lim_{y \rightarrow \infty} g(y) = (2^\gamma \beta C_2)^{\frac{1}{1+\gamma}}. \quad (5.14)$$

Similarly,

$$\lim_{y \rightarrow -\infty} g(y) = (2^\gamma \beta C_1)^{\frac{1}{1+\gamma}}. \quad (5.15)$$

Since  $\bar{u}$  is unique, the solution of (5.13)-(5.15) is also unique. From the relation  $u_\alpha(0, t) \rightarrow t^{\frac{1}{1+\gamma}} g(0)$ , we see that

$$\frac{u(a, \alpha^2 t)}{\alpha^{2/(1+\gamma)}} \longrightarrow t^{\frac{1}{1+\gamma}} g(0) \quad \text{as } \alpha \rightarrow \infty.$$

In particular,

$$\frac{u(0, t)}{t^{1/(1+\gamma)}} \longrightarrow g(0) \quad \text{as } t \rightarrow \infty.$$

Using Lemma 5.2, we get that for any  $y \in [-M, M]$

$$\begin{aligned} \frac{|u(y, t) - u(0, t)|}{t^{1/(1+\gamma)}} &\leq \frac{|u_y(\lambda y, t)| |y|}{t^{1/(1+\gamma)}} \\ &\leq \frac{CM}{\sqrt{t}}. \end{aligned}$$

We conclude that for any  $M > 0$

$$\frac{u(y, t)}{t^{1/(1+\gamma)}} \longrightarrow g(0) \quad \text{uniformly in } y \in [-M, M] \text{ as } t \rightarrow \infty.$$

Using (5.3) and the fact that  $w = P + Q$ , we have established

**Theorem 5.4** If  $F$  satisfies the assumption (5.1), then for any  $M > 0$

$$\lim_{t \rightarrow \infty} \frac{P(x, t)}{t^{1/(1+\gamma)}} = \lim_{t \rightarrow \infty} \frac{Q(x, t)}{t^{1/(1+\gamma)}} = \frac{1}{2}g(0),$$

uniformly in  $x \in [-M, M]$ ; here  $g$  is the unique solution of (5.13)-(5.15).

**Acknowledgement.** I am grateful to Dr. David S. Ross from Eastman Kodak company who introduce the problem to us at an IMA industrial problem seminar; I would also thank Professor Avner Friedman for several stimulating conversations.

## References

- [1] Friedman A., *Mathematics in industrial problems, Part II*, to be published by Springer-Verlag
- [2] Ladyzenskaja O.A., Solonnikov V.A. & Ural'ceva N.N., *Linear and Quasilinear Equations of Parabolic Type*, Amer. Math. Soc.(1968).

Recent IMA Preprints

#	Author/s	Title
495	Lucas Hsu, Niky Kamran and Peter J. Olver,	Equivalence of Higher Order Lagrangians II. The Cartan Form for Particle Lagrangians
496	D.J. Kaup and Peter J. Olver,	Quantization of BiHamiltonian Systems
497	Metin Arik, Fahrünisa Neyzi, Yavuz Nutku, Peter J. Olver and John M. Verosky	Multi-Hamiltonian Structure of the Born-Infeld Equation
498	David H. Wagner,	Detonation Waves and Deflagration Waves in the One Dimensional ZND Model for High Mach Number Combustion
499	Jerrold R. Griggs and Daniel J. Kleitman,	Minimum Cutsets for an Element of a Boolean Lattice
500	Dieter Jungnickel,	On Affine Difference Sets
501	Pierre Leroux,	Reduced Matrices and q-log Concavity Properties of q-Stirling Numbers
502	A. Narain and Y. Kizilyalli,	The Flow of Pure Vapor Undergoing Film Condensation Between Parallel Plates
503	Donald A. French,	On the Convergence of Finite Element Approximations of a Relaxed Variational Problem
504	Yisong Yang,	Computation, Dimensionality, and Zero Dissipation Limit of the Ginzburg-Landau Wave Equation
505	Jürgen Sprekels,	One-Dimensional Thermomechanical Phase Transitions with Non-Convex Potentials of Ginzburg-Landau Type
506	Yisong Yang,	A Note On Nonabelian Vortices
507	Yisong Yang,	On the Abelian Higgs Models with Sources
508	Chjan. C. Lim,	Existence of Kam Tori in the Phase Space of Vortex Systems
509	John Weiss,	Bäcklund Transformations and the Painlevé Property
510	Pu Fu-cho and D.H. Sattinger,	The Yang-Baxter Equation for Integrable Systems
511	E. Bruce Pitman and David G. Schaeffer,	Instability and Ill-Posedness in Granular Flow
512	Brian A. Coomes,	Polynomial Flows on $C^n$ *
513	Bernardo Cockburn, Suchung Hou and Chi-Wang Shu,	The Runge-Kutta Local Projection Discontinuous Galerkin Finite Element Method for Conservation Laws IV: The Multidimensional Case
514	Peter J. Olver,	Invariant Theory, Equivalence Problems, and the Calculus of Variations
515	Daniel D. Joseph and Thomas S. Lundgren with an appendix by R. Jackson and D.A. Saville,	Ensemble Averaged and Mixture Theory Equations
516	P. Singh, Ph. Caussignac, A. Fortes, D.D. Joseph and T. Lundgren,	Stability of Periodic Arrays of Cylinders Across the Stream by Direct Simulation
517	Daniel D. Joseph,	Generalization of the Foscolo-Gibilaro Analysis of Dynamic Waves
518	A. Narain and D.D. Joseph,	Note on the Balance of Energy at a Phase Change Interface
519	Daniel D. Joseph,	Remarks on inertial radii, persistent normal stresses, secondary motions , and non-elastic extensional viscosities
520	D. D. Joseph,	Mathematical Problems Associated with the Elasticity of Liquids
521	Henry C. Simpson and Scott J. Spector,	Some Necessary Conditions at an Internal Boundary for Minimizers in Finite Elasticity
522	Peter Gritzmann and Victor Klee,	On the 0-1 Maximization of Positive Definite Quadratic Forms
523	Fu-Cho Pu and D.H. Sattinger,	The Yang-Baxter Equations and Differential Identities
524	Avner Friedman and Fernando Reitich,	A Hyperbolic Inverse Problem Arising in the Evolution of Combustion Aerosol
525	E.G. Kalnins, Raphael D. Levine and Willard Miller, Jr.,	Conformal Symmetries and Generalized Recurrences for Heat and Schrödinger Equations in One Spatial Dimension
526	Wang Jinghua and Gerald Warnecke,	On Entropy Consistency of Large Time Step Godunov and Glimm Schemes
527	C. Guillopé and J.C. Saut,	Existence Results for the Flow of Viscoelastic Fluids with a Differential Constitutive Law
528	H.L. Bodlaender, P. Gritzmann, V. Klee and J. Van Leeuwen	Computational Complexity of Norm-Maximization
529	Li Ta-tsien (Li Da-qian) and Yu Xin,	Life-Span of Classical Solutions to Fully Nonlinear Wave Equations
530	Jong-Shenq Guo,	A Variational Inequality Associated with a Lubrication Problem
531	Jong-Shenq Guo,	On the Semilinear Elliptic Equation $\Delta w - \frac{1}{2}y \cdot \nabla w + \lambda w - w^{-\beta} = 0$ in $R^n$
532	Andrew E. Yagle,	Inversion of the Bloch transform in magnetic resonance imaging using asymmetric two-component inverse scattering

## Recent IMA Preprints (Continued)

#	Author/s	Title
533	Bei Hu,	A Fiber Tapering Problem
534	Peter J. Olver,	Canonical Variables for BiHamiltonian Systems
535	Michael Renardy,	A Well-Posed Boundary Value Problem for Supercritical Flow of Viscoelastic Fluids of Maxwell Type
536	Michael Renardy,	Ill-Posedness Resulting from Slip As a Possible Explanation of Melt Fracture
537	Michael Renardy,	Compatibility Conditions at Corners Between Walls and Inflow Boundaries for Fluids of Maxwell Type
538	Rolf Rees,	The Spectrum of Restricted Resolvable Designs with $r = 2$
539	D. Lewis and J.C. Simo,	Nonlinear stability of rotating pseudo-rigid bodies
540	Robert Hardt and David Kinderlehrer,	Variational Principles with Linear Growth
541	San Yih Lin and Yisong Yang,	Computation of Superconductivity in Thin Films
542	A. Narain,	Pressure Driven Flow of Pure Vapor Undergoing Laminar Film Condensation Between Parallel Plates
543	P.J. Vassiliou,	On Local Equivalence for Vector Field Systems
544	Brian A. Coomes,	On Conditions Sufficient for Injectivity of Maps
545	Yanchun Zhao,	A Class of Global Smooth Solutions of the One Dimensional Gas Dynamics System
546	H. Holden, L. Holden and N.H. Risebro,	Some Qualitative Properties of $2 \times 2$ Systems of Conservation Laws of Mixed Type
547	M. Slemrod,	Dynamics of Measured Valued Solutions to a Backward-Forward Heat Equation
548	Avner Friedman and Jürgen Sprekels,	Steady States of Austenitic–Martensitic–Domains in the Ginzburg–Landau Theory of Shape Memory Alloys
549	Avner Friedman and Bei Hu,	Degenerate Hamilton–Jacobi–Bellman Equations in a Bounded Domain
550	E.G. Kalnins, Willard Miller, Jr., and M.V. Tratnik,	Families of Orthogonal and Biorthogonal Polynomials on the N-Sphere
551	Heinrich Freistühler,	On Compact Linear Degeneracy
552	Matthew Witten,	Quantifying the Concepts of Rate and Acceleration/Deceleration of Aging
553	J.P. Albert and J.L. Bona,	Total Positivity and the Stability of Internal Waves in Stratified Fluids of Finite Depth
554	Brian Coomes and Victor Zurkowski,	Linearization of Polynomial Flows and Spectra of Derivations
555	Yuriko Renardy,	A Couette-Poiseuille Flow of Two Fluids in a Channel
556	Michael Renardy,	Short wave instabilities resulting from memory slip
557	Daniel D. Joseph and Michael Renardy,	Stokes' first problem for linear viscoelastic fluids with finite memory
558	Xiayi Ding,	Superlinear Conservation Law with Viscosity
559	J.L. Ericksen,	Liquid Crystals with Variable Degree of Orientation
560	F. Robert Ore, Jr. and Xinfu Chen,	Electro-Optic Modulation in an Arbitrary Cross-Section Waveguide
561	M.V. Tratnik,	Multivariable biorthogonal continuous-discrete Wilson and Racah polynomials
562	Yisong Yang,	Existence of Solutions for a Generalized Yang-Mills Theory
563	Peter Gritzmann, Laurent Habsieger and Victor Klee,	Good and Bad Radii of Convex Polygons
564	Martin Golubitsky, Martin Krupa and Chjan. C. Lim,	Time-Reversibility and Particle Sedimentation
565	G. Yin,	Recent Progress in Parallel Stochastic Approximations
566	G. Yin,	On H-Valued SA: Finite Dimensional Approximations
567	Chien-Cheng Chang,	Accurate Evaluation of the Effect of Diffusion and Conductivity in Certain Equations
568	Chien-Cheng Chang and Ruey-Ling Chern,	The Effect of Viscous Diffusion in Discrete Vortex Dynamics for Slightly Viscous Flows
569	Li Ta-Tsien (Li Da-qian) and Zhao Yan-Chun,	Global Existence of Classical Solutions to the Typical Free Boundary Problem for General Quasilinear Hyperbolic Systems and its Applications
570	Thierry Cazenave and Fred B. Weissler,	The Structure of Solutions to the Pseudo-Conformally Invariant Nonlinear Schrödinger Equation
571	Marshall Slemrod and Athanasios E. Tzavaras,	A Limiting Viscosity Approach for the Riemann Problem in Isentropic Gas Dynamics
572	Richard D. James and Scott J. Spector,	The Formation of Filamentary Voids in Solids
573	P.J. Vassiliou,	On the Geometry of Semi-Linear Hyperbolic Partial Differential Equations in the Plane Integrable by the Method of Darboux
574	Jerome V. Moloney and Alan C. Newell,	Nonlinear Optics
575	Keti Tenenblat,	A Note on Solutions for the Intrinsic Generalized Wave and Sine-Gordon Equations
576	P. Szmolyan,	Heteroclinic Orbits in Singularly Perturbed Differential Equations
577	Wenxiong Liu,	A Parabolic System Arising In Film Development