

**HIGHER ORDER NONLINEAR
DEGENERATE PARABOLIC EQUATIONS**

By

Francisco Bernis

and

Avner Friedman

IMA Preprint Series # 474

February 1989

HIGHER ORDER NONLINEAR DEGENERATE PARABOLIC EQUATIONS*

FRANCISCO BERNIS† AND AVNER FRIEDMAN‡

Abstract. This paper is concerned with nonlinear degenerate parabolic equations of the form $u_t + (-1)^{m-1} D(f(u)D^{2m+1}u) = 0$ with $f(u) \sim |u|^n$ ($n \geq 1$) near $u = 0$ and $D = \partial/\partial x$. Under appropriate boundary conditions it is shown that there exists a weak solution u . Some of the main results of the paper are that $u \geq 0$ if $u_0 \geq 0$, and that the support of $u(\cdot, t)$ (when $u_0 \geq 0$) increases with t (for the last property we require that $n \geq 2$ and $m = 1$).

§1. Introduction. In this paper we consider higher order nonlinear degenerate parabolic equations of the form

$$(1.1) \quad \frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left(f(u) \frac{\partial^3 u}{\partial x^3} \right) = 0$$

and, more generally,

$$(1.2) \quad \frac{\partial u}{\partial t} + (-1)^{m-1} \frac{\partial}{\partial x} \left(f(u) \frac{\partial^{2m+1} u}{\partial x^{2m+1}} \right) = 0$$

where

$$(1.3) \quad f(u) = |u|^n f_0(u), \quad f_0(u) > 0$$

and n is a real number, $n \geq 1$.

Equation (1.1) arises in modeling the motion of viscous droplets spreading over a solid surface: in [3] [4] [5] and [9] the authors take $f(u) = |u|^3$ (but also $f(u) = |u|^3 + \beta|u|$ in [4], and $f(u) = |u|^3 + \beta u^2$ in [9]); further, since they assume, on physical grounds, that $u \geq 0$, they replace $|u|$ by u . Equation (1.2) with $m = 2$ and $f(u) = |u|^n$, $n = 3$ appears in recent work by King [6] [7] [8] and Tayler and King [13] in a model of oxidation of silicon in semiconductor devices. Some explicit solutions and heuristic asymptotic analysis with respect to n , for (1.2) with $f(u) = |u|^n$, is given in a recent work by Smyth and Hill [11].

In this paper we shall consider first (1.1) in a bounded interval in the x -space, with appropriate boundary conditions, and we shall prove in §§2,3 the existence of a weak solution. We next establish (in §4) the remarkable phenomena that

$$(1.4) \quad \begin{array}{l} \text{if the initial data are } \geq 0 \\ \text{then the solution is } \geq 0. \end{array}$$

*This paper is partially supported by National Science Foundation Grant DMS-86-12880 and U.S. Army Grant No. DAJA-86-C-0040.

†Departamento de Matematica Aplicada, Universidad Complutense, 28040 Madrid, Spain

‡Institute for Mathematics and its Applications, University of Minnesota, Minneapolis, Minnesota 55455

It is well known that this positivity result is false for solutions of the linear equation $u_t + \partial^4 u / \partial x^4 = 0$. In the process of proving (1.4) we establish some estimates. These estimates provide additional regularity of the weak solution and, in particular, allow us to assert that the weak solution satisfies:

$$\iint u \phi_t = \iint f(u) u_{xx} \phi_{xx} + \iint f'(u) u_x u_{xx} \phi_x$$

for any test function ϕ .

In §5 we show that if $n \geq 4$ then the support of the solution $u(\cdot, t)$ increases with t . A slightly weaker result is established in §6 in case $2 \leq n < 4$.

Finally, in §7 we extend all the results of §§2-4 to equations of the form (1.2) with $m \geq 2$.

§2. The approximating problems. In §§2-5 we study the equation

$$(2.1) \quad u_t + (f(u) u_{xxx})_x = 0 \quad \text{in} \quad Q_{T_0} \equiv \Omega \times (0, T_0)$$

where $T_0 > 0$, Ω is a bounded interval, say

$$\Omega = \{-a < x < a\}.$$

with initial conditions

$$(2.2) \quad u(x, 0) = u_0(x) \quad , \quad u_0 \in H^1(\Omega)$$

and boundary conditions

$$(2.3) \quad u_x = u_{xxx} = 0 \quad \text{on} \quad x = \pm a .$$

We assume that

$$(2.4) \quad f(u) = |u|^n f_0(u), \quad f_0 \in C^{1+\alpha}(\mathbb{R}^1), \quad f_0 > 0$$

where $\alpha \in (0, 1)$, and take

$$(2.5) \quad n > 1 ;$$

the case $n = 1$ will be considered at the end of Section 4.

Since (2.1) is degenerate at $u = 0$, we begin by approximating it by a family of non-degenerate diffusions:

$$(2.6) \quad u_t + ((f(u) + \epsilon) u_{xxx})_x = 0 \quad \text{in} \quad Q_{T_0} ,$$

where $\epsilon > 0$.

We also approximate u_0 in the $H^1(\Omega)$ -norm by $C^{4+\alpha}$ functions $u_{0\epsilon}$ satisfying (2.3), and replace (2.2) by

$$(2.7) \quad u(x, 0) = u_{0\epsilon}(x) .$$

Using the parabolic Schauder estimates [1] [2] [12] one can prove that (2.6), (2.7), (2.3) has a unique solution in a small time interval, say in Q_σ for some small $\sigma > 0$. The derivatives

$$u_t, u_x, u_{xx}, u_{xxx}, u_{xxxx}$$

are all Hölder continuous in $\overline{Q_\sigma}$. Later on we shall prove an a priori Hölder estimate for the solution u_ϵ of (2.6), (2.7), (2.3) in every domain Q_σ , independently of σ . This allows us to extend the solution u_ϵ step-by-step to all of Q_{T_0} .

We shall now assume that u_ϵ is a solution in Q_σ , for some $0 < \sigma < T_0$ and derive various estimates.

Setting $u = u_\epsilon$, we begin with

$$\begin{aligned} & \int_{\Omega} [u_x(x, t+h)^2 - u_x(x, t)^2] dx \\ &= \int_{\Omega} [u_x(x, t+h) + u_x(x, t)][u_x(x, t+h) - u_x(x, t)] dx \\ &= - \int_{\Omega} [u_{xx}(x, t+h) + u_{xx}(x, t)][u(x, t+h) - u(x, t)] dx \end{aligned}$$

since $u_x = 0$ on the boundary. Dividing by h and letting $h \rightarrow 0$ we get, for any $0 < t_1 < t_2 < \sigma$,

$$- \int_{t_1}^{t_2} \int_{\Omega} u_t u_{xx} dx dt = \frac{1}{2} \left[\int_{\Omega} u_x(x, t)^2 dx \right]_{t=t_1}^{t=t_2} .$$

Multiplying (2.6) by $u_{\epsilon, xx}$ and integrating over Q_T ($0 < T < \sigma$) and using the last identity, we get

$$(2.8) \quad \frac{1}{2} \int_{\Omega} u_{\epsilon, x}^2(x, T) dx + \int_0^T \int_{\Omega} (f(u_\epsilon) + \epsilon) u_{\epsilon, xxx}^2 dx dt = \frac{1}{2} \int_{\Omega} u_{0\epsilon, x}^2 dx .$$

Hence

$$(2.9) \quad \int_{\Omega} u_{\epsilon, x}^2(x, T) dx \leq \int_{\Omega} u_{0\epsilon, x}^2 dx .$$

Integrating (2.6) over Ω_T we also have

$$(2.10) \quad \int_{\Omega} u_{\epsilon,x}(x, T) dx = \int_{\Omega} u_{0\epsilon,x} dx .$$

Notice that

$$(2.11) \quad \int_{\Omega} u_{0\epsilon,x}^2 \leq (1 + \eta(\epsilon)) \int_{\Omega} u_{0,x}^2 \quad (\eta(\epsilon) \rightarrow 0 \text{ if } \epsilon \rightarrow 0).$$

Hence from (2.9), (2.10) we deduce, by the Poincare inequality, that

$$(2.12) \quad |u_{\epsilon}(x, t)| \leq A \quad \text{in } Q_{\sigma}$$

where A is a constant independent of ϵ, σ .

From (2.9), (2.11) and Sobolev's inequality we also deduce that

$$(2.13) \quad |u_{\epsilon}(x_2, t) - u_{\epsilon}(x_1, t)| \leq K|x_2 - x_1|^{1/2} \quad \text{in } Q_{\sigma}$$

where K a constant independent of ϵ, σ .

Setting

$$(2.14) \quad h_{\epsilon} = (f(u_{\epsilon}) + \epsilon)u_{\epsilon,xxx}$$

we see from (2.8), (2.12) that

$$(2.15) \quad |h_{\epsilon}|_{L^2(Q_{\sigma})} \leq A_1, \quad A_1 \text{ independent of } \sigma, \epsilon.$$

LEMMA 2.1. *There exists a constant M independent of σ, ϵ such that*

$$(2.16) \quad |u_{\epsilon}(x, t_2) - u_{\epsilon}(x, t_1)| \leq M|t_2 - t_1|^{1/8}$$

for all $x \in \Omega$, t_1 and t_2 in $(0, \sigma)$.

Proof. We suppose that

$$|u_{\epsilon}(x_0, t_2) - u_{\epsilon}(x_0, t_1)| > M|t_2 - t_1|^{1/8}$$

for some x_0 and t_2, t_1 and derive an upper bound for M which is independent of σ, ϵ ; for simplicity we suppose that $u_{\epsilon}(x_0, t_2) > u_{\epsilon}(x_0, t_1)$ and that $t_2 > t_1$; thus

$$(2.17) \quad u_{\epsilon}(x_0, t_2) - u_{\epsilon}(x_0, t_1) > M|t_2 - t_1|^{\beta}, \quad 0 < t_1 < t_2 < \sigma$$

where $\beta = 1/8$.

We shall use the relation

$$(2.18) \quad \iint u_\epsilon \phi_t = - \iint h_\epsilon \phi_x$$

which is valid for any "reasonable" test-function. Since $u_{\epsilon,t}$ is continuous in \overline{Q}_σ and $h_\epsilon = 0$ on the lateral boundary, we may take any ϕ such that

$$\phi \in Lip(\overline{Q}_\sigma), \quad \phi = 0 \quad \text{near } t = 0 \text{ and near } t = \sigma;$$

ϕ need not vanish on the lateral boundary. We shall construct a test function ϕ of the form

$$(2.19) \quad \phi(x, t) = \xi(x) \theta_\delta(t)$$

where ξ and θ_δ are defined as follows:

Definition of ξ .

$$\xi(x) = \xi_0 \left(\frac{x - x_0}{\frac{M^2}{16K^2} (t_2 - t_1)^{2\beta}} \right)$$

where M is from (2.17) and K from (2.13), and $\xi_0(x) = \xi_0(-x)$, $\xi_0 \in C_0^\infty$, $\xi_0(x) = 1$ if $0 \leq x < \frac{1}{2}$, $\xi_0(x) = 0$ if $x \geq 1$ and $\xi_0'(x) \leq 0$ if $x \geq 0$. Thus

$$(2.20) \quad \xi(x) = \begin{cases} 0 & \text{if } |x - x_0| \geq \frac{M^2}{16K^2} (t_2 - t_1)^{2\beta} \\ 1 & \text{if } |x - x_0| \leq \frac{1}{2} \frac{M^2}{16K^2} (t_2 - t_1)^{2\beta}. \end{cases}$$

Definition of θ_δ . We take

$$\theta_\delta(t) = \int_{-\infty}^t \theta'_\delta(s) ds$$

where

$$\theta'_\delta(t) = \begin{cases} 1/\delta & \text{if } |t - t_2| < \delta \\ -1/\delta & \text{if } |t - t_1| < \delta \\ 0 & \text{elsewhere,} \end{cases}$$

where $\delta < \frac{1}{2}(t_2 - t_1)$. Notice that θ_δ is Lipschitz continuous and $|\theta_\delta| \leq 1$; $\theta_\delta = 0$ near $t = 0$ and near $t = \sigma$, if δ is small enough.

Inserting (2.19) into (2.18) we get

$$(2.21) \quad \iint u_\epsilon \xi(x) \theta'_\delta(t) = - \iint h_\epsilon \xi'(x) \theta_\delta(t).$$

The left-hand side satisfies

$$\iint u_\epsilon \xi \theta'_\delta(t) \longrightarrow \int \xi(x) (u_\epsilon(x, t_2) - u_\epsilon(x, t_1)) dx \quad \text{as } \delta \longrightarrow 0.$$

We shall estimate the last expression from below. In view of (2.20) we only need to consider values of x such that

$$(2.22) \quad |x - x_0| \leq \frac{M^2}{16K^2} (t_2 - t_1)^{2\beta}.$$

For such values,

$$\begin{aligned} u_\epsilon(x, t_2) - u_\epsilon(x, t_1) &= [u_\epsilon(x, t_2) - u_\epsilon(x_0, t_2)] \\ &\quad + [u_\epsilon(x_0, t_2) - u_\epsilon(x_0, t_1)] + [u_\epsilon(x_0, t_1) - u_\epsilon(x, t_1)] \\ &\geq -2K|x - x_0|^{1/2} + M(t_2 - t_1)^\beta \quad \text{by (2.13), (2.17),} \\ &\geq \frac{M}{2}(t_2 - t_1)^\beta \quad \text{by (2.22).} \end{aligned}$$

Hence, if we assume that the set $\{\xi = 1\}$ is included in Ω (otherwise, very minor modifications are necessary),

$$\int \xi(x) (u_\epsilon(x, t_2) - u_\epsilon(x, t_1)) dx \geq \frac{M}{2} (t_2 - t_1)^\beta \frac{M^2}{16K^2} (t_2 - t_1)^{2\beta}.$$

On the other hand, the right-hand side of (2.21) is bounded from above by

$$\left| \iint h_\epsilon \xi'(x) \theta_\delta \right| \leq \frac{C_1}{\frac{M^2}{16K^2} (t_2 - t_1)^{2\beta}} \left(\iint h_\epsilon^2 \right)^{1/2} \frac{\sqrt{2}M}{4K} (t_2 - t_1)^\beta (t_2 - t_1 + 2\delta)^{1/2}.$$

We thus conclude, after letting $\delta \rightarrow 0$, that

$$M^3 (t_2 - t_1)^{3\beta} \leq C_2 \frac{1}{M} (t_2 - t_1)^{\frac{1}{2} - \beta}$$

where C_2 is a constant independent of ϵ, M and σ . Since $\beta = \frac{1}{8}$, we find that $M \leq C_2^{1/4}$, and the lemma follows.

From Lemma 2.1 and (2.13), (2.12), it follows that there is an upper bound on the $C_{x,t}^{\frac{1}{2}, \frac{1}{8}}$ -norm of u_ϵ in Q_σ , which is independent of σ, ϵ . This a priori bound allows us to conclude that u_ϵ can be extended step-by-step to a solution of (2.6), (2.7), (2.3) in all of Q_{T_0} , and that

$$(2.23) \quad \{u_\epsilon\} \text{ is a uniformly bounded and equi-continuous family in } \overline{Q}_{T_0}.$$

§3. Existence of weak solution. By (2.23), every sequence $\epsilon \rightarrow 0$ has a subsequence such that

$$(3.1) \quad u_\epsilon \longrightarrow u \quad \text{uniformly in } \overline{Q}_{T_0}.$$

THEOREM 3.1. *Any function u obtained as in (3.1) satisfies the following properties:*

$$(3.2) \quad \begin{aligned} u \in C(\overline{Q}_{T_0}), \quad & \text{in fact } u \text{ is uniformly Hölder continuous} \\ & \text{(exponent } \frac{1}{2}) \text{ in } x \text{ and Hölder continuous (exponent } \frac{1}{8}) \text{ in } t, \end{aligned}$$

$$(3.3) \quad u_t, u_x, u_{xx}, u_{xxx}, u_{xxxx} \text{ belong to } C(P)$$

where $P = \overline{Q}_{T_0} \setminus (\{u = 0\} \cup \{t = 0\})$, and

$$(3.4) \quad f(u)u_{xxx} \in L^2(P);$$

u satisfies (2.1) in the following sense:

$$(3.5) \quad \iint_{Q_{T_0}} u \phi_t + \iint_P f(u)u_{xxx} \phi_x = 0$$

for all $\phi \in \text{Lip}(\overline{Q}_{T_0})$, $\phi = 0$ near $t = 0$ and near $t = T_0$,

$$(3.6) \quad u(x, 0) = u_0(x), \quad x \in \overline{\Omega},$$

$$(3.7) \quad u_x(\cdot, t) \rightarrow u_{0x} \quad \text{strongly in } L^2(\Omega) \text{ as } t \rightarrow 0,$$

and

$$(3.8) \quad u \text{ satisfies (2.3) at all points of the lateral boundary where } u \neq 0.$$

Proof. The assertions (3.2), (3.6) are obvious. For ϕ as asserted in (3.5) we have

$$(3.9) \quad \iint_{Q_{T_0}} u_\epsilon \phi_t + \iint_{Q_{T_0}} f(u_\epsilon) u_{\epsilon,xxx} \phi_x + \epsilon \iint_Q u_{\epsilon,xxx} \phi_x = 0.$$

From (2.8), $\epsilon \iint u_{\epsilon,xxx}^2 \leq C$; hence, by Hölder's inequality,

$$(3.10) \quad \epsilon \iint u_{\epsilon,xxx} \phi_x \rightarrow 0 \quad \text{if } \epsilon \rightarrow 0.$$

From (2.15) it follows that, for a subsequence,

$$(3.11) \quad h_\epsilon \rightarrow h \quad \text{weakly in } L^2(Q_{T_0}).$$

Next, by regularity theory of uniformly parabolic equations and the uniform Hölder continuity of the u_ϵ we deduce that

$$(3.12) \quad u_{\epsilon,t}, u_{\epsilon,x}, u_{\epsilon,xx}, u_{\epsilon,xxx}, u_{\epsilon,xxxx} \quad \text{are uniformly convergent in any compact subset of } P.$$

It follows that

$$(3.13) \quad f(u) u_{xxx} = h \quad \text{on } P,$$

that (3.3), (3.8) hold and (recalling (3.11)) that (3.4) holds; further, for any $\delta > 0$,

$$(3.14) \quad \iint_{|u|>\delta} f(u_\epsilon) u_{\epsilon,xxx} \phi_x \rightarrow \iint_{|u|>\delta} f(u) u_{xxx} \phi_x.$$

On the other hand, if ϵ is sufficiently small, depending on δ , then by (2.4)

$$(3.15) \quad \left| \iint_{|u|\leq\delta} f(u_\epsilon) u_{\epsilon,xxx} \phi_x \right| \leq C \delta^{n/2} \left\{ \iint f(u_\epsilon) u_{\epsilon,xxx}^2 \right\}^{1/2} \leq C_1 \delta^{n/2}$$

where (2.8) has been used in the last inequality.

To prove (3.7) notice that from $u_{0\epsilon} \rightarrow u_0$ in $H^1(\Omega)$ and (2.9) we get

$$\limsup_{t \rightarrow 0} \int_{\Omega} u_x^2(x, t) dx \leq \int_{\Omega} u_{0x}^2 dx.$$

Since also

$$u_x(\cdot, t) \rightarrow u_{0x} \quad \text{weakly in } L^2(\Omega)$$

as $t \rightarrow 0$, the assertion (3.7) follows.

Taking $\epsilon \rightarrow 0$ in (3.9) and using (3.10), (3.14), (3.15) we deduce, since δ is arbitrary, that (3.5) is satisfied, and the proof of the theorem is complete.

REMARK 3.1. Since $u_t + h_x = 0$ in the sense of weak derivatives in Q_{T_0} , we have:

$$u_t \in L^2(0, T_0; H^{-1}(\Omega)).$$

REMARK 3.2. From (2.10) we deduce that

$$(3.16) \quad \int_{\Omega} u(x, t) dx = \text{const.} = \int_{\Omega} u_0(x) dx.$$

REMARK 3.3. If ϕ is smooth and $\phi_x = 0$ on the lateral boundary then from (3.9) we obtain

$$(3.17) \quad \begin{aligned} \iint_{Q_T} u_{\epsilon} \phi_t + \epsilon \iint_{Q_T} u_{\epsilon, xxx} \phi_x &= \iint_{Q_T} f(u_{\epsilon}) u_{\epsilon, xx} \phi_{xx} \\ &+ \iint_{Q_T} f'(u_{\epsilon}) u_{\epsilon, x} u_{\epsilon, xx} \phi_x \end{aligned}$$

This equation will be used in §4.

REMARK 3.4. Let $u_0 \geq 0$. In general the solution of (2.6) may take negative values. Let, for example, $\phi_{\epsilon}(x)$ be the solution of

$$\begin{aligned} (|\phi|^n + \epsilon) \phi''' &= h(x) \quad \text{where } h'(0) = 1 \\ \phi(0) = \phi'(0) &= 0, \quad \phi''(0) = 1. \end{aligned}$$

If we take $u_0 = \phi_{\epsilon}$ then $u_0 \geq 0$ near $x = 0$ and the solution u_{ϵ} of (2.6) with $f(u) = |u|^n$ satisfies

$$\frac{\partial u_{\epsilon}(0, 0)}{\partial t} = -((|\phi_{\epsilon}(x)|^n + \epsilon) \phi_{\epsilon}'''(x))' \Big|_{x=0} = -1$$

Thus $u_{\epsilon}(0, t)$ takes negative values near $t = 0$. In §4. we shall prove, however, that if $u_0 \geq 0$ then $u \geq 0$.

Definition 3.1. The solution u satisfying the assertions of Theorem 3.1 will be called a *weak solution*.

This concept is very weak; it includes stationary solution with compact support of the form

$$(x - b)^+(c - x)^+, \quad -a < b < c < a.$$

Such solutions will be excluded in §4 when we shall prove that for u_0 satisfying a certain positivity condition, the weak solution constructed in Theorem 3.1 satisfies: $u_{xx} \in L^2(Q_{T_0})$.

§4. **Nonnegative solutions.** In this section we assume, in addition to $u_0 \in H^1(\Omega)$, that

$$(4.1) \quad u_0 \geq 0$$

and prove that the weak solution $u(x, t)$ constructed in Theorem 3.1 satisfies:

$$(4.2) \quad u(x, t) \geq 0 \quad \text{a.e.}$$

Under some additional positivity assumptions on u_0 (depending on n) we shall prove additional positivity and regularity properties for u .

We introduce the functions

$$(4.3) \quad g_\epsilon(s) = - \int_s^A \frac{dr}{f(r) + \epsilon}, \quad G_\epsilon(s) = - \int_s^A g_\epsilon(r) dr$$

where $A > \max |u_\epsilon|$ for all small ϵ . Then

$$(4.4) \quad G'_\epsilon(s) = g_\epsilon(s), \quad G''_\epsilon(s) = g'_\epsilon(s) = \frac{1}{f(s) + \epsilon},$$

$$(4.5) \quad g_\epsilon(s) \leq 0, \quad G_\epsilon(s) \geq 0 \quad \text{if } s \leq A,$$

$$(4.6) \quad G_\epsilon(s) \leq G_0(s) \quad \text{for all } s$$

where $G_0 = \lim_{\epsilon \rightarrow 0} G_\epsilon$ and, for $0 < s \leq A$,

$$(4.7) \quad G_0(s) = \begin{cases} A_0 + O(s^{2-n}) & \text{if } 1 < n < 2, A_0 > 0, \\ C_2 \log \frac{1}{s} + O(1) & \text{if } n = 2, C_2 > 0, \\ C_1 s^{2-n} + R(s) & \text{if } n > 2, \end{cases}$$

$$C_1 > 0, \quad R(s) = \begin{cases} O(s^{3-n}) & \text{if } n > 3 \\ O(\log \frac{1}{s}) & \text{if } n = 3 \\ O(1) & \text{if } n < 3; \end{cases}$$

the constants are positive and depend on $f_0(0)$, where f_0 is the function appearing in (2.4).

Denote by $\tilde{G}_0(s)$ the function $G_0(s)$ corresponding to $f(u) = |u|^n$, i.e.,

$$(4.8) \quad \tilde{G}_0(s) = \begin{cases} \frac{A^{2-n}}{2-n} + \frac{s A^{1-n}}{n-1} - \frac{s^{2-n}}{(2-n)(n-1)} & \text{if } 1 < n < 2 \\ \log \frac{A}{s} + \frac{s}{A} - 1 & \text{if } n = 2 \\ \frac{s^{2-n}}{(n-2)(n-1)} - \frac{A^{2-n}}{n-2} + \frac{s A^{1-n}}{n-1} & \text{if } n > 2 \end{cases}$$

Then, for $0 \leq s \leq A$,

$$(4.9) \quad c_1 \tilde{G}_0(s) \leq G_0(s) \leq c_2 \tilde{G}_0(s), \quad c_1 > 0;$$

indeed this follows from (see (2.4))

$$k_1 |s|^n \leq f(s) \leq k_2 |s|^n \quad (k_1 > 0, 0 \leq s \leq A).$$

From (4.8), (4.9) we deduce, in particular, that

$$(4.10) \quad G_0(0) = \begin{cases} \infty & \text{if } n \geq 2 \\ A_0 & \text{if } 1 < n < 2. \end{cases}$$

If we formally multiply (2.1) by $G_0(u)$, integrate over Q_T , and use the relations

$$(4.11) \quad g_0'(s) = G_0''(s) = \frac{1}{f(s)},$$

we obtain

$$(4.12) \quad \int_{\Omega} G(u(x, T)) dx + \int_0^T \int_{\Omega} u_{xx}^2 dx dt = \int_{\Omega} G(u_0(x)) dx.$$

In order to proceed rigorously we assume, in addition to (4.1), that

$$(4.13) \quad \begin{aligned} \int_{\Omega} |\log u_0| &< \infty & \text{if } n = 2, \\ \int_{\Omega} u_0^{2-n} dx &< \infty & \text{if } 2 < n < 4, \\ u_0 &> 0 & \text{in } \bar{\Omega} & \text{if } n \geq 4. \end{aligned}$$

REMARK 4.1. If $n \geq 4$ then the conditions $u_0 \in H^1(\Omega)$ and $\int |u_0|^{2-n} < \infty$ imply that $u_0 \neq 0$ in $\bar{\Omega}$; see the proof of Theorem 4.1 (iii) below. If $2 \leq n < 4$ then (4.13) implies that

the set of zeros of u_0 must have zero measure; if $1 < n < 2$ then u_0 may have compact support.

THEOREM 4.1 (NONNEGATIVITY). *Under the assumption (4.1), (4.13),*

- (i) *if $1 < n < 2$ then the solution u is ≥ 0 in Q_{T_0} ,*
- (ii) *if $2 \leq n < 4$ then again $u \geq 0$; further, the set $\{u = 0\}$ has zero measure and, in fact,*

$$(4.14) \quad \int_{\Omega} |\log u(x, t)| dx \leq C < \infty \quad \forall t \in [0, T_0] \quad \text{if } n = 2,$$

$$(4.15) \quad \int_{\Omega} u(x, t)^{2-n} dx \leq C < \infty \quad \forall t \in [0, T_0] \quad \text{if } 2 < n < 4;$$

- (iii) *if $n \geq 4$ then $u > 0$ in \overline{Q}_{T_0} ; such a solution is unique.*

From (iii) it follows that, when $u_0 > 0$ and $n \geq 4$, the weak solution u is a classical solution and all the derivatives

$$u_t, u_x, u_{xx}, u_{xxx}, u_{xxxx}$$

are continuous in $\overline{Q}_{T_0} \setminus \{t = 0\}$.

Proof of Theorem 4.1. We can choose the smooth approximation $u_{0\epsilon}$ of u_0 such that $u_{0\epsilon} \geq u_0$. Then, from (4.1), (4.13) and (4.9) we have

$$(4.16) \quad \int_{\Omega} G_{\epsilon}(u_{0\epsilon}(x)) dx \leq C, \quad C \text{ independent of } \epsilon.$$

Multiplying equation (2.6) by $g_{\epsilon}(u_{\epsilon})$ and integrating over Q_T , $T \in (0, T_0)$, we get, after performing an integration by parts and using the boundary condition for u_{ϵ} and (4.4),

$$(4.17) \quad \int_{\Omega} G_{\epsilon}(u_{\epsilon}(x, T)) dx + \int_0^T \int_{\Omega} u_{\epsilon, xx}^2 dx dt = \int_{\Omega} G_{\epsilon}(u_{0\epsilon}(x)) dx.$$

Consequently, by (4.16),

$$(4.18) \quad \int_{\Omega} G_{\epsilon}(u_{\epsilon}(x, T)) dx \leq C, \quad G_{\epsilon}(u_{\epsilon}) \geq 0$$

and

$$(4.19) \quad \iint_{Q_{T_0}} u_{\epsilon,xx}^2 dx dt \leq C.$$

We proceed to prove that

$$(4.20) \quad u \geq 0 \quad \text{in} \quad Q_{T_0}.$$

If this is not true then there is a point $(x_0, t_0) \in Q_{T_0}$ such that $u(x_0, t_0) < 0$. Since $u_\epsilon \rightarrow u$ uniformly, there exist $\delta > 0$ and $\epsilon_0 > 0$ such that

$$u_\epsilon(x, t_0) < -\delta \quad \text{if} \quad |x - x_0| < \delta, \quad x \in \Omega, \quad \epsilon < \epsilon_0.$$

But for such x ,

$$\begin{aligned} G_\epsilon(u_\epsilon(x, t_0)) &= - \int_{u_\epsilon(x, t_0)}^A g_\epsilon(s) ds \geq - \int_{-\delta}^0 g_\epsilon(s) ds \\ &\rightarrow - \int_{-\delta}^0 g_0(s) ds \quad \text{as} \quad \epsilon \rightarrow 0, \end{aligned}$$

by the monotone convergence theorem where $g_0(s) = \lim_{\epsilon \rightarrow 0} g_\epsilon(s)$, and the integral on the right-hand side is equal to $+\infty$ for $n \geq 1$ since $g_0(s) = -\infty$ if $s < 0$, by (4.3). It follows that

$$\lim_{\epsilon \rightarrow 0} \int G_\epsilon(u_\epsilon(x, t_0)) dx = \infty,$$

a contradiction to (4.18).

Having proved (4.20), we now specialize to $n \geq 2$ and prove that, for each $T \in (0, T_0)$,

$$(4.21) \quad \text{the set } \{u(\cdot, T) = 0\} \text{ has measure zero}$$

If the assertion (4.21) is not true then for some $t_0 \in (0, T_0)$ the set $E = \{u(\cdot, t_0) = 0\}$ has positive measure. Since $u_\epsilon \rightarrow u$ uniformly, there exists a modulus of continuity $\sigma(\epsilon)$ such that

$$u_\epsilon(x, t_0) < \sigma(\epsilon) \quad \text{for all } x \in E.$$

Now, for any $x \in E$ and for any $\delta > 0$,

$$G_\epsilon(u_\epsilon(x, t_0)) \geq - \int_{\sigma(\epsilon)}^A g_\epsilon(s) ds \geq - \int_{\delta}^A g_\epsilon(s) ds \rightarrow - \int_{\delta}^A g_0(s) ds$$

if ϵ is small enough (so that $\sigma(\epsilon) < \delta$), and

$$\int_{\delta}^A g_0(s) ds \geq \begin{cases} c \delta^{2-n} & \text{if } n > 2 \quad (c > 0) \\ c \log \frac{1}{\delta} & \text{if } n = 2 \quad (c > 0). \end{cases}$$

Hence

$$\overline{\lim}_{\epsilon \rightarrow 0} \int G_{\epsilon}(u_{\epsilon}(x, t_0)) dx \geq \begin{cases} c \delta^{2-n} (\text{meas } E) \rightarrow \infty & \text{if } n > 2 \\ c \log \frac{1}{\delta} (\text{meas } E) \rightarrow \infty & \text{if } n = 2 \end{cases}$$

if $\delta \rightarrow 0$, a contradiction to (4.18).

Let $n \geq 2$. At the points (x, t) where $u(x, t) > 0$,

$$(4.22) \quad G_{\epsilon}(u_{\epsilon}(x, t)) \rightarrow G_0(u(x, t)),$$

where G_0 satisfies (4.8), (4.9). Since the set $\{u(\cdot, t) = 0\}$ has measure zero for any t , it follows that, for any t , (4.22) holds for almost all x . From (4.18) and Fatou's lemma we then deduce that

$$\int_{\Omega} G_0(u(x, t)) dx \leq C,$$

which, in view of (4.8), (4.9), yields the assertions (4.14), (4.15), for all $n \geq 2$.

In order to complete the proof of Theorem 4.1 it remains to prove (iii). If u is not positive everywhere in \overline{Q}_{T_0} then there exists a point (x_0, t_0) in Q_{T_0} such that $u(x_0, t_0) = 0$. By the Hölder continuity of u ,

$$u(x, t_0) < K|x - x_0|^{1/2}$$

and thus

$$\int_{\Omega} u(x, t_0)^{2-n} dx \geq c \int_{\Omega} |x - x_0|^{(2-n)/2} dx = \infty \quad \text{if } n \geq 4,$$

which is a contradiction to (4.15) (which was proved for all $n > 2$).

To prove uniqueness of positive solutions (if $n \geq 4$) suppose v is another positive solution. Then for any $0 < T < T_0$,

$$(4.23) \quad 0 < C_1 \leq u(x, t), v(x, t) \leq C_2 \quad \text{for all } x \in \Omega, 0 \leq t \leq T.$$

Set $w = u - v$. Subtracting the differential equations for u, v and multiplying by w_{xxx} , and then integrating over $\Omega \times (t_0, t)$ and letting $t_0 \rightarrow 0$, we obtain

$$\frac{1}{2} \int_{\Omega} w_x(x, t)^2 dx + \int_0^t \int_{\Omega} (f(u)u_{xxx} - f(v)v_{xxx})w_{xxx} = 0;$$

here we have used the fact that

$$u_x(\cdot, t) \rightarrow u_{0x} \quad \text{strongly in } L^2(\Omega) \text{ as } t \rightarrow 0,$$

and the same for v . Writing

$$f(u)u_{xxx} - f(v)v_{xxx} = f(u)w_{xxx} + (f(u) - f(v))v_{xxx}$$

and noting that

$$|f(u) - f(v)| \leq C_3|w|,$$

we get

$$\sup_{0 < \tau < t} \frac{1}{2} \int_{\Omega} w_x(x, \tau)^2 dx + C_4 \int_0^t \int_{\Omega} w_{xxx}^2 \leq C_3 \int_0^t \int_{\Omega} |w v_{xxx} w_{xxx}|.$$

Since the right-hand side is bounded by

$$\frac{1}{2} C_4 \int_0^t \int_{\Omega} w_{xxx}^2 + C_5 \int_0^t \int_{\Omega} w^2 v_{xxx}^2,$$

we get

$$(4.24) \quad \sup_{0 < \tau < t} \int_{\Omega} w_x(x, \tau)^2 dx + \int_0^t \int_{\Omega} w_{xxx}^2 \leq C_6 \int_0^t \int_{\Omega} w^2 v_{xxx}^2.$$

From (2.8) with $\epsilon \rightarrow 0$ and (4.23) we see that

$$(4.25) \quad \int_0^T \int_{\Omega} v_{xxx}^2 \quad \text{is finite.}$$

Next

$$\int_{\Omega} u(x, t) dx = \int_{\Omega} v(x, t) dx = \int_{\Omega} u_0 dx,$$

so that $\int_{\Omega} w(x, t) dx = 0$ and, by Poincaré's inequality,

$$(4.26) \quad \sup_{\Omega \times (0, t)} w^2 \leq C_7 \sup_{0 < \tau < t} \int_{\Omega} w_x^2(x, \tau) dx.$$

Using (4.25), (4.26) in (4.24) we get

$$\sup_{0 < \tau < t} \int_{\Omega} w_x(x, \tau)^2 dx \leq C_8 \left(\int_0^t \int_{\Omega} v_{xxx}^2 \right) \sup_{0 < \tau < t} \int_{\Omega} w_x(x, \tau)^2 dx$$

and then, for small t , $w_x \equiv 0$. This yields the asserted uniqueness

THEOREM 4.2. Let u_0 be as in Theorem 4.1 and let $n > 1$. Then the solution u satisfies:

$$(4.27) \quad u_x \in L^2(0, T_0; H_0^1(\Omega)),$$

and (2.1) holds in the following sense:

$$(4.28) \quad \iint_{Q_{T_0}} u \phi_t = \iint_{Q_{T_0}} f(u) u_{xx} \phi_{xx} + \iint_{Q_{T_0}} f'(u) u_x u_{xx} \phi_x$$

for all $\phi \in C^2(\overline{Q}_{T_0})$ with $\phi = 0$ near $t = 0$ and near $t = T_0$, and $\phi_x = 0$ on $\partial\Omega \times (0, T_0)$.

REMARK 4.2. (4.27) implies that $u_x = 0$ on $\partial\Omega \times (0, T_0)$ for almost all t ($u_{xxx} = 0$ on $\partial\Omega \times (0, T_0)$ where $u \neq 0$).

REMARK 4.3. Notice that in view of (4.27) and the continuity of u , all the integrals in (4.28) make sense.

Proof. The assertion (4.27) follows from (4.17). To prove (4.28) we shall let $\epsilon \rightarrow 0$ in (3.17). But first we establish:

LEMMA 4.3. As $\epsilon \rightarrow 0$

$$(4.29) \quad u_{\epsilon,x} \rightarrow u_x \quad \text{in } L^2(Q_{T_0}) \quad \text{strongly.}$$

Proof. From (4.19) it follows that

$$(4.30) \quad u_{\epsilon,x} \rightarrow u_x \quad \text{weakly in } L^2(0, T_0; H_0^1(\Omega)).$$

Next, recalling that $u_{\epsilon,t} = -h_{\epsilon,x}$, where h_ϵ is defined by (2.14) and satisfies (2.15), we have

$$u_{\epsilon,xt} = -h_{\epsilon,xx}$$

in the distribution sense, and thus

$$(4.31) \quad u_{\epsilon,xt} \quad \text{are uniformly bounded in } L^2(0, T_0; H^{-2}(\Omega)).$$

We shall now use a compactness lemma of Lions [10; p. 58]:

Let E_0, E and E_1 be reflexive Banach spaces such that $E_0 \subset E \subset E_1$, the imbedding $E_0 \rightarrow E$ is compact and the imbedding $E \rightarrow E_1$ is continuous. Assume also that $1 < p_0, p_1 < \infty$. If $\{v_k\}$ is a bounded sequence in $L^{p_0}(0, T_0; E_0)$ and $\{dv_k/dt\}$ is a bounded sequence in $L^{p_1}(0, T_0; E_1)$, then there exists a subsequence of $\{v_k\}$ which converges strongly both in $L^{p_0}(0, T_0; E)$ and in $C([0, T_0]; E_1)$.

Taking $E_0 = H_0^1(\Omega)$, $E = L^2(\Omega)$, $E_1 = H^{-2}(\Omega)$, $p_0 = p_1 = 2$ and $v_k = u_{\epsilon_k, x}$, the assertion (4.29) then follows (using (4.30), (4.31)).

Having proved Lemma 4.3, we now let $\epsilon \rightarrow 0$ in (3.17); using the uniform convergence $u_\epsilon \rightarrow u$ and (4.29), (4.30), the relation (4.28) follows.

From (4.17) we see that

$$t \rightarrow \int_{\Omega} G_\epsilon(u_\epsilon(x, t)) dx \quad \text{is monotone decreasing.}$$

Taking $\epsilon \rightarrow 0$ we get:

COROLLARY 4.4. *The function*

$$t \rightarrow \int_{\Omega} G_0(u(x, t)) dx$$

is monotone decreasing.

If, in particular, $f(u) = |u|^n$ then the function $G_0(s)$ is given by (4.8) for $s > 0$; since $\int_{\Omega} u(x, t) dx$ is constant, we conclude that

$$\int_{\Omega} u(x, t)^{2-n} dx \quad \text{increases in } t \text{ if } 1 < n < 2$$

$$\text{and decreases in } t \text{ if } n > 2,$$

$$\int_{\Omega} \log \frac{1}{u(x, t)} dx \quad \text{decreases in } t \text{ if } n = 2.$$

COROLLARY 4.5. *Let u_0 be as in Theorem 4.1 and suppose $n \geq \frac{8}{3}$. Then the set*

$$\{t \in (0, T_0); \exists x \in \bar{\Omega} \text{ with } u(x, t) = 0\}$$

has zero measure and, consequently, the boundary condition $u_{xxx} = 0$ holds for almost all t .

Proof. From (4.27) it follows that

$$(4.32) \quad u_{xx}(\cdot, t_0) \in L^2(\Omega), \quad u_x(\pm a, t_0) = 0$$

for almost all t_0 . Thus it suffices to show that (4.32) implies $u(x, t_0) \neq 0$ for all $x \in \bar{\Omega}$. Suppose $u(x_0, t_0) = 0$ for some $x_0 \in \bar{\Omega}$. Since (4.32) implies $u(\cdot, t_0) \in C^{1,1/2}$, and $u(\cdot, t_0) \geq$

0, we have that $u_x(x_0, t_0) = 0$ if $x_0 \in \text{int } \Omega$; also (by (4.32)) $u_x(x_0, t_0) = 0$ if $x_0 \in \partial\Omega$. It follows that

$$u(x, t_0) \leq C|x - x_0|^{3/2}$$

where C is a constant (depending on t_0). Consequently

$$\int u(x, t_0)^{2-n} dx \geq c \int |x - x_0|^{3(2-n)/2} dx = \infty \quad (c > 0)$$

if $n \geq 8/3$, a contradiction to (4.15).

We conclude this section by considering the case where $u_0(x) \geq 0$ without the additional condition (4.13). If we define

$$\tilde{u}_{0\delta}(x) = u_0(x) + \delta$$

and denote by $\tilde{u}_\delta(x, t)$ the solution u constructed in Theorem 3.1 for the initial data $\tilde{u}_{0\delta}$, which then satisfies all the properties asserted in Theorems 4.1, 4.2, then \tilde{u}_δ satisfies the estimates

$$\begin{aligned} \int_{\Omega} |\tilde{u}_{\delta,x}|^2 dx &\leq C, \quad |\tilde{u}_\delta| \leq A, \quad \iint_{Q_{x_0}} f(\tilde{u}_\delta) \tilde{u}_{\delta,xxx}^2 dxdt \leq C, \\ |\tilde{u}_\delta(x_1, t_1) - \tilde{u}_\delta(x_2, t_2)| &\leq K(|x_1 - x_2|^{1/2} + |t_1 - t_2|^{1/8}) \end{aligned}$$

with constants C, A, K independent of δ . Taking a subsequence

$$\tilde{u}_\delta \rightarrow u$$

we obtain:

THEOREM 4.6. *For any $u_0 \geq 0$ there exists a weak solution in the sense of Theorem 3.1 such that $u \geq 0$.*

REMARK 4.4. Most of the results of the previous sections remain true if $n = 1$, although some of the arguments require minor modifications. In particular, Lemma 2.1, Theorem 3.1 and Theorem 4.1 remain valid if $n = 1$. Furthermore, Theorem 4.2 still holds provided that we use the positive approximations of Section 6 in order to prove (4.28).

§5. Expansion of the support. In this section we continue to assume that

$$(5.1) \quad u_0 \in H^1(\Omega), \quad u_0 \geq 0$$

and assume also that $n \geq 4$. We consider the weak solution u constructed in Theorem 4.6; then $u = \lim_{\delta \rightarrow 0} \tilde{u}_\delta$ where \tilde{u}_δ is the classical positive solution of (2.1), (2.3) with initial data

$$\tilde{u}_\delta(x, 0) = u_0(x) + \delta, \quad \delta > 0.$$

THEOREM 5.1. *The support of the function $t \rightarrow u(\cdot, t)$ is increasing with t .*

Proof. Set $v = \tilde{u}_\delta$. Let $\xi(x)$ be a smooth nonnegative function such that

$$(5.2) \quad \xi'(\pm a) = 0 ,$$

$$(5.3) \quad \int_{-a}^a \xi(x) u_0^{2-n}(x) dx \leq C < \infty .$$

Introduce also the function

$$G_0(s) = \int_A^s d\tau \int_A^\tau \frac{d\tau_1}{f(\tau_1)} , \text{ as in §4,}$$

where $A > \max \tilde{u}_\delta$. Multiplying the equation (2.1) for $u = \tilde{u}_\delta = v$ by $\xi G_0'(\tilde{u}_\delta)$ and integrating over Q_T , we get

$$(5.4) \quad \int_{\Omega} \xi(x) G_0(v(x, T)) dx - \int_{\Omega} \xi(x) G_0(v_0(x)) dx - \iint_{Q_T} [v_{xxx} v_x \xi + v_{xxx} h(v) \xi'] dx dt = 0$$

where

$$h(v) = -f(v) \int_v^A \frac{ds}{f(s)} .$$

Also

$$(5.5) \quad - \iint_{Q_T} v_{xxx} v_x \xi = \iint_{Q_T} [v_{xx}^2 \xi + v_x v_{xx} \xi']$$

since $v_x = 0$ on $x = \pm a$, and

$$(5.6) \quad \iint_{Q_T} v_{xxx} h(v) \xi' = - \iint_{Q_T} [h'(v) v_x v_{xx} \xi' + h(v) v_{xx} \xi'']$$

by (5.2)

Substituting (5.5), (5.6) into (5.4) and using the relations

$$h(v) \sim C_1 v, \quad h'(v) \sim C_2 \quad \text{for } v \text{ near } 0 \quad (C_1 > 0, C_2 > 0)$$

where C_1, C_2 depend on $f_0(0)$, we deduce that, since $v < A$,

$$(5.7) \quad \int_{-a}^a \xi(x)v(x, T)^{2-n} + \iint_{Q_T} v_{xx}^2 \xi \leq C \iint_{Q_T} [|\xi' v_x v_{xx}| + |v_{xx} \xi''|] + C.$$

We now choose ξ to have the form $\xi = \zeta^s$ where ζ is a smooth nonnegative function and $s \geq 4$. Then

$$|\xi'| \leq C \zeta^{s-1}, \quad |\xi''| \leq C \zeta^{s-2}.$$

Hence

$$\begin{aligned} \iint |\xi' v_x v_{xx}| &\leq C \left\{ \iint \zeta^s v_{xx}^2 \quad \iint \zeta^{s-2} v_x^2 \right\}^{1/2} \\ &\leq C_1 \left\{ \iint \xi v_{xx}^2 \right\}^{1/2}, \end{aligned}$$

since $\iint v_x^2$ is bounded (independently of δ), and

$$\begin{aligned} \iint |v_{xx} \xi''| &\leq C \iint \zeta^{s-2} |v_{xx}| \leq C \left\{ \iint \zeta^s v_{xx}^2 \quad \iint \zeta^{s-4} \right\}^{1/2} \\ &\leq C_1 \left\{ \iint \xi v_{xx}^2 \right\}^{1/2} \end{aligned}$$

since ζ^{s-4} is bounded (recalling that $s \geq 4$). Substituting these estimates in (5.7) we conclude that

$$(5.8) \quad \int_{-a}^a \xi(x)v(x, T)^{2-n} dx \leq C'$$

where C' is a constant independent of T, δ .

Suppose $u_0(x) > 0$ in an interval $\{\lambda \leq x \leq \mu\}$ and choose $\zeta(x)$ smooth in \mathbf{R}^1 , positive in $\lambda < x < \mu$ and vanishing on $\{-\infty < x \leq \lambda\} \cup \{\mu \leq x < \infty\}$. Then the function $\xi = \zeta^4$ satisfies (5.2), (5.3) and consequently (5.8) must hold and, in particular,

$$\int_{\lambda+\epsilon}^{\mu-\epsilon} v(x, T)^{2-n} dx \leq C''(\epsilon) \quad \forall \epsilon > 0 \quad (v = \tilde{u}_\delta)$$

where $C''(\epsilon)$ is a constant independent on T and δ . Letting $\delta \rightarrow 0$ we get

$$\int_{\lambda+\epsilon}^{\mu-\epsilon} u(x, T)^{2-n} dx \leq C''(\epsilon).$$

Since $u(\cdot, T) \in C^{1/2}$ and $n \geq 4$, this inequality implies that $u(x, T) > 0$ if $\lambda + \epsilon < x < \mu - \epsilon$ (cf. the proof of Theorem 4.1 (iii)). Recalling that ϵ is arbitrary, it follows that $u(x, T) > 0$ for all x in Ω for which $u_0(x) > 0$. This is also true if $x = \pm a$, by choosing $\zeta(\pm a) > 0$, $\zeta'(\pm a) = 0$ in the above proof. Hence the support of $u(\cdot, T)$ contains the support of $u_0(\cdot)$.

Similarly one can show that if $0 < t_1 < t_2 < T_0$ then the support of $u(\cdot, t_2)$ contains the support of $u(\cdot, t_1)$.

§6. Approximation by positive u_ϵ . We can construct a weak solution using also other approximations to $f(s)$ and u_0 . In this section we shall use the approximations

$$(6.1) \quad f_\epsilon(s) = \frac{s^4 f(s)}{\epsilon f(s) + s^4},$$

$$(6.2) \quad u_{0\epsilon}(x) = u_0(x) + \epsilon^\theta \quad (0 < \theta < \frac{1}{2})$$

(this choice of θ is needed in the proof of (6.7) below) in order to show that, if $2 \leq n < 4$, the resulting weak solution, u , is such that the "weak support" of $t \rightarrow u(\cdot, t)$ is monotone increasing.

The solution u_ϵ of the approximating system satisfies

$$(6.3) \quad u_{\epsilon,t} + (f_\epsilon(u_\epsilon)u_{\epsilon,xxx})_x = 0,$$

$$(6.4) \quad u_\epsilon(x, 0) = u_{0\epsilon}(x).$$

Since $\lim_{s \rightarrow 0} \frac{f_\epsilon(s)}{s^4} = \frac{1}{\epsilon}$ if $1 < n < 4$ while $f_\epsilon(s)$ has the form (2.4) if $n \geq 4$,

$$\lim_{s \rightarrow 0} \frac{f_\epsilon(s)}{s^4} = \begin{cases} \frac{f_0(0)}{\epsilon f_0(0) + 1} & \text{if } n = 4 \\ 0 & \text{if } n > 4 \end{cases}$$

and $u_{0\epsilon}(x) > 0$, Theorem 4.1 (iii) implies that there exists a unique positive (and smooth) solution u_ϵ of (6.3), (6.4), (2.3) for all $t > 0$. Let u be any limit of a subsequence of u_ϵ , $\epsilon \rightarrow 0$. One can easily modify the arguments in §§2-5 to show that u is a weak solution satisfying all the properties derived above. Let us for instance establish (4.18), (4.19). To do this we use the functions

$$g_\epsilon(s) = - \int_s^A \frac{dr}{f_\epsilon(r)}, \quad G_\epsilon(r) = - \int_s^A g_\epsilon(r) dr$$

and establish, analogously to (4.17), that

$$(6.5) \quad \int_{\Omega} G_\epsilon(u_\epsilon(x, T)) dx + \int_0^T \int_{\Omega} u_{\epsilon,xx}^2 dx dt = \int_{\Omega} G_\epsilon(u_0(x) + \epsilon^\theta) dx.$$

We compute

$$G'_\epsilon(s) - G'(s) = \frac{1}{f_\epsilon(s)} - \frac{1}{f(s)} = \frac{\epsilon}{s^4}$$

and consequently

$$(6.6) \quad G_\epsilon(s) - G(s) = \epsilon \int_s^A \frac{r-s}{r^4} dv = \epsilon \left(\frac{1}{6s^2} + \frac{s}{3A^3} - \frac{1}{2A^2} \right).$$

It follows that

$$(6.7) \quad |G_\epsilon(u_0 + \epsilon^\theta) - G(u_0 + \epsilon^\theta)| \leq C\epsilon^{1-2\theta} \rightarrow 0 \quad \text{if } \epsilon \rightarrow 0$$

and therefore

$$(6.8) \quad \int_{\Omega} G_\epsilon(u_0 + \epsilon^\theta) dx \rightarrow \int_{\Omega} G(u_0) dx.$$

From (6.5) and (6.8) we obtain (4.18), (4.19), upon which Theorem 4.1 is based.

The fact that the u_ϵ are positive smooth functions will enable us to extend the proof of Theorem 5.1 to the case $2 \leq n < 4$:

THEOREM 6.1. *If $2 \leq n < 4$ and $u_0 \geq 0$, then any weak solution obtained by the approximations (6.1)–(6.4) has the following property:*

*if $u_0(x) > 0$ in an interval $\omega \subset \Omega$, then
 $u(\cdot, t) > 0$ a.e. on ω , for all $0 < t < T_0$.*

Proof. The proof is based on extending the ideas which occur in the proof of Theorem 5.1. Take $\zeta \in C^2(\bar{\Omega})$, $\text{supp } \zeta$ in ω , $\zeta \geq 0$, $(\zeta^4)' = 0$ on $\partial\Omega$, and let $\xi = \zeta^4$. Multiplying (6.3) by $\xi(x)g_\epsilon(u_\epsilon(x, t))$ and integrating, we obtain after several integrations by parts,

$$(6.9) \quad \int_{\Omega} \xi(x)G_\epsilon(u_\epsilon(x, T)) dx + \int_0^T \int_{\Omega} \xi u_{\epsilon,xx}^2 = R + \int_{\Omega} \xi(x)G_\epsilon(u_0(x) + \epsilon^\theta) dx,$$

where

$$R = - \iint \xi' u_{\epsilon,xx} u_{\epsilon,x} - \iint \xi' u_{\epsilon,xx} u_{\epsilon,x} (1 + f'_\epsilon(u_\epsilon)g_\epsilon(u_\epsilon)) \\ - \iint \xi'' u_{\epsilon,xx} f_\epsilon(u_\epsilon)g_\epsilon(u_\epsilon).$$

We easily estimate, if $0 < s < A$,

$$|g_\epsilon(s)| \leq K \frac{\epsilon s^n + s^4}{s^{n+3}},$$

$$f_\epsilon(s) \leq K \frac{s^{n+4}}{\epsilon s^n + s^4}.$$

Also, by explicitly computing $f'_\epsilon(s)$ we find that

$$|f'_\epsilon(s)| \leq K_2 \frac{s^{n+3}}{\epsilon s^n + s^4}.$$

Hence

$$|f_\epsilon(u_\epsilon)g_\epsilon(u_\epsilon)| \leq K_3|u_\epsilon|,$$

$$|f'_\epsilon(u_\epsilon)g_\epsilon(u_\epsilon)| \leq K_3.$$

Using these estimates we can estimate $|R|$ from above:

$$|R| \leq C \iint \zeta^2 |u_{\epsilon,xx}| \cdot \zeta |u_{\epsilon,x}| + C \iint \zeta^2 |u_{\epsilon,xx}| \cdot u_\epsilon.$$

If we estimate the right-hand side by the Schwarz inequality and use the result in (6.9), we get,

$$(6.10) \quad \int_{\Omega} \zeta^4 G_\epsilon(u_\epsilon(x, T)) dx + \int_0^T \int_{\Omega} \zeta^4 u_{\epsilon,xx}^2 \leq C \int \zeta^4 G_\epsilon(u_0 + \epsilon^\theta) dx$$

$$+ C \int_0^T \int_{\Omega} (\zeta^2 u_{\epsilon,x} + u_\epsilon^2).$$

Letting $\epsilon \rightarrow 0$ and noting (cf. (6.7), (6.8)) that

$$\int \zeta^4 G_\epsilon(u_0 + \epsilon^\theta) \rightarrow \int \zeta^4 G(u_0)$$

and that the last term on the right-hand side of (6.10) is bounded independently of ϵ , we conclude that

$$(6.11) \quad \int_{\Omega} \zeta^4 G(u_0) dx < \infty \text{ implies } \int_{\Omega} \zeta^4 G(u(\cdot, t)) < \infty$$

for all $0 < t < T_0$. But this property can be used to establish Theorem 6.1, by the same arguments used in the proof of Theorem 4.1 for $2 \leq n < 4$.

§7. **Equations of order ≥ 6 .** In this section we shall extend some of the results of §§2–4 to equation (1.2) with $m \geq 2$; $f(u)$ is assumed to satisfy (2.4). We shall take the initial condition (2.2) with

$$(7.1) \quad u_0 \in H^m(\Omega)$$

and the boundary conditions

$$(7.2) \quad Du = D^3u = D^5u = \dots = D^{2m+1}u = 0 \quad \text{on} \quad \partial\Omega \times (0, T_0)$$

where $D = \frac{\partial}{\partial x}$. We begin by introducing the approximating equations

$$(7.3) \quad u_t + (-1)^{m-1} D((f(u) + \epsilon)D^{2m+1}u) = 0$$

with the boundary conditions (7.2) and the initial conditions

$$(7.4) \quad u(x, 0) = u_{0\epsilon}(x)$$

where $u_{0\epsilon}$ are smooth ($C^{2m+2+\alpha}$) and satisfy (7.2), and $u_{0\epsilon} \rightarrow u_0$ in $H^m(\Omega)$ as $\epsilon \rightarrow 0$

Denote the solution of (7.2)–(7.4) by u_ϵ . Multiplying (7.3) by $D^{2m}u_\epsilon$ and integrating over Q_T , we get

$$(7.5) \quad \frac{1}{2} \int_{\Omega} |D^m u_\epsilon(x, T)|^2 + \iint_{Q_T} (f(u_\epsilon) + \epsilon) |D^{2m+1} u_\epsilon|^2 = \frac{1}{2} \int_{\Omega} |D^m u_{0\epsilon}(x)|^2 ;$$

hence

$$(7.6) \quad \int_{\Omega} |D^m u_\epsilon(x, T)|^2 \leq C ,$$

$$(7.7) \quad \iint_{Q_{T_0}} f(u_\epsilon) |D^{2m+1} u_\epsilon|^2 \leq C$$

where C is a constant independent of T, ϵ . By integrating (7.3) over Q_T we also have

$$(7.8) \quad \int_{\Omega} u_\epsilon(x, T) dx = \text{const.} = \int_{\Omega} u_{0\epsilon}(x) dx.$$

Using (7.6), (7.2) and (7.8) we deduce, by the Poincaré inequality, that

$$(7.9) \quad \int_{\Omega} |D^j u_\epsilon(x, T)|^2 \leq C \quad \forall \quad 0 \leq j \leq m-1$$

by the boundary conditions (7.2), $D^2u, D^4u, \dots, D^m u$ have at least one zero for each T and this is enough to apply the Poincaré inequality. By the Sobolev inequality,

$$(7.10) \quad |u_\epsilon| \leq A,$$

$$(7.11) \quad |u_\epsilon(x_1, t) - u_\epsilon(x_2, t)| \leq K|x_1 - x_2| \quad \forall x_1, x_2 \in \Omega, t \in (0, T)$$

where A, K are constants independent of ϵ ; in fact

$$\begin{aligned} D^j u_\epsilon(\cdot, t) & \text{ is Lipschitz in } x, \quad 0 \leq j \leq m-2 \\ D^{m-1} u_\epsilon(\cdot, t) & \text{ is Hölder continuous in } x \text{ (exponent } \frac{1}{2}) \end{aligned}$$

uniformly in t, ϵ .

Consider the function

$$(7.12) \quad h_\epsilon = (-1)^{m-1} (f(u_\epsilon) + \epsilon) D^{2m+1} u_\epsilon.$$

By (7.7) and (7.10) it follows that

$$(7.13) \quad \iint_{Q_{T_0}} |h_\epsilon|^2 \leq C, \quad C \text{ independent of } \epsilon.$$

Using this fact and the relation

$$(7.14) \quad \frac{\partial}{\partial t} u_\epsilon + D h_\epsilon = 0$$

we can now repeat the argument of Lemma 2.1. Using however (7.11) (instead of (2.13)) we get

$$(7.15) \quad |u_\epsilon(x, t_1) - u_\epsilon(x, t_2)| \leq M|t_1 - t_2|^{1/5}.$$

We can now proceed as in §3 and establish:

THEOREM 7.1. *There exists a function u , a uniform limit of the u_ϵ (for a sequence $\epsilon \rightarrow 0$) in Q_{T_0} , such that u is Lipschitz continuous in x and Hölder continuous (exponent $1/5$) in t ,*

$$(7.16) \quad \int_{\Omega} |D^j u(x, t)| dx \leq C \quad \forall \quad 1 \leq j \leq m, \quad 0 < t < T_0,$$

$$(7.17) \quad f(u) D^{2m+1} u \in L^2(P) \quad \text{where } P = \overline{Q_{T_0}} \setminus (\{u = 0\} \cup \{t = 0\});$$

u is a classical solution of (1.2) in P , and

$$(7.18) \quad \iint_{Q_{T_0}} u \phi_t + \iint_P f(u) D^{2m+1} u \cdot \phi_x = 0$$

for any $\phi \in Lip(\overline{Q}_{T_0})$, $\phi = 0$ near $t = 0$ and near $t = T_0$; further,

$$u(x, 0) = u_0(x), \quad x \in \Omega$$

and

u satisfies the boundary conditions (7.2)

at all points of the lateral boundary where $u \neq 0$.

REMARK 7.1: The argument of Lemma 2.1 also gives that

$D^j u$ is Hölder continuous in \overline{Q}_{T_0} , for all $0 \leq j \leq m - 1$.

In fact, if $m \geq 1$,

$$D^{m-1} u \text{ belongs to } C_{x,t}^{\frac{1}{2}, \frac{1}{4(m+1)}}$$

and if $m \geq 2$ then

$$D^j u \text{ belongs to } C_{x,t}^{1, \frac{1}{5+2j}} \text{ for } 0 \leq j \leq m - 2.$$

We next consider the case where $u_0 \geq 0$ and

$$(7.19) \quad \begin{aligned} \int_{\Omega} |\log u_0| &< \infty && \text{if } n = 2, \\ \int_{\Omega} u_0^{2-n} dx &< \infty && \text{if } 2 < n < n^*, \\ u_0(x) &> 0 \text{ in } \overline{\Omega} && \text{if } n \geq n^* \end{aligned}$$

where $n^* = 8/3$ if $m = 2$ and $n^* = 5/2$ if $m \geq 3$.

REMARK 7.2. If $u_0 \geq 0$ and $\int_{\Omega} u_0^{2-n} dx < \infty$, $n \geq n^*$ then $u_0(x)$ must be strictly positive (cf. Remark 4.1). Indeed if $m = 2$ then $u_{0,xx} \in L^2(\Omega)$ and we argue as in the proof of Corollary 4.5; if $m \geq 3$ then $u_{0,xx}$ is continuous and thus if $u_0(x_0) = 0$ then $u_{0,x}(x_0) = 0$ and $u_0(x) \leq C|x - x_0|^2$, which implies

$$\int_{\Omega} u_0(x)^{2-n} dx \geq c \int_{\Omega} |x - x_0|^{2(2-n)} dx = \infty \quad \text{if } n \geq \frac{5}{2} \quad (c > 0),$$

a contradiction.

THEOREM 7.2. *If $u_0 \geq 0$ and (6.19) holds then $u \geq 0$ and, in fact, all the assertions of Theorem 4.1 hold; further*

$$(7.20) \quad u(x, t) > 0 \quad \text{in } \overline{Q_{T_0}} \quad \text{if } n \geq n^* .$$

Proof. Multiplying (7.3) by $g_\epsilon(u_\epsilon)$ (as in §4) and integrating over Q_T we easily get after some integrations by parts,

$$(7.21) \quad \int_{\Omega} G_\epsilon(u_\epsilon(x, T)) + \iint_{Q_T} |D^{m+1}u_\epsilon|^2 = \int_{\Omega} G_\epsilon(u_{0\epsilon}(x)) .$$

This relation allows us to deduce all the assertions of Theorem 4.1. Finally, if $n \geq n^*$, $u(x, t)$ must be strictly positive in $\overline{Q_{T_0}}$; for, if $u(x_0, t_0) = 0$, then we get (cf. Remark 7.2)

$$\int_{\Omega} u^{2-n}(x, t_0) dx = \infty$$

which is a contradiction

Recall that for $n \geq n^*$ the solution u is positive and classical.

For any $u_0 \geq 0$ we can construct (by Theorem 7.2) solutions $u = \tilde{u}_\delta$ corresponding to the initial data $u_0(x) + \delta$. Taking $\delta \rightarrow 0$ we obtain a limiting function $u = \lim \tilde{u}_\delta$ which is ≥ 0 ; u is a weak solution in the sense of Theorem 7.1. Thus, for any $u_0 \geq 0$ there exists a weak solution which is ≥ 0 (cf. Theorem 4.6).

We summarize:

THEOREM 7.3. *For any $u_0 \in H^m(\Omega)$, $u_0 \geq 0$ there exists a weak solution (in the sense of Theorem 7.1) which is ≥ 0 .*

Acknowledgement. The first author wishes to thank Professor M. Bertsch for useful discussions.

REFERENCES

- [1] S.D. EIDELMAN, *Parabolic systems*, North-Holland, Amsterdam, 1969.
- [2] A. FRIEDMAN, *Interior estimates for parabolic systems of partial differential equations*, J. Math. and Mech., 7 (1958), 393-418.
- [3] H.P. GREENSPAN, *On the motion of a small viscous droplet that wets a surface*, J. Fluid Mech., 84 (1978), 125-143.
- [4] H.P. GREENSPAN AND B.M. MCCAY, *On the wetting of a surface by a very viscous fluid*, Stud. Appl. Math., 64 (1981), 95-112.

- [5] L.M. HOCKING, *Sliding and spreading of thin two-dimensional drops*, Quart. J. Mech. and Appl. Math., 34 (1981), 37-55.
- [6] J.R. KING, *Thesis*, Oxford University, 1986.
- [7] J.R. KING, *The isolation oxidation of silicon*, SIAM J. Appl. Math., to appear.
- [8] J.R. KING, *The isolation oxidation of silicon: the reaction-controlled case*, SIAM J. Appl. Math., to appear.
- [9] A.A. LACEY, *The motion with slip of a thin viscous droplet over a solid surface*, Stud. Appl. Math., 67 (1982), 217-230.
- [10] J.L. LIONS, *Quelques méthodes de résolution des problèmes aux limites non linéaires*, Dunod, Paris, 1969.
- [11] N.F. SMITH AND J.M. HILL, *High order nonlinear diffusion*, IMA J. Appl. Math., 40 (1988), 73-86.
- [12] V.A. SOLONNIKOV, *On boundary value problems for linear parabolic systems of differential equations of general form*, Trudy Mat. Inst. Steklov, 83 (1965), 3-163. English translation: Proc. Steklov Inst. Math. 83 (1965), 1-184.
- [13] A.B. TAYLER AND J.R. KING, *Free boundaries in semi-conductor fabrication*, International Colloquium on Free Boundary Problems, Irsee, Bavaria (West Germany), 1987. To appear in Pitman, Research Notes in Mathematics.