

A STEFAN PROBLEM FOR A REACTION-DIFFUSION SYSTEM

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Abstract. The paper deals with a Stefan problem for a system of three weakly coupled semilinear parabolic equations. The system describes dissolution of a spherical particle in solution. The dissolved species A reacts chemically with species B already in the solution, thereby forming species C . Species C diffuses in the solution and some of it adsorbs to the particle's boundary and gradually shuts down the dissolution. It is shown that the mathematical model has a unique solution with finite shut-down time. When the reaction rate K increases to infinity, the limit model should exhibit phase separation between A and B and it thus has two free boundaries: the particle's boundary, and the $A - B$ interface. It is proved, in the case in which A and B diffuse at the same rate, that the solution with finite K converges to the solution of the limit problem, and that the A phase in the limit problem disappears in finite time.

§1. The model. Consider a solid spherical particle composed of chemical A with uniform concentration A^* . The particle is in a solution of chemical B . As the particle dissolves, the A that enters the solution reacts with B to form chemical C . Then C diffuses in the solution and some of it reaches the solid particle and adsorbs to its surface. The presence of the adsorbed C inhibits the dissolution, and ultimately shuts it down entirely.

Assuming radially symmetric data and radially symmetric functions A, B, C , we denote by $r = R(t)$ the radius of the solid sphere at time t . Then the equations

$$(1.1) \quad \frac{\partial A}{\partial t} = D_A \Delta A - K A B ,$$

$$(1.2) \quad \frac{\partial B}{\partial t} = D_B \Delta B - K A B ,$$

$$(1.3) \quad \frac{\partial C}{\partial t} = D_C \Delta C + K A B ,$$

hold in $\{r > R(t)\}$, where K is the reaction rate and D_A, D_B, D_C are the diffusion coefficients. These equations indicate that A and B are lost in a second-order reaction in which C is formed, and all three species diffuse. In the standard mass-action model of chemical kinetics, the concentrations are all expressed in moles/liter, and the coefficient K , the second-order reaction rate, is expressed in liters/(mole-sec). Then $K A B$ is the number of moles per liter per second that undergo reaction; in our case, A and B are consumed, C is created, the same number of moles of A and B are lost, and this number of moles of C is created. A nice reference for this material is the book by Erdi and Toth [4].

Next,

$$(1.4) \quad \frac{dR}{dt} = \alpha \frac{\partial A}{\partial r} \quad \text{on} \quad r = R(t)$$

where α is a positive constant, i.e., the rate at which the radius of the particle decreases is proportional to the flux of species A away from the particle. We also have

$$(1.5) \quad \frac{\partial B}{\partial r} = 0 \quad \text{on} \quad r = R(t) ,$$

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i.e., there is no flux of B through the particle's surface and B does not undergo any surface reaction.

The adsorption of C to the surface is proportional to the local saturation; it is given by an empirical law $D_C \partial C / \partial r = \gamma C^n$ for some positive constants γ, n (see [13, pp. 104–105]); for definiteness we take $n = 4$, that is,

$$(1.6) \quad D_C \frac{\partial C}{\partial r} = \gamma C^4 .$$

However, all the results of this paper remain valid with minor changes if we replace γC^4 by any other monotone increasing function $f(C)$ with $f(0) = 0$, $f(C) > 0$ for $C > 0$.

The boundary conditions at $r = \infty$ are

$$(1.7) \quad A(\infty, t) = 0, \quad B(\infty, t) = B^* , \quad C(\infty, t) = 0 ,$$

where B^* is a positive constant.

We now impose initial conditions. First,

$$(1.8) \quad R(0) = R_0 > 0 .$$

Next we assume that

$$(1.9) \quad A(r, 0) = A_0(r) , \quad B(r, 0) = B_0(r) , \quad C(r, 0) = C_0(r)$$

for $r > R_0$ where A_0, B_0, C_0 are approximately 0, $B^*, 0$ (i.e., initially mostly only B is present in the solution and its concentration is nearly uniform). We also assume that the initial conditions are smooth and fit smoothly with the boundary conditions:

$$(1.10) \quad \begin{cases} A_0 \in C^{2+\nu}[R_0, \infty), & A_0'(r) \leq 0, & A_0(R_0) = A^* , \\ A_0(r) = 0 & \text{if } r > R_0 + \delta_1 \text{ for some } \delta_1 > 0 ; \end{cases}$$

$$(1.11) \quad \begin{cases} B_0 \in C^{2+\nu}[R_0, \infty), & B_0(r) \geq 0, & B_0'(r) \geq 0 , \\ B_0(r) = B^* & \text{if } r \geq R_0 + \delta_2 \text{ for some } \delta_2 > 0 ; \end{cases}$$

$$(1.12) \quad \begin{cases} C_0(r) \in C^2[R_0, \infty), & C_0(r) \geq 0 , \\ C_0(r) = 0 & \text{if } r > R_0 + \delta_3 \text{ for some } \delta_3 > 0, \\ D_C \frac{\partial C_0}{\partial r} = \gamma C_0^4 & \text{at } r = R_0 , \end{cases}$$

for some $0 < \nu < 1$.

We finally need to determine the boundary condition for A at the particle's surface. The flux of A from the particle depends on the amount of C that is adsorbed to the surface. On a portion of $\{r = R\}$ where there is no adsorbed C , $A = A^*$, the saturation concentration of A ; local thermodynamic equilibrium is established instantaneously. On portion which is

fully covered by C , $\partial A/\partial r = 0$, i.e., the dissolution shuts down. This is actually a microscopic statement, which we shall now “average.” We shall use the “weighted average”

$$(1.13) \quad \zeta(t) = \frac{\beta \int_0^t R^2(s) D_C \frac{\partial C}{\partial r} \Big|_{r=R(s)} ds + \delta}{R^2(t)}$$

where β is a positive empirical parameter, and δ is a small positive parameter such that $\delta/R_0^2 < 1$. We then impose the boundary condition

$$(1.14) \quad -\zeta_0(t) D_A \frac{\partial A}{\partial r} + (1 - \zeta_0(t))(A - A^*) = 0 \quad \text{on} \quad r = R(t)$$

where

$$(1.15) \quad \zeta_0(t) = \min\{\zeta(t), 1\} .$$

Thus the dissolution shuts down as soon as $\zeta(t)$ becomes equal to 1. This boundary condition has the basic properties demanded by the physical problem; it reduces to the Dirichlet condition in the absence of adsorbed C , it reduces to the Neumann condition when the surface is covered, and it makes a continuous monotone transition between these two conditions as a function of the fraction of surface area that is covered.

REMARK 1.1. The parameter δ in (1.13) ensures that $\zeta(0) > 0$ and therefore the boundary condition in (1.14) does not degenerate at $t = 0$. All the results of this paper, however, except uniqueness, extend to the case $\delta = 0$ by simply going to the limit with $\delta \rightarrow 0$. If $\delta = 0$, the solution is not smooth at $(R_0, 0)$ and our proof of uniqueness (for the case $\delta > 0$) does not carry through.

For additional information on the model see [9; Chap. 18].

§2. The main results.

DEFINITION 2.1. We refer to the system (1.1)–(1.15) as Problem (P) . By a solution to Problem (P) we mean (A, B, C, R) satisfying (1.1)–(1.15) in the classical case; in particular, $R(t)$ is continuously differentiable for all $t \geq 0$.

DEFINITION 2.2. Suppose T^* is such that

$$\zeta(t) < 1 \quad \text{if} \quad t < T^* , \quad \zeta(t) \geq 1 \quad \text{if} \quad t \geq T^* .$$

Then we call T^* the *shut-down time*.

Note that (1.4) and (1.14) reduce to

$$(2.1) \quad R(t) = R(T^*), \quad A_r(R(T^*), t) = 0 \quad \text{for} \quad t > T^* .$$

THEOREM 2.1. *There exists a unique solution of (P) , and it has the following properties: (i) $R(t) > 0$, $R'(t) \leq 0$ and $\zeta'(t) > 0$ for all $t > 0$; (ii) it has a finite shut-down time T^* and $R'(t) < 0$ if $0 < t < T^*$, and (iii) R and ζ belong to $C^{1+\mu}[0, \infty) \cap C^\infty(0, T^*)$ for any $0 < \mu < 1$.*

We are interested in the case of fast reaction, that is, $K \gg 1$. This motivates the study of the solution

$$(A_K, B_K, C_K, R_K) \quad \text{of} \quad (P)$$

as $K \rightarrow \infty$. It can be shown that

$$\int_0^T \int_{R_K(t)}^{\infty} A_K B_K \leq \frac{\text{const.}}{K} ,$$

and that

$$\frac{\partial}{\partial r} A_K \leq 0, \quad \frac{\partial}{\partial r} B_K \geq 0 .$$

It follows, formally, that the limits

$$(2.2) \quad A = \lim A_K, \quad B = \lim B_K, \quad C = \lim C_K ,$$

$$R = \lim R_K \quad \text{and} \quad \zeta = \lim \zeta_K$$

are such that $AB = 0$, i.e.,

$$(2.3) \quad \begin{aligned} A(r, t) &> 0 \quad \text{if} \quad R(t) < r < S(t) , \\ &= 0 \quad \text{if} \quad r > S(t) , \end{aligned}$$

$$(2.4) \quad \begin{aligned} B(r, t) &> 0 \quad \text{if} \quad r > S(t) \\ &= 0 \quad \text{if} \quad R(t) < r < S(t) \end{aligned}$$

for some function $S(t)$; furthermore,

$$(2.5) \quad \frac{\partial A}{\partial t} = D_A \Delta A \quad \text{if} \quad R(t) < r < S(t), \quad t > 0 ,$$

$$(2.6) \quad \frac{\partial B}{\partial t} = D_B \Delta B \quad \text{if} \quad S(t) < r < \infty, \quad t > 0 ,$$

$$(2.7) \quad \frac{\partial C}{\partial t} = D_C \Delta C - \left(D_A \frac{\partial A}{\partial r} \Big|_{r=S(t)} \right) \delta(r - S(t)) \quad \text{if} \quad R(t) < r < \infty, \quad t > 0 ,$$

$$(2.8) \quad \frac{dR}{dt} = \alpha \frac{\partial A}{\partial r} \Big|_{r=R(t)} ,$$

$$(2.9) \quad -\zeta_0 D_A \frac{\partial A}{\partial r} + (1 - \zeta_0)(A - A^*) = 0 \quad \text{at} \quad r = R(t)$$

where ζ_0 is defined by (1.13), (1.15),

$$(2.10) \quad D_C \frac{\partial C}{\partial r} = \gamma C^4 \quad \text{on} \quad r = R(t) ,$$

$$(2.11) \quad A(S(t), t) = 0 ,$$

$$(2.12) \quad B(S(t), t) = 0 ,$$

$$(2.13) \quad D_A \frac{\partial A}{\partial r} = -D_B \frac{\partial B}{\partial r} \quad \text{on} \quad r = S(t) ;$$

since $A+B \rightarrow C$, the generation of C which occurs at $r = S(t)$ is at the same rate as $D_A \partial A / \partial r$, or $-D_B \partial B / \partial r$, which explains both (2.13) and the source term in (2.7).

We finally have the initial conditions:

$$(2.14) \quad \begin{aligned} A(r, 0) &= A_0(r) \quad \text{if} \quad R(0) < r < S(0) , \\ B(r, 0) &= B_0(r) \quad \text{if} \quad S(0) < r < \infty , \\ C(r, 0) &= C_0(r) \quad \text{if} \quad R(0) < r < \infty , \end{aligned}$$

$$(2.15) \quad S(0) = r_0 \quad \text{where } r_0 \text{ is such that } A_0(r_0) = B_0(r_0)$$

(r_0 is uniquely determined), and the conditions at $r = \infty$:

$$(2.16) \quad B(\infty, t) = B^*, \quad C(\infty, t) = 0 .$$

DEFINITION 2.3. We shall denote the problem (P) by (P_K) and refer to the system (2.5)–(2.16) as Problem (P_∞) . By a solution to Problem (P_∞) we mean (A, B, C, R, S, T_f) such that all the equations are satisfied in the classical case for $0 < t < T_f$, and

$$\begin{aligned} R(t) < S(t) < \infty & \quad \text{if} \quad t < T_f , \\ S(t) - R(t) & \rightarrow 0 \quad \text{if} \quad t \rightarrow T_f ; \end{aligned}$$

in particular, $R(t)$ and $S(t)$ are continuously differentiable for $0 < t < T_f$, and continuous for $0 \leq t \leq T_f$, and $A_r(r, t)$ is continuous for $R(t) \leq r \leq S(t)$.

The curve $r = S(t)$ is the interface between the separated phases A and B , and T_f is the *final time*, i.e., the time at which phase A has totally disappeared.

THEOREM 2.2. Assume that $D_A = D_B$ and $C_0(r) \not\equiv 0$. Then there exists a unique solution to Problem (P_∞) , and it has the following properties: (i) $R(t) > 0$, $R'(t) \leq 0$ and $\zeta'(t) > 0$ for all $0 < t < T_f$, and $R'(t) < 0$ as long as $\zeta(t) < 1$; (ii) $T_f < \infty$, and (iii) R, S and ζ belong to $C^0[0, T_f] \cap C^{1+\mu}[0, T_f] \cap C^\infty(0, T_f)$.

THEOREM 2.3. Assume that $D_A = D_B$. Then, as $K \rightarrow \infty$, the limits in (2.2) exist where (A, B, C, R, S, T_f) is the solution to Problem (P_∞) ; the convergence of A_K, B_K, C_K is uniform in any compact subset of

$$(2.17) \quad \{R(t) < r < S(t), 0 \leq t < T_f\} \cup \{S(t) < r < \infty, 0 \leq t < T_f\} ,$$

and the convergence of R_K and ζ_K is uniform for $0 \leq t < T_f$.

Theorem 2.1 is proved in Sections 3–6, and Theorems 2.2, 2.3 are proved in Sections 7–10.

REMARK 1. Reaction-diffusion systems of the form

$$(2.18) \quad A_t = D_A \Delta A - KAB, \quad B_t = D_B \Delta B - KAB$$

with $K \rightarrow \infty$ have been studied in [1] [12]. Evans [5] considered (2.18) in a fixed cylinder $\Omega \times (0, T)$ under the assumptions

$$(2.19) \quad \frac{\partial A}{\partial n} = 0 \quad \text{on} \quad \partial\Omega, \quad \frac{\partial B}{\partial n} = 0 \quad \text{on} \quad \partial\Omega, \quad A_0 B_0 = 0$$

where A_0, B_0 are the nonnegative initial values for A and B respectively. He proved that, as $K \rightarrow \infty$, $A \rightarrow u^+$ and $B \rightarrow u^-$ where u is the solution of

$$u_t = \operatorname{div}(a(u)\nabla u) = 0 \quad \text{in} \quad \Omega \times (0, T)$$

with $\partial u / \partial n = 0$ on $\partial\Omega$ and $u|_{t=0} = A_0 - B_0$; here

$$a(u) = D_A \quad \text{if} \quad u > 0, \quad a(u) = D_B \quad \text{if} \quad u < 0.$$

(Uniqueness for the limit problem was established in [2].) His proof relies heavily on the assumptions in (2.19) and does not seem to extend to the present case where we have a moving boundary $r = R(t)$ and different boundary conditions than in (2.19).

REMARK 2. For the general study of the Stefan problem in n -dimensions we refer to [8] [10] and the references therein. In the case of radially symmetric solutions, existence, uniqueness and regularity have been established by several methods (see [6] [7] and the references therein). In the standard Stefan problem one assumes that A vanishes on the free boundary. The condition (1.14) is called a “kinetic” condition. A Stefan problem for one heat equation in one-dimension with kinetic condition, was studied by Visintin [14] and Xie [15].

§3. Local existence for (P). In this section we prove:

THEOREM 3.1. *There exists a solution (A, B, C, R) of Problem (P) for $0 < t < T$, where T is some small positive number.*

Proof. The proof is based on a fixed-point argument. Set, for any $N_0 > 0$,

$$K_R = \{R(t) \in C^{0,1}[0, T], R(0) = R_0, -M \leq \dot{R}(t) \leq 0 \text{ a.e.}\},$$

$$K_\zeta = \{\zeta(t) \in C^{0,1}[0, T], \frac{\delta}{R_0^2} \leq \zeta(t) \leq N_0, 0 \leq \dot{\zeta}(t) \leq N \text{ a.e.}\}$$

where

$$M = \frac{\alpha}{D_A} \frac{A^* R_0^2}{\delta}$$

and N is a positive constant to be determined. We endow K_R and K_ζ with the $C^0[0, T]$ -norm; then $K_R \times K_\zeta$ is a compact set in $C^0[0, T] \times C^0[0, T]$.

For each $(R(t), \zeta(t)) \in K_R \times K_\zeta$ there exists a unique solution $(A(r, t), B(r, t))$ of (1.1), (1.2), (1.5), (1.14), (1.7) with the initial conditions as in (1.9); since the parabolic system is weakly coupled, such a solution exists for any given time T . By the maximum principle,

$$(3.1) \quad 0 \leq A(r, t) \leq A^*, \quad 0 \leq B(r, t) \leq B^*.$$

We next prove that

$$(3.2) \quad A_r(r, t) \leq 0, \quad B_r(r, t) \leq 0.$$

If we differentiate (1.1), (1.2) with respect to r , we get a coupled system of parabolic equations

$$\begin{aligned}\frac{\partial}{\partial t}A_r - \mathcal{L}A_r &= -KAB_r, \\ \frac{\partial}{\partial t}B_r - \mathcal{L}B_r &= -KBA_r\end{aligned}$$

where \mathcal{L} is an elliptic operator. On $t = 0$ and on the boundary $r = R(t)$ we have $A_r \leq 0$, $B_r \geq 0$. We approximate A_r, B_r by solutions $A_r^\varepsilon, B_r^\varepsilon$ satisfying the same parabolic system, with initial and boundary conditions given by

$$A_r^\varepsilon = A_r - \varepsilon, \quad B_r^\varepsilon = B_r + \varepsilon.$$

Then $A_r^\varepsilon < 0$, $B_r^\varepsilon > 0$ for $R(t) \leq r < \infty, 0 \leq t \leq T$. Indeed, otherwise there is a smallest t_0 such that $A_r^\varepsilon \leq 0$, $B_r^\varepsilon \geq 0$ for $R(t) \leq r < \infty, 0 \leq t \leq t_0$, and $A_r^\varepsilon = 0$ or $B_r^\varepsilon = 0$ at some point (r_0, t_0) . This is a contradiction to the strong maximum principle applied to A_r^ε or to B_r^ε .

If we now let $\varepsilon \rightarrow 0$, we obtain the assertion (3.2).

Motivated by (1.4), (1.14), we now define

$$(3.3) \quad \bar{R}(t) = R_0 + \frac{\alpha}{D_A} \int_0^t \frac{1 - \zeta_0(s)}{\zeta_0(s)} [A(R(s), s) - A^*] ds.$$

Next we consider the parabolic equation (1.3) in $\bar{R}(t) < r < \infty, 0 < t < T$ with boundary conditions (1.6) and $C(\infty, t) = 0$, and with initial conditions $C(r, 0) = C_0(r)$. Since the functions $KA^*B^*t + \hat{C}$ ($\hat{C} = \sup C_0$) and 0 are supersolution and sub-solution, respectively, the existence of a solution can be established by a fixed-point argument (cf. [6; Chap. 7, Sec. 5]). Uniqueness follows by a comparison principle [6].

We now define

$$(3.4) \quad \bar{\zeta}(t) = \frac{\beta\gamma \int_0^t \bar{R}^2(s) C^4(\bar{R}(s), s) + \delta}{\bar{R}^2(t)}$$

and consider the mapping W :

$$W(R(t), \zeta(t)) = (\bar{R}(t), \bar{\zeta}(t)).$$

If we show that W has a fixed point in $K_R \times K_\zeta$, then this yields a solution to Problem (P).

LEMMA 3.2. *W maps $K_R \times K_\zeta$ into itself.*

Proof. From (3.3), (3.4) we get

$$(3.5) \quad \frac{d}{dt} \bar{R}(t) = \frac{\alpha}{D_A} \frac{1 - \zeta_0(t)}{\zeta_0(t)} (A - A^*),$$

$$(3.6) \quad \frac{d}{dt} \bar{\zeta}(t) = \beta\gamma C^4(\bar{R}(t), t) - \frac{2}{\bar{R}(t)} \bar{\zeta}(t) \frac{d\bar{R}(t)}{dt}.$$

From (3.5) we see that

$$(3.7) \quad -M \leq \frac{d\bar{R}(t)}{dt} \leq 0$$

Since $C \leq KA^*B^*t + \widehat{C}$ we find from (3.4) that $\bar{\zeta}(t) \leq N_0$ if T is small enough. From (3.6) and (3.7) it then follows that

$$0 < \frac{d}{dt}\bar{\zeta}(t) \leq \widehat{N}$$

provided \widehat{N} is a sufficiently large number independent of N (see the definition of K_ζ). Hence, choosing $N = \widehat{N}$ we see that $(\bar{R}, \bar{\zeta})$ belongs to $K_R \times K_\zeta$.

LEMMA 3.3. *W is continuous (when $K_R \times K_\zeta$ is endowed with the $C^0[0, T] \times C^0[0, T]$ norm).*

Proof. If we use the transformation $\hat{r} = r - R(t)$ in order to flatten the boundary $r = R(t)$, we get a new parabolic equation for A where a new term $\dot{R}A_r$ is added to the heat operator. On the lateral boundary $\hat{r} = 0$

$$-A_r + a(t)A = b(t)$$

where a, b are Lipschitz continuous and their Lipschitz constants are bounded independently of (R, ζ) in $K_R \times K_\zeta$. We can therefore apply L^p -estimates [10] to deduce that

$$(3.8) \quad \|A_r\|_{L^p(\Omega_T)} + \|A_{rr}\|_{L^p(\Omega_T)} + \|A_t\|_{L^p(\Omega_T)} \leq C_1$$

for any $1 < p < \infty$, where

$$\Omega_T = \{R(t) \leq r < \infty, 0 \leq t \leq T\}$$

and C_1 is a constant depending on p but not on (R, ζ) and T . By Sobolev's imbedding [11] we then have the Hölder estimate

$$(3.9) \quad \|A_r\|_{C^\mu(\Omega_T)} \leq C_2$$

for some $0 < \mu < 1$ and C_2 independent of T and (R, ζ) .

We now proceed to prove that W is continuous.

Suppose (R_n, ζ_n) and (R, ζ) belong to $K_R \times K_\zeta$ and $R_n \rightarrow R$, $\zeta_n \rightarrow \zeta$ in the $C^0[0, T]$ -norm. We need to prove that

$$W(R_n, \zeta_n) \rightarrow W(R, \zeta) .$$

Define $A_n, B_n, \bar{\zeta}_n, \bar{R}_n$ and C_n corresponding to R_n, ζ_n , so that $W(R_n, \zeta_n) = (\bar{R}_n, \bar{\zeta}_n)$. Applying the estimates (3.8), (3.9) to A_n and similar estimates to B_n and C_n , we can easily show that any subsequence of n 's has a subsequence for which

$$A_n \rightarrow \tilde{A}, B_n \rightarrow \tilde{B}$$

where \tilde{A}, \tilde{B} satisfy the same parabolic system in $R(t) < r < \infty$, $0 < t < T$ which A, B satisfy. By uniqueness it follows that $\tilde{A} = A$ and $\tilde{B} = B$. Therefore

$$(3.10) \quad A_n \rightarrow A, B_n \rightarrow B \quad \text{uniformly in } \rho_n(t) \leq r < \infty, 0 \leq t \leq T$$

where $\rho_n(t) = \max\{R_n(t), R(t)\}$.

We proceed to prove that

$$(3.11) \quad \begin{aligned} \overline{R}_n(t) - \overline{R}(t) &= \frac{\alpha}{D_A} \left\{ \int_0^t \frac{1 - \zeta_{n,0}(s)}{\zeta_{n,0}(s)} [A_n(R_n(s), s) - A^*] ds \right. \\ &\quad \left. - \int_0^t \frac{1 - \zeta_0(s)}{\zeta_0(s)} [A(R(s), s) - A^*] ds \right\} \rightarrow 0 \end{aligned}$$

where $\zeta_{n,0} = \min\{\zeta_n, 1\}$. By (3.9)

$$\begin{aligned} |A_n(R_n(s), s) - A_n(\rho_n(s), s)| &\leq C_1 |R_n(s) - R(s)|^\mu, \\ |A(R(s), s) - A(\rho_n(s), s)| &\leq C_1 |R_n(s) - R(s)|^\mu. \end{aligned}$$

Using also (3.10) we conclude that

$$\left| \frac{\alpha}{D_A} \int_0^t \frac{1 - \zeta_{n,0}(s)}{\zeta_{n,0}(s)} [A_n(R_n(s), s) - A(R(s), s)] ds \right| \leq C_2 \int_0^t |R_n(s) - R(s)|^\mu ds + \varepsilon_n$$

where $\varepsilon_n \rightarrow 0$ uniformly in t , as $n \rightarrow \infty$. Since $\zeta_n(t) \rightarrow \zeta(t)$ uniformly, we deduce from the expression for $\overline{R}_n - \overline{R}$ in (3.11) that

$$|\overline{R}_n(t) - \overline{R}(t)| \rightarrow 0 \quad \text{uniformly in } t \in [0, T].$$

Similarly we can prove that $C_n \rightarrow C$ for $\overline{R}(t) < r < \infty$, $0 \leq t \leq T$ and $\overline{\zeta}_n(t) \rightarrow \overline{\zeta}(t)$ uniformly in $0 \leq t \leq T$, and this completes the proof of Lemma 3.3.

From Lemmas 3.2, 3.3 we have that W maps the compact set $K_R \times K_\zeta$ into itself and is continuous. By the Schauder fixed-point theorem, W has a fixed point (R, ζ) and this completes the proof of Theorem 3.1.

§4. Uniqueness. In this section we prove:

THEOREM 4.1. *For any $T > 0$, Problem (P) has at most one solution.*

Proof. Suppose there are two solutions $(A_i, B_i, C_i, R_i, \zeta_i)$ ($i = 1, 2$). Set

$$\hat{A}_i(x, t) = A_i(r, t), \hat{B}_i(x, t) = B_i(r, t), \hat{C}_i(x, t) = C_i(r, t)$$

where $x = r - R_i(t)$. Then

$$(4.1) \quad \begin{aligned} \hat{A}_{i,t} &= D_A \left(\hat{A}_{i,xx} + \frac{2}{x + R_i} \hat{A}_{i,x} \right) + \dot{R}_i \hat{A}_{i,x} - K \hat{A}_i \hat{B}_i \\ &\quad \text{in } Q_T = \{x > 0, t > 0\}, \\ -D_A \zeta_{i,0} \hat{A}_{i,x} + (1 - \zeta_{i,0})(\hat{A}_i - A^*) &= 0 \quad \text{for } x = 0, t > 0, \\ \hat{A}_i(x, 0) &= A_0(x + R_0) \quad \text{for } x > 0, \\ \lim_{x \rightarrow \infty} \hat{A}_i(x, t) &= 0 \quad \text{for } t > 0. \end{aligned}$$

Similar systems can be written for \widehat{B}_i and \widehat{C}_i . Set

$$u = \widehat{A}_1 - \widehat{A}_2, \quad v = \widehat{B}_1 - \widehat{B}_2, \quad w = \widehat{C}_1 - \widehat{C}_2$$

and

$$\rho = R_1 - R_2, \quad \eta = \zeta_1 - \zeta_2.$$

Then

$$\dot{\rho}(t) = \frac{\alpha}{D_A} \left[\frac{1 - \zeta_{1,0}(t)}{\zeta_{1,0}(t)} (\widehat{A}_1(x, 0) - A^*) - \frac{1 - \zeta_{2,0}(t)}{\zeta_{2,0}(t)} (\widehat{A}_2(x, 0) - A^*) \right]$$

so that

$$(4.2) \quad |\dot{\rho}(t)| \leq C[|u(0, t)| + |\eta(t)|].$$

Next we easily estimate from (1.13),

$$(4.3) \quad |\eta(t)| \leq C \left[\|\rho\|_{L^\infty(0, T)} + \|w(0, t)\|_{L^\infty(0, T)} \right].$$

Since $\rho(0) = 0$, (4.2) yields

$$(4.4) \quad \|\rho\|_{L^\infty(0, T)} \leq CT \left[\|u(0, t)\|_{L^\infty(0, T)} + \|\eta\|_{L^\infty(0, T)} \right].$$

Substituting this into (4.3), we get

$$(4.5) \quad \|\eta\|_{L^\infty(0, T)} \leq \frac{CT}{1 - CT} \left[\|u(0, t)\|_{L^\infty(0, T)} + \|w(0, t)\|_{L^\infty(0, T)} \right]$$

if T is such that $CT < 1$. Finally, using (4.5) in (4.4) we get

$$(4.6) \quad \|\rho\|_{L^\infty(0, T)} \leq CT \left[\|u(0, t)\|_{L^\infty(0, T)} + \|w(0, t)\|_{L^\infty(0, T)} \right].$$

From (4.1) we see that u satisfies a parabolic system

$$(4.7) \quad \begin{aligned} u_t &= D_A u_{xx} + \left(\frac{2D_A}{x + R_1} + \dot{R}_1 \right) u_x - K \widehat{B}_1 u + F \quad \text{in } Q_T, \\ &-D_A \zeta_{1,0} u_x + (1 - \zeta_{1,0})u = G \quad \text{if } x = 0, t > 0, \\ u(x, 0) &= 0, \quad x > 0 \\ \lim_{x \rightarrow \infty} u(x, t) &= 0, \quad t > 0. \end{aligned}$$

where

$$\begin{aligned} |F| &\leq C \left[\|\dot{\rho}\|_{L^\infty(0, T)} + \|v\|_{L^\infty(Q_T)} \right] \\ &\leq C \left[\|u\|_{L^\infty(Q_T)} + \|v\|_{L^\infty(Q_T)} + \|w\|_{L^\infty(Q_T)} \right] \end{aligned}$$

by (4.2), (4.5), and

$$|G| \leq C|\eta(t)| \leq CT \left[\|u\|_{L^\infty(Q_T)} + \|w\|_{L^\infty(Q_T)} \right].$$

It is easy to see that the function

$$C(t+T) \cdot \left[\|u\|_{L^\infty(Q_T)} + \|v\|_{L^\infty(Q_T)} + \|w\|_{L^\infty(Q_T)} \right]$$

is a supersolution to (4.7), so that

$$\|u\|_{L^\infty(Q_T)} \leq CT \left[\|u\|_{L^\infty(Q_T)} + \|v\|_{L^\infty(Q_T)} + \|w\|_{L^\infty(Q_T)} \right] .$$

The same estimate can similarly be established for v and for w . Hence if T is small enough then $u = v = w = 0$ in Q_T . We can now proceed step-by-step to prove that $u = v = w = 0$ for all t , as long as the two solutions exist.

§5. Global existence. In this section we prove that there exists a solution to Problem (P) for all time. We first recall that as long as $\zeta(t) \leq 1$

$$(5.1) \quad \dot{R}(t) = \frac{\alpha}{D_A} \frac{1 - \zeta(t)}{\zeta(t)} (A - A^*) ,$$

$$(5.2) \quad \dot{\zeta}(t) = \beta\gamma C^4(R(t), t) - \frac{2\zeta(t)}{R(t)} \dot{R}(t) .$$

Hence

$$(5.3) \quad \dot{R}(t) < 0 , \quad \dot{\zeta}(t) > 0 .$$

LEMMA 5.1. *There exists a positive constant R_* such that*

$$(5.4) \quad R(t) \geq R_* \quad \text{as long as} \quad \zeta(t) \leq 1 .$$

Proof. Take λ and ε_1 positive and small such that

$$C^4(R(t), t) \leq \varepsilon_1 \quad \text{and} \quad \frac{1}{2} R_0 < R(t)$$

for all $\lambda \leq t \leq 2\lambda$. For $t > 2\lambda$

$$\zeta(t) \geq \frac{\beta\gamma\varepsilon_1\lambda + \delta}{R^2(t)}$$

Hence $R(t) > R_*$ as long as $\zeta(t) \leq 1$, where $R_* = (\beta\gamma\varepsilon_1\lambda)^{1/2}$.

REMARK 5.1. Note that the lower bound R_* is independent of the regularizing parameter δ ; cf. Remark 1.1.

THEOREM 5.2. *There exists a global solution to Problem (P), and $R \in C^{1+\mu}[0, \infty) \cap C^\infty(0, T^*)$, for any $0 < \mu < 1$.*

Proof. Suppose we already have a solution for $0 \leq t \leq T$ where T is a positive number, not necessarily small. By Lemma 5.1, $R(t) \geq R_*$ for $0 \leq t \leq T$.

Since $R(t) \geq R_* > 0$ (R_* independent of t), a review of the proof of local existence shows that the solution can be extended to $0 \leq t \leq T + T_0$ provided T_0 is a small positive constant depending only on an a priori bound on $\sup |\dot{R}(t)|$. By (5.1) $|\dot{R}(t)|$ is indeed uniformly bounded

(independently on t), and therefore the solution to Problem (P) can be extended step-by-step to all $t > 0$.

To prove the a priori regularity of R and ζ , we perform a change of variables

$$\widehat{A}(x, t) = A(r, t), \quad \widehat{B}(x, t) = B(r, t), \quad \widehat{C}(x, t) = C(r, t)$$

where $x = r - R(t)$. Then \widehat{A} satisfies, for any $T > 0$,

$$\begin{aligned} \widehat{A}_t &= D_A \left(\widehat{A}_{xx} + \frac{2}{x+R} \widehat{A}_x \right) + \dot{R}A_x - K\widehat{A}\widehat{B} \quad \text{in } Q_T, \\ -D_A\zeta_0\widehat{A}_x + (1-\zeta_0)(\widehat{A} - A^*) &= 0 \quad \text{for } x=0, \quad 0 < t < T, \\ \widehat{A}(x, 0) &= A_0(x + R_0), \quad x > 0, \\ \lim_{x \rightarrow \infty} \widehat{A}(x, t) &= 0, \quad 0 < t < T. \end{aligned} \tag{5.5}$$

Since $\dot{R}(t)$ is uniformly bounded, the same is true of $\dot{\zeta}(t)$, by (5.2) (Recall that $C(R(t), t)$ is bounded by $KA^*B^*T + \widehat{C}$). We can therefore apply the L^p parabolic estimates [10] to \widehat{A} and conclude that

$$\int_{Q_T} (|\widehat{A}_r| + |\widehat{A}_{rr}| + |\widehat{A}_t|)^p \leq C_{p,T}$$

for any $p > 1$. This implies that

$$\widehat{A} \in C^{\mu, \mu/2}(Q_T) \quad \text{for } 0 < \mu < 1,$$

and yields the $C^{1+\mu/2}[0, T]$ regularity of $R(t)$.

Similar L^p estimates can be established for \widehat{B} and \widehat{C} and then, from (5.2), one deduces the $C^{1+\mu/2}[0, T]$ regularity of $\zeta(t)$.

The above arguments can be used to prove step-by-step the $C^{1+\mu}[0, T]$ and $C^\infty(0, T)$ regularity of $R(t)$ and $\zeta(t)$ for any $0 < T \leq T^*$

§6. Finite shut-down.

LEMMA 6.1. *If $\zeta(t) < 1$ for all $t > 0$ then*

$$\lim_{t \rightarrow \infty} \zeta(t) = 1 \tag{6.1}$$

and

$$\lim_{t \rightarrow \infty} \dot{R}(t) = 0. \tag{6.2}$$

Proof. The idea of the proof is to show by comparison (with linear problems) that both A and B do not go to zero as $t \rightarrow \infty$, and the same is then true of C ; this implies that $\zeta(t) \rightarrow \infty$ as $t \rightarrow \infty$, which is a contradiction. Set

$$\rho = \lim_{t \rightarrow \infty} R(t) > 0, \quad \sigma = \lim_{t \rightarrow \infty} \zeta(t).$$

We shall assume that $\sigma < 1$ and derive a contradiction. We begin by introducing the solution \tilde{A}, \tilde{B} to

$$(6.3) \quad \begin{aligned} \tilde{A}_t &= D_A \Delta \tilde{A} - K B^* \tilde{A} \quad \text{for } r > R(t), t > 0, \\ -D_A \zeta(t) \tilde{A}_r + (1 - \zeta(t))(\tilde{A} - A^*) &= 0 \quad \text{if } r = R(t), t > 0, \\ \tilde{A}(r, 0) &= A_0(r) \quad \text{if } r > R_0, \\ \lim_{r \rightarrow \infty} \tilde{A}(r, t) &= 0 \quad \text{for any } t > 0, \end{aligned}$$

and

$$(6.4) \quad \begin{aligned} \tilde{B}_t &= D_B \Delta \tilde{B} - K A^* \tilde{B} \quad \text{for } r > R(t), t > 0, \\ \tilde{B}_r &= 0 \quad \text{if } r = R(t), t > 0, \\ \tilde{B}(r, 0) &= B_0(r) \quad \text{if } r > R_0, \\ \lim_{r \rightarrow \infty} \tilde{B}(r, t) &= B^* \quad \text{for any } t > 0. \end{aligned}$$

By comparison, $\tilde{A} \leq A$, $\tilde{B} \leq B$.

Let $\tilde{A}_\infty, \tilde{B}_\infty$ denote the solutions to the limit problems ($t \rightarrow \infty$) (6.3) and (6.4):

$$(6.5) \quad \begin{aligned} D_A \Delta \tilde{A}_\infty - K B^* \tilde{A}_\infty &= 0, \quad \rho < r < \infty, \\ -D_A \sigma \tilde{A}_{\infty, r} + (1 - \sigma)(\tilde{A}_\infty - A^*) &= 0, \quad r = \rho, \\ \lim_{r \rightarrow \infty} \tilde{A}_\infty(r) &= 0 \end{aligned}$$

and

$$(6.6) \quad \begin{aligned} D_B \Delta \tilde{B}_\infty - K A^* \tilde{B}_\infty &= 0, \quad \rho < r < \infty, \\ \tilde{B}_{\infty, r}(\rho) &= 0, \\ \lim_{r \rightarrow \infty} \tilde{B}_\infty(r) &= B^*. \end{aligned}$$

Note that, since $\sigma < 1$, $\tilde{A}_\infty(r) \geq 0$ for all $r > \rho$; $\tilde{B}_\infty(r)$ is positive for all $r \geq \rho$.

We shall prove that the functions

$$\tilde{u}(r, t) = \tilde{A}(r, t) - \tilde{A}_\infty(r), \quad \tilde{v}(r, t) = \tilde{B}(r, t) - \tilde{B}_\infty(r)$$

converge to zero as $t \rightarrow \infty$.

The function \tilde{u} satisfies

$$(6.7) \quad \begin{aligned} \tilde{u}_t &= D_A \Delta \tilde{u} - K B^* \tilde{u} \quad \text{if } r > R(t), t > 0 \\ -D_A \sigma \tilde{u}_r + (1 - \sigma) \tilde{u} &= \tilde{\varepsilon}(t) \quad \text{if } r = R(t), t > 0 \end{aligned}$$

where $\tilde{\varepsilon}(t) \rightarrow 0$ as $t \rightarrow \infty$. For any $\varepsilon > 0$,

$$\frac{1}{\sigma} |\tilde{\varepsilon}(t)| < \varepsilon \quad \text{if } t > t_\varepsilon,$$

and the function

$$(6.8) \quad \tilde{U}(r, t) = A^* e^{-KB^*(t-t_\varepsilon)} + \frac{\varepsilon}{r} (R(t_\varepsilon))^2$$

is then a supersolution of (6.7) in $r > R(t)$, $t > t_\varepsilon$, which majorizes $|\tilde{u}|$ both on $t = t_\varepsilon$ and as $r \rightarrow \infty$. Hence, by comparison,

$$|\tilde{u}(r, t)| \leq \tilde{U}(r, t) .$$

Taking $t \rightarrow \infty$ and recalling that ε is arbitrary, we conclude that

$$(6.9) \quad \sup_{r \geq R(t)} |\tilde{A}(r, t) - \tilde{A}_\infty(r)| \rightarrow 0 \quad \text{if } t \rightarrow \infty .$$

Similarly, the function v satisfies

$$\begin{aligned} \tilde{v}_t &= D_B \Delta \tilde{v} - KA^* \tilde{v} \quad \text{for } r > R(t), t > 0 , \\ v_r(R(t), t) &\rightarrow 0 \quad \text{if } t \rightarrow \infty \end{aligned}$$

and, with a supersolution similar to (6.8), we conclude that

$$(6.10) \quad \sup_{r \geq R(t)} |\tilde{B}(r, t) - \tilde{B}_\infty(r)| \rightarrow 0 \quad \text{if } t \rightarrow \infty .$$

Since $\tilde{A}_\infty(r) > 0$, $\tilde{B}_\infty(r) > 0$ for $r > \rho$, it follows from (6.9), (6.10) that

$$(6.11) \quad \tilde{A}(r, t) \tilde{B}(r, t) \geq c_0 \chi_{[a, b]} \quad \text{for } t \geq t_*$$

where c_0 is a positive constant, $\rho < a < b < \infty$, and t_* is sufficiently large.

Consider the parabolic problem

$$(6.12) \quad \begin{aligned} \tilde{C}_t &= D_C \Delta \tilde{C} + Kc_0 \chi_{[a, b]} \quad \text{if } r > R(t), t > t_* , \\ D_C \tilde{C}_r &= \gamma \tilde{C}^4 \quad \text{if } r = R(t), t > t_* , \\ \tilde{C}(r, t_*) &= 0 \quad \text{if } r > R(t_*) , \\ \lim_{r \rightarrow \infty} \tilde{C}(r, t) &= 0 \quad \text{if } t > t_* . \end{aligned}$$

Since $AB \geq \tilde{A}\tilde{B}$, it follows from (6.11), by comparison, that

$$(6.13) \quad C(r, t) \geq \tilde{C}(r, t) .$$

Denote by \tilde{C}_∞ the solution to

$$(6.14) \quad \begin{aligned} D_C \Delta \tilde{C}_\infty + Kc_0 \chi_{[a, b]} &= 0 , \quad r > \rho , \\ D_C \tilde{C}_{\infty, r} &= \gamma \tilde{C}_\infty^4 , \quad r = \rho , \\ \tilde{C}_\infty(\infty) &= 0 ; \end{aligned}$$

clearly $\tilde{C}_\infty(r) > 0$ for $r \geq \rho$. We shall evaluate the difference

$$(6.15) \quad \tilde{w}(r, t) = \tilde{C}(r, t) - \tilde{C}_\infty(r) .$$

It satisfies

$$\begin{aligned} \tilde{w}_t &= D_C \Delta w \quad \text{if } r > R(t), t > t_* , \\ -D_C w_r + a(t)w &= \varepsilon(t) \quad \text{if } r = R(t), t > 0 , \\ \lim_{r \rightarrow \infty} w(r, t) &= 0 \quad \text{if } t > t_* , \end{aligned}$$

where

$$\varepsilon(t) = D_C \tilde{C}_{\infty, r}(R(t)) - \gamma \tilde{C}_\infty^4(R(t)) \rightarrow 0 \quad \text{if } t \rightarrow \infty$$

and

$$a(t) = (\gamma \tilde{C}^4 - \gamma \tilde{C}_\infty^4) / (\tilde{C} - \tilde{C}_\infty) > 0 .$$

As in the proof of (6.9) we now construct, for any $\varepsilon > 0$, a supersolution (cf. (6.8))

$$(6.16) \quad \tilde{W}(r, t) = \frac{M_\varepsilon}{t^{3/2}} e^{-\frac{r^2}{4t\sqrt{D_C}}} + \frac{\varepsilon}{r} (R(t_\varepsilon))^2 \quad \text{for } t > t_\varepsilon$$

where $D_A |\varepsilon(t)| < \varepsilon$ if $t > t_\varepsilon$. By comparison,

$$|w(r, t)| \leq \tilde{W}(r, t) \quad \text{if } r > R(t), t > t_\varepsilon .$$

Recalling (6.15) we conclude that

$$\tilde{C}(R(t), t) \rightarrow \tilde{C}_\infty(\rho) \quad \text{if } t \rightarrow \infty$$

and, by (6.13),

$$(6.17) \quad C(R(t), t) \geq \nu \quad \text{if } t \geq \tilde{t} \quad (\nu = \frac{1}{2} \tilde{C}_\infty(\rho) > 0)$$

for some $\tilde{t} > 0$.

Since $D_C C_r = \gamma C^4$ at $(R(t), t)$, and since $0 < R_* \leq R(t) < R_0$, we can use (6.17) in (1.13) to deduce that $\zeta(t) \rightarrow \infty$ if $t \rightarrow \infty$; this is a contradiction to the assumption that $\zeta(t) < 1$ for all t . We have thus proved that $\sigma = 1$, i.e., (6.1) holds. The assertion (6.2) then follows from (5.1) and (6.1).

To prove that there is a finite shut-down time we may proceed by contradiction as before, namely, we assume that $\zeta(t) < 1$ for all $t > 0$, and then wish to derive a contradiction by using (1.13). This however requires a good lower bound on $C(R(t), t)$ in case $\sigma = \lim_{t \rightarrow \infty} \zeta(t)$ is equal to 1 (Note that if $\sigma = 1$ then $\tilde{A} = 0$ and the previous proof does not provide a useful lower bound on $C(R(t), t)$.) The desired bound is provided in the following lemma.

LEMMA 6.2. *If $\zeta(t) < 1$ for all $t > 0$ then there exists a positive constant M such that*

$$(6.18) \quad C_r(R(t), t) \geq \frac{M}{t^6}$$

for all t large.

Proof. For any $t_0 > 0$, $C(r, t_0) > 0$ for all $r \geq R(t_0)$. Choose $R(t_0) < a < b < \infty$ such that

$$(6.19) \quad C(r, t_0) > \varepsilon_0 \chi_{[a, b]}(r) \equiv c_0(r)$$

for some $\varepsilon_0 > 0$, and define

$$(6.20) \quad \tilde{C}(r, t) = \int_{\mathbb{R}^3} \frac{e^{-\frac{|x-y|^2}{4\sqrt{D_C}(t-t_0)}}}{(t-t_0)^{3/2}} \varepsilon c_0(|y|) dy, \quad 0 < \varepsilon < 1$$

for $r = |x| > R(t)$, $t > t_0$. Since $\tilde{C}(R(t_0), t_0) = 0$,

$$C(R(t), t) > \tilde{C}(R(t), t) \quad \text{if } t_0 \leq t \leq t_1$$

for some $t_1 > t_0$. On the other hand, for $r = R(t)$, $t \geq t_1$ we have

$$\begin{aligned} -D_C \tilde{C}_r + \gamma \tilde{C}^4 &\leq -D_C \varepsilon \int_{\{a < |y| < b\}} \frac{|y| - R(t)}{(4\pi)^{3/2} 2\sqrt{D_C} (t-t_0)^{5/2}} e^{-\frac{|R(t)-y|^2}{4\sqrt{D_C}(t-t_0)}} \\ &\quad + C_1 \varepsilon^4 < 0 \end{aligned}$$

if ε is sufficiently small. Recalling also (6.19), (6.20) and noting that \tilde{C} is a subsolution of (1.3), we conclude, by comparison, that if ε is small enough then

$$C(r, t) > \tilde{C}(r, t) \quad \text{if } t > t_0,$$

from which (6.18) follows upon using (1.6).

We finally prove:

THEOREM 6.3. *There is a finite shut-down time T^* .*

Proof. Suppose $\zeta(t) < 1$ for all $t > 0$. We can write, for $r = R(t)$,

$$\begin{aligned} \dot{\zeta}(t) &= \beta D_C C_r - \zeta(t) \frac{2\dot{R}(t)}{R(t)} = \beta D_C C_r - \frac{2\alpha}{R(t)} \zeta(t) A_r \\ &= \beta D_C C_r - \frac{2\alpha}{R(t)} (1 - \zeta(t)) (A - A^*). \end{aligned}$$

Setting

$$\eta(t) = \zeta(t) - 1, \quad q(t) = \frac{2\alpha}{R(t)} (A^* - A(R(t), t))$$

we can write

$$(6.21) \quad \dot{\eta} + q(t)\eta = \beta D_C C_r(R(t), t) \equiv p(t),$$

so that,

$$(6.22) \quad \eta(t) = e^{-\int_0^t q(\tau)d\tau} \eta(0) + \int_0^t p(\tau)e^{-\int_\tau^t q(s)ds} d\tau .$$

Consider first the case that there exists a positive constant α_0 and a sequence $t_n \rightarrow \infty$ such that

$$(6.23) \quad \frac{1}{t_n} \int_0^{t_n} (A^* - A(R(\tau), \tau))d\tau \geq \alpha_0 \quad \text{as } t_n \rightarrow \infty ;$$

then,

$$\int_0^{t_n} q(\tau)d\tau \geq \frac{2\alpha}{R_0} \alpha_0 t_n \quad \text{as } t_n \rightarrow \infty .$$

By using (6.18) in (6.22), we get

$$\begin{aligned} \eta(t_n) &\geq \eta(0)e^{-\frac{2\alpha\alpha_0}{R_0} t_n} + \int_{t_0}^{t_n} \frac{M}{\tau^6} e^{-\frac{2\alpha A^*}{R_*}(t_n - \tau)} d\tau \\ &\geq \eta(0)e^{-\frac{2\alpha\alpha_0}{R_0} t_n} + \frac{1}{2} \frac{R_*}{2\alpha A^*} \frac{M}{t_n^6} \end{aligned}$$

for some $t_0 > 0$ and all t_n sufficiently large. This is a contradiction since $\eta(t) = \zeta(t) - 1 < 0$ for all $t > 0$.

It remains to consider the case where (6.23) is not satisfied for any $\alpha_0 > 0$, that is,

$$(6.24) \quad \frac{1}{t} \int_0^t A(R(s), s)ds \rightarrow A^* \quad \text{as } t \rightarrow \infty .$$

Since

$$A_r(R(t), t) = \frac{1}{\alpha} \dot{R}(t) \rightarrow 0$$

by Lemma 6.1, we can derive, by comparison,

$$A(r, t) \leq \widetilde{W}(r, t)$$

for any $\varepsilon > 0$ and $t > t_\varepsilon$, where \widetilde{W} is the function introduced in (6.16). It follows that

$$A(R(t), t) \rightarrow 0 \quad \text{if } t \rightarrow \infty ,$$

a contradiction to (6.24).

§7. **Asymptotic estimates as $K \rightarrow \infty$.** We now study the behavior of the solution $(A_K, B_K, C_K, R_K, \zeta_K)$ as $K \rightarrow \infty$, assuming that

$$(7.1) \quad D_B = D_A .$$

Recall that

$$(7.2) \quad 0 \leq A_K \leq A^* , \quad 0 \leq B_K \leq B^*$$

and

$$(7.3) \quad \frac{\partial}{\partial r} A_K \leq 0 , \quad \frac{\partial}{\partial r} B_K \geq 0 , \quad \dot{R}_K(t) \leq 0 .$$

LEMMA 7.1. *There exists a positive constant M such that*

$$(7.4) \quad -M \leq \frac{\partial}{\partial r} A_K , \quad \frac{\partial}{\partial r} B_K \leq M , \quad -\alpha M \leq \dot{R}_K(t)$$

for all K .

Proof. Consider the function $u = A_K - B_K$. It satisfies

$$(7.5) \quad \begin{aligned} u_t &= D_A u \quad \text{if } r > R_K(t), t > 0 \\ -\zeta_{K,0}(t) D_A u_r &= (1 - \zeta_{K,0}(t))(A^* - A) \quad \text{if } r = R_K(t), t > 0 , \\ u(r, 0) &= A_0(r), \quad r > R_0 , \\ u(\infty, t) &= 0 , \quad t > 0 \end{aligned}$$

where $\zeta_{K,0} = \min\{\zeta_K, 1\}$. Applying the maximum principle to u_r we deduce that

$$(7.6) \quad -M \leq u_r \leq 0$$

where M is a constant independent of K . But then, upon recalling also (7.3),

$$(A_{K,r})^2 = A_{K,r} u_r + A_{K,r} B_{K,r} \leq A_{K,r} u_r \leq M |A_{K,r}|$$

so that $|A_{K,r}| \leq M$. The proof that $|B_{K,r}| \leq M$ is similar. Finally, $\dot{R}_K = \alpha A_{K,r} \leq -\alpha M$.

LEMMA 7.2. *There exist positive constants N_0, N independent of K such that, for any $T > 0$,*

$$(7.7) \quad \int_0^T \int_{R_K(t)}^\infty K A_K B_K r^2 dr dt \leq N_0 + NT$$

for all K .

Proof. Integrating the equation

$$(7.8) \quad K A_K B_K = D_A \Delta A_K - \frac{\partial}{\partial t} A_K$$

over $R_K(t) < r < \infty$, $0 < t < T$ and using Lemma 7.1, (7.7) readily follows.

In order to obtain uniform Hölder estimates on A_K, B_K , we consider the function $v = A_K B_K$. It satisfies:

$$\begin{aligned}
(7.9) \quad & v_t = D_A \Delta v - 2A_{K,r} B_{K,r} - K(A_K + B_K)v, \quad r > R_K(t), \\
& -\zeta_{K,0}(t) D_A v_r = (1 - \zeta_{K,0}(t))(A^* - A_K)B_K, \quad r = R_K(t), \\
& v(r, 0) = A_0(r)B_0(r), \quad r > R_0, \\
& v(\infty, t) = 0, \quad t > 0.
\end{aligned}$$

LEMMA 7.3. *For any $T > 0$ and any compact set Ω_T in $\{R_K(t) < r < \infty, 0 \leq t \leq T\}$ whose distance to $r = R_K(t)$ is $\geq c_* > 0$, there exists a constant M depending on T and c_* (but not on K) such that*

$$(7.10) \quad \int_{\Omega_T} [K A_K B_K (A_K + B_K)]^2 r^2 dr dt \leq M.$$

Proof. Let $\xi(r, t)$ be a cutoff function such that $\xi = 1$ in Ω_T and $\xi = 0$ outside $(\frac{1}{2} c_*)$ -neighborhood of Ω_T . Multiplying the differential equation for v by $\xi^2 K v r^2$ and integrating, we obtain

$$\begin{aligned}
& \int_{R_K(t)}^{\infty} \xi^2 K \frac{v^2}{2} r^2 dr + \int_0^T \int_{R_K(t)}^{\infty} \xi^2 K (v_r)^2 r^2 dr dt \\
& + \int_0^T \int_{R_K(t)}^{\infty} \xi^2 K^2 v^2 (A_K + B_K) r^2 dr dt \leq I
\end{aligned}$$

where

$$I = \int_0^T \int_{R_K(t)}^{\infty} \left[|\xi \xi_t| K v^2 + 2\xi |\xi_r| |v_r| K v + 2|A_{K,r} B_{K,r}| \xi^2 K v \right] r^2 dr dt$$

is bounded independently of K , by Lemmas 7.1 and 7.2; this implies the assertion (7.10).

LEMMA 7.4. *Let Ω_T be as in Lemma 7.3. Then there exists a constant M depending on T and c_* (but not on K) such that*

$$(7.11) \quad \|A_K\|_{C^{1/4,1/8}(\Omega_T)} \leq M,$$

$$(7.12) \quad \|B_K\|_{C^{1/2,1/8}(\Omega_T)} \leq M.$$

That means that the A_K and B_K are uniformly Hölder continuous (in Ω_T) with exponent $1/4$ in r and exponent $1/8$ in t .

Proof. By Lemmas 7.2, 7.3,

$$v_t - D_A \Delta v \quad \text{is in} \quad L^2(\Omega_T),$$

uniformly in K . By L^2 estimates it then follows that

$$\|v\|_{W_2^{2,1}(\Omega_T)} \leq M$$

with a slightly smaller set Ω_T and a larger constant M , both independent however of K . By Sobolev's imbedding [11] we then have

$$(7.13) \quad \|v\|_{C^{1/2,1/4}(\Omega_T)} \leq M$$

with yet another constant M .

The function $u = A_K - B_K$ satisfies (7.5), and we can apply L^p estimates to deduce u also satisfies the estimate (7.13). Thus both $A_K B_K$ and $A_K - B_K$ belong to $C^{1/2,1/4}(\Omega_T)$, uniformly in K . Since

$$A_K + B_K = (4v + u^2)^{1/2}$$

it follows that $A_K + B_K$ is in $C^{1/4,1/8}(\Omega_T)$, and the same then holds for A_K and B_K .

THEOREM 7.5. *Let $\tilde{\Omega}_T$ be any compact which is contained in $R_K(t) < r < \infty$, $0 \leq t < T$ for all K sufficiently large. Then*

$$(7.14) \quad \lim_{K \rightarrow \infty} A_K(r, t) B_K(r, t) = 0$$

uniformly in $(r, t) \in \tilde{\Omega}_T$.

Proof. By Lemma 7.2

$$\iint_{\tilde{\Omega}_T} A_K B_K r^2 dr dt \leq \frac{M}{K} \rightarrow 0$$

if $K \rightarrow \infty$. Since $A_K B_K$ is uniformly Hölder continuous in $\tilde{\Omega}_T$ with exponent and coefficient independent of K , (7.14) follows.

§8. Asymptotic limits as $K \rightarrow \infty$. The estimates of §7 show that for any sequence $K'_n \rightarrow \infty$ there is a subsequence K_n such that, as $K_n \rightarrow \infty$,

$$(8.1) \quad R_{K_n}(t) \rightarrow R(t) \in Lip[0, T] \quad \text{in} \quad C^0[0, T],$$

$$(8.2) \quad A_{K_n}(r, t) \rightarrow A(r, t) \in C^{1/4,1/8}(Q_T) \quad \text{in} \quad C^0(Q_T),$$

$$(8.3) \quad B_{K_n}(r, t) \rightarrow B(r, t) \in C^{1/4,1/8}(Q_T) \quad \text{in} \quad C^0(Q_T),$$

$$(8.4) \quad K_n A_{K_n} B_{K_n} \rightarrow f \quad \text{in the sense of weak convergence of measures,}$$

for any $0 < T < \infty$, where

$$Q_T = \{(r, t); r > R(t), 0 \leq t < T\},$$

and f is a measure.

From Theorem 7.5 we have

$$(8.5) \quad A(r, t) B(r, t) = 0.$$

The functions A, B both satisfy the equation

$$(8.6) \quad w_t - D_A \Delta w + f = 0 \quad \text{in} \quad \mathcal{D}'(Q_T),$$

whereas the function $u = A - B$ satisfies

$$(8.7) \quad u_t - D_A \Delta u = 0 \quad \text{in} \quad Q_T$$

since each of the functions $A_{K_n} - B_{K_n}$ satisfies this equation.

By (7.3) it follows that $u_r \leq 0$ in Q_T and then, by the strong maximum principle,

$$(8.8) \quad u_r < 0 \quad \text{in} \quad Q_T.$$

It follows that there exists a curve $r = S(t)$ with $S(t) \in C^{1+\nu}[0, T] \cap C^\infty(0, T)$ such that

$$(8.9) \quad \begin{aligned} u(r, t) &> 0 \quad \text{if} \quad r < S(t), \\ u(r, t) &< 0 \quad \text{if} \quad r > S(t); \end{aligned}$$

here ν is as in (1.10), (1.11). Since $A_0 - B_0$ is positive at $r = R_0$ and negative at $r = \infty$,

$$(8.10) \quad R_0 < S(0) < \infty.$$

Take T such that

$$(8.11) \quad S(t) > R(t) \quad \text{for all} \quad 0 \leq t \leq T.$$

For $r < S(t)$ we have $u > 0$, or $A > B$. Since $AB = 0$, it follows that $B = 0$. Similarly $A = 0$ if $r > S(t)$; thus

$$(8.12) \quad \begin{aligned} A(r, t) &= 0 \quad \text{if} \quad r > S(t), \\ B(r, t) &= 0 \quad \text{if} \quad r < S(t). \end{aligned}$$

In any closed domain in $\{r < S(t), t < T\}$ we have $B = 0$ and then, by (8.6) with $w = B$, $f = 0$. Similarly $f = 0$ if $r > S(t)$. It follows that f is a measure supported on $r = S(t)$, $0 < t < T$. In particular,

$$A_t - D_A \Delta A = 0 \quad \text{if} \quad r < S(t), \quad 0 < t < T.$$

Since also $A(S(t), t) = 0$ and $S(t)$ is smooth, regularity results for the heat equation imply that A is in $C^{1+\nu}$ in $R(t) \leq r \leq S(t)$, $0 \leq t \leq T$ and in C^∞ in $R(t) \leq r \leq S(t)$, $0 < t \leq T$.

Next, from (8.6) for $w = A$,

$$(8.13) \quad - \int \int A(-\varphi_t - D_A \Delta \varphi) r^2 dr dt = \int \int f \varphi r^2 dr dt$$

for any test function φ in Q_T . Using the fact that $A = 0$ if $r > S(t)$, and integrating by parts in (8.13) we find that

$$\int \int f \varphi r^2 dr dt = - \int_0^T S^2(t) D_A A_r(S(t) - 0, t) \varphi(S(t), t) dt.$$

It follows that

$$(8.14) \quad f(r, t) = D_A A_r(S(t) - 0, 0)\delta(r - S(t)) .$$

THEOREM 8.1. *For any sequence $K'_n \rightarrow 0$ there is a subsequence $K_n \rightarrow \infty$ such that the solutions of (P_{K_n}) converge to a solution of (P_∞) uniformly in compact subsets of $\{r > R(t), 0 \leq t \leq T\}$*

Proof. We have already proved most of the theorem. Since $S(t) > R(t)$ for $0 \leq t \leq T$, it follows that A satisfies the two boundary conditions (2.8), (2.9), where $\zeta = \lim_{K \rightarrow \infty} \zeta_K$. Using the $C^{1+\mu}$ -regularity of A (which one obtains by the same argument as for Problem (P_K)) as well as the $C^{1+\mu}$ -regularity of $R_K(t)$ and of the A_K near $r = R_K(t)$, we can deduce that $C_{K_n} \rightarrow C$ uniformly near $r = R(t)$ and that C satisfies the boundary condition (2.10). The remaining assertions of Theorem 8.1 have already been proved.

Denote by T_f be the supremum of all T 's for which (8.11) holds. We claim:

$$(8.15) \quad \text{if } T_f < \infty \text{ then } S(t) \rightarrow R(T_f) \text{ if } t \rightarrow T_f .$$

Indeed, if $\overline{\lim} S(t) > R(T_f)$, then (8.11) holds for $T > T_f$ (since u is smooth and $u_r < 0$ in $\{r > R(t), t > 0\}$). On the other hand, if the limit $S(T_f - 0)$ does not exist then $u_r(r, T_f)$ will vanish on a nonempty interval, contradicting the inequality $u_r < 0$.

In §9 we shall prove uniqueness for Problem (P_∞) ; this implies that the convergence asserted in Theorem 8.1 is not just for a subsequence $K_n \rightarrow \infty$ but for all $K \rightarrow \infty$.

In §10 we shall prove that $T_f < \infty$, and this will conclude the proof of all the assertions made in Theorems 2.2 and 2.3.

§9. Uniqueness for (P_∞) . In this section we prove:

THEOREM 9.1. *Assume that $D_A = D_B$. Then there exists at most one solution to Problem (P_∞) .*

Proof. We begin with some remarks on the regularity of any solution (A, B, C, R, S) . Consider the function

$$u = \begin{cases} A & \text{if } R(t) \leq r \leq S(t) \\ -B & \text{if } r > S(t) \end{cases}$$

for all $t < T$, where $T < T_f$, the final time of the solution. Then

$$u_t = D_A \Delta u ,$$

and $u = A$ in a neighborhood of $\{(R(t), t), 0 < t < T\}$. From this and the boundary conditions for A and C at $r = R(t)$ we easily deduce, as in earlier sections, that $R(t)$ and $\zeta(t)$ belong to $C^{1+\mu}[0, T]$ and $C(r, t)$ belongs to $C^{1+\mu}$ in $R(t) \leq r < S(t)$, $0 \leq t \leq T$, for some $0 < \mu < 1$, and

$$(9.1) \quad -M < u_r < 0 \quad \text{for } r > R(t), 0 \leq t \leq T ;$$

furthermore, the function $S(t)$ defined by

$$r = S(t) \quad \text{if } u(r, t) = 0$$

is in $C^{1+\mu}[0, T]$ with

$$\dot{S}(t) = -\frac{u_t(S(t), t)}{u_r(S(t), t)}.$$

(Note that $u_t \in C^\mu$ with $\mu = \nu$ up to $t = 0$ since A_0 and B_0 belong to $C^{2+\nu}$). The function $C(r, t)$ is C^∞ off $r = S(t)$. Near $r = S(t)$ the regularity of $C(r, t)$ is the same as the regularity of the special solution of (2.7):

$$(9.2) \quad \tilde{C}(|x|, t) = \int_0^t \int_{|y|=R(\tau)} \frac{e^{-\frac{|x-y|^2}{\sqrt{D_C} 4(t-\tau)}}}{(4\pi^2 D_C (t-\tau))^{3/2}} u_r(S(\tau), \tau) d\sigma_y d\tau.$$

This is a single layer potential and, since $u_r(S(\tau), \tau)$ is in C^μ , this potential is in C^β (across $r = S(t)$) for any $0 < \beta < 1$, and its derivative from each side of $r = S(t)$ is uniformly continuous (the proof is similar to the proof of Lemma 1 in [6; p. 217]). Thus in particular,

$$(9.3) \quad |C_r|_{L^\infty} \leq M', \quad M' \text{ constant.}$$

Suppose now that $(A_i, B_i, C_i, R_i, S_i)$ are solutions of Problem (P_∞) for $i = 1, 2$, and set

$$u_i(r, t) = \begin{cases} A_i(r, t) & \text{if } R_i(t) \leq r \leq S_i(t) \\ -B_i(r, t) & \text{if } S_i(t) < r < \infty \end{cases}$$

for $0 \leq t \leq T$ where T is such that

$$(9.4) \quad S_i(t) - R_i(t) \geq \lambda > 0 \quad \text{for } 0 \leq t \leq T \text{ and } i = 1, 2$$

where $x = r - R_i(t)$, and

$$\begin{aligned} \tilde{u}(x, t) &= \hat{u}_1(x, t) - \hat{u}_2(x, t), \\ \tilde{C}(x, t) &= \hat{C}_1(x, t) - \hat{C}_2(x, t). \end{aligned}$$

Set

$$\Omega_\tau = \{0 < x < \infty, 0 < t < \tau\} \quad \text{for any } \tau \leq T.$$

The functions \tilde{u}, \tilde{C} satisfy (for simplicity we take $D_A = D_C = 1$):

$$(9.5) \quad \begin{aligned} \tilde{u}_t &= \tilde{u}_{xx} + \left(\frac{2}{x+R_1} + \dot{R}_1\right) \tilde{u}_x + \left[\frac{2}{x+R_1} + \dot{R}_1 - \frac{2}{x+R_2} - \dot{R}_2\right] \tilde{u}_{2,x} \quad \text{in } \Omega_T, \\ -\zeta_{1,0} \tilde{u}_x + (1 - \zeta_{1,0}) \tilde{u} &= (\zeta_{1,0} - \zeta_{2,0}) [\tilde{u}_{2,x} + \tilde{u}_2 - A^*] \quad \text{if } x = 0, 0 < t < T, \\ \tilde{u}(x, 0) &= 0, \quad x > 0, \\ \tilde{u}(\infty, t) &= 0, \quad 0 < t < T, \end{aligned}$$

and

$$\begin{aligned}
(9.6) \quad & \tilde{C}_t = \tilde{C}_{xx} + \left(\frac{2}{x + R_1} + \dot{R}_1 \right) \tilde{C}_x + \left[\frac{2}{x + R_1} + \dot{R}_1 - \frac{2}{x + R_2} - \dot{R}_2 \right] \hat{C}_{2,x} \\
& + \hat{u}_{2,x}(S_2 - R_2, t) \delta(x + R_2 - S_2) - \hat{u}_{1,x}(S_1 - R_1, t) \delta(x + R_1 - S_1) \quad \text{in } \Omega_T, \\
& -\tilde{C}_x + \gamma(\hat{C}_1 + \hat{C}_2)(\hat{C}_1^2 + \hat{C}_2^2) \tilde{C} = 0 \quad \text{if } x = 0, \quad 0 < t < T, \\
& \tilde{C}(x, 0) = 0, \quad x > 0, \\
& \tilde{C}(\infty, t) = 0, \quad 0 < t < T.
\end{aligned}$$

By the maximum principle,

$$(9.7) \quad \|\tilde{u}\|_{L^\infty(\Omega_\tau)} \leq \|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)} + N\tau \left[\|R_1 - R_2\|_{L^\infty(0,\tau)} + \|\dot{R}_1 - \dot{R}_2\|_{L^\infty(0,\tau)} \right]$$

for any $0 < \tau < T$, where N is a constant independent of τ ; in the sequel we shall denote by N any such constant.

Differentiating in x the differential equation in (9.5), we obtain a parabolic equation for \tilde{u}_x . Using the maximum principle, we find that

$$(9.8) \quad \|\tilde{u}_x\|_{L^\infty(\Omega_\tau)} \leq N \|\zeta_1 - \zeta_2\|_{L^\infty(\Omega_\tau)} + N\tau \left[\|R_1 - R_2\|_{L^\infty(0,\tau)} + \|\dot{R}_1 - \dot{R}_2\|_{L^\infty(0,\tau)} \right].$$

Since $R_1(0) = R_2(0)$ we also have

$$(9.9) \quad \|R_1 - R_2\|_{L^\infty(0,\tau)} \leq \tau \|\dot{R}_1 - \dot{R}_2\|_{L^\infty(0,\tau)}.$$

From the boundary conditions for \tilde{u}_i at $x = 0$,

$$\begin{aligned}
|\dot{R}_1 - \dot{R}_2| &= \alpha |\hat{u}_{1,x}(0, t) - \hat{u}_{2,x}(0, t)| \\
&= \alpha \left| \frac{1 - \zeta_{1,0}}{\zeta_{1,0}} (\hat{u}_1 - A^*) - \frac{1 - \zeta_{2,0}}{\zeta_{2,0}} (\hat{u}_2 - A^*) \right|,
\end{aligned}$$

so that

$$(9.10) \quad \|\dot{R}_1 - \dot{R}_2\| \leq N \left[\|\tilde{u}\|_{L^\infty(\Omega_\tau)} + \|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)} \right].$$

Substituting (9.9), (9.10) into (9.7), (9.8) and choosing $\tau < 1/(2N)$, we get

$$(9.11) \quad \|\tilde{u}\|_{L^\infty(\Omega_\tau)} + \|\tilde{u}_x\|_{L^\infty(\Omega_\tau)} \leq N \|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)}.$$

Next, from the definition (1.14), (1.15) of $\zeta_{i,0}, \zeta_i$ we deduce that

$$(9.12) \quad \|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)} \leq N\tau \|\tilde{C}(0, t)\|_{L^\infty(0,\tau)} + N \|R_1 - R_2\|_{L^\infty(0,\tau)},$$

and we need to estimate $\tilde{C}(0, t)$.

Denote by $G(x, y, t, s)$ the Green function for the problem (see [3])

$$\begin{aligned}
w_t &= w_{xx} + \left(\frac{2}{x + R_1} + \dot{R}_1 \right) w_x \quad \text{in } \Omega_T, \\
-w_x + \gamma(\hat{C}_1 + \hat{C}_2)(\hat{C}_1^2 + \hat{C}_2^2) w &= 0 \quad \text{if } x = 0, \quad 0 < t < T, \\
w(\infty, t) &= 0, \quad 0 < t < T.
\end{aligned}$$

Then (using (9.6)) we can represent \tilde{C} in the form

$$\begin{aligned}
\tilde{C}(|x|, t) &= \int_0^t \int_0^\infty G(x, y, t, s) \left[\frac{2}{y + R_1(s)} + \dot{R}_1(s) - \frac{2}{y + R_2(s)} - \dot{R}_2(s) \right] \hat{C}_{2,y}(y, s) dy ds \\
&\quad + \int_0^t G(x, S_2(s) - R_2(s), t, s) \hat{u}_{2,x}(S_2(s) - R_2(s), s) ds \\
&\quad - \int_0^t G(x, S_1(s) - R_1(s), t, s) \hat{u}_{1,x}(S_1(s) - R_1(s), t - s) ds . \\
&\equiv J_1 + J_2 + J_3 .
\end{aligned}$$

We can write

$$\begin{aligned}
|J_2(0, t) + J_3(0, t)| &\leq \int_0^t |G(0, S_2 - R_2, t, s)| |\hat{u}_{2,x}(S_2 - R_2, s) - \hat{u}_{1,x}(S_1 - R_1, s)| \\
&\quad + \int_0^t |\hat{u}_{1,x}(S_1 - R_1, s)| |G(0, S_2 - R_2, t, s) - G(0, S_1 - R_1, t, s)|
\end{aligned}$$

where $S_i = S_i(s)$, $R_i = R_i(s)$. The difference in the first integral is equal to

$$[\hat{u}_{2,x}(S_2 - R_2, s) - \hat{u}_{2,x}(S_1 - R_1, s)] + \tilde{u}_x(S_1 - R_1, s) .$$

Therefore, if $0 < t < \tau$,

$$|J_2(0, t) - J_3(0, t)| \leq N \left[\|S_1 - S_2\|_{L^\infty(0,\tau)} + \|R_1 - R_2\|_{L^\infty(0,\tau)} + \|\tilde{u}_x\|_{L^\infty(\Omega_\tau)} \right] .$$

Using the regularity result (9.3), we can immediately estimate also $|J_1(0, t)|$ by the L^∞ norm of $|R_1 - R_2| + |\dot{R}_1 - \dot{R}_2|$. Hence

$$(9.13) \quad |\tilde{C}(0, t)| \leq N \left[\|\dot{R}_1 - \dot{R}_2\|_{L^\infty(0,\tau)} + \|S_1 - S_2\|_{L^\infty(0,\tau)} + \|\tilde{u}_x\|_{L^\infty(\Omega_\tau)} \right] .$$

Next we need to estimate $S_1 - S_2$. By the mean value theorem,

$$\hat{u}_1(S_1 - R_1, t) - \hat{u}_1(S_2 - R_2, t) = \hat{u}_{1,x}(y, t) \cdot (S_1 - R_1 - S_2 + R_2)$$

where y is a point between $S_1 - R_1$ and $S_2 - R_2$. Since $|u_{1,x}| \geq \nu_0 > 0$ (ν_0 constant which depends on the λ in (9.4)), we get

$$\begin{aligned}
|S_1(t) - S_2(t)| &\leq |R_1 - R_2| + \frac{1}{\nu_0} |\hat{u}_1(S_1 - R_1, t) - \hat{u}_1(S_2 - R_2, t)| \\
&= |R_1 - R_2| + \frac{1}{\nu_0} |\hat{u}_2(S_2 - R_2, t) - \hat{u}_1(S_2 - R_2, t)| \\
&\leq \|R_1 - R_2\|_{L^\infty(0,\tau)} + \frac{1}{\nu_0} \|\tilde{u}\|_{L^\infty(\Omega_\tau)} .
\end{aligned}$$

Substituting this into (9.13) and using the result in (9.12), we find after using also (9.10) and (9.11), that (9.12) becomes

$$\|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)} \leq N\tau \|\zeta_1 - \zeta_2\|_{L^\infty(0,\tau)} .$$

Hence $\zeta_1(t) = \zeta_2(t)$ if $0 \leq t \leq \tau$, τ small. This implies (by (9.11) and (9.10)) that $u_1 \equiv u_2$, $R_1 \equiv R_2$ and the two solutions coincide if $0 \leq t \leq \tau$. We similarly proceed step-by-step to complete the proof of uniqueness.

§10. $T_f < \infty$.

In this section we prove:

THEOREM 10.1. *If $D_A = D_B$ and $C_0(r) \not\equiv 0$ then $T_f < \infty$.*

We shall denote by T_∞^* the smallest t (if existing) such that $\zeta(t) \geq 1$, and first prove:

LEMMA 10.2. *If $T_f = \infty$ then $T_\infty^* < \infty$.*

Proof. Suppose the assertion is not true; then $\zeta(t) < 1$ for all $t \geq 0$. Since also $T_f = \infty$, $A(r, t) = u(r, t)$ for all (r, t) in a neighborhood of $\{(R(t), t), 0 \leq t < \infty\}$ and, consequently,

$$(10.1) \quad D_A \zeta u_r + (1 - \zeta)(u - A^*) = 0 \quad \text{if } r = R(t), \quad t > 0 ,$$

$$(10.2) \quad u_r = \alpha \dot{R} \quad \text{if } r = R(t), \quad t > 0 .$$

where $u = \lim_{K \rightarrow \infty} (A_K - B_K)$. Set

$$\sigma = \lim_{t \rightarrow \infty} \zeta(t); \quad \text{then } \sigma \leq 1 .$$

If $\sigma < 1$ then we can use the comparison argument as in the proof of (6.7) to deduce that

$$(10.3) \quad |u(r, t) - u_\infty(r)| \leq \widetilde{W}(r, t) \quad \text{for all } t > 0 ,$$

where \widetilde{W} is defined in (6.16) and $u_\infty(r)$ is the harmonic function in $\rho < r < \infty$ where ($\rho = \lim_{t \rightarrow \infty} R(t)$) satisfying

$$(10.4) \quad -D_A \sigma u_{\infty, r}(\rho) + (1 - \sigma)(u_\infty(\rho) - A^*) = 0 , \quad u_\infty(\infty) = -B^* .$$

Next we argue as in Lemma 6.2 (with $t_0 = 0$ in (6.19)) and deduce that

$$(10.5) \quad C_{K,r}(R_K(t), t) \geq \frac{M}{t^6} \quad (M > 0, \quad t \geq 1) .$$

where M is independent of K . Hence

$$(10.6) \quad C_r(R(t), t) \geq \frac{M}{t^6} \quad (M > 0, \quad t \geq 1) .$$

Using this estimate and (10.1), (10.2), we can proceed as in the proof of Theorem 6.3, deriving (6.22) for $\eta(t) = \zeta(t) - 1$, and concluding that if

$$(10.7) \quad \frac{1}{t_n} \int_0^{t_n} (A^* - u(R(\tau), \tau)) d\tau \geq \alpha_0 > 0$$

for a sequence $t_n \rightarrow \infty$, then $\eta(t_n) > 0$ for t_n large enough, which is a contradiction to the assumption that $\zeta(t) < 1$ for all $t > 0$. Hence

$$(10.8) \quad \frac{1}{t} \int_0^t u(R(\tau), \tau) d\tau \rightarrow A^* \quad \text{as } t \rightarrow \infty .$$

This together with (10.3) implies that $u_\infty(\rho) = A^*$. Hence u_∞ takes its maximum at $r = \rho$ and, by the maximum principle, $u_{\infty,r}(\rho) < 0$. However, since $u_\infty(\rho) = A^*$, we must also have $u_{\infty,r}(\rho) = 0$ by (10.4), which is a contradiction. We conclude that $\sigma = 1$, i.e., $\zeta(\infty) = 1$ and then, by (10.1),

$$(10.9) \quad u_r(R(t), t) \rightarrow 0 \quad \text{if } t \rightarrow \infty .$$

By (10.9) and a comparison argument (as in the proof of Lemma 6.1) we then get

$$u(r, t) \leq \widetilde{W}(r, t)$$

where \widetilde{W} is defined in (6.16); therefore

$$(10.10) \quad u(R(t), t) \rightarrow 0 \quad \text{if } t \rightarrow \infty .$$

Observe next that the argument that led from (10.5) to (10.8) is independent of whether $\sigma < 1$ or $\sigma = 1$. Thus (10.8) still holds and this is a contradiction to (10.10). The proof that T_∞^* is finite is thereby complete.

Proof of Theorem 10.1. Suppose $T_f = \infty$. Then, by Lemma 10.2, $T_\infty^* < \infty$. Therefore

$$u_r(\overline{R}, t) = 0 \quad \text{if } t > T_\infty^* \quad (\overline{R} = R(T_\infty^*)) .$$

Since also $u_t = D_A \Delta u$ and $u(\infty, t) = -B^*$, it follows by the comparison argument that

$$|u(r, t) + B^*| \leq M \frac{e^{-\frac{r^2}{\sqrt{D_A} 4t}}}{t^{3/2}} \quad \text{if } t > T_\infty^* .$$

Consequently, $u(r, t) < 0$ for all $r \geq R(t)$, $t \geq T$ provided T is sufficiently large, and this is a contradiction to the assumption that $T_f = \infty$ (which implies that u remains positive for all $t > 0$).

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