

**First, thanks for the organizers for inviting me here**

---

and **thanks to Keith**

...for hiring me here in 1992 for my 2<sup>nd</sup> post doc and the first one “truly away from home”.

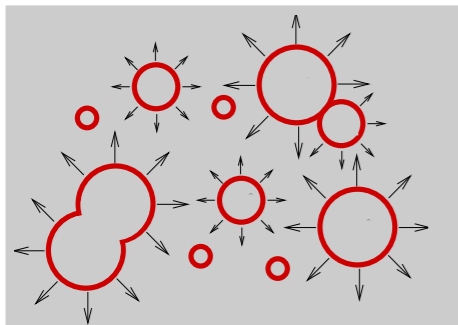
It was with Keith that I first started working on **Baryogenesis** (protecting primordial BAU), that eventually led me to work on **EWBG**

This was (is) a great place with great people, **Keith** obviously, but also **Arkadi, Larry, Mischa S and Mischa V**,... and a bunch of great post docs including **Jim Cline** (who sends his regards!).



# Baryogenesis from Dark Sector

Kimmo Kainulainen,  
Olivefest 2017 / 17.5.2017



with:

Tommi Alanne, **Jim Cline**, Pat Scott,  
Mike Trott, Kimmo Tuominen,  
Ville Vaskonen, Christoph Weniger,  
Dave Tucker-Smith, Sami Nurmi,  
Tommi Tenkanen ...

**SM almost UV-complete:**

- simplicity as a guiding principle, solve issues with **DM** and **BAU**

**EWBG:** Obvious models are ~dead:  
~~MSM, MSSM, 2HDM...~~

- Transition strength (TS)?
- CP-violation ?

**EWBG goes Dark Side**

- Higgs portal/singlet models (DM,TS)
- CP-portal

# UV-completion

May we have just (almost) SM all the way to the Planck scale?

Espinosa, Giudice, Riotto, JCAP 0805 (2008) 002  
 Degrassi et al, JHEP 1208 (2012) 098

Possible running to negative coupling can be cured for example by a **singlet S**

$$\frac{1}{2} \lambda_{hs} |H|^2 S^2 \Rightarrow \beta(\lambda) \rightarrow \beta(\lambda)_{SM} + \frac{1}{2} \lambda_{hs}^2 \quad (\lambda_{hs} \approx 0.7)$$

~~FT. UNIFICATION~~ →

problem of gauge Landau poles

~~HIERARCHY PROBLEM~~ →

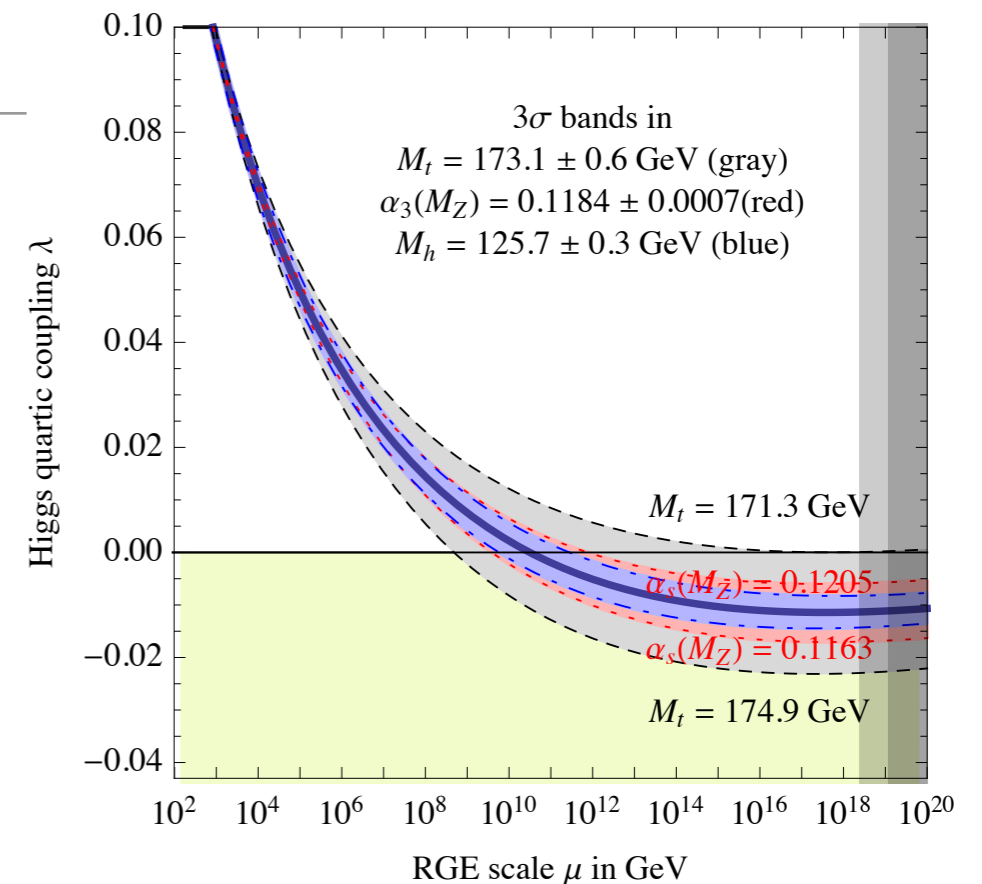
behaviour of relevant operators

**Asymptotic safety** ✓

with gravity corrections, or...

**No intermediate scales** ✓

(technically)



Robinson and Wilczek, PRL 96, 231601 (2006)  
 Wetterich and Shaposhnikov, Phys.Lett. B683 (2010)

Remain the problems of

**BARYON ASYMMETRY**  
**DARK MATTER**

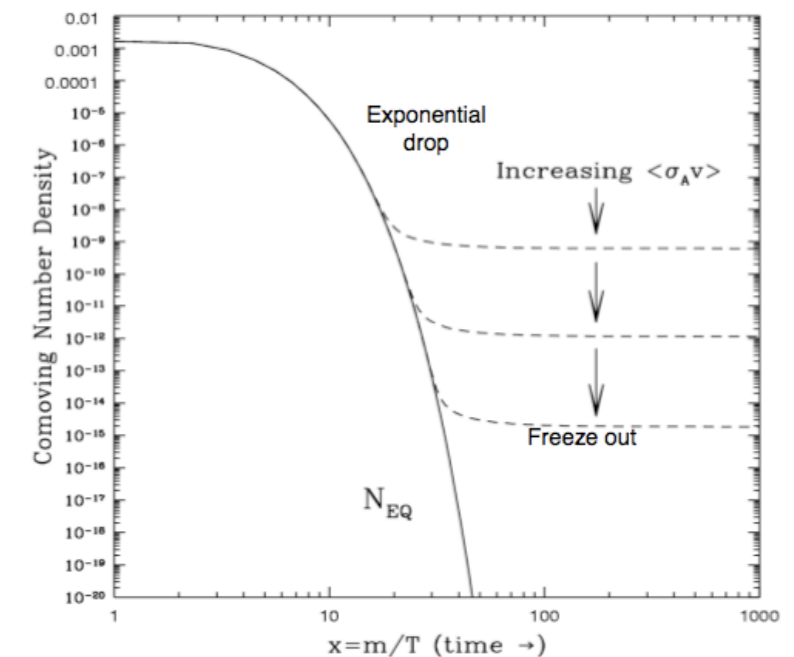
If **minimality** is the guiding principle, these **should** rather be solved at **low E**

# Dark Matter (Thermal) **WIMP** (Freeze-out)

- ◆ Cold
- ◆ Non-baryonic
- ◆ Neutral, or weakly interacting
- ◆ **Stable** (by some discrete symmetry)

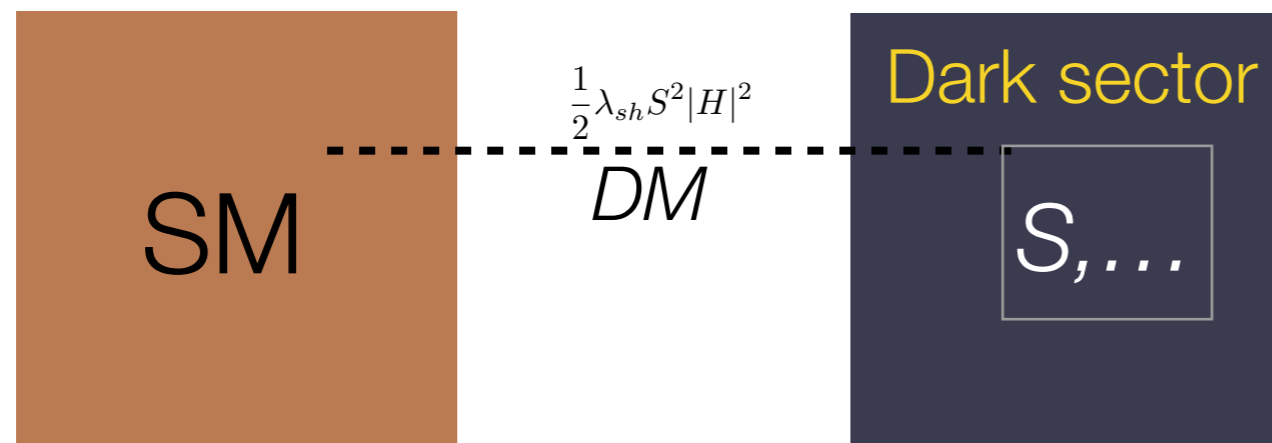
DM DM  $\rightarrow$  SM SM  
DM  $\not\rightarrow$  SM + ...

$$\Omega_{DM} \sim \frac{1}{\langle v\sigma \rangle}$$



## Popular context:

Simple **Portals to** (perhaps complex) **Dark sector**

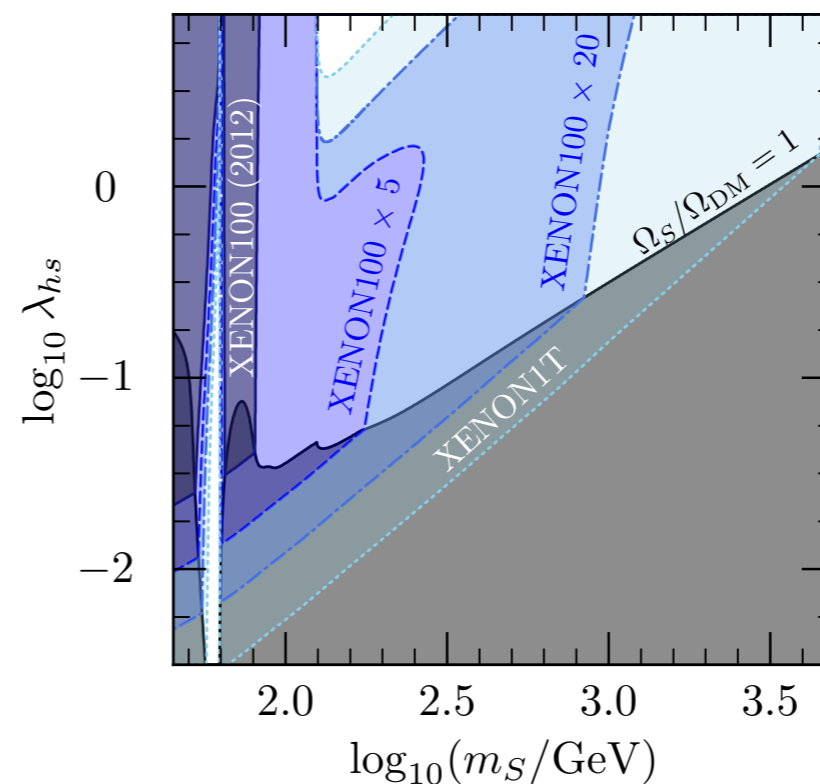
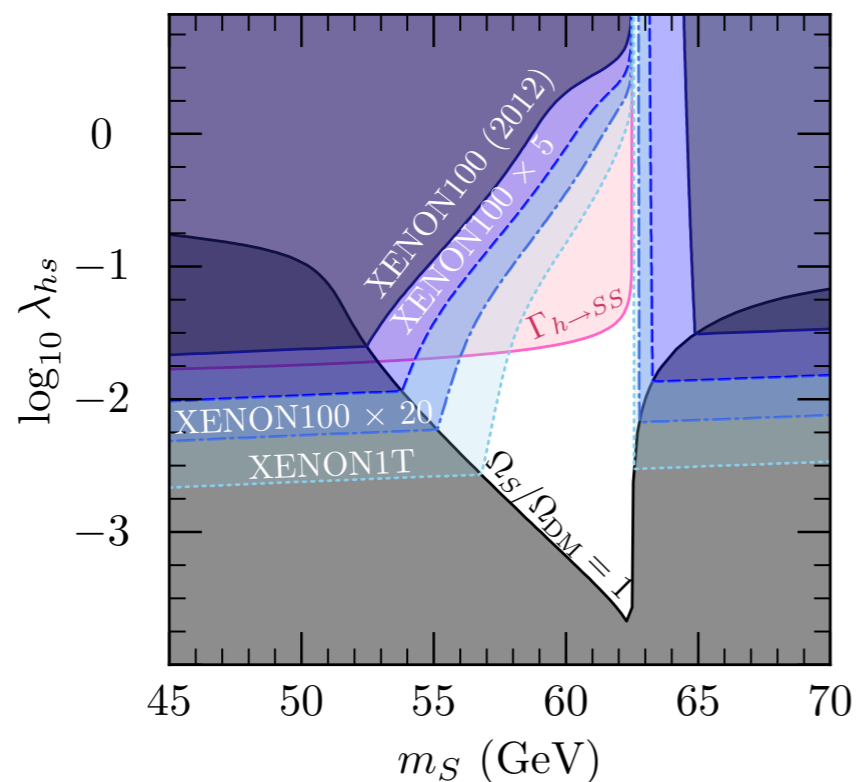


- J.M.Cline, KK,  
JHEP 1111 (2011) 089  
JCAP 1301 (2013) 012  
Phys.Rev. D87 (2013)7,071701
- J.M.Cline, KK, P.Scott, C.Weniger  
PRD88 (2013) 055025
- J.M.Cline, KK. D.Tucker-Smith,  
in progress.
- KK, K.Tuominen and V.Vaskonen  
Phys.Rev. D93 (2016) 7 ,015016
- T.Alanne, KK, K.Tuominen, V.Vaskonen  
JCAP 1608 (1016) no08, 057,
- KK, S.Nurmi, T.Tenkanen,  
K.Tuominen and V.Vaskonen  
JCAP 1606 (2016) no.06, 022

# Solution: DM (Given $Z_2$ -symmetry) DM can be singlet scalar

The (simplest possible) model:

$$V = V_{\text{MSM}} + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{sh} S^2 |H|^2 + \frac{1}{4}\lambda_s S^4$$



Xenon bounds account for the fact that  $f_{\text{rel}} \leq 1$ .

J.M.Cline, KK, P.Scott, C.Weniger  
PRD88 (2013) 055025

$$\Omega \sim \frac{1}{\langle v_{\text{Mol}} \sigma \rangle} \sim \frac{1}{\lambda_{hs}^2}$$

MANY (more complicated) alternatives also work



# EWBG

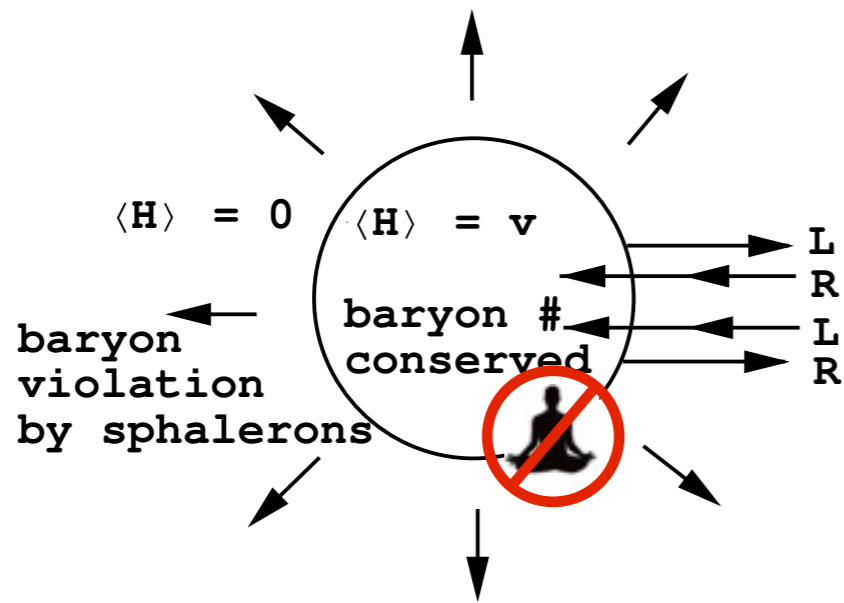
\* Out-of-eq => BSM

$$H \sim 10^{-14} T_{100}^2 \text{ GeV}$$

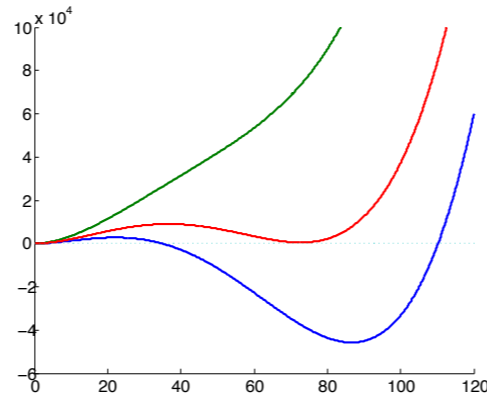
$$\Gamma \sim 10^{-5} T_{100} \text{ GeV}$$



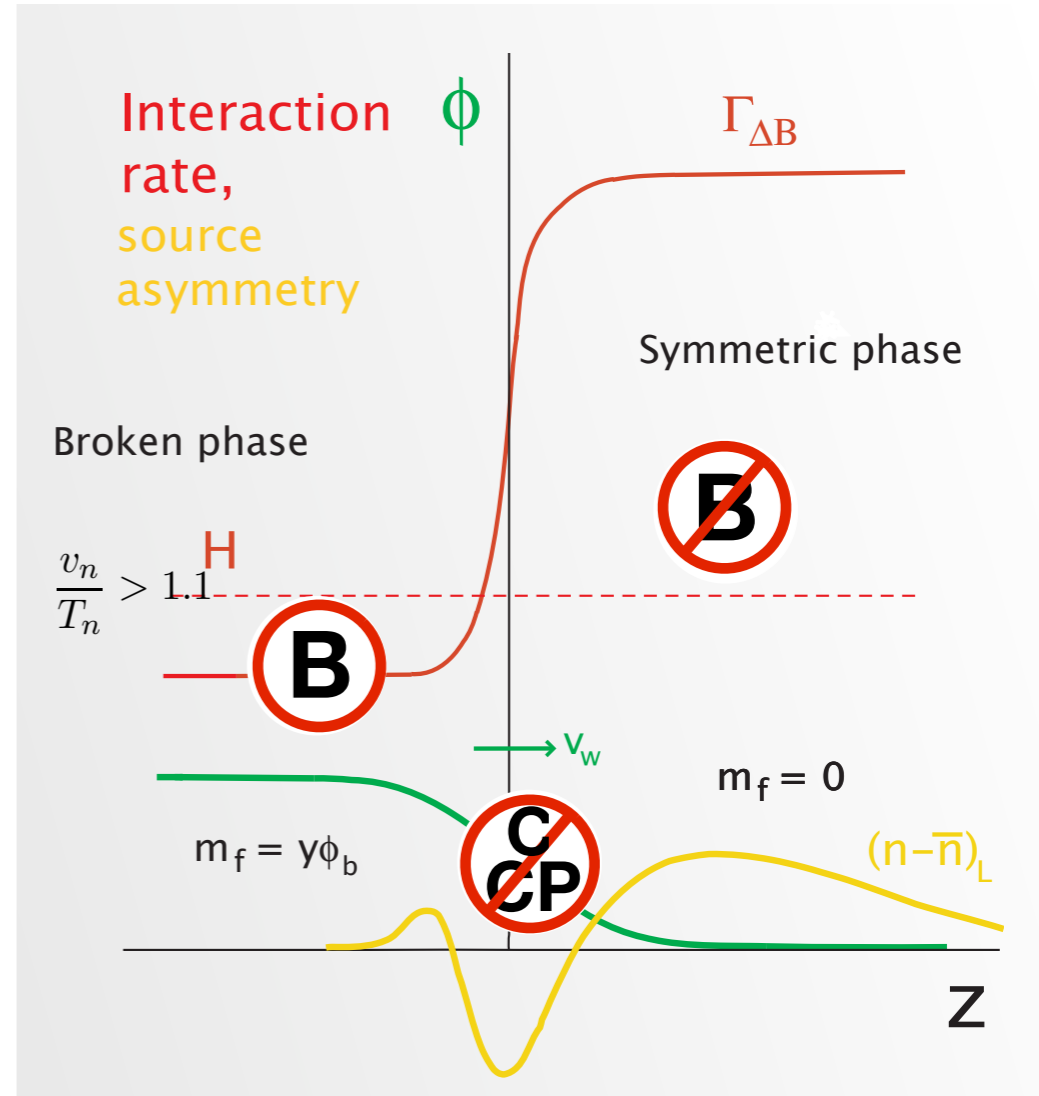
=> need a bubble



$$\frac{1}{t_w} \sim \frac{v_w}{L_w} \sim 10^{-5} T_{100}$$



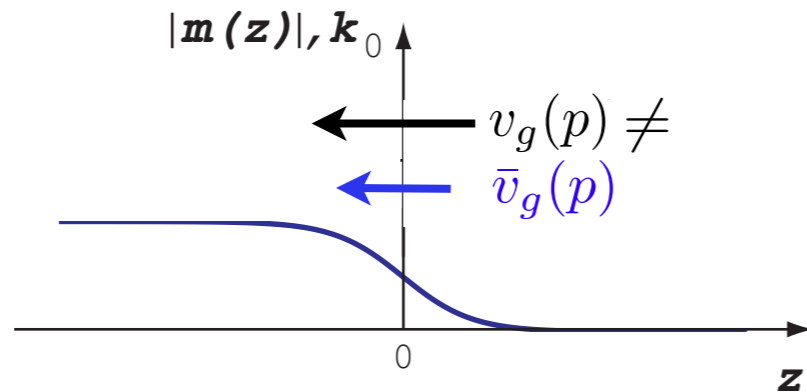
Sakharov:



May be killed by lack of

- \* Strong transition
- \* CP-violation

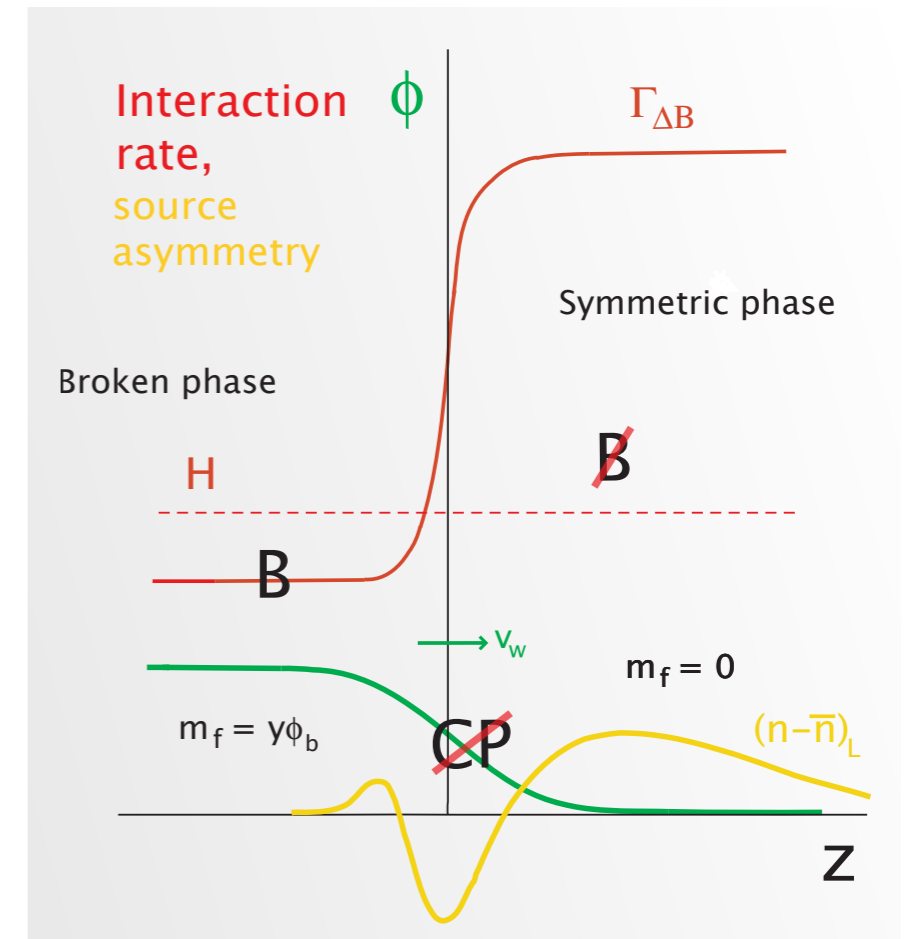
# Creating baryons / CP-violation / Semiclassical limit



$$(\partial_t + \mathbf{v}_g \cdot \partial_{\mathbf{x}} + \mathbf{F} \cdot \partial_{\mathbf{p}}) f_i = C[f_i, f_j, \dots].$$

$$v_g = \frac{p_0}{\omega} \left( 1 + s_{CP} \frac{s|m|^2 \theta'}{2p_0^2 \omega} \right) \quad m = |m(z)| e^{i\theta(z)}$$

$$F = -\frac{|m||m'|}{\omega} + s_{CP} \frac{s(|m|^2 \theta)'}{2\omega^2}. \quad \text{CP-force}$$



NEEDS LARGE, CHANGING CP-VIOLATING PHASES AT THE TRANSITION REGION

**WKB** J.M. Cline, M. Joyce and K. Kainulainen. PLB417 (1998) 79; JHEP 0007 (2000) 018  
 J.M. Cline and K. Kainulainen, PRL85 (2000) 5519.  
 M. Joyce, T. Prokopec and N. Turok, PRD53, 2958 (1996); PRL75, 1695 (1995); PRD53, 2930 (1996).

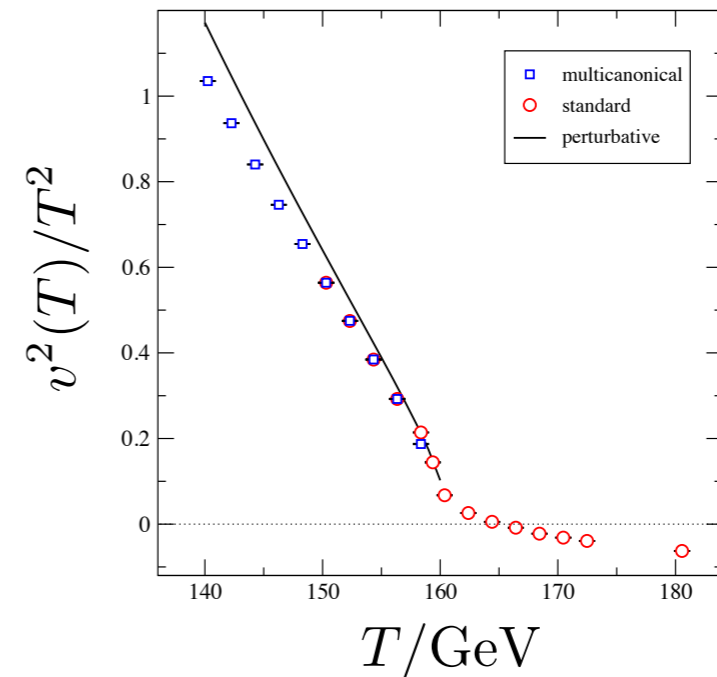
**CTP** K. Kainulainen, T. Prokopec, M.G. Schmidt and S. Weinstock, JHEP 0106, 031 (2001); PRD66 (2002) 043502.  
 T. Prokopec, M.G. Schmidt and S. Weinstock, Annals Phys. 314, 208 (2004), Annals Phys. 314, 267 (2004)

# MSM Transition strength / B-violation rate

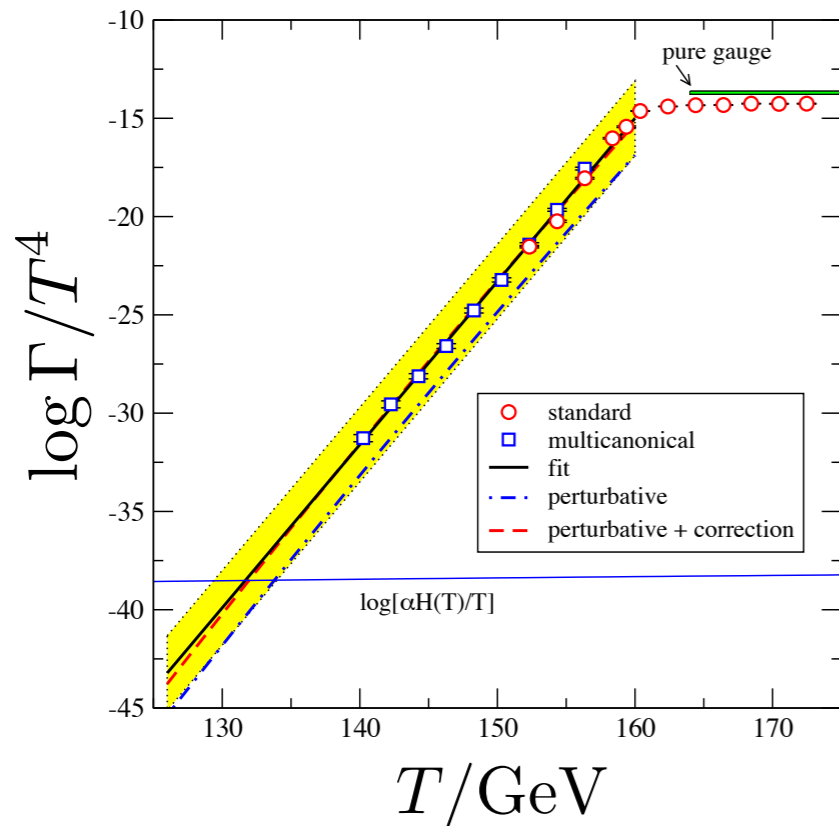
PT in SM, is a **cross-over** at

$$T = 159.6 \pm 0.1 \pm 1.5 \text{ GeV}$$

M.d'Onofrio, K.Rummukainen, A.Tranberg,  
Phys.Rev.Lett. 113 (2014) no.14, 141602



**no  
bubble**



## Sphaleron rate in SM

$$\Gamma_{\text{Symm.}}/T^4 = (8.0 \pm 1.3) \times 10^{-7} \approx (18 \pm 3)\alpha_W^5$$

$$\log \frac{\Gamma_{\text{Broken}}}{T^4} = (0.83 \pm 0.01) \frac{T}{\text{GeV}} - (147.7 \pm 1.9)$$

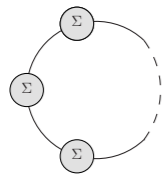
Sphalerons in equilibrium until  $T \approx 132 \text{ GeV} < T_{PT} \Rightarrow$  **washout !**



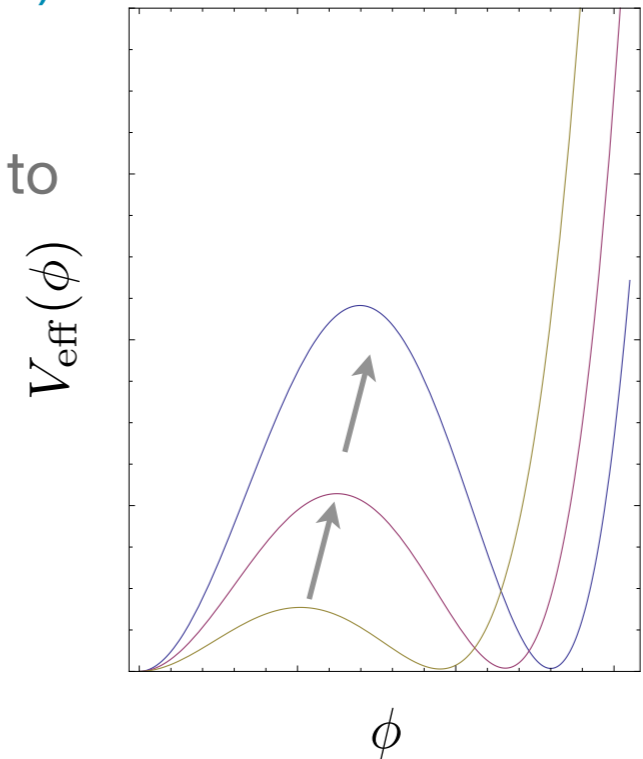
# MSSM LSS - transition strength

Traditionally: **increase the strength by (effective cubic) loop corrections**

Need new **light** ( $m_i < T$ ) **bosonic** fields strongly coupled to Higgs



$$\delta V_{\text{eff}} = - \sum_i \frac{T m_i^3(\phi, T)}{12\pi} + \dots$$



=> **Light Stop Scenario** in **the MSSM** and NMSSM

[Carena, Quiros, Wagner (1996), ... Espinosa, Quiros, Zwirner, Laine, Cline, KK, Losada, ...]

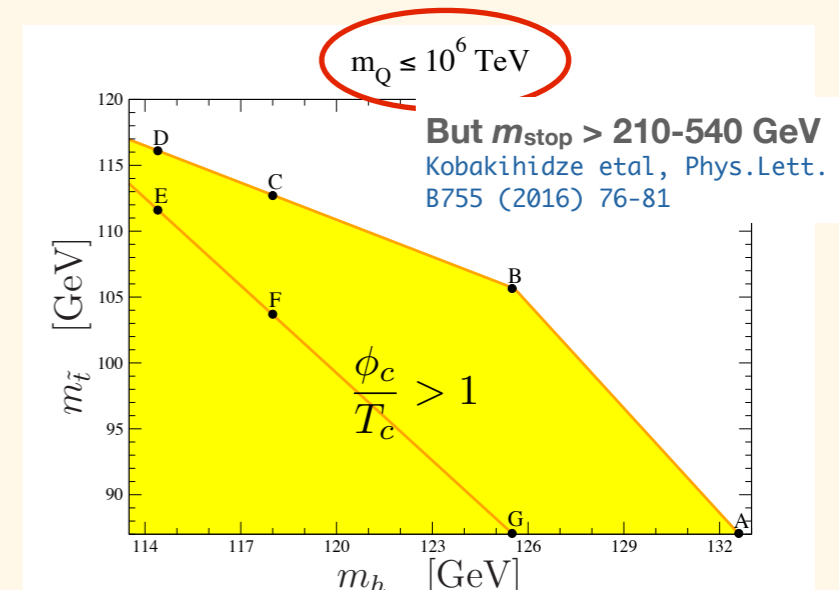
However, higgs mass mostly from squarks:

$$m_h^2 \sim y_t^2 \log \frac{m_{t_R}^2 m_{t_L}^2}{m_t^4}$$

**Tension:** light  $t_R$  => **very heavy  $t_L$**

.. and this just has gotten too tight ...

Carena, et al. 2009, 2013

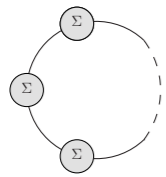


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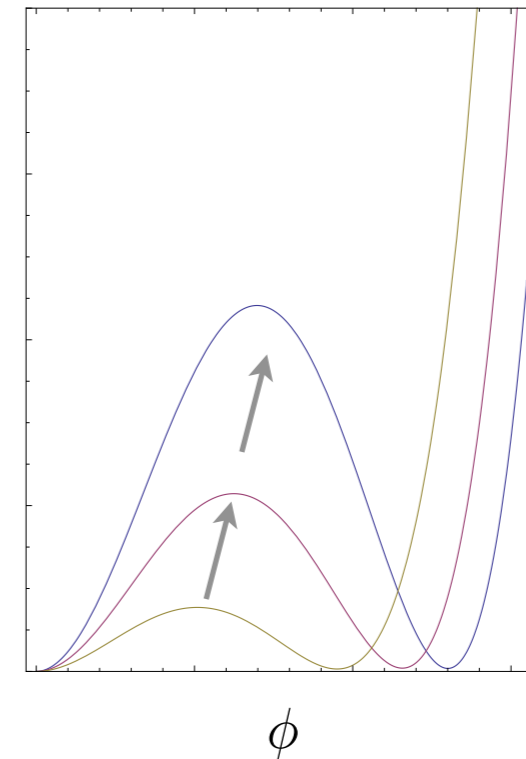
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$V_{\text{eff}}(\phi)$



Leszek Roszkowski (Nordita, June 2015):

“SUSY cannot be disproved, only abandoned”

(the actual content misunderstood here in purpose)

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[Carena, Quiros, Wagner (1996), ... Espinosa, Quiros, Zwirner, Laine, Cline, KK, Losada, ...]

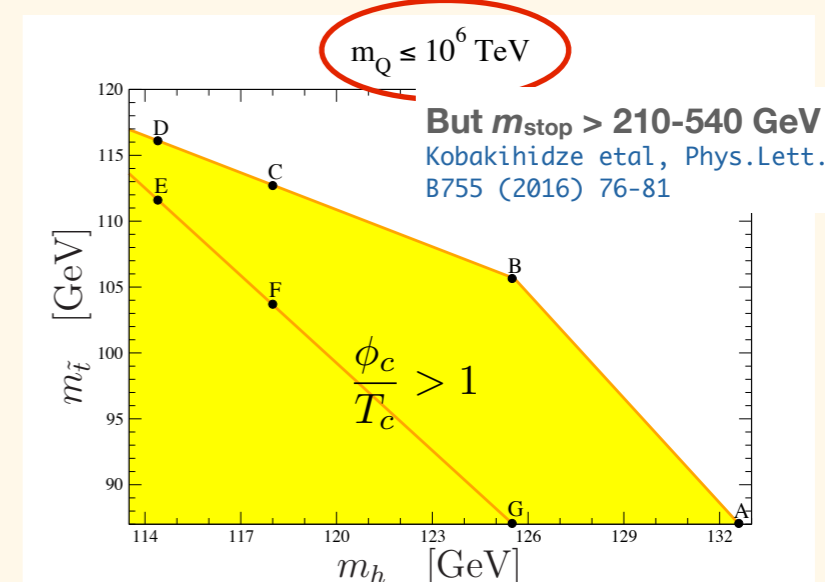
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# 2HD models (no DM )

many new CP-violating pms.

Full GL(2,C)-reparametrization invariant potential

J.Cline, KK, M.Trott, JHEP 1111 (2011) 089

$$\begin{aligned}
 V = & \frac{\lambda}{4} \left( H^{\dagger i} H_i - \frac{v^2}{2} \right)^2 + m_1^2 (S^{\dagger i} S_i) + (m_2^2 H^{\dagger i} S_i + \text{h.c.}) \\
 & + \lambda_1 (H^{\dagger i} H_i) (S^{\dagger j} S_j) + \lambda_2 (H^{\dagger i} H_j) (S^{\dagger j} S_i) + [\lambda_3 H^{\dagger i} H^{\dagger j} S_i S_j + \text{h.c.}] , \\
 & + [\lambda_4 H^{\dagger i} S^{\dagger j} S_i S_j + \lambda_5 S^{\dagger i} H^{\dagger j} H_i H_j + \text{h.c.}] + \lambda_6 (S^{\dagger i} S_i)^2 , \\
 & + y_t \bar{t}_L (H^{0*} \delta_{ti} + (\eta_U \delta_{ti} + \eta'_U \gamma_{tb}^* V_{bi})) S^{0*} q_R^i
 \end{aligned}$$

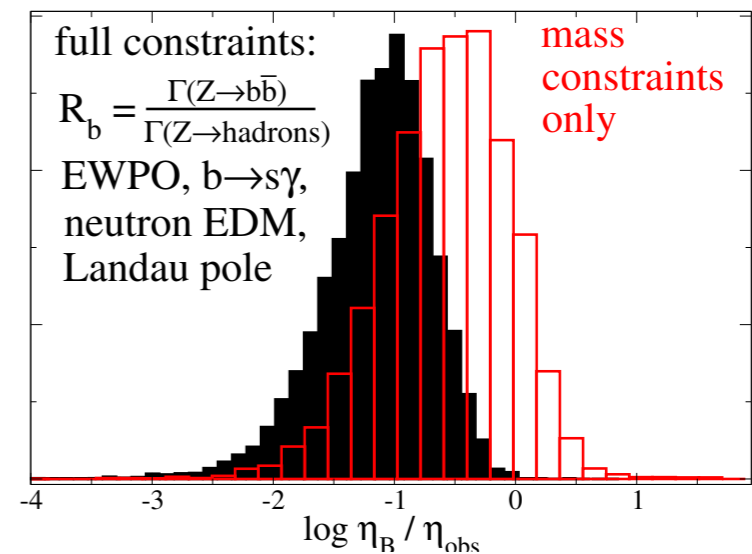
MFV to avoid FCNC (Yukawa sector)

G.C.Branco, W.Grimus & L.Lavoura, PLB380 (1996) 119

Strong EWPT and large BAU points were rare then:  $<1/10^4$ .

**Strongest bounds always from EDM's, so ACME killed these models**

ACME Collaboration, Science 343 (2014) 269-272



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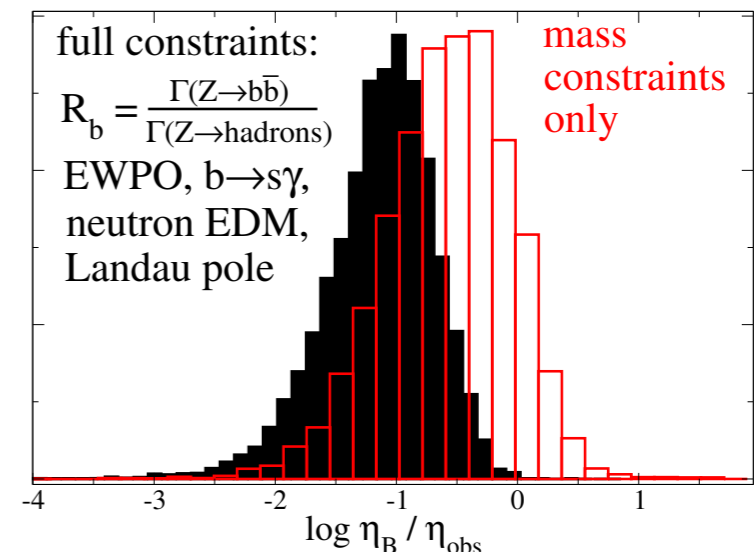
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## Strong transition and strong CP-violation

Generic issue:

Strong transition > large  $|\lambda_i$ 's  
Large B > large phases

-> large EDM's

Need something to alleviate burden on 2HDM  $\lambda_i$ 's

# Singlet model Strong transition

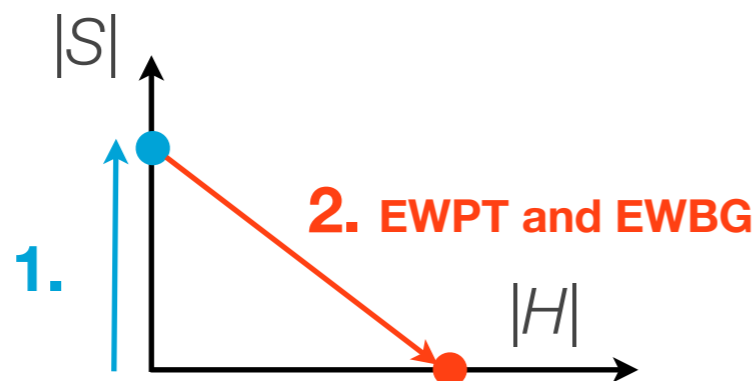
Anderson, Hall, PRD45, 2685 (1992)  
Profumo, Ramsey-Musolf, Shaughnessy,  
JHEP 0708 (2007) 010

## Use tree level barrier:

J.R.Espinosa, T.Konstandin, F.Riva, NPB854 (2012) 592

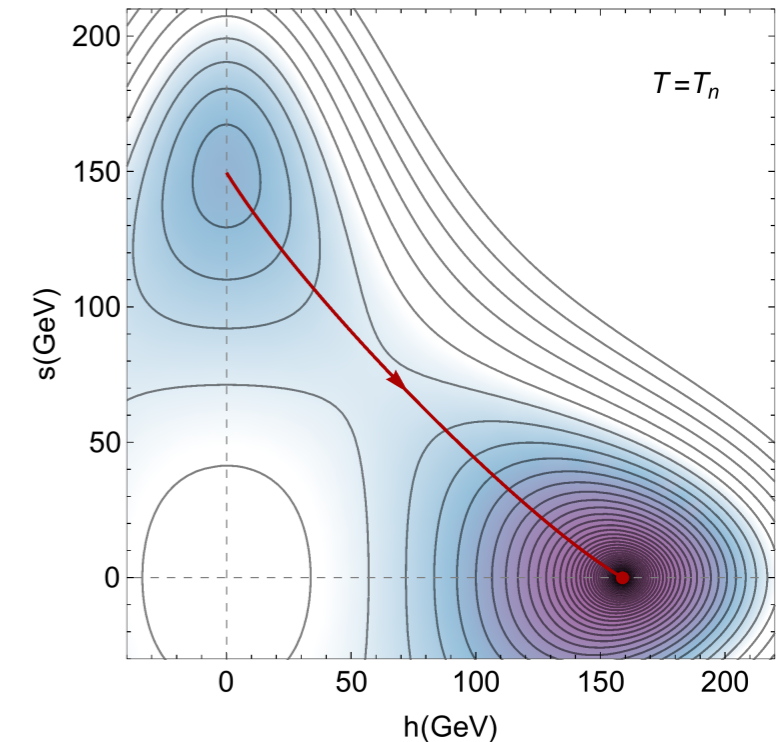
$$V = \frac{1}{2} \lambda_{hs} h^2 s^2 - (\mu_s^2 - c_s T^2) s^2 - (\mu_h^2 - c_h T^2) h^2 + \dots$$

only the leading high-T



Arrange  $c_s$ ,  $c_h$ ,  $\mu_s$  and  $\mu_h$  so that transition goes in *two steps*  
→ large barrier at  $T_c$  → strong transition. This WORKS

- large barrier requires largish  $\lambda_{hs}$
- 2-step mechanism needs small  $m_s$



### Variants of the scheme:

Inoue, Ovanessian, Ramsey-Musolf,  
PRD93 (2016) 015013, etc...

**Idea actually present** in the  
MSSM ~color breaking

Laine, Rummukainen,  
Cline, Moore, Quiros ...



# Singlet model Strong transition *and* S-DM?

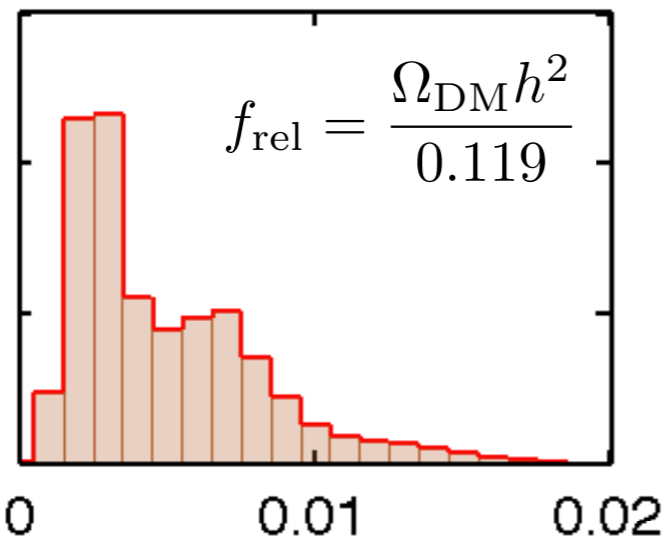
However, large  $\lambda_{hs}$  gives small  $\Omega$ :

$$\Omega \sim \frac{1}{\langle v_{Mol}\sigma \rangle} \sim \frac{1}{\lambda_{hs}^2}$$

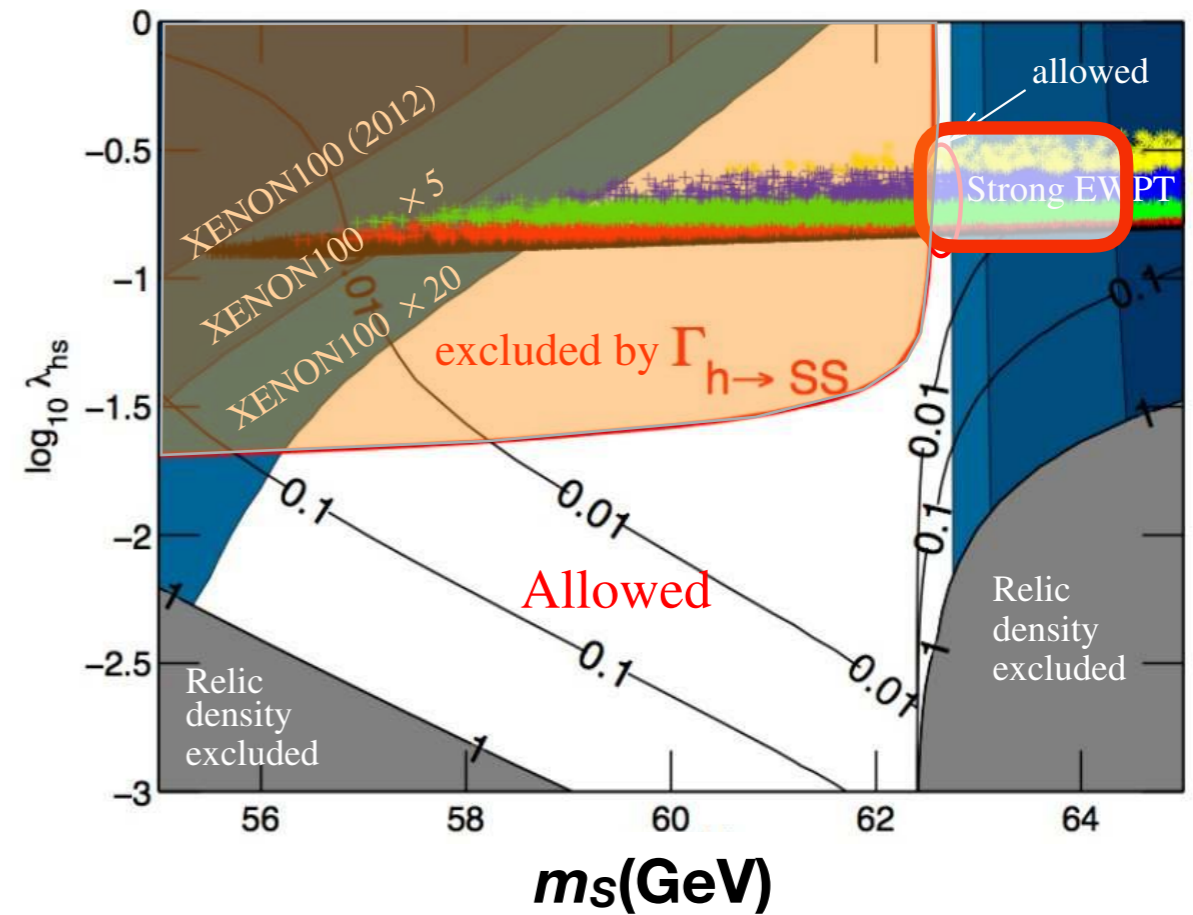
Strong transition implies a **subdominant DM**



BAU acceptable  $v/T > 1$  models



J.M. Cline, KK, JCAP 1301 (2013) 012



J.M.Cline, KK, P.Scott and C.Weniger, PRD88 (2013) 055025



**Further extensions with more singlets, or new fermions, or ...**

KK, K.Tuominen and V.Vaskonen, PRD93(2016) 7,015016



# Singlet model Strong transition *and* BAU?

## BAU from top transport

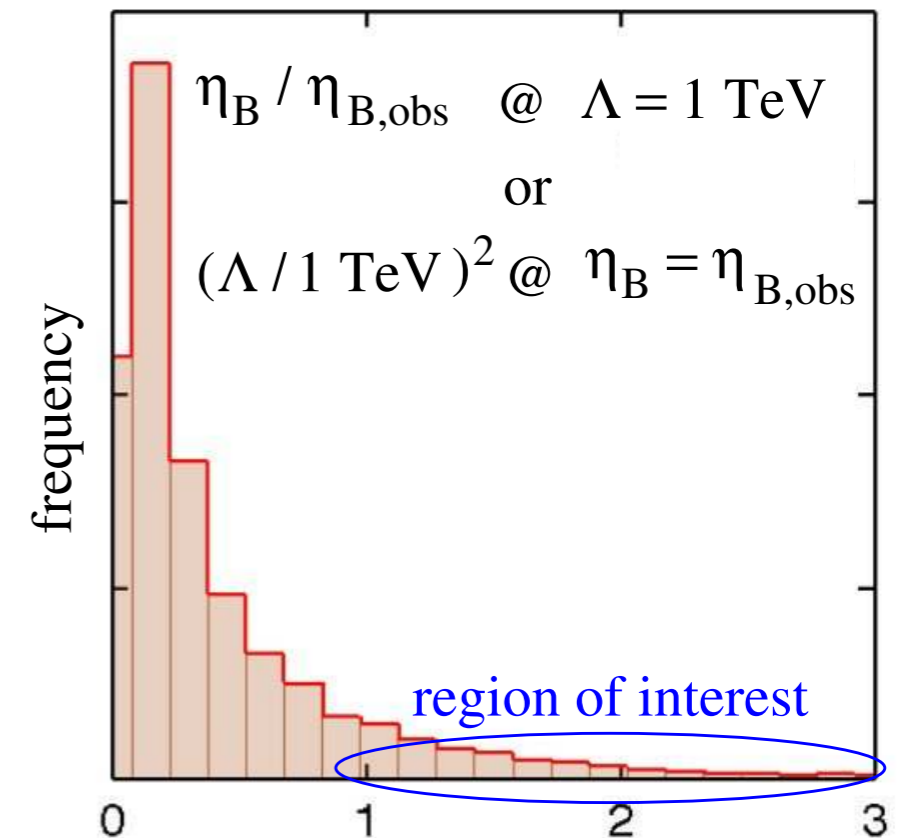
DM stability =>  $Z_2$  symmetry:  $\langle S \rangle_{T=0} = 0$

**CP violation, NEW PHYSICS:** Dim-6 operator  
(If not DM could take Dim-5 as well) Espinosa, etal

$$y_t \bar{Q}_L H \left( 1 + \frac{\eta}{\Lambda^2} S^2 \right) t_R + \text{h.c.} \quad (\eta \equiv i)$$

$$m_t(z) = \frac{y_t}{\sqrt{2}} h(z) \left( 1 + i \frac{S^2(z)}{\Lambda^2} \right)$$

J.M. Cline, KK, JCAP 1301 (2013) 012



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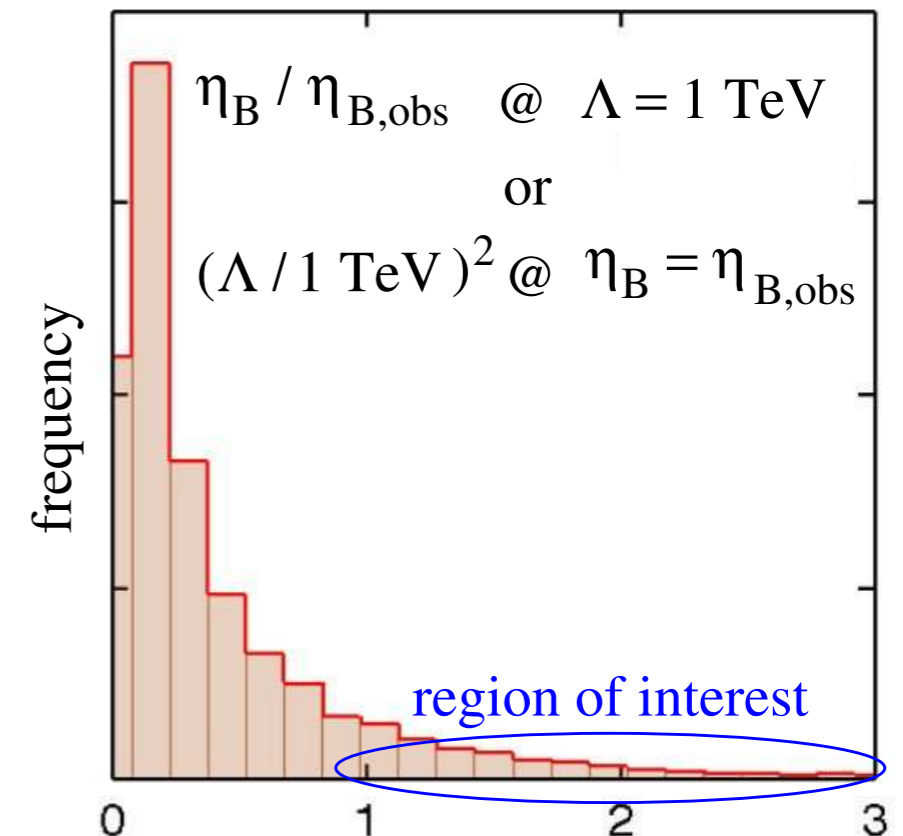
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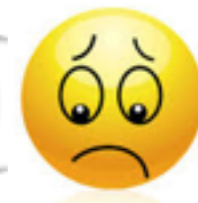
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J.M. Cline, KK, JCAP 1301 (2013) 012



However, there goes the UV-completion



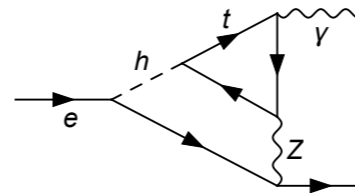


$$\mathcal{L}_{\text{scalar}} = Z^{ij} (D^\mu H_i)^\dagger D_\mu H_j + \frac{1}{2} (\partial_\mu S)^2 - V(H_1, H_2)_{2\text{HDM}} - \left[ \frac{1}{4} \lambda_S S^4 + \frac{1}{2} \lambda_{S1} S^2 |H_1|^2 + \frac{1}{2} \lambda_{S2} S^2 |H_2|^2 + \left( \frac{1}{2} \lambda_{S12} S^2 H_2^\dagger H_1 + \text{h.c.} \right) \right]$$

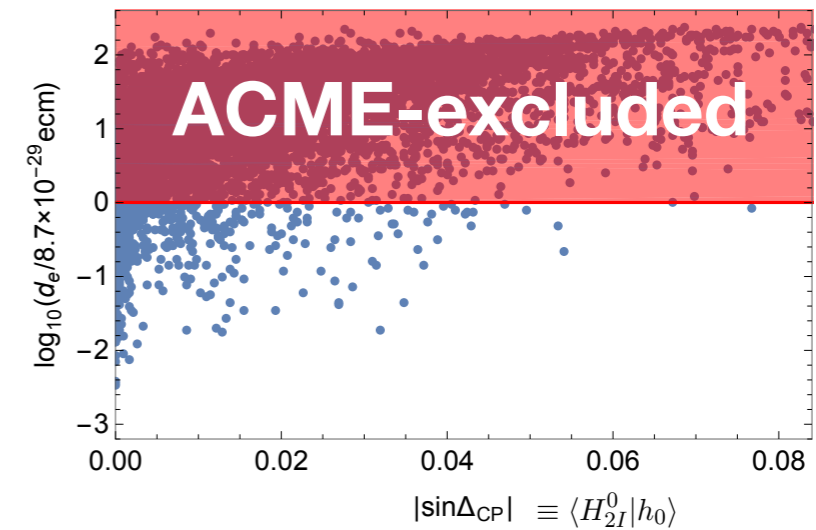
# 2HD+S model ~~CP~~ from 2HDM - strength from S / constraints

MCMC scan of parameter space with:

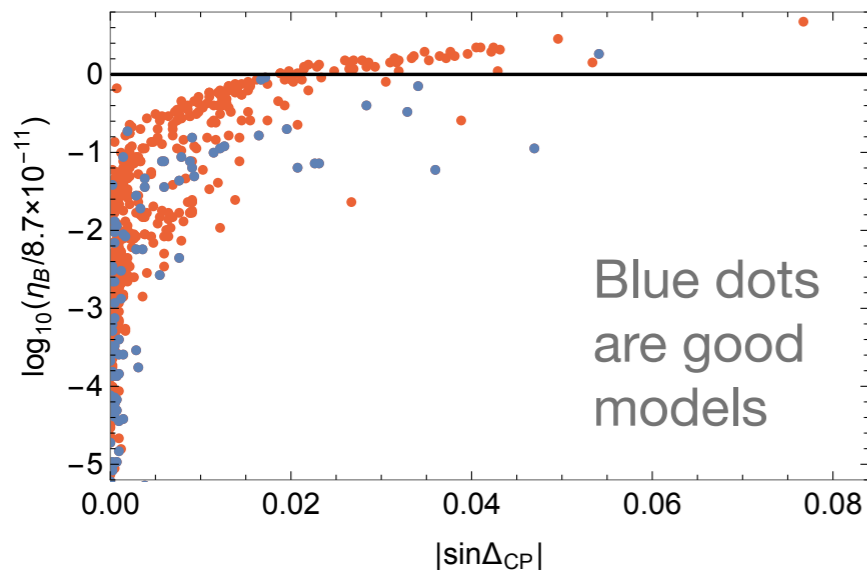
- Accelerator constraints
- EWP-data
- **EDM's (electron, ACME)**
- **Perturbativity up to 1.5 TeV**
- **LUX-limits on DM (singlet s)**
- **Strong transition (=>s is subleading DM)**



ACME Collaboration, Science 343 (2014) 269-272



If all above ok, compute BAU



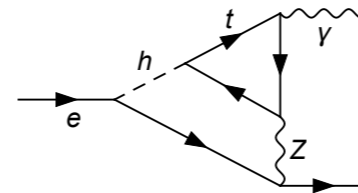
T.Alanne, KK, K.Tuominene and V.Vaskonen JCAP 1608 (2016) no.08, 057

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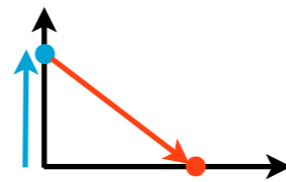
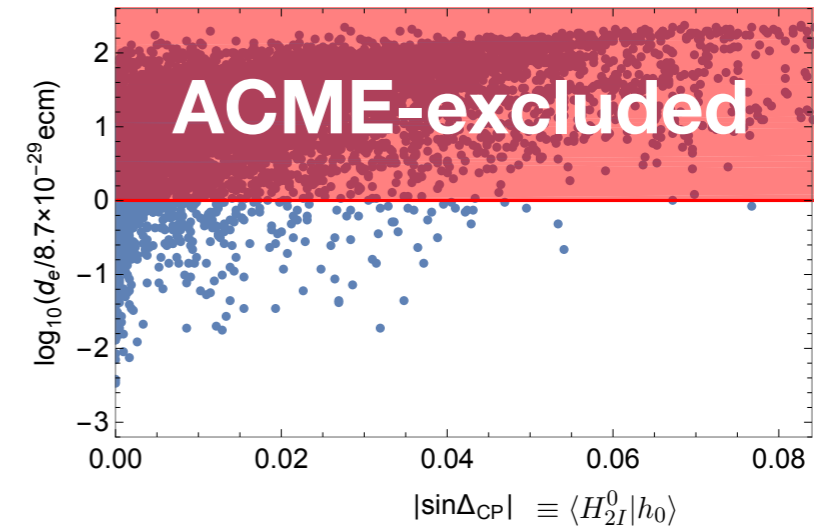
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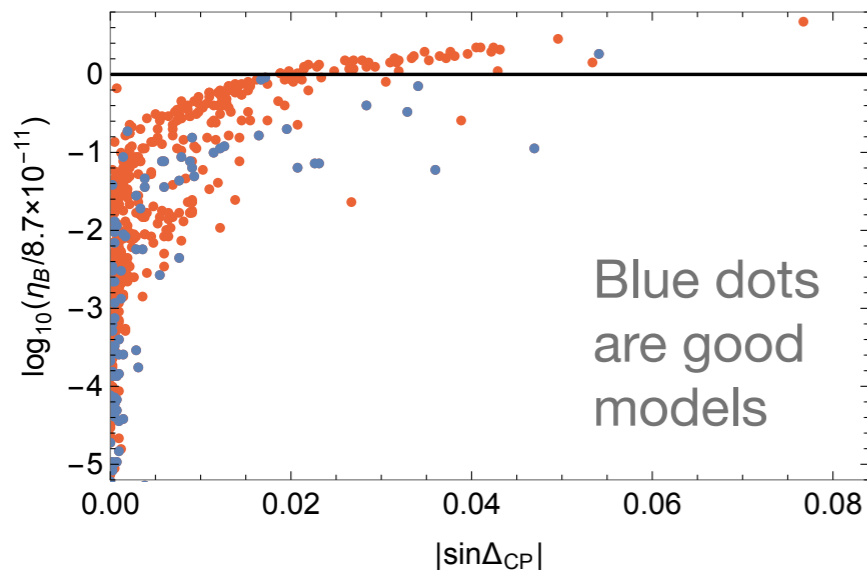
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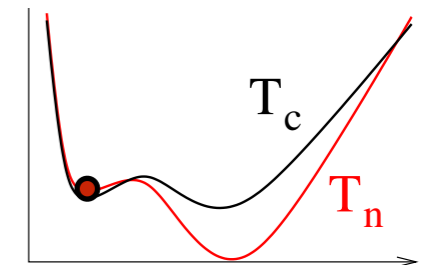


If all above ok, compute BAU



The bubble nucleation rate

$$\Gamma \sim T^4 \left( \frac{S_3(T)}{2\pi T} \right)^{3/2} \exp \left( -\frac{S_3(T)}{T} \right)$$



is slow if  $T_n \ll T_c$  (lot of supercooling) =>

**Danger of being trapped => No EW-breaking**

T.Alanne, KK, K.Tuominene and V.Vaskonen JCAP 1608 (2016) no.08, 057

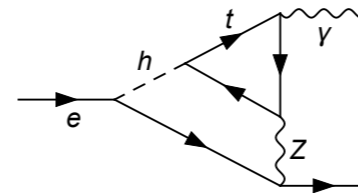
Indeed: Models in red: no  $T_n$

$$\mathcal{L}_{\text{scalar}} = Z^{ij} (D^\mu H_i)^\dagger D_\mu H_j + \frac{1}{2} (\partial_\mu S)^2 - V(H_1, H_2)_{2\text{HDM}} - \left[ \frac{1}{4} \lambda_S S^4 + \frac{1}{2} \lambda_{S1} S^2 |H_1|^2 + \frac{1}{2} \lambda_{S2} S^2 |H_2|^2 + \left( \frac{1}{2} \lambda_{S12} S^2 H_2^\dagger H_1 + \text{h.c.} \right) \right]$$

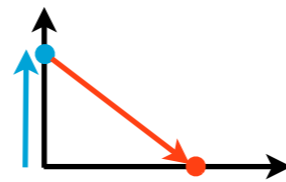
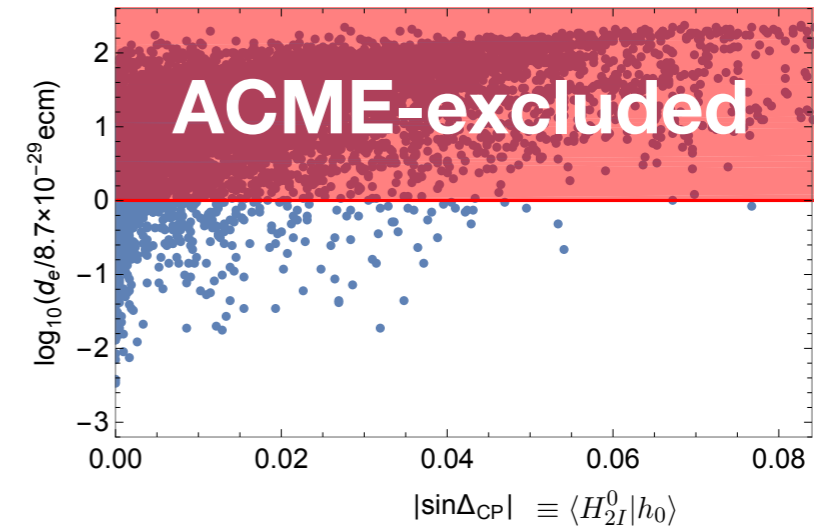
# 2HD+S model ~~CP~~ from 2HDM - strength from S / constraints

MCMC scan of parameter space with:

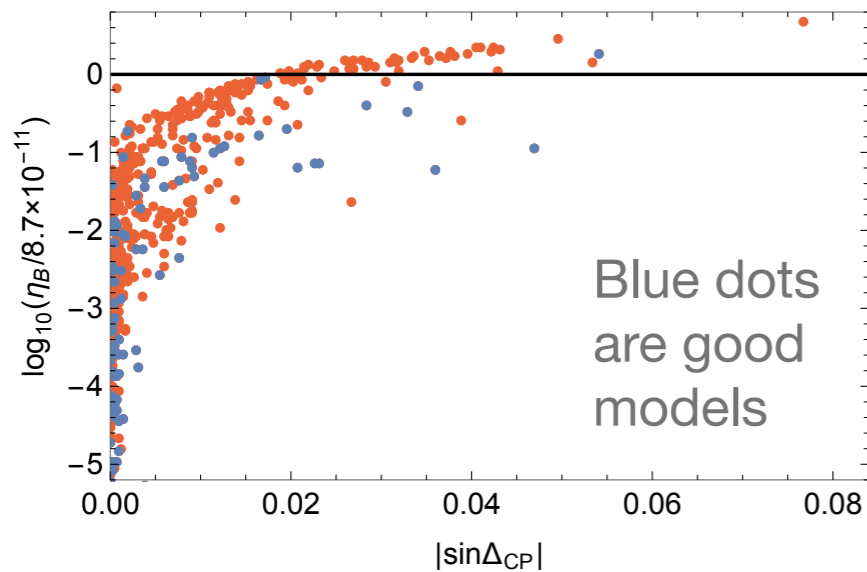
- Accelerator constraints
- EWP-data
- **EDM's (electron, ACME)**
- **Perturbativity up to 1.5 TeV**
- **LUX-limits on DM (singlet s)**
- **Strong transition** (=>s is subleading DM)



ACME Collaboration, Science 343 (2014) 269-272



If all above ok, compute BAU

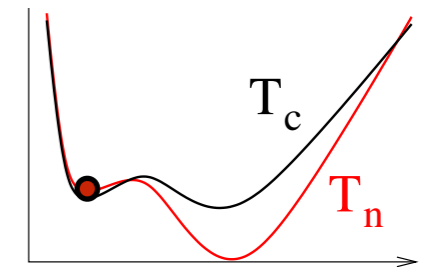


T.Alanne, KK, K.Tuominene and V.Vaskonen JCAP 1608 (2016) no.08, 057

Indeed: Models in red: no  $T_n$

The bubble nucleation rate

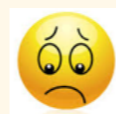
$$\Gamma \sim T^4 \left( \frac{S_3(T)}{2\pi T} \right)^{3/2} \exp \left( -\frac{S_3(T)}{T} \right)$$



is slow if  $T_n \ll T_c$  (lot of supercooling) =>

**Danger of being trapped => No EW-breaking**

+ **low energy Landau poles < 10-100 TeV**



- No UV-completion => solving B => need for BSM
- **Problem is with the CP-violation and DM**

# Dark sector EWBG

**Generic issue:** Difficulty of obtaining large CP-violation consistent with observations (and UV-completeness idea)

**Idea:** Move CP-violation to dark sector

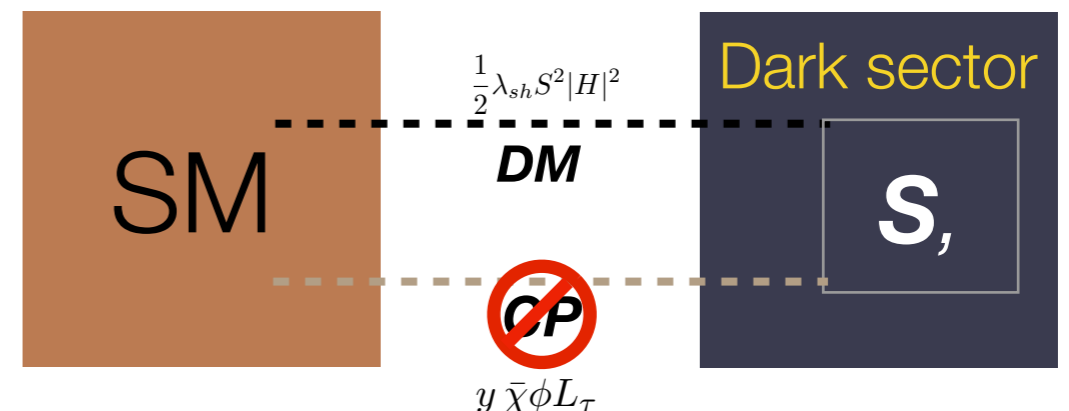
**Gain:** Get rid of most constraints, EDM's in particular

**Problem:** How to transfer the asymmetry to the visible sector / sphalerons?

**Solution:** a **CP-portal**

“Electroweak baryogenesis from a dark sector”

James M. Cline, KK and David Tucker-Smith  
CERN-TH-2017-050, arXiv:1702.08909 [hep-ph]  
accepted for publication in PRD



# Dark sector EWBG main features (IDM + S + $\chi$ )

1. **CP-asymmetry source** involving a sterile neutrino  $\chi$

$$\mathcal{L} \ni \frac{1}{2} \bar{\chi} [m_\chi + \overset{\text{singlet scalar}}{S}(\eta P_L + \eta^* P_R)] \chi$$

spontaneous CP-violation  $\Rightarrow$   $\chi$ -helicity asymmetry

2. **CP-portal:** involving an inert doublet  $\phi$

$$\mathcal{L} \ni y \bar{\chi} \phi L_\tau$$

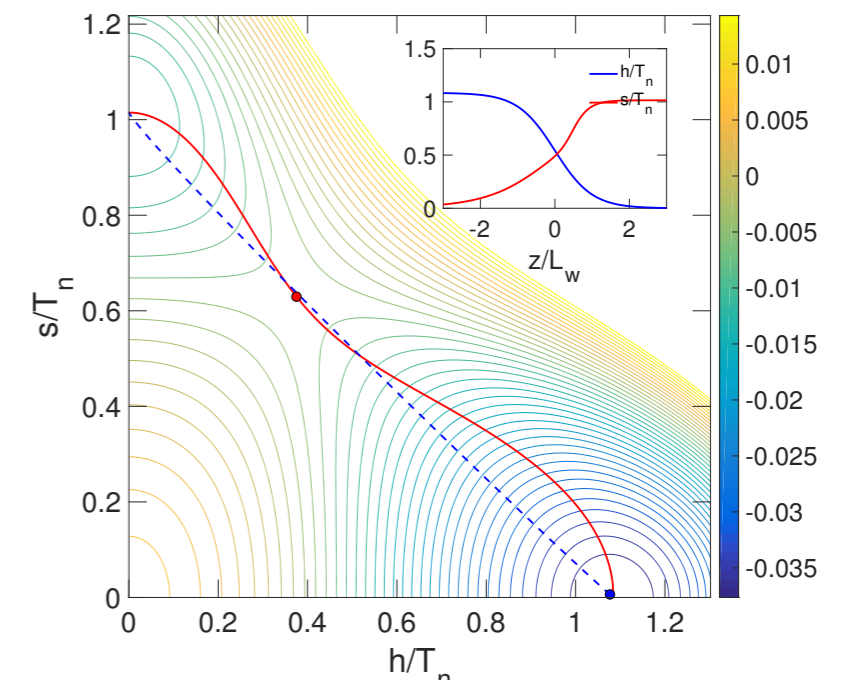
Asymmetry transported to  $\tau$ 's mainly via decays

$$\chi \bar{L}_\tau \rightarrow \phi, \quad \phi \rightarrow \bar{L}_\tau \chi$$

3.  **$\tau$ 's bias sphalerons** to produce baryon asymmetry

This can be made to work with reasonable parameters:

$$m_\chi \sim 50 \text{ GeV}, \quad m_\phi \sim 150 \text{ GeV}, \quad y \cong 0.65 \quad \eta \cong 0.1$$



PT is 2-step as in the the singlet model  
Blue line: tunneling path at  $T_n$   
Red line: path over the expanding bubble wall.

4. **Bonus:**  $\chi$  is a natural **DM** candidate

# Dark sector EWBG fluid equations sample solutions

Fluid equations are generically of the form:

$$A_\chi \begin{pmatrix} \mu'_\chi \\ u'_\chi \end{pmatrix} = C_\chi - v_w m_\chi^{2'} \begin{pmatrix} K_{2,\chi} \mu_\chi \\ K_{6,\chi} u_\chi \end{pmatrix} + \begin{pmatrix} 0 \\ S_\chi^h \end{pmatrix}$$

$$A_i \begin{pmatrix} \mu'_i \\ u'_i \end{pmatrix} = C_i, \quad \text{for } i = \phi, \tau$$

where

$$A_i \equiv \begin{pmatrix} v_w K_{1,i} & 1 \\ -K_{4,i} & v_w K_{5,i} \end{pmatrix}, \quad C_i \equiv \begin{pmatrix} C_i^\mu \\ -\tilde{\Gamma}_{\text{el},i} u_i \end{pmatrix},$$

and (model specific) collision terms are

$$C_\chi^\mu = 2\tilde{\Gamma}_{\text{hf}} \mu_\chi + 2\tilde{\Gamma}_d (\mu_\chi + c\mu_\tau - c\mu_\phi)$$

$$C_\phi^\mu = \tilde{\Gamma}_d (\mu_\phi - \mu_\tau - c\mu_\chi) + 2\tilde{\Gamma}_{\times,\phi} (\mu_\phi - \mu_\tau)$$

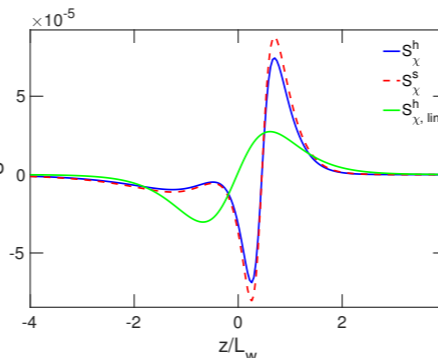
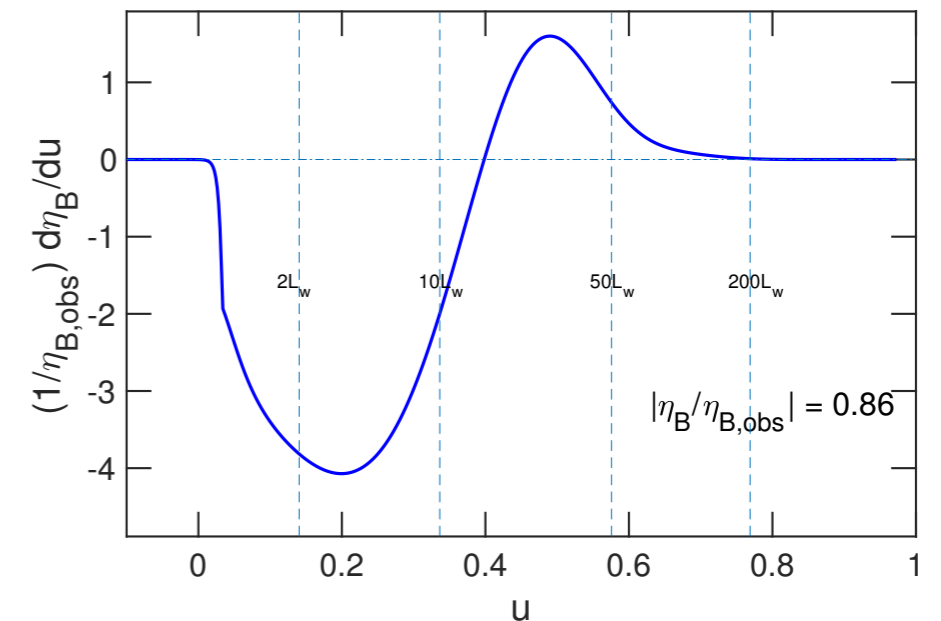
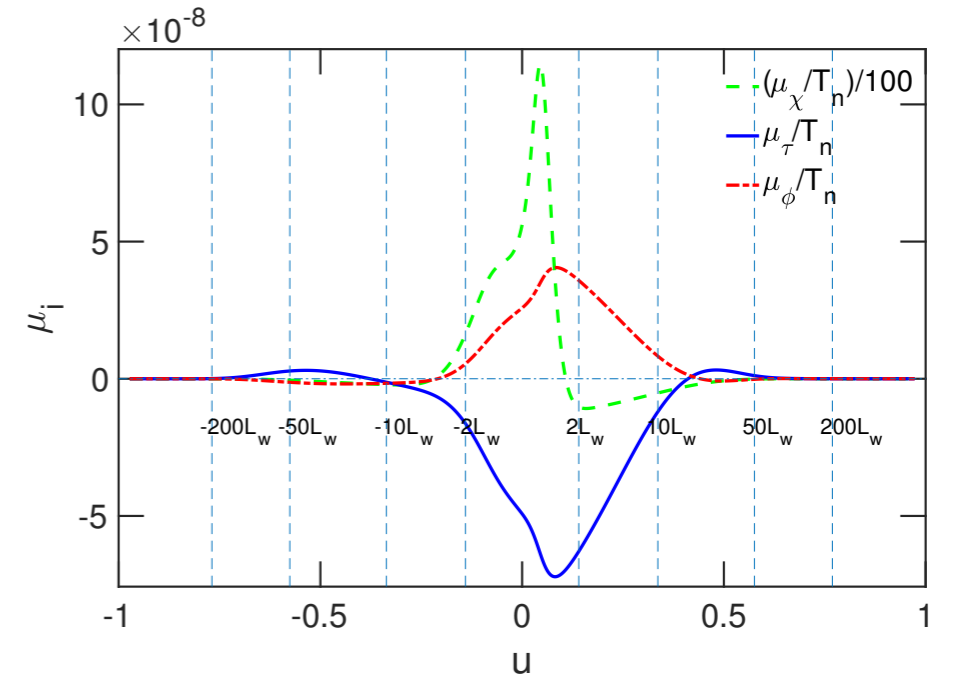
$$C_\tau^\mu = \tilde{\Gamma}_d (\mu_\tau + c\mu_\chi - \mu_\phi) + 2\tilde{\Gamma}_{\times,\tau} (\mu_\tau - \mu_\phi)$$

and the source

$$S_\chi^h = v_w \left( -K_8^h (m_\chi^2 \theta')' + K_9^h \theta' m_\chi^2 m_\chi^{2'} \right)$$

with kinematic functions:

$$K_8^h(x) = \left\langle \frac{k_z^2 f_0'}{2\omega_0 \omega_{0\parallel}^2 |\mathbf{k}|} \right\rangle \quad K_9^h(x) = \left\langle \frac{k_z^2}{4\omega_0^2 \omega_{0\parallel}^2 |\mathbf{k}|} \left( \frac{f_0'}{\omega_0} - f_0'' \right) \right\rangle$$



$$\eta_B = \frac{405 \Gamma_{\text{sph}}}{4\pi^2 v_w g_* T} \int_{-\infty}^{\infty} dz \mu_\tau f_{\text{sph}} e^{-\nu z}$$

# Dark sector EWBG fluid equations sample solutions

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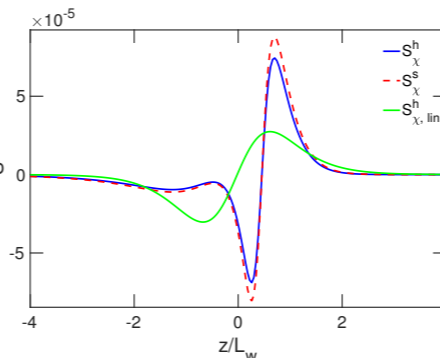
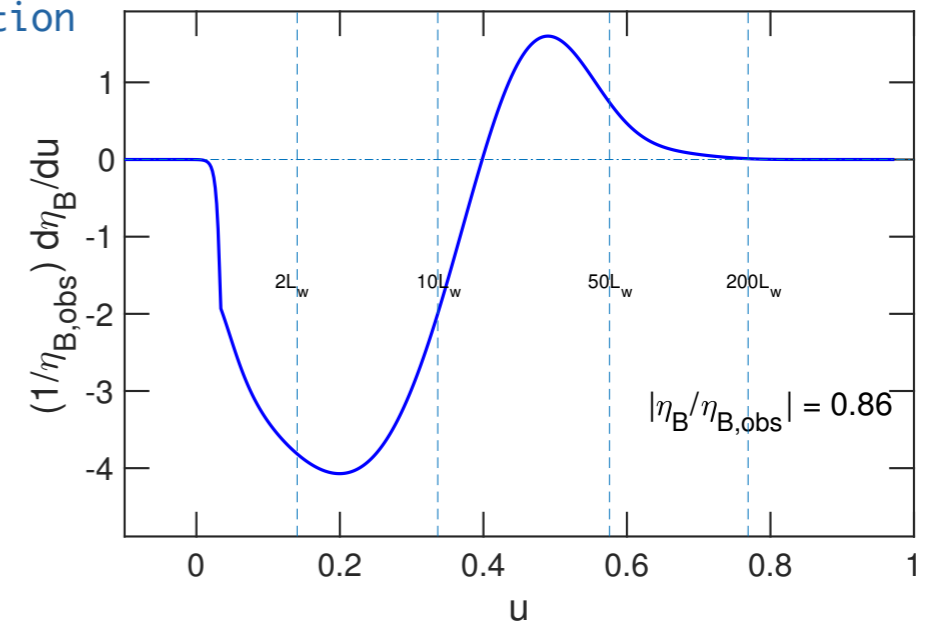
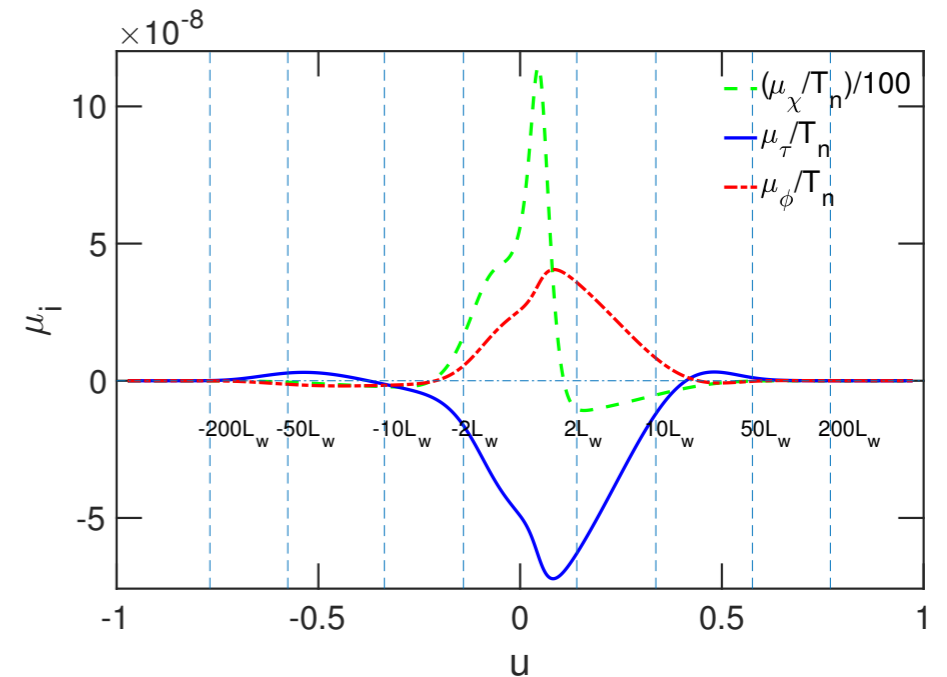
CP-potential interaction  
From, JM. Cline,  
KK and K.Olive,  
PRD49 (1994)  
6394-6409

and the source

$$S_\chi^h = v_w \left( -K_8^h (m_\chi^2 \theta')' + K_9^h \theta' m_\chi^2 m_\chi^{2'} \right)$$

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$$\eta_B = \frac{405 \Gamma_{\text{sph}}}{4\pi^2 v_w g_* T} \int_{-\infty}^{\infty} dz \mu_\tau f_{\text{sph}} e^{-\nu z}$$

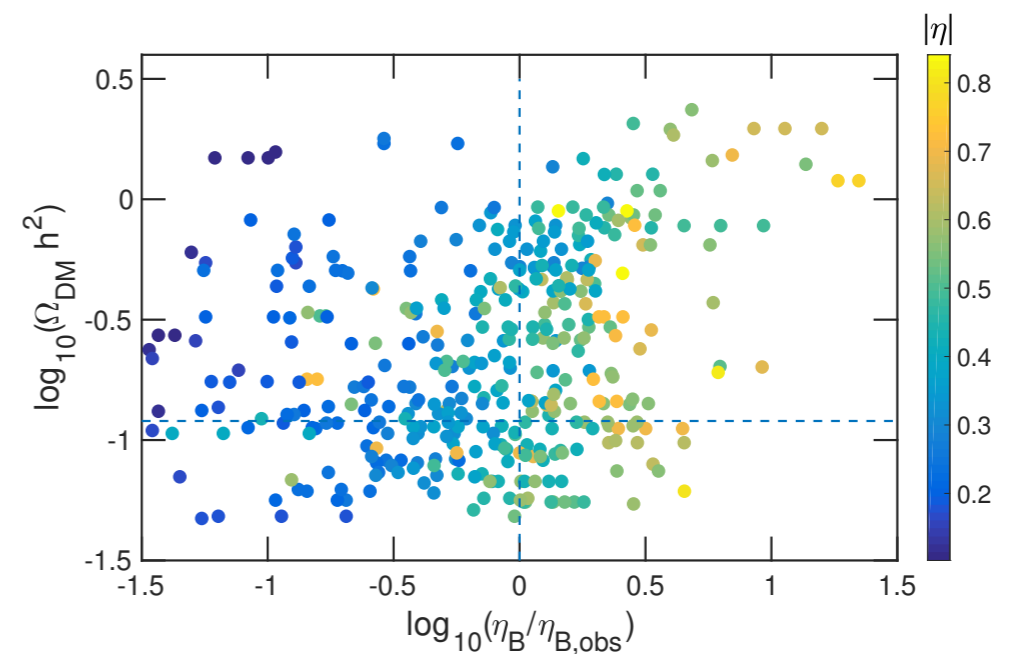
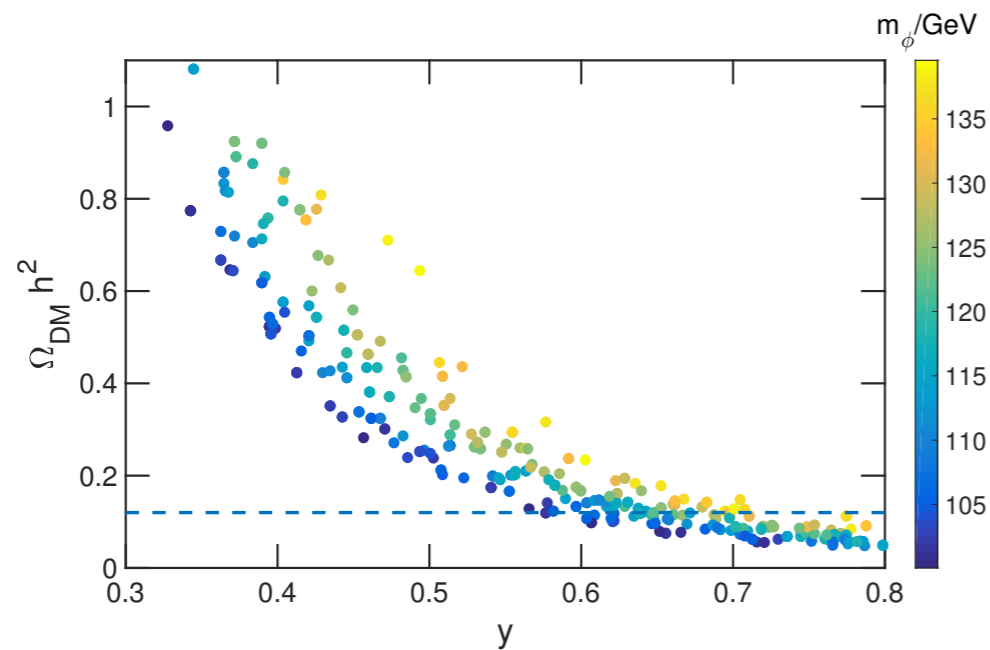
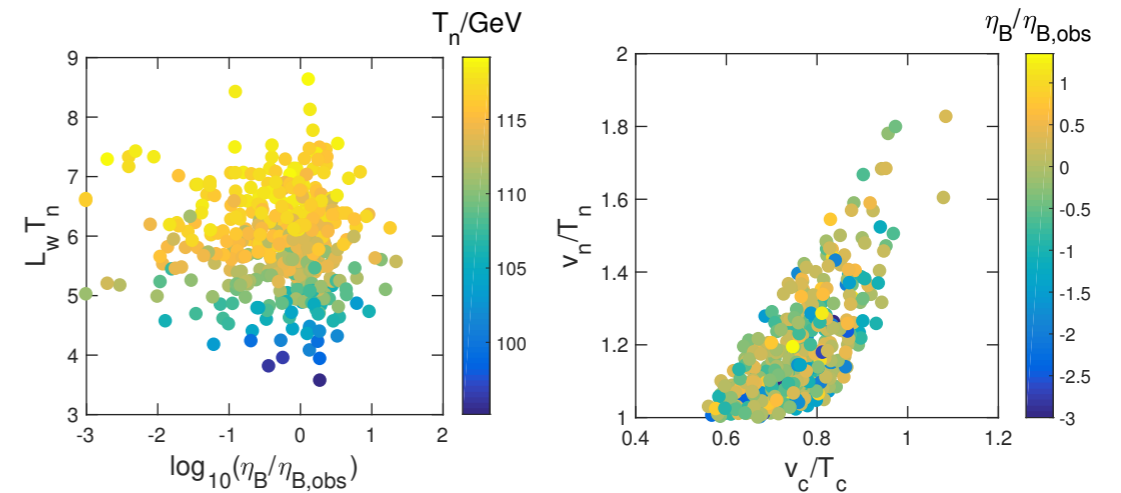
# Dark sector EWBG successful EWBG and DM

## Parameter scan:

Region of parameters that gives the BAU and DM of the right magnitude:

$$y \in [0.5, 0.9], \eta \in [0.025, 0.25], \lambda_m \in [0.5, 0.8]$$

$$m_\chi \in [40, 60], \quad m_\phi \in [100, 140],$$



**Couplings reasonably small, so should easily be made UV-complete (not tested yet)**



# Dark sector EWBG main (collider) constraint

Main collider signature of the singlet  $\chi$  and inert doublet  $\phi$  model is the Drell-Yan production of  $\phi^+\phi^-$  followed by  $\phi_{\pm} \rightarrow \tau^{\pm}\chi$ , exactly as  $pp \rightarrow \tau\tilde{\tau}^*$ ,  $\tau\tilde{\tau} \rightarrow \tau\chi^0_1$  in the MSSM.

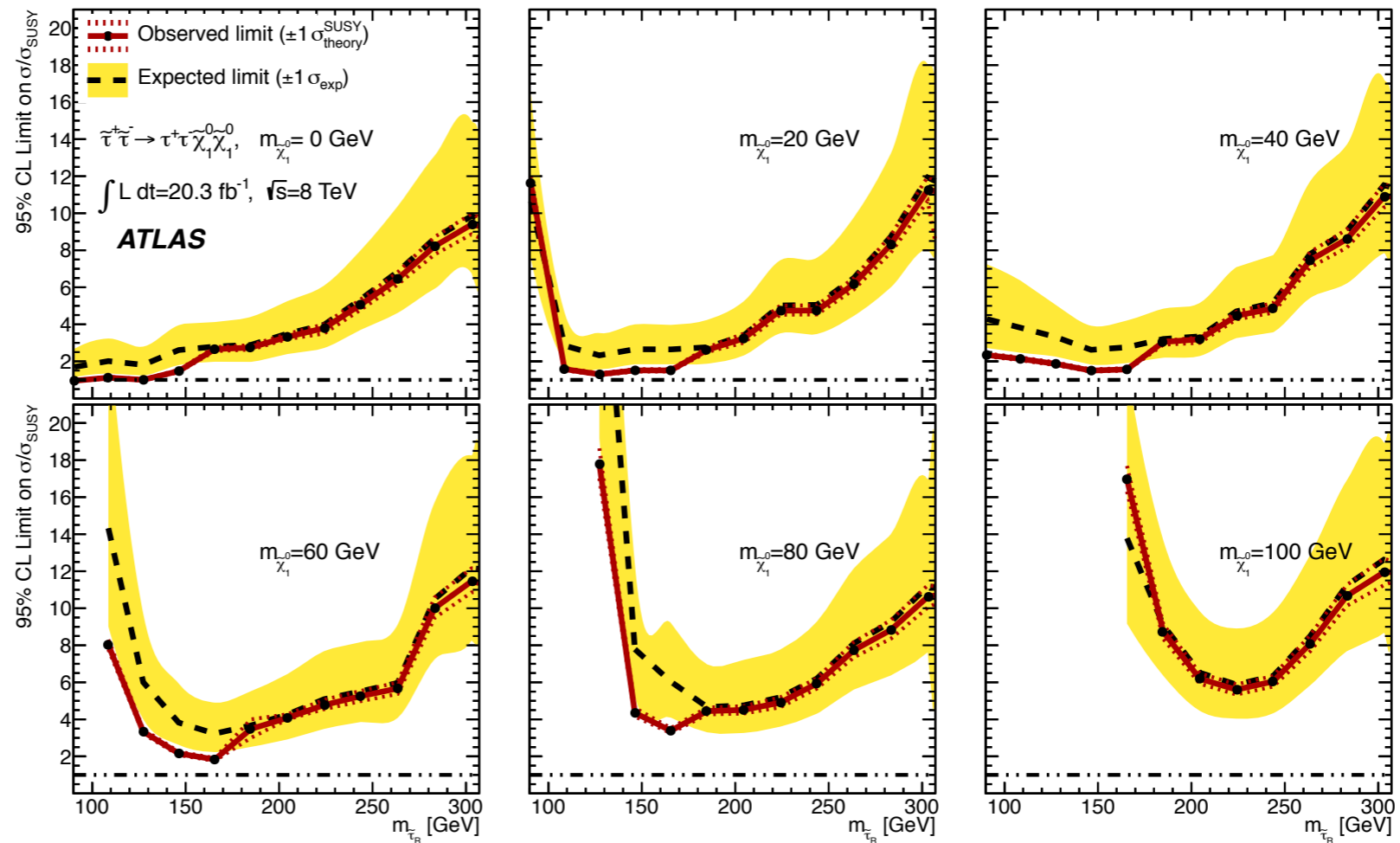
ATLAS: 1407.0350

JHEP 1410 (2014) 096

$$m_{\chi} \lesssim 20 \text{ GeV}$$

$$m_{\phi} < 130 \text{ GeV}$$

~ Excluded



Run 2 should probe the model in interesting parameter range  $m_{\chi} \cong 40 \text{ GeV}$  and  $m_{\phi} < 170 \text{ GeV}$

# Conclusions

---

## DM and BAU require BSM physics

- “Complete” solutions appear to be in decline (no SUSY found)

**Simplicity:** an interesting paradigm to follow

- Unification and Hierarchy Problem maybe not relevant

Some aspects are easily realized with help of singlets

- DM
- Strong EWPT

EWBG challenged by need for new, unconstrained CP-violation

**Dark sector EWBG with a CP-portal, can fulfill all expectations:  
DM, BAU and (likely) UV-completeness**

Will be probed by LHC run II

**Thank you for your patience**

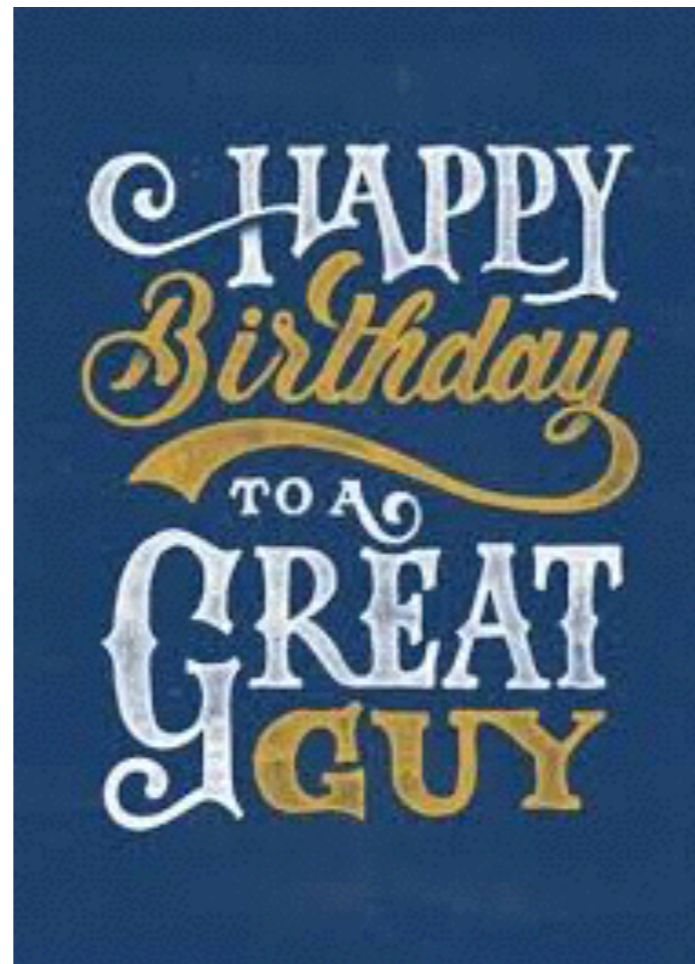
---

and **thanks again to**  
**Keith**

**Thank you for your patience**

---

and **thanks again to  
Keith**



**Extra slides. Please  
remove these from the  
web-version ...**

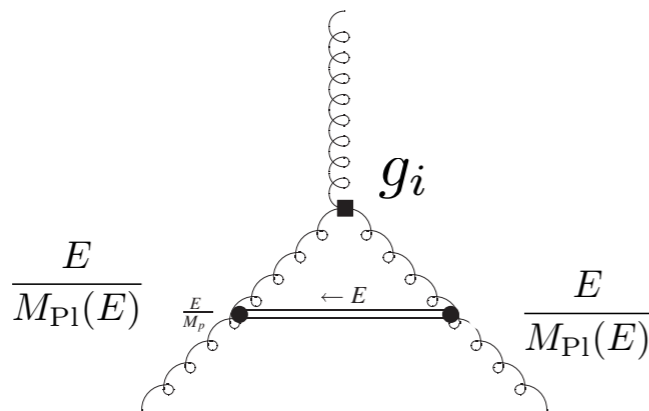
# Asymptotic safety by Gravity, maybe no need for FT unification

Running couplings (incl. yukawas):

$$\mu \frac{\partial g_i}{\partial \mu} = \beta_{\text{SM}}(g_i) + \beta_{\text{grav}}(g_i)$$

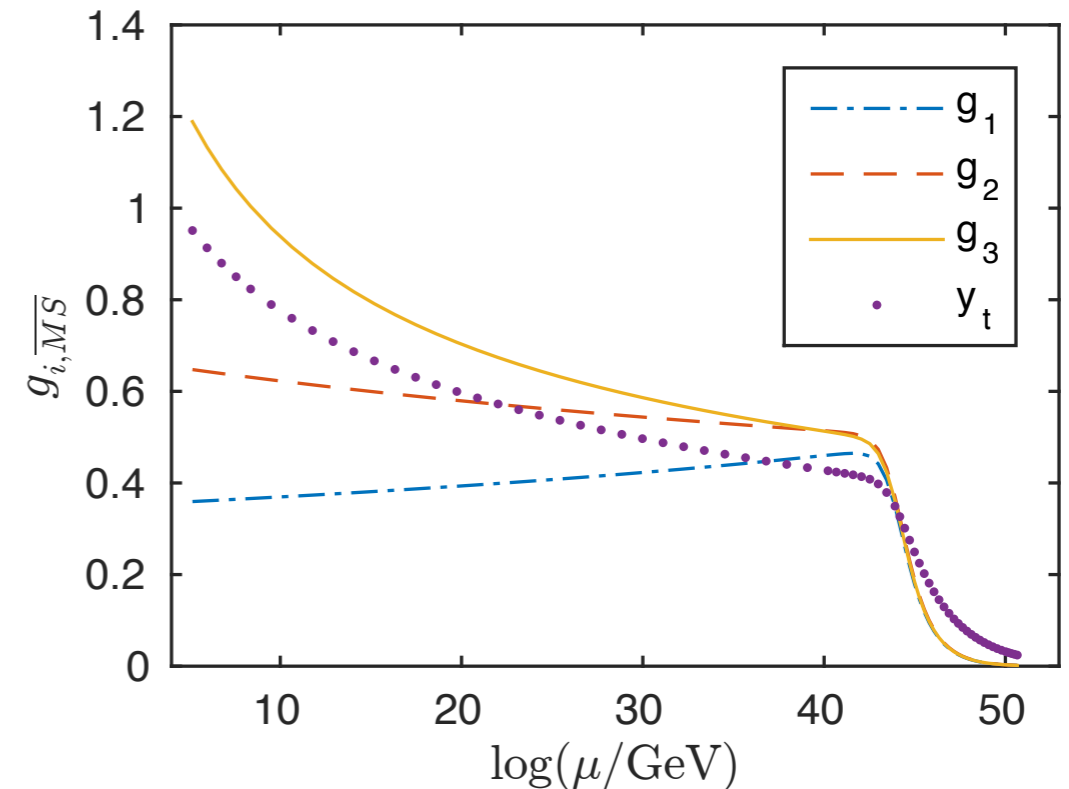
Where gravity correction is **parametrically**:

$$\beta(g_i)_{\text{grav}} = \frac{a_i \mu^2}{M_{\text{Pl}}^2 + \xi \mu^2} g_i$$



S.P. Robinson and F.Wilczek, PRL96 (2006) 231601

Laura Laulumaa,  
MSc Thesis, JyU 2015



$$a_g \approx -1$$

$$a_{y_t} \approx -0.5$$

$$\xi \approx 0.024$$

Wetterich and Shaposhnikov,  
Phys.Lett. B683 (2010) ...

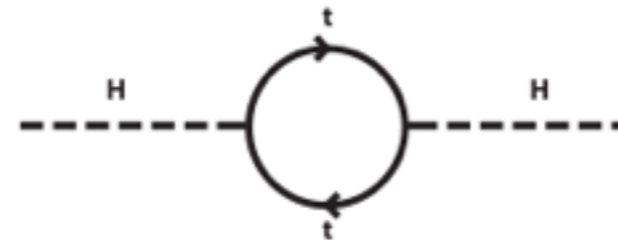
**Gravity may remain non-renormalizable for all E**

# Hierarchy problem

Excessive running of Higgs mass:

Naive **renormalization with a cutoff**:

$$\Delta m_H^2 = -\frac{3y_f^2}{8\pi^2} \Lambda_{UV}^2$$



Same also from **Wilsonian RGE-flow** argument

$$\left. \begin{array}{l} \lambda\phi^4, F_{\mu\nu}F^{\mu\nu}, \bar{\psi}\phi\psi \\ g\phi^5, (F_{\mu\nu})^2, \bar{\psi}\phi^2\psi \\ m^2\phi^2, g\phi^3 \end{array} \right\} \begin{array}{l} \sim \text{constant} \\ \text{die} \\ \text{grow} \end{array} \Rightarrow \begin{array}{l} \text{hence FT:s} \\ \text{are so simple!} \\ \text{problem} \end{array}$$

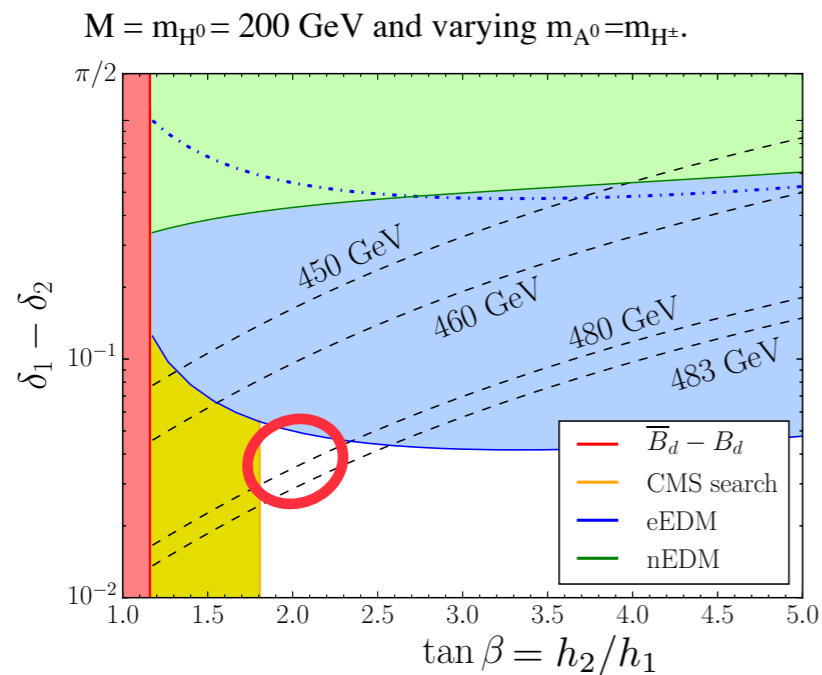
**However**, both arguments introduce an arbitrary scale, *breaking conformal invariance*. Eg. in dimensional regularization, instead:

$$\Delta m_H^2 \sim -\frac{3y_f^2}{8\pi^2} \frac{m_f^2}{\epsilon}$$

- 1) same as for fermions
- 2) vanishes if no new mass scales

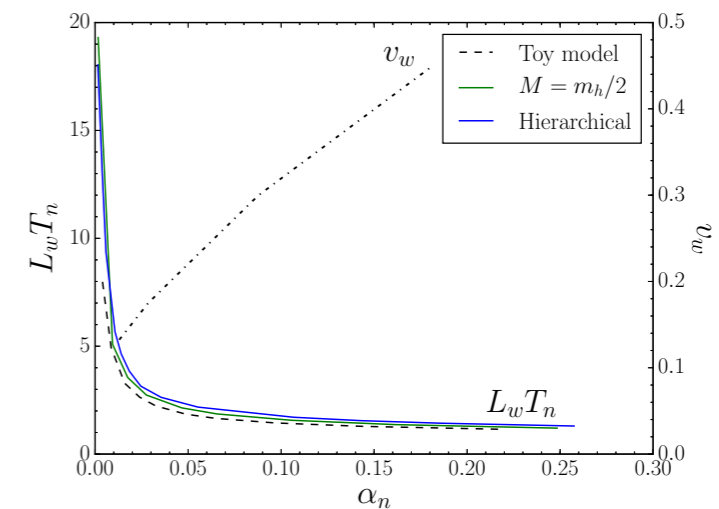
# 2HDM bounces back (not likely)

Recent paper claim 2HDM (with  $Z_2$  symmetry) is still OK (and may also give lots of GW's)

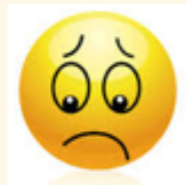


Dorsh, Huber, Konstandin and No,  
arXiv:1611.05874

But model requires  
 very small  $L_w T < 2$ ,  
 very large  $v_n/T_n \approx 3-5$   
 very large couplings



=> Neither the semiclassical BEq's, nor  
 perturbative effective action reliable.



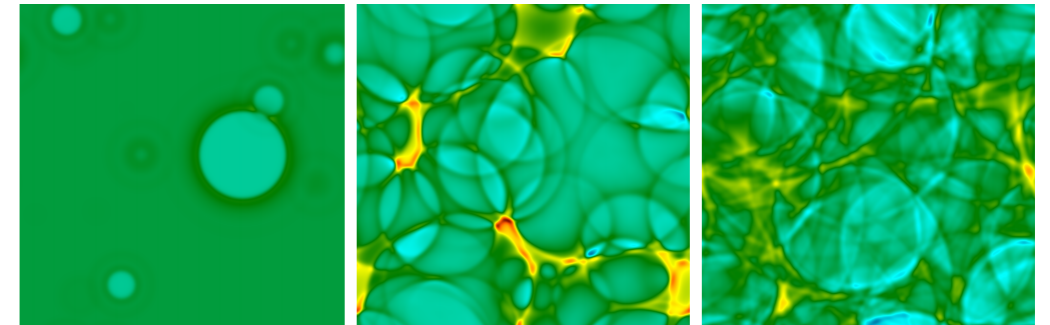
**Very low scale Landau poles at  $\sim$ TeV =>**

- No UV-completion
- *Problem is with the CP-violation*

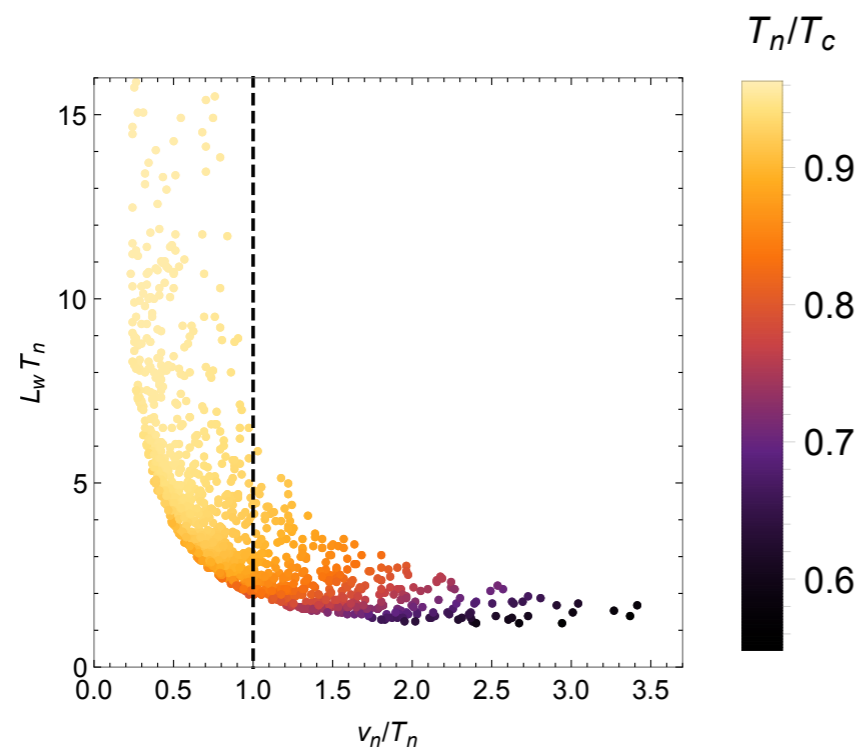


# Singlet model: BAU and gravitational waves

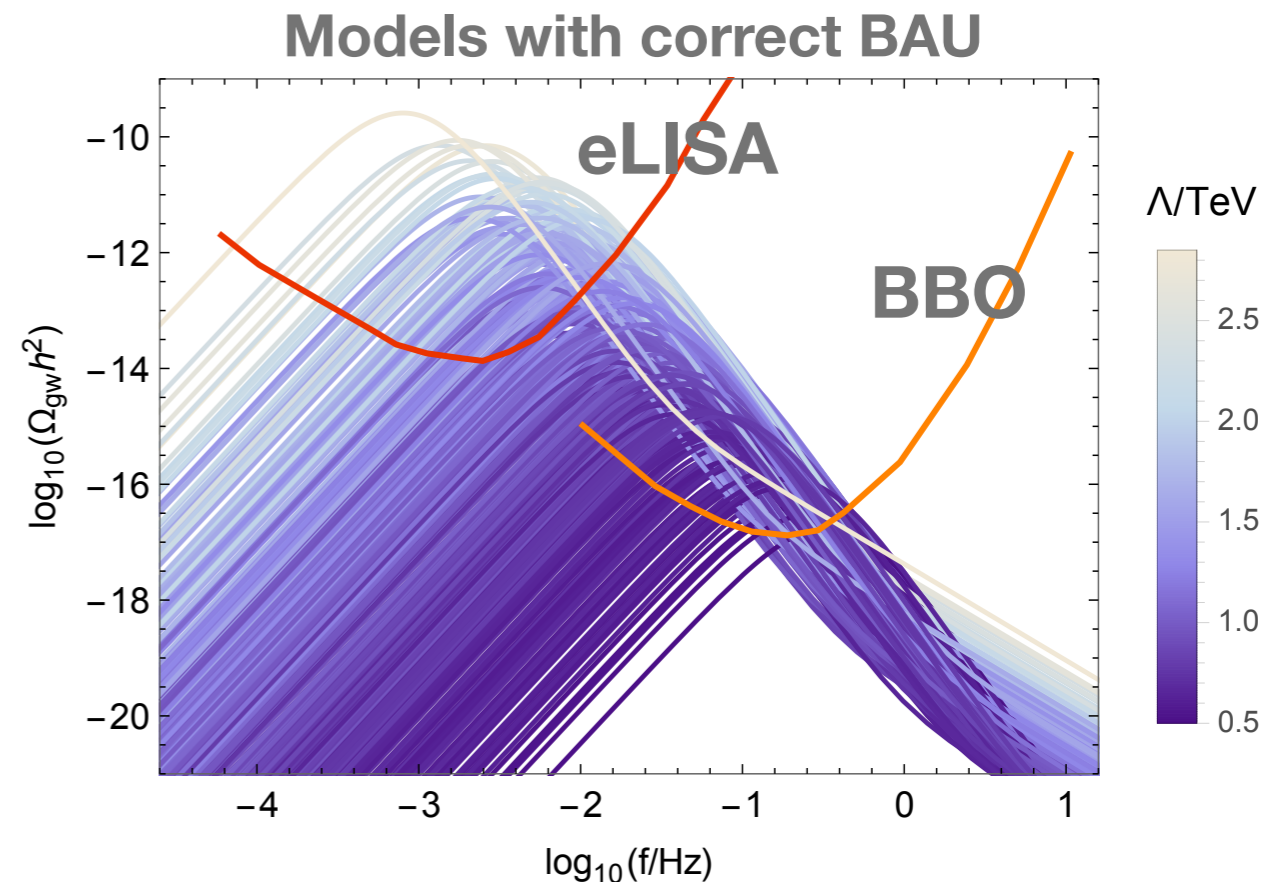
This singlet model can also give rise to an observable gravitational wave signature:



Kari Rummukainen, Physics Days 2016

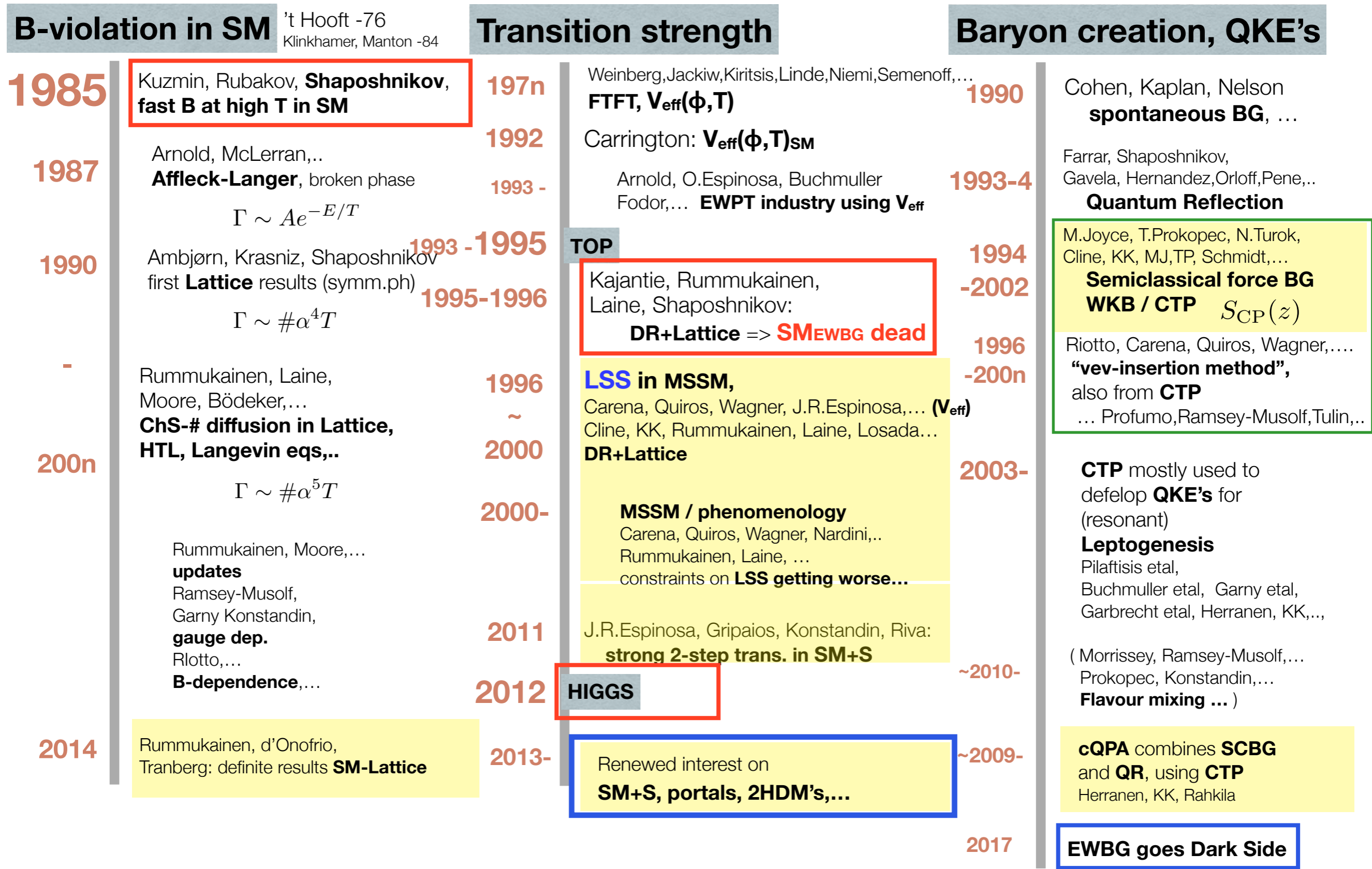


Large GW-signal correlates with strong transition (large supercooling) and a **very thin wall** ...



Ville Vaskonen, arXiv:1611.02073

# Short history of the EWBG



# Dark energy ...

---

It may require:

“Extended” gravity

or

“Alternative” gravity

# Dark energy ...

---

It may require:

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or

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gravity = truth, facts

# Dark energy ...

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# Dark energy ...

---

It may require:

“Extended” truths

or

“Alternative” facts



gravity = truth, facts

...or it is an illusion. Anyway, I will not consider it here...

# EWBG

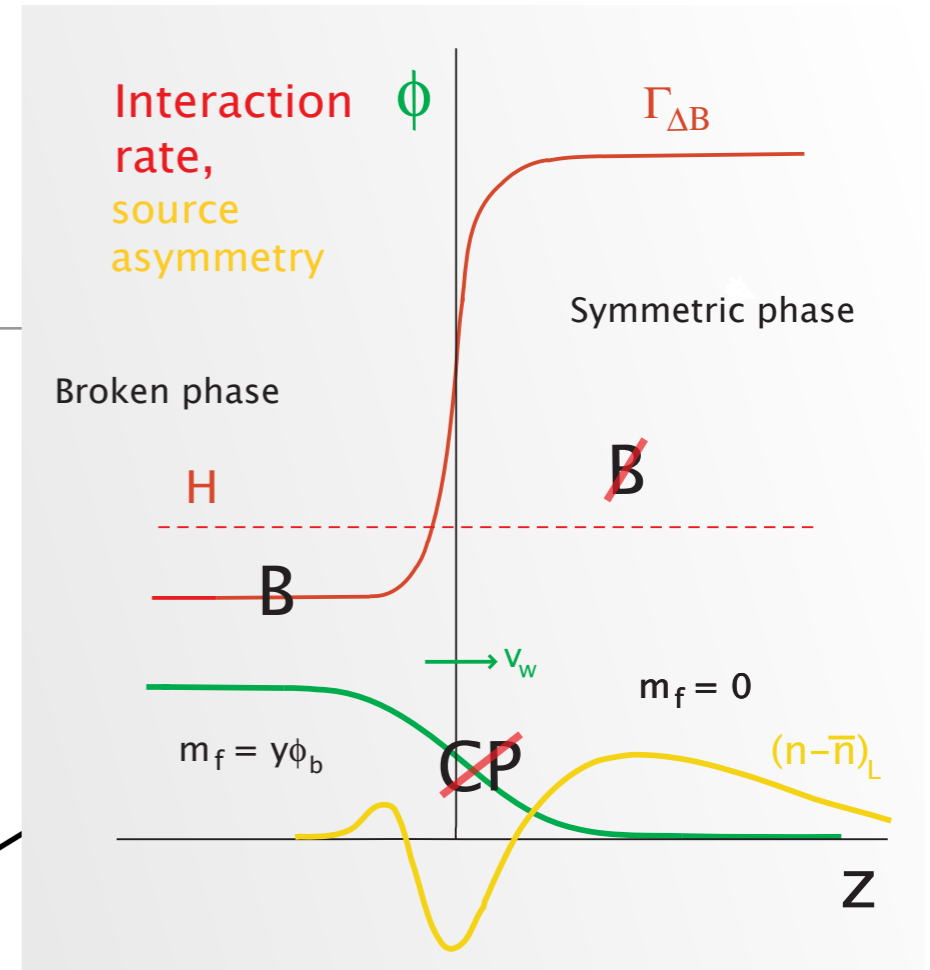
$$H \sim 10^{-14} T_{100}^2 \text{ GeV}$$

$$\Gamma \sim 10^{-5} T_{100} \text{ GeV}$$



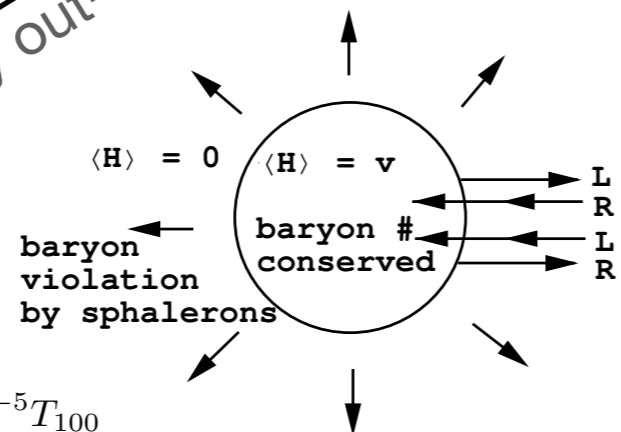
PT must be 1st order

To keep BA



Equilibrium / Nonperturbative / Gauge issues

Mostly out-of-equilibrium / quantum



$$\frac{1}{t_w} \sim \frac{v_w}{L_w} \sim 10^{-5} T_{100}$$

To make BA



# EWBG

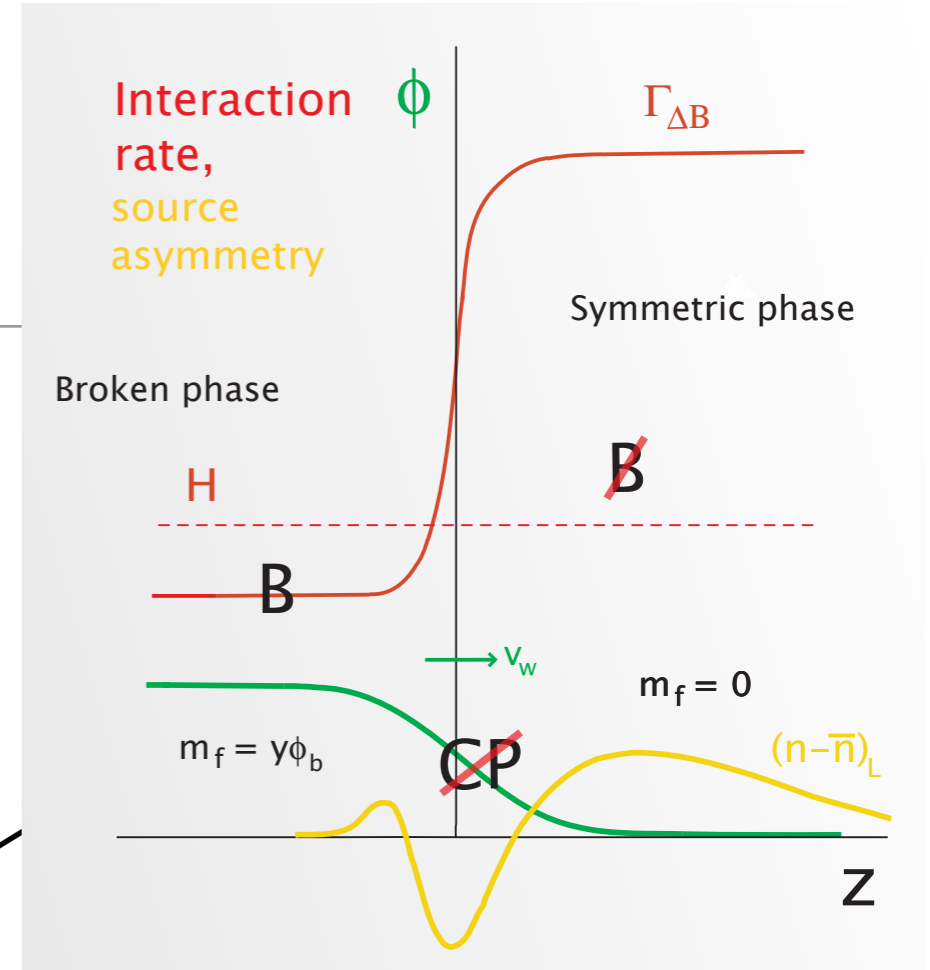
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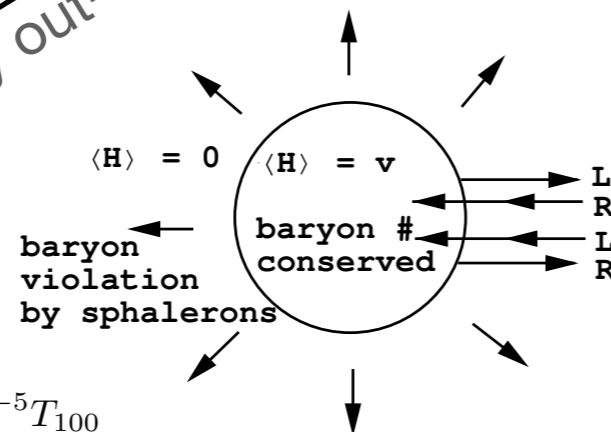


PT must be 1st order

To keep BA



Equilibrium / Nonperturbative / Gauge issues  
 Mostly out-of-equilibrium / quantum



$$\frac{1}{t_w} \sim \frac{v_w}{L_w} \sim 10^{-5} T_{100}$$

Sphaleron rate in the unbroken phase (equilibrium)  
 Ambjorn et al, ... Moore; Rummukainen et al, ...

 To make BA

# EWBG

$$H \sim 10^{-14} T_{100}^2 \text{ GeV}$$

$$\Gamma \sim 10^{-5} T_{100} \text{ GeV}$$

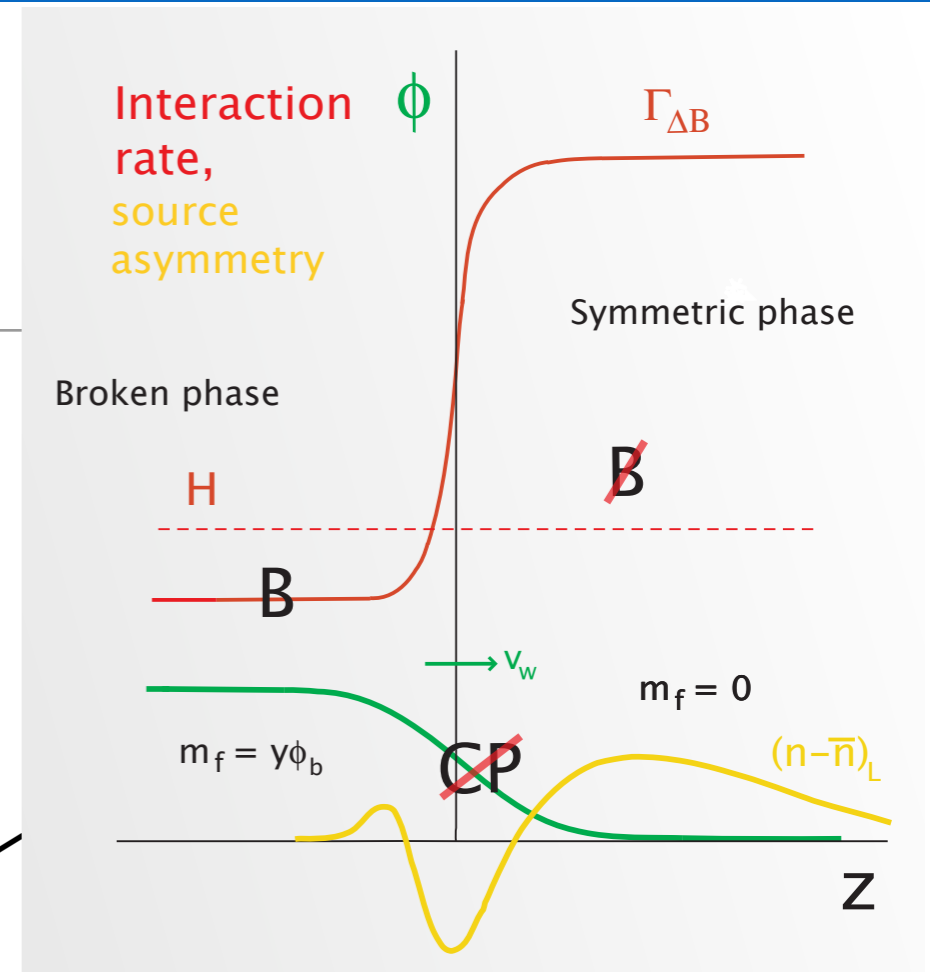


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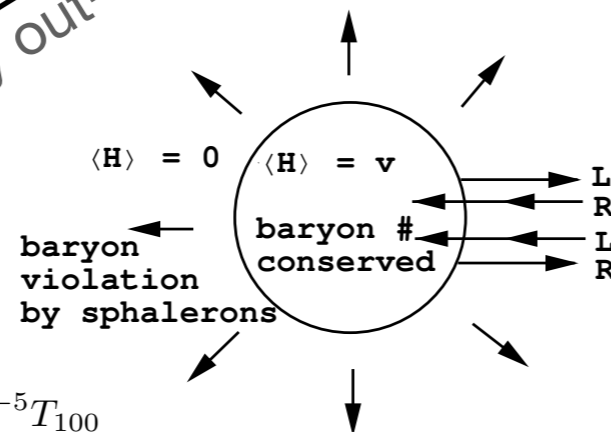
## To keep BA

(Small) sphaleron rate in the broken phase

Kuzmin, Rubakov & Shapshonkinov, Arnold & McLerran, Moore, Rummukainen...;



Equilibrium / Nonperturbative / Gauge issues  
 Mostly out-of-equilibrium / quantum



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To make BA

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## To keep BA

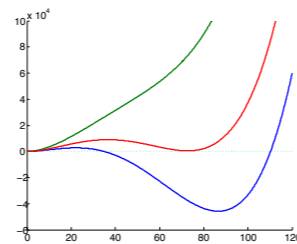
### (Small) sphaleron rate in the broken phase

Kuzmin, Rubakov & Shapshonkinov, Arnold & McLerran, Moore, Rummukainen...;

- $V_{\text{eff}}$  in Landau gauge

$$\frac{\phi_n}{T_n} > 1$$

H.H.Patel, M.J.Ramsey-Musolf, C.Wainwright, S.Profumo  
 JHEP 07 (2011) 029; PRD84 (2011) 023521; PRD86 (2012) 083537.  
 M.Garny and T.Konstandin, JHEP1207 (2012) 189, ....



- Dim. reduction to a 3D-Higgs-gauge theory simulated in Lattice

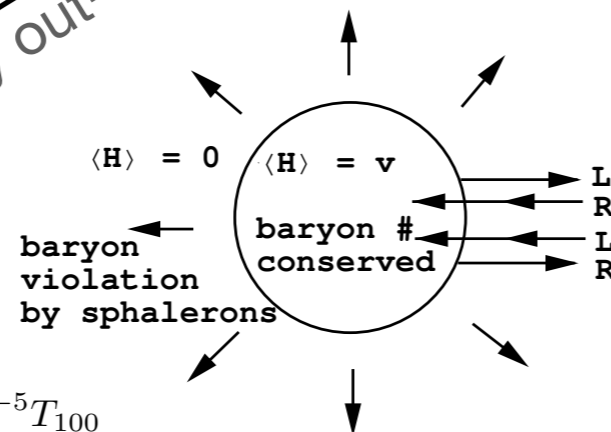
K.Kajantie, M.Laine, K.Rummukainen and M.E.Shaposhnikov,  
 NPB458 (1996) 90; NPB466 (1996) 189; PRL77, 2887 (1996)....

### 2-loop $V_{\text{eff}}$ in LG

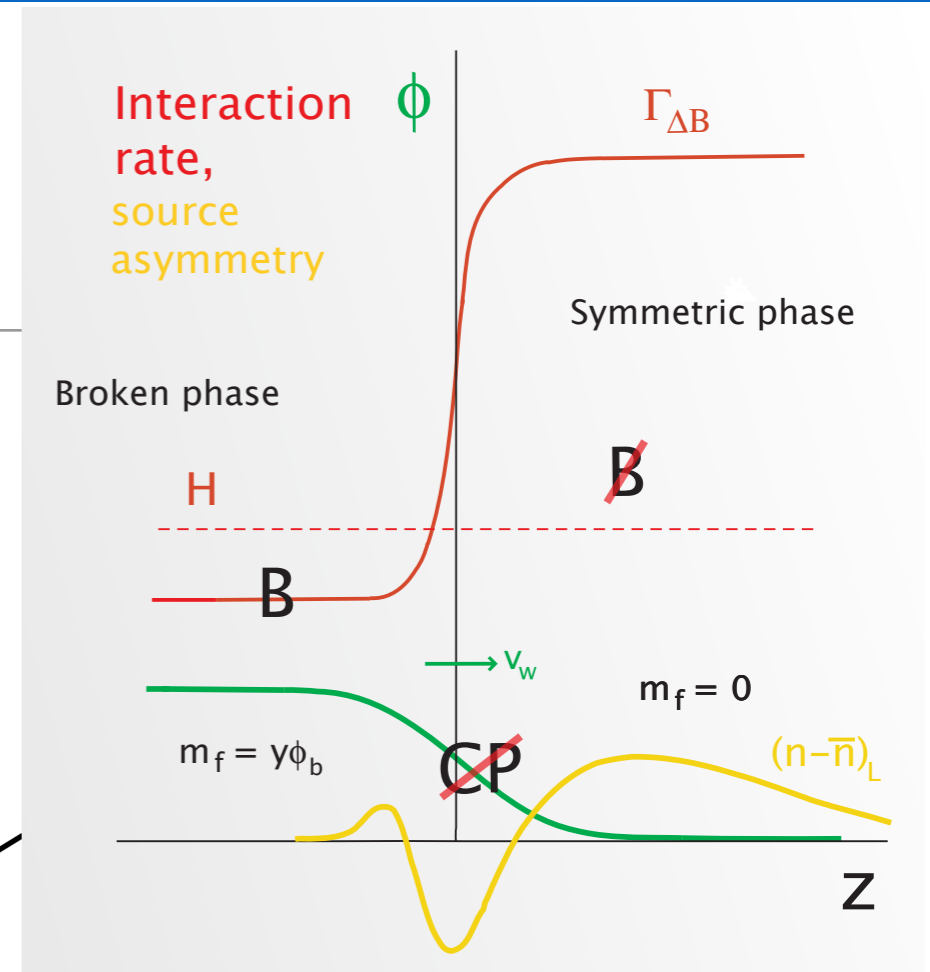
~OK

M.Laine, G.Nardini and K.Rummukainen,  
 JCAP 1301 (2013) 011...

Equilibrium / Nonperturbative / Gauge issues  
 Mostly out-of-equilibrium / quantum



$$\frac{1}{t_w} \sim \frac{v_w}{L_w} \sim 10^{-5} T_{100}$$



### Sphaleron rate in the unbroken phase (equilibrium)


Ambjorn etal, ... Moore; Rummukainen etal, ...



To make BA

# EWBG

$$H \sim 10^{-14} T_{100}^2 \text{ GeV}$$

$$\Gamma \sim 10^{-5} T_{100} \text{ GeV}$$


PT must be 1st order

## To keep BA

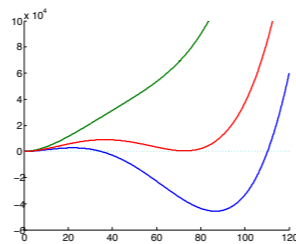
### (Small) sphaleron rate in the broken phase

Kuzmin, Rubakov & Shapshonkinov, Arnold & McLerran, Moore, Rummukainen...; 

- $V_{\text{eff}}$  in Landau gauge

$$\frac{\phi_n}{T_n} > 1$$

H.H.Patel, M.J.Ramsey-Musolf, C.Wainwright, S.Profumo  
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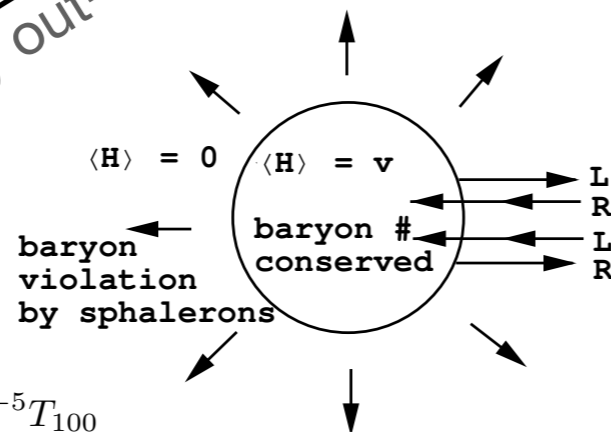
- Dim. reduction to a 3D-Higgs-gauge theory simulated in Lattice

K.Kajantie, M.Laine, K.Rummukainen and M.E.Shaposhnikov,  
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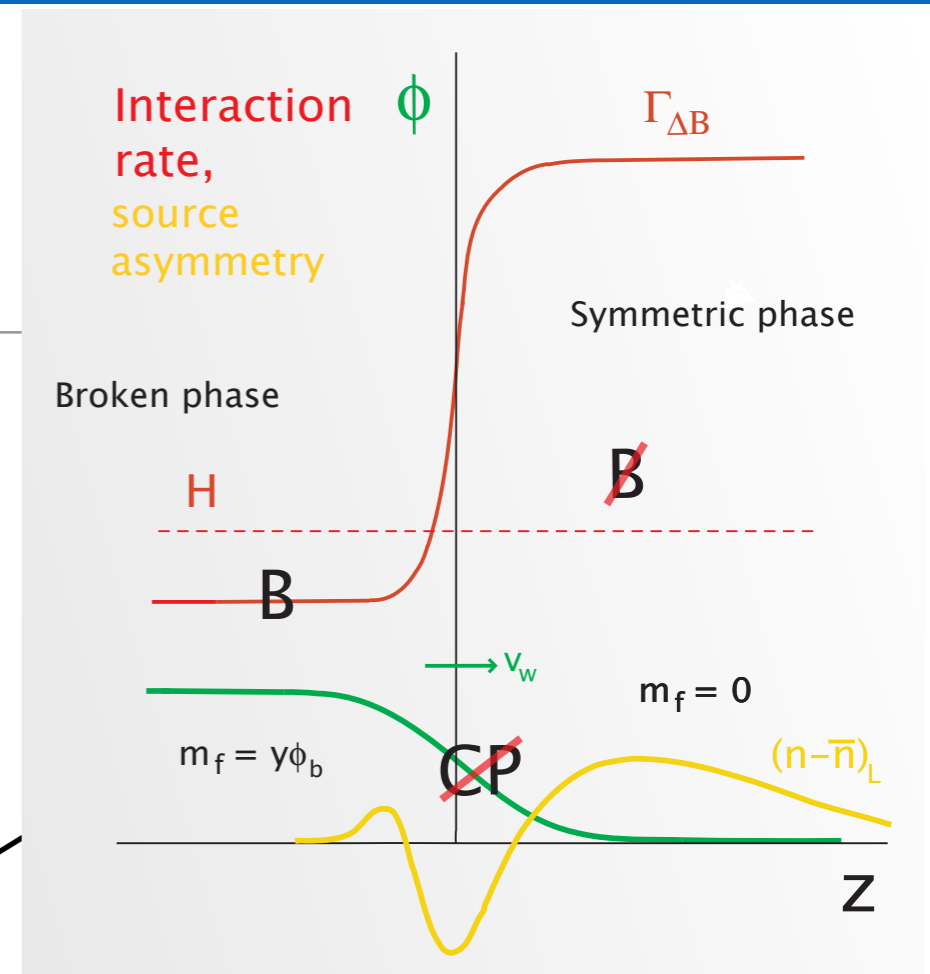
### 2-loop $V_{\text{eff}}$ in LG ~OK

M.Laine, G.Nardini and K.Rummukainen,  
 JCAP 1301 (2013) 011...

Equilibrium / Nonperturbative / Gauge issues  
 Mostly out-of-equilibrium / quantum



$$\frac{1}{t_w} \sim \frac{v_w}{L_w} \sim 10^{-5} T_{100}$$



### CP-violating source in transport eqs.

- Thin wall: *quantum*

- Thick wall SC:



**SC force** Joyce, Prokopec, Turok, Cline, KK, Schmidt, Weinstock, Konstandin, ...

**vev-insertion expansion**  
 Riotto, Carena, Quiros, Wagner, ...


### Sphaleron rate in the unbroken phase (equilibrium)

Ambjorn et al, ... Moore; Rummukainen et al, ...

To make BA 

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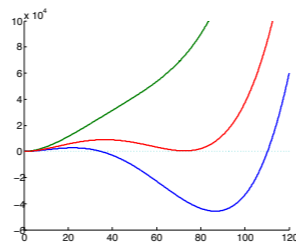
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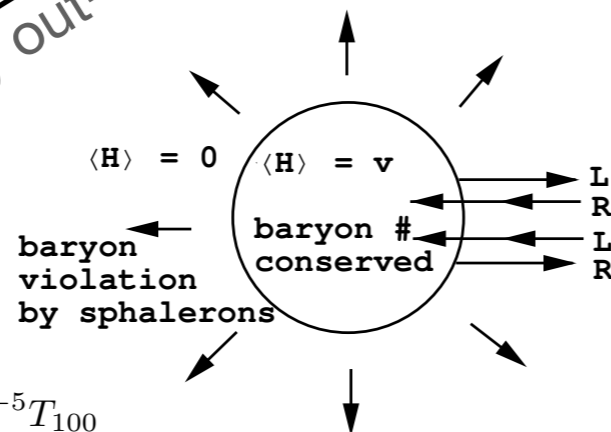
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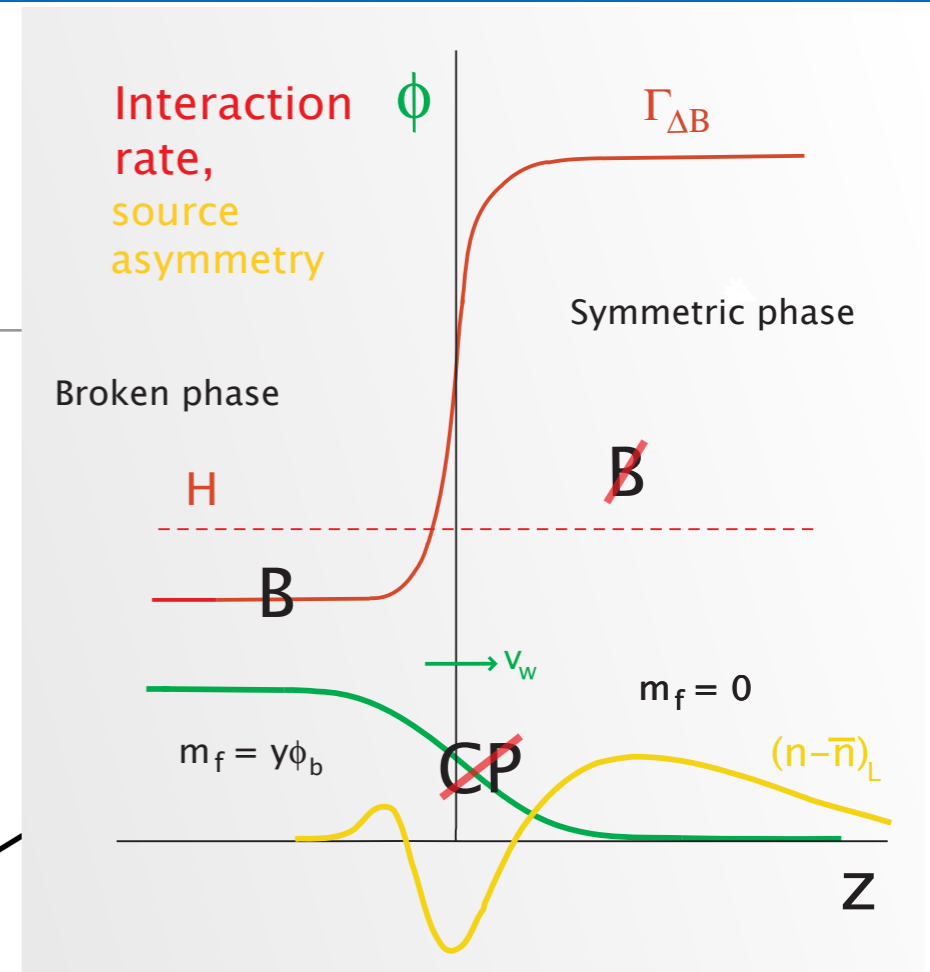
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### (CP-even) dynamics of the expanding wall

Parametrized by  $v_w$  and  $\phi(z)$

Kajantie et al, Prokopec & Moore. John & Smith Espinosa, Konstandin, No & Servant (2010)....

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To make BA

# B-violation I/EW chiral anomaly

B perturbatively conserved in SM => Proton stability!

Baryon + (and) lepton current

Chern-Simons current

$$\partial_\mu j_{B+L}^\mu = N_f \frac{g^2}{16\pi^2} * F_{\mu\nu} F^{\mu\nu} = N_f \partial_\mu j_{CS}^\mu [A_\mu^a]$$

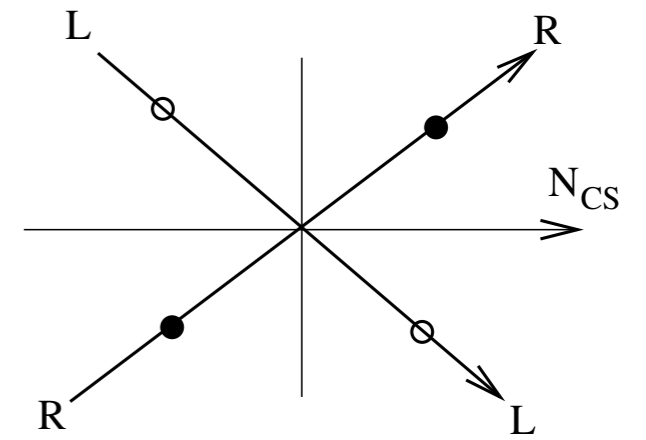
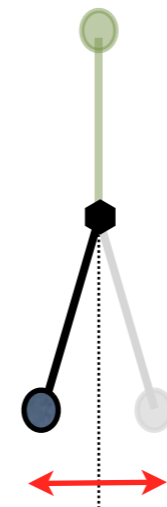
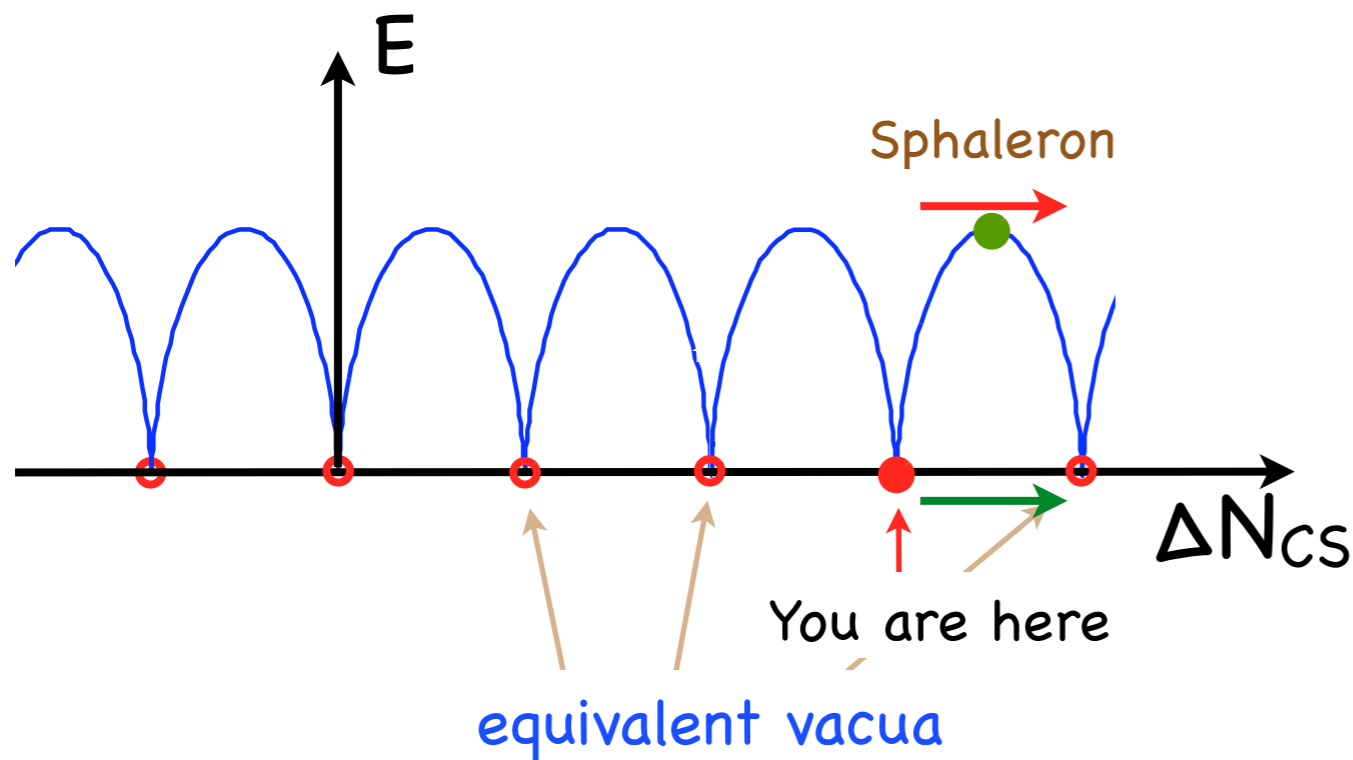
↑  
gauge fields

**NOETHER:**

$$\Delta N_{B+L} = N_f \Delta N_{CS}$$

Nonperturbative fluctuations can lead to “large” gauge transformations that to change baryon and lepton numbers.

# B-violation II / G-vacuum / CS#-diffusion / level crossing

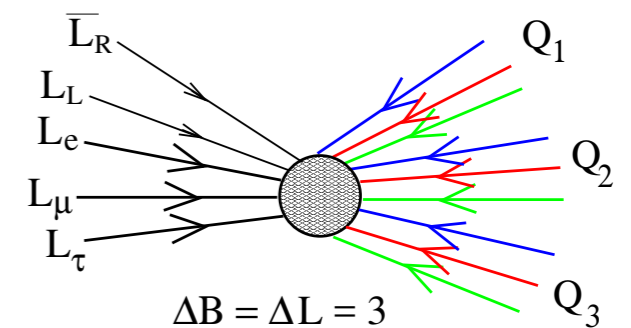


Fermion zero modes kicked from vacuum

- SLOW at  $T=0$  (tunnelling) 'tHooft -76  
 →

- FAST at high  $T$  (thermal activation)  
 →

Kuzmin, Rubakov and Shaposhnikov, -85



$$\Gamma \sim e^{-E_{Sph}(T)/T}$$

$$\frac{E_{Sph}(T)}{T} \approx 40 \frac{\phi}{T}$$

**B** , if  $\phi/T < 1.1$

**B** , if  $\phi/T > 1.1$

NEEDS A BIG JUMP IN THE ORDER PARAMETER AT TRANSITION