

**AXIS-SYMMETRIC BOUNDARY-VALUE PROBLEMS
FOR NEMATIC LIQUID CRYSTALS WITH VARIABLE
DEGREE OF ORIENTATION**

By

Vincenzo M. Tortorelli

and

Epifanio G. Virga

IMA Preprint Series # 764

February 1991

AXIS-SYMMETRIC BOUNDARY-VALUE PROBLEMS FOR NEMATIC LIQUID CRYSTALS WITH VARIABLE DEGREE OF ORIENTATION

VINCENZO M. TORTORELLI* AND EPIFANIO G. VIRGA†

Abstract. We employ a special class of orientations for the optical axis to minimize an energy functional for nematic liquid crystals whose degree of orientation varies in space. When the boundary data possess a certain symmetry with respect to an axis, we prove that the energy minimizers share that symmetry, and we establish a criterion to decide about uniqueness of minimizers within the class we employ here.

Table of Contents

1. Introduction
2. Longitudinal minimizers and uniqueness
3. An example

1. Introduction. Liquid crystals are oriented materials whose microscopic state is described at a macroscopic scale by the *optical axis*. The regions in space where the optical axis suffers discontinuities are called *defects*. ERICKSEN [1] has recently proposed a new mathematical model to describe defects in liquid crystals. ERICKSEN's model differs from OSEEN and FRANK's classical model [2,3] in that it allows the *degree of orientation* of liquid crystals to vary in space.

The interpretation of both the optical axis and the degree of orientation in terms of microscopic parameters can be found in Section 2 of [1]. Here we regard both as macroscopic, phenomenological quantities.

We represent the optical axis by a unit vector \mathbf{n} and the degree of orientation by a scalar s that ranges in the interval $[-1/2, 1]$. A liquid crystal becomes isotropic wherever s vanishes: there \mathbf{n} cannot be defined.

Let \mathcal{B} be the region in space occupied by a liquid crystal. The subset of \mathcal{B} defined by

$$(1.1) \quad \mathcal{S}(s) := \{p \in \mathcal{B} | s(p) = 0\}$$

is called *singular set*, because there \mathbf{n} may exhibit defects.

*Scuola Normale Superiore, 56100 Pisa, Italy

†Dipartimento di Matematica, Università di Pavia, 27100 Pavia, Italy

In ERICKSEN's model the free energy of a nematic liquid crystal is a functional \mathcal{F} that depends on both s and \mathbf{n} . A simple formula for \mathcal{F} compatible with ERICKSEN's general formula is

$$(1.2) \quad \mathcal{F}[s, \mathbf{n}] := \kappa \int_{\mathcal{B}} \{k|\nabla s|^2 + s^2|\nabla \mathbf{n}|^2 + \sigma_0(s)\} ,$$

where κ and k are both positive constants and σ_0 is a given positive function (cf. Section 5 of [1]). The fields s and \mathbf{n} in (1.2) are defined thus:

$$s : \mathcal{B} \rightarrow [-1/2, 1] \quad , \quad \mathbf{n} : \mathcal{B} \setminus \mathcal{S}(s) \rightarrow \mathcal{S}^2 ,$$

where \mathcal{S}^2 is the unit sphere of \mathcal{V} , the translation space of the three-dimensional space.

If \mathcal{B} is an open bounded set with Lipschitzian boundary and if both s and \mathbf{n} are prescribed on the whole of $\partial\mathcal{B}$ or on an open subset of $\partial\mathcal{B}$, say \mathcal{D} , \mathcal{F} attains its minimum in the following class of admissible functions:

$$(1.3) \quad \mathcal{A} := \{(s, \mathbf{n}) | s \in W^{1,2}(\mathcal{B}; [-1/2, 1]), \mathbf{u} \in W^{1,2}(\mathcal{B}; \mathcal{V}), \mathbf{u} := s\mathbf{n}\} .$$

If the boundary data are prescribed on the whole of $\partial\mathcal{B}$, this result is proven in [4] (cf. also [5]); essentially the same arguments apply if the boundary data are prescribed only on \mathcal{D} . For all members of \mathcal{A} we agree to set $\mathbf{u} = 0$ in $\mathcal{S}(s)$.

The regularity of the minimizers of \mathcal{F} has been examined in [6] and [7]: in particular, we know that any pair that minimizes \mathcal{F} is such that s is continuous on the whole of \mathcal{B} and, if σ_0 is of class C^m with m an integer, both s and \mathbf{n} are of class $C^{m,\alpha}$, for all $\alpha < 1$, away from the singular set $\mathcal{S}(s)$. Furthermore, reasoning as in Section 5 of [7], one sees that if \mathcal{D} is a surface of class $C^{1,\gamma}$, with $\gamma > 0$, and if both s and \mathbf{n} are prescribed on \mathcal{D} as Lipschitzian functions, then for all minimizers of \mathcal{F} both s and \mathbf{u} are Hölder continuous in \mathcal{B} up to \mathcal{D} . As to the singular set of a minimizing pair, another theorem of [7] (cf. Section 6) shows that its Hausdorff dimension is at most 2 for all $k \leq 1$, while it is at most 1 for all $k > 1$.

Also the uniqueness of the minimizers of \mathcal{F} has been addressed in [7]. In particular, if $\sigma_0 = 0$ and $k \in]0, 1[$, we know that there is only one pair (s, \mathbf{n}) that minimizes \mathcal{F} . If $k \geq 1$, this conclusion may in general be false.

Several boundary-value problems for \mathcal{F} have been solved so far; a review can be found in [8]. In some of these problems both the region \mathcal{B} and the data prescribed on its boundary are axis-symmetric (see [9], for example). Thus, the solution is assumed to share the same symmetry. Though this is quite natural, an explanation still wants. In this paper we get closer to such an explanation.

In Section 2 we introduce the class \mathcal{R} of all fields \mathbf{n} whose longitude on \mathcal{S}^2 can be defined as a smooth function on the whole of $\mathcal{B} \setminus \mathcal{S}(s)$, except where \mathbf{n} is directed along the

polar axis of \mathcal{S}^2 . A subclass of \mathcal{R} plays a special rôle in our work: its members have the same longitude on all planes through a given axis; we call \mathcal{R}_0 this class not *longitudinal* orientations its members. Assuming no special symmetry for \mathcal{B} , we show that when the boundary data for \mathbf{n} are longitudinal, the minimum of \mathcal{F} in \mathcal{R} is the same as the minimum in \mathcal{R}_0 , though the minimizers need not be all in \mathcal{R}_0 . Furthermore, we give a criterion to decide whether all minimizers of \mathcal{F} in \mathcal{R} belong indeed to \mathcal{R}_0 . Of course, when the minimizer of \mathcal{F} happens to be unique in \mathcal{R}_0 , and our criterion applies, we get uniqueness in the whole of \mathcal{R} .

In Section 3 we take \mathcal{B} as a circular cylinder and we prescribe \mathbf{n} on the lateral boundary of \mathcal{B} so as to be constant on all the generatrices of \mathcal{B} and symmetric about its axis. We also prescribe s on the lateral boundary of \mathcal{B} as a positive constant. In [9] the minimizers of \mathcal{F} subject to these conditions have been studied within the class \mathcal{R}_0 relative to the axis of \mathcal{B} . Here we prove that the variational problem studied in [9] has a unique minimizer in \mathcal{R} , *modulo* a reflection. We think that for this problem each minimizer of \mathcal{F} belongs indeed to \mathcal{R} , and so there would be only one minimizer among all admissible orientations.

2. Longitudinal minimizers and uniqueness. Let \mathcal{B} have Lipschitzian boundary and let \mathcal{D} , an open subset of $\partial\mathcal{B}$, be a connected surface of class $C^{1,\gamma}$, for some $\gamma > 0$. Let 0 be a point of the three-dimensional space. We call \mathbf{b} the axis through 0 in the direction of a given unit vector \mathbf{e}_z :

$$(2.1) \quad \mathbf{b} := \{0 + z\mathbf{e}_z | z \in \mathbb{R}\} .$$

We call *longitudinal plane* of \mathcal{B} each plane containing \mathbf{b} that intersects \mathcal{B} . The axis \mathbf{b} itself may or may not intersect \mathcal{B} . In either case we can represent \mathcal{B} in the cylindrical co-ordinates (r, θ, z) whose axis is \mathbf{b} . Away from \mathbf{b} we employ the local orthonormal frame $\{\mathbf{e}_r, \mathbf{e}_\theta, \mathbf{e}_z\}$, where \mathbf{e}_r is the radial unit vector.

Let a pair (s, \mathbf{n}) be given in $C(\mathcal{B}; [-1/2, 1]) \times C^1(\mathcal{B} \setminus \mathcal{S}(s); \mathcal{S}^2)$. We define a closed subset of \mathcal{B} by

$$(2.2) \quad \mathcal{B}_0 := \{p \in \mathcal{B} \setminus \mathcal{S}(s) | (\mathbf{n} \cdot \mathbf{e}_z)^2 = 1\} \cup \mathcal{S}(s) .$$

DEFINITION. We say that \mathbf{n} belongs to the class \mathcal{R} if there are two functions of class C^1 , $\alpha : \mathcal{B} \setminus (\mathcal{B}_0 \cup \mathbf{b}) \rightarrow \mathbb{R}$ and $\beta : \mathcal{B} \setminus \mathcal{S}(s) \rightarrow [0, \pi]$, such that in $\mathcal{B} \setminus (\mathcal{B}_0 \cup \mathbf{b})$ \mathbf{n} can be represented by the following formula

$$(2.3) \quad \mathbf{n} = \sin \beta \cos \alpha \mathbf{e}_r + \sin \beta \sin \alpha \mathbf{e}_\theta + \cos \beta \mathbf{e}_z .$$

We say that \mathbf{n} belongs to \mathcal{R}_0 , if it belongs to \mathcal{R} and $\alpha = 0$ in (2.3).

REMARK 1. Locally, each field \mathbf{n} that minimizes \mathcal{F} in \mathcal{A} is represented by (2.3). Nevertheless, while β is well defined in the whole of $\mathcal{B} \setminus \mathcal{S}(s)$ for all smooth fields \mathbf{n} , this is not

the case for α , whose determination need not be unique along all closed curves in $\mathcal{B} \setminus \mathcal{B}_0$, though it is well defined on all simply connected subset of $\mathcal{B} \setminus \mathcal{B}_0$.

We assume that \mathcal{F} attains minimum in \mathcal{R} .

In our opinion this is not a severe assumption. We think that the existence of a minimizer of \mathcal{F} in \mathcal{R} is to be proved along the same lines of though pursued in [4]. Here, as is made plain by (2.6) below, $\sin \beta$ is to be treated the same way as s and \mathcal{B}_0 is to play the same rôle as $\mathcal{S}(s)$ in the existence theorem of [4]. In so doing, one should achieve a weak existence result in a suitable subclass of \mathcal{A} . Then, a regularity result for α, β , and s , away from \mathcal{B}_0 , should be at hand. We think so since both α and β are locally determined in terms of \mathbf{n} , to within an additive constant.

THEOREM. *Let σ_0 be a function of class C^1 . If both s and \mathbf{n} are prescribed on \mathcal{D} as Lipschitz continuous functions such that*

$$(2.4) \quad s|_{\mathcal{D}} \geq s_0 > 0, \quad \mathbf{n} \cdot \mathbf{e}_r|_{\mathcal{D} \setminus \mathcal{b}} > 0, \quad \mathbf{n} \cdot \mathbf{e}_\theta|_{\mathcal{D} \setminus \mathcal{b}} = 0,$$

then the minimum of \mathcal{F} in \mathcal{R} is the same as the minimum in \mathcal{R}_0 . Furthermore, if the set $\mathcal{B} \setminus \mathcal{B}_0$ is connected for all minimizers of \mathcal{F} in \mathcal{R}_0 , then there is no other minimizer of \mathcal{F} in \mathcal{R} .

Proof. Let the pair (s, \mathbf{n}) , with s a continuous function and \mathbf{n} in \mathcal{R} , be a minimizer of \mathcal{F} . It follows from (2.3) that

$$(2.5) \quad \int_{\mathcal{B} \setminus \mathcal{B}_0} s^2 |\nabla \mathbf{n}|^2 = \int_{\mathcal{B} \setminus \mathcal{B}_0} s^2 \left\{ \beta_r^2 + \sin^2 \beta \alpha_r^2 + \frac{1}{r^2} (\beta_\theta^2 + \sin^2 \beta (\alpha_\theta + 1)^2) + \beta_z^2 + \sin^2 \beta \alpha_z^2 \right\} r dr d\theta dz =: \mathcal{G}[\alpha, \beta],$$

where partial derivatives with respect to co-ordinates are denoted by subscripts. If the pair (s, \mathbf{u}) minimizes \mathcal{F} subject to (2.4), then the pair (α, β) minimizes \mathcal{G} subject to

$$(2.6) \quad \alpha|_{\mathcal{D} \setminus \mathcal{b}} = 0$$

and to the approximate boundary condition for β on $\mathcal{D} \setminus \mathcal{b}$. In particular, the function $g : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ defined by

$$(2.7) \quad g(\lambda) := \mathcal{G}[\lambda \alpha, \beta]$$

attains its minimum at $\lambda = 1$. It is easily seen that

$$(2.8) \quad g(\lambda) = a\lambda^2 + b\lambda + c,$$

where

$$(2.9) \quad a := \int_{\mathcal{B} \setminus \mathcal{B}_0} s^2 \sin^2 \beta |\nabla \alpha|^2, \quad b := 2 \int_{\mathcal{B} \setminus \mathcal{B}_0} \frac{s^2}{r^2} \sin^2 \beta \alpha_\theta, \quad c := \int_{\mathcal{B} \setminus \mathcal{B}_0} |\nabla \beta|^2.$$

We show now that $g(1)$ is the minimum of g only if both a and b vanish. We see from (2.9)₁ that $a \geq 0$, and so $b \leq 0$, otherwise we would have $g(\lambda) < g(1)$ for all $\lambda \in [0, 1[$. Were $a > 0$, the minimum of g would be $g(-\frac{b}{2a})$; since

$$g\left(-\frac{b}{2a}\right) - g(1) = -\frac{(b + 2a)^2}{4a^2} < 0,$$

we conclude that $a = 0$. Thus, g is a linear function, which attains minimum at $\lambda = 1$ only if it is constant. Likewise, α minimizes \mathcal{G} only if it is constant in $\mathcal{B} \setminus \mathcal{B}_0$. Furthermore, the value of α does not affect the value of \mathcal{G} , and so the minimum of \mathcal{F} in \mathcal{R} is the same as in \mathcal{R}_0 where $\alpha \equiv 0$.

If a pair (s, \mathbf{n}_0) , with $\mathbf{n}_0 \in \mathcal{R}_0$, minimizes \mathcal{F} and the set $\mathcal{B} \setminus \mathcal{B}_0$ corresponding to \mathbf{n}_0 is not connected, then infinitely many minimizers of \mathcal{F} in \mathcal{R} can be produced by merely rotating \mathbf{n}_0 about \mathbf{e}_z in each component of $\mathcal{B} \setminus \mathcal{B}_0$ whose boundary does not include \mathcal{D} . If, on the contrary, the set $\mathcal{B} \setminus \mathcal{B}_0$ is connected for all minimizers of \mathcal{F} in \mathcal{R}_0 , there is no other orientation of \mathcal{R} with constant α that minimizes \mathcal{F} . \square

The following Corollary is an immediate consequence of the Theorem.

COROLLARY. *If \mathbf{n}_0 is the only minimizer of \mathcal{F} in \mathcal{R}_0 and the set $\mathcal{B} \setminus \mathcal{B}_0$ is connected, then \mathbf{n}_0 is the only minimizer of \mathcal{F} in \mathcal{R} .*

REMARK 2. It is easily seen that the Theorem above holds also if $\mathbf{n} \cdot \mathbf{e}_r|_{\mathcal{D} \setminus \mathfrak{b}} < 0$, but if $\mathbf{n} \cdot \mathbf{e}_r$ has not the same sign on $\mathcal{D} \setminus \mathfrak{b}$, explicit counterexamples have been constructed in [10].

REMARK 3. At each point of \mathcal{B} away from \mathfrak{b} all fields of \mathcal{R}_0 lie in a longitudinal plane of \mathcal{B} . Thus, the Theorem tells us that \mathcal{F} possess in \mathcal{R} a *longitudinal minimizer*. Then, the uniqueness criterion established in the Corollary can be rephrased as follows: If there is only one longitudinal minimizer of \mathcal{F} and for it $\mathcal{B} \setminus \mathcal{B}_0$ is connected, then the longitudinal minimizer of \mathcal{F} is its unique minimizer in the whole of \mathcal{R} .

3. An example. In this section we apply the Theorem proved in Section 2 to the special case in which \mathcal{B} is a circular cylinder:

$$(3.1) \quad \mathcal{B} = \{0 + r\mathbf{e}_r + z\mathbf{e}_z | z \in]0, H[, r \in [0, R[, \vartheta \in [0, 2\pi[\}.$$

We denote by $\partial_R \mathcal{B}$ the lateral boundary of \mathcal{B} . Here, of course, \mathfrak{b} intersects \mathcal{B} .

THEOREM. Let σ_0 be of class C^1 . If

$$(3.2) \quad \mathbf{n}|_{\partial_R \mathcal{B}} = \cos \varphi_0 \mathbf{e}_r + \sin \varphi_0 \mathbf{e}_z \quad \text{and} \quad s|_{\partial_R \mathcal{B}} = s_0 ,$$

where $\varphi_0 \in]-\frac{\pi}{2}, \frac{\pi}{2}[$ and $s_0 > 0$ are given constants, then all minimizers of \mathcal{F} in \mathcal{R}_0 have the form

$$(3.3) \quad \mathbf{n} = \cos \varphi(r) \mathbf{e}_r + \sin \varphi(r) \mathbf{e}_z$$

and all minimizing s depend only on r .

Proof. Away from $\mathfrak{S}(s) \cup \mathfrak{b}$ any field of \mathcal{R}_0 can be represented as follows

$$(3.4) \quad \mathbf{n} = \cos \varphi(r, \vartheta, z) \mathbf{e}_r + \sin \varphi(r, \vartheta, z) \mathbf{e}_z ,$$

where φ is a function of class C^1 into $]-\frac{\pi}{2}, \frac{\pi}{2}[$. To prove (3.3) it remains to see that φ is indeed a function of r only. If \mathbf{n} is as in (3.4), an easy computation shows that

$$(3.5) \quad \begin{aligned} \mathcal{F}[s, \mathbf{n}] = & \kappa \int_0^{2\pi} d\theta \int_0^H dz \int_0^R dr r \left\{ k \left(s_r^2 + s_z^2 + \frac{s_\theta^2}{r^2} \right) + \right. \\ & \left. + s^2 \left(\varphi_r^2 + \varphi_z^2 + \frac{\cos^2 \varphi + \varphi_\theta^2}{r^2} \right) + \sigma_0(s) \right\} . \end{aligned}$$

Thus,

$$(3.6) \quad \mathcal{F}[s, \mathbf{n}] \geq 2\pi H \kappa \inf_A F[t, \psi]$$

where

$$(3.7) \quad F[t, \psi] := \int_0^R dr r \left\{ kt'^2 + t^2 \left(\psi'^2 + \frac{\cos^2 \psi}{r^2} \right) + \sigma_0(t) \right\}$$

and

$$(3.8) \quad A := \{(t, \psi) | t \in AC(0, R), \psi \in AC_{\text{loc}}(]0, R[\setminus S(t)) : t(R) = s_0, \psi(R) = \varphi_0\} .$$

In (3.7) a prime denotes differentiation with respect to r . In (3.8) the set $S(t)$ is defined as

$$(3.9) \quad S(t) := \{r \in [0, R] | t(r) = 0\} .$$

In (3.6) the equality holds if, and only if, both s and φ depend only on r and there is a pair $(t_0, \psi_0) \in A$ such that

$$(3.10) \quad F[t_0, \psi_0] = \inf_A F[t, \psi] .$$

Since the infimum of F is indeed attained in A , as one sees in Section 4 of [11], the proof of the Theorem is complete. \square

REMARK 1. The problem of finding the longitudinal minimizers of \mathcal{F} has been solved in [9] under the assumption that $\varphi_0 = 0$. Apart from the reflection that changes φ into $-\varphi$, there is only one longitudinal minimizer of \mathcal{F} for all $k > 0$. Furthermore, $\mathcal{B}_0 = \mathcal{B} \cap \mathfrak{b}$ for all $k > 0$, and so the longitudinal minimizer is the only minimizer of \mathcal{F} in \mathcal{R} , *modulo* a reflection.

REMARK 2. We conjecture that if \mathcal{B} is axis-symmetric and both s and \mathbf{n} are prescribed on the lateral boundary of \mathcal{B} as in (3.2), then all orientations that minimize \mathcal{F} belong to \mathcal{R} , and so the uniqueness criterion stated in Section 2 holds indeed among all admissible fields.

Acknowledgement. We are indebted to R.M. HARDT for an enlighten discussion about the longitudinal minimizers discussed here. An early version of this paper was written while E. VIRGA was visiting the Institute for Mathematics and its Applications at the University of Minnesota. The hospitality and the secretarial support of this Institute is gratefully acknowledged.

REFERENCES

- [1] J.L. ERICKSEN, *Liquid crystals with variable degree of orientation*, IMA Preprint Series No. 559 (1989).
- [2] C.W. OSEEN, *The theory of liquid crystals*, Trans. Faraday Soc., **29** (1933), 883–899.
- [3] F.C. FRANK, *On the theory of liquid crystals*, Discuss. Faraday Soc., **25** (1958), 19–28.
- [4] L. AMBROSIO, *Existence of minimal energy configurations of nematic liquid crystals with variable degree of orientation*, Manuscripta Math., **68** (1990), 215–228.
- [5] F.H. LIN, *Nonlinear theory of defects in nematic liquid crystals: phase transition and flow phenomena*, Comm. Pure Appl. Math., **42** (1989), 789–814.
- [6] L. AMBROSIO, *Regularity of solutions of a degenerate elliptic variational problem*, Manuscripta Math., **68** (1990), 309–326.
- [7] F.H. LIN, *On nematic liquid crystals with variable degree of orientation*, to appear in Comm. Pure Appl. Math. (1990).
- [8] E.G. VIRGA, *New variational problems in the statics of liquid crystals*, Preprint, 1990.
- [9] D. ROCCATO, V.J. MIZEL and E.G. VIRGA, *A variational problem for nematic liquid crystals with variable degree of orientation*, Preprint, 1990.
- [10] R.M. HARDT, *Point and line singularities in liquid crystals*, to appear in the Proceedings of the Conference on Variational Problems, École Normale Supérieure, Paris, 1988.
- [11] L. AMBROSIO and E.G. VIRGA, *A boundary-value problem for nematic liquid crystals with variable degree of orientation*, to appear in Arch. Rational Mech. Anal., (1990).

Recent IMA Preprints

#	Author/s	Title
693	A. Eden, A.J. Milani and B. Nicolaenko	Finite dimensional exponential attractors for semilinear wave equations with damping
694	A. Eden, C. Foias, B. Nicolaenko & R. Temam	Inertial sets for dissipative evolution equations
695	A. Eden, C. Foias, B. Nicolaenko & R. Temam	Hölder continuity for the inverse of Mañé's projection
696	Michel Chipot and Charles Collins	Numerical approximations in variational problems with potential wells
697	Huanan Yang	Nonlinear wave analysis and convergence of MUSCL schemes
698	László Gerencsér and Zsuzsanna Vágó	A strong approximation theorem for estimator processes in continuous time
699	László Gerencsér	Multiple integrals with respect to L -mixing processes
700	David Kinderlehrer and Pablo Pedregal	Weak convergence of integrands and the Young measure representation
701	Bo Deng	Symbolic dynamics for chaotic systems
702	P. Galdi, D.D. Joseph, L. Preziosi, S. Rionero	Mathematical problems for miscible, incompressible fluids with Korteweg stresses
703	Charles Collins and Mitchell Luskin	Optimal order error estimates for the finite element approximation of the solution of a nonconvex variational problem
704	Peter Gritzmann and Victor Klee	Computational complexity of inner and outer j -radii of polytopes in finite-dimensional normed spaces
705	A. Ronald Gallant and George Tauchen	A nonparametric approach to nonlinear time series analysis: estimation and simulation
706	H.S. Dumas, J.A. Ellison and A.W. Sáenz	Axial channeling in perfect crystals, the continuum model and the method of averaging
707	M.A. Kaashoek and S.M. Verduyn Lunel	Characteristic matrices and spectral properties of evolutionary systems
708	Xinfu Chen	Generation and Propagation of interfaces in reaction diffusion systems
709	Avner Friedman and Bei Hu	Homogenization approach to light scattering from polymer-dispersed liquid crystal films
710	Yoshihisa Morita and Shuichi Jimbo	ODEs on inertial manifolds for reaction-diffusion systems in a singularly perturbed domain with several thin channels
711	Wenxiong Liu	Blow-up behavior for semilinear heat equations: multi-dimensional case
712	Hi Jun Choe	Hölder continuity for solutions of certain degenerate parabolic systems
713	Hi Jun Choe	Regularity for certain degenerate elliptic double obstacle problems
714	Fernando Reitich	On the slow motion of the interface of layered solutions to the scalar Ginzburg–Landau equation
715	Xinfu Chen and Fernando Reitich	Local existence and uniqueness of solutions of the Stefan problem with surface tension and kinetic undercooling
716	C.C. Lim, J.M. Pimbley, C. Schmeiser and D.W. Schwendeman	Rotating waves for semiconductor inverter rings
717	W. Balsler, B.L.J. Braaksma, J.-P. Ramis and Y. Sibuya	Multisummability of formal power series solutions of linear ordinary differential equations
718	Peter J. Olver and Chehrzad Shakiban	Dissipative decomposition of partial differential equations
719	Clark Robinson	Homoclinic bifurcation to a transitive attractor of Lorenz type, II
720	Michelle Schatzman	A simple proof of convergence of the QR algorithm for normal matrices without shifts
721	Ian M. Anderson, Niky Kamran and Peter J. Olver	Internal, external and generalized symmetries
722	C. Foias and J.C. Saut	Asymptotic integration of Navier–Stokes equations with potential forces. I
723	Ling Ma	The convergence of semidiscrete methods for a system of reaction-diffusion equations
724	Adelina Georgescu	Models of asymptotic approximation
725	A. Makagon and H. Salehi	On bounded and harmonizable solutions on infinite order arma systems
726	San-Yih Lin and Yan-Shin Chin	An upwind finite-volume scheme with a triangular mesh for conservation laws
727	J.M. Ball, P.J. Holmes, R.D. James, R.L. Pego & P.J. Swart	On the dynamics of fine structure
728	KangPing Chen and Daniel D. Joseph	Lubrication theory and long waves
729	J.L. Ericksen	Local bifurcation theory for thermoelastic Bravais lattices
730	Mario Taboada and Yuncheng You	Some stability results for perturbed semilinear parabolic equations
731	A.J. Lawrance	Local and deletion influence
732	Bogdan Vernescu	Convergence results for the homogenization of flow in fractured porous media
733	Xinfu Chen and Avner Friedman	Mathematical modeling of semiconductor lasers
734	Yongzhi Xu	Scattering of acoustic wave by obstacle in stratified medium
735	Songmu Zheng	Global existence for a thermodynamically consistent model of phase field type
736	Heinrich Freistühler and E. Bruce Pitman	A numerical study of a rotationally degenerate hyperbolic

- system part I: the Riemann problem
- 737 **Epifanio G. Virga**, New variational problems in the statics of liquid crystals
- 738 **Yoshikazu Giga and Shun'ichi Goto**, Geometric evolution of phase-boundaries
- 739 **Ling Ma**, Large time study of finite element methods for 2D Navier–Stokes equations
- 740 **Mitchell Luskin and Ling Ma**, Analysis of the finite element approximation of microstructure in micromagnetics
- 741 **M. Chipot**, Numerical analysis of oscillations in nonconvex problems
- 742 **J. Carrillo and M. Chipot**, The dam problem with leaky boundary conditions
- 743 **Eduard Harabetian and Robert Pego**, Efficient hybrid shock capturing schemes
- 744 **B.L.J. Braaksma**, Multisummability and Stokes multipliers of linear meromorphic differential equations
- 745 **Tae Il Jeon and Tze-Chien Sun**, A central limit theorem for non-linear vector functionals of vector Gaussian processes
- 746 **Chris Grant**, Solutions to evolution equations with near-equilibrium initial values
- 747 **Mario Taboada and Yuncheng You**, Invariant manifolds for retarded semilinear wave equations
- 748 **Peter Rejto and Mario Taboada**, Unique solvability of nonlinear Volterra equations in weighted spaces
- 749 **Hi Jun Choe**, Holder regularity for the gradient of solutions of certain singular parabolic equations
- 750 **Jack D. Dockery**, Existence of standing pulse solutions for an excitable activator-inhibitory system
- 751 **Jack D. Dockery and Roger Lui**, Existence of travelling wave solutions for a bistable evolutionary ecology model
- 752 **Giovanni Alberti, Luigi Ambrosio and Giuseppe Buttazzo**, Singular perturbation problems with a compact support semilinear term
- 753 **Emad A. Fatemi**, Numerical schemes for constrained minimization problems
- 754 **Y. Kuang and H.L. Smith**, Slowly oscillating periodic solutions of autonomous state-dependent delay equations
- 755 **Emad A. Fatemi**, A new splitting method for scalar conservation laws with stiff source terms
- 756 **Hi Jun Choe**, A regularity theory for a more general class of quasilinear parabolic partial differential equations and variational inequalities
- 757 **Haitao Fan**, A vanishing viscosity approach on the dynamics of phase transitions in Van Der Waals fluids
- 758 **T.A. Osborn and F.H. Molzahn**, The Wigner–Weyl transform on tori and connected graph propagator representations
- 759 **Avner Friedman and Bei Hu**, A free boundary problem arising in superconductor modeling
- 760 **Avner Friedman and Wenxiong Liu**, An augmented drift-diffusion model in semiconductor device
- 761 **Avner Friedman and Miguel A. Herrero**, Extinction and positivity for a system of semilinear parabolic variational inequalities
- 762 **David Dobson and Avner Friedman**, The time-harmonic Maxwell equations in a doubly periodic structure
- 763 **Hi Jun Choe**, Interior behaviour of minimizers for certain functionals with nonstandard growth
- 764 **Vincenzo M. Tortorelli and Epifanio G. Virga**, Axis-symmetric boundary-value problems for nematic liquid crystals with variable degree of orientation
- 765 **Nikan B. Firoozye and Robert V. Kohn**, Geometric parameters and the relaxation of multiwell energies
- 766 **Haitao Fan and Marshall Slemrod**, The Riemann problem for systems of conservation laws of mixed type
- 767 **Joseph D. Fehribach**, Analysis and application of a continuation method for a self-similar coupled Stefan system
- 768 **C. Foias, M.S. Jolly, I.G. Kevrekidis and E.S. Titi**, Dissipativity of numerical schemes
- 769 **D.D. Joseph, T.Y.J. Liao and J.-C. Saut**, Kelvin–Helmholtz mechanism for side branching in the displacement of light with heavy fluid under gravity
- 770 **Chris Grant**, Solutions to evolution equations with near-equilibrium initial values
- 771 **B. Cockburn, F. Coquel, Ph. LeFloch and C.W. Shu**, Convergence of finite volume methods
- 772 **N.G. Lloyd and J.M. Pearson**, Computing centre conditions for certain cubic systems
- 773 **João Palhoto Matos**, Young measures and the absence of fine microstructures in the $\alpha - \beta$ quartz phase transition
- 774 **L.A. Peletier & W.C. Troy**, Self-similar solutions for infiltration of dopant into semiconductors
- 775 **H. Scott Dumas and James A. Ellison**, Nekhoroshev's theorem, ergodicity, and the motion of energetic charged particles in crystals
- 776 **Stathis Filippas and Robert V. Kohn**, Refined asymptotics for the blowup of $u_t - \Delta u = u^p$.
- 777 **Patricia Bauman, Nicholas C. Owen and Daniel Phillips**, Maximum principles and a priori estimates for an incompressible material in nonlinear elasticity
- 778 **Patricia Bauman, Nicholas C. Owen and Daniel Phillips**, Maximal smoothness of solutions to certain Euler–Lagrange equations from nonlinear elasticity
- 779 **Jack Carr and Robert Pego**, Self-similarity in a coarsening model in one dimension
- 780 **J.M. Greenberg**, The shock generation problem for a discrete gas with short range repulsive forces
- 781 **George R. Sell and Mario Taboada**, Local dissipativity and attractors for the Kuramoto–Sivashinsky equation in thin 2D domains
- 782 **T. Subba Rao**, Analysis of nonlinear time series (and chaos) by bispectral methods