

TIME AVERAGING AND TURBULENCE TERMS IN METEOROLOGY

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ABSTRACT. We discuss averaging in time for the planetary boundary layer in the Boussinesq approximation. We introduce a notion of instantaneous turbulent kinetic energy (ITKE) and then derive a balance equation for ITKE.

1. INTRODUCTION

It is important to have a rigorous treatment of averaging (or smoothing) processes applied to the Euler and Navier-Stokes Equations, especially since averaged values are registered by meteorological devices, they are also used for numerical simulations (ARPS, etc.).

2. VARIOUS NOTIONS OF TIME AVERAGING

Any (scalar) flow $u(x, t)$ may be represented as (cf. [1, 2])

$$(2.1) \quad u(x, t) = \sum_{k=1}^K u_k(x, t),$$

where each $u_k(x, t)$ ($k = 1, 2, \dots, K$) describes a specific feature of the flow, where x is a spatial variable and t is time

In atmospheric science it is customary to view (2.1) as

$$(2.2) \quad u(x, t) = \overline{u(x, t)} + u'(x, t),$$

where $\overline{u(x, t)}$ is a **slowly varying average** component and $u'(x, t)$ is a **rapidly fluctuating** or **turbulent** component.

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Standard (Reynolds) Time Averaging. We assume that $u(x, t)$ is a C^1 function of its variables. We assume that there is a typical averaging period T that includes a “large” number of oscillations in $u'(x, t)$, then $\overline{u(x, t)}$ is independent of time and is given by

$$(2.3) \quad \overline{u(x, t)} \equiv \frac{1}{T} \int_0^T u(x, s) ds,$$

so in this case $\overline{u(x, t)}$ is not slowly varying but a constant (in time) function for $t \in [0, T]$. It is easy to show that

$$(2.4) \quad \overline{u'(x, t)} = 0, \quad \overline{\overline{u(x, t)}v(x, t)} = \overline{u(x, t)}\overline{v(x, t)}, \quad \overline{\overline{u(x, t)}v'(x, t)} = 0$$

and

$$(2.5) \quad \overline{\overline{u(x, t)}} = \overline{u(x, t)}$$

Using (2.4), we have

$$(2.6) \quad \overline{u(x, t)v(x, t)} = \overline{u(x, t)}\overline{v(x, t)} + \overline{u'(x, t)v'(x, t)}$$

and

$$(2.7) \quad \begin{aligned} \overline{u(x, t)v(x, t)w(x, t)} &= \overline{u(x, t)}\overline{v(x, t)}\overline{w(x, t)} + \overline{u(x, t)v'(x, t)w'(x, t)} \\ &+ \overline{v(x, t)u'(x, t)w'(x, t)} + \overline{w(x, t)v'(x, t)u'(x, t)} \\ &+ \overline{u'(x, t)v'(x, t)w'(x, t)} \end{aligned}$$

Also for spatial derivatives under appropriate smoothness assumptions we have

$$(2.8) \quad \overline{\frac{\partial}{\partial x}u(x, t)} = \frac{\partial}{\partial x}\overline{u(x, t)}, \quad \overline{\frac{\partial}{\partial x}u'(x, t)} = \frac{\partial}{\partial x}\overline{u'(x, t)} = 0$$

However, temporal derivatives and time averages do not commute, i.e.

$$(2.9) \quad \begin{aligned} \overline{\frac{\partial}{\partial t}u(x, t)} &= \overline{\frac{\partial}{\partial t}u'(x, t)} \\ &= \frac{1}{T}(u(x, T) - u(x, 0)) = \frac{1}{T}(u'(x, T) - u'(x, 0)) \\ &\neq 0 = \frac{\partial}{\partial t}\overline{u(x, t)} = \frac{\partial}{\partial t}\overline{u'(x, t)}. \end{aligned}$$

The problem with (2.9) might be helped by assuming that there exists an appropriate averaging period T so that

$$(2.10) \quad \overline{u(x, t)} = u(x, 0) = u(x, T), \quad u'(x, 0) = u'(x, T) = 0.$$

However, this is a superficial remedy since $\overline{u(x, t)}$ is a constant (in time) function and temporal derivative is zero.

Running Averages. We define

$$(2.11) \quad \overline{u(x, t)} \equiv \frac{1}{2T} \int_{t-T}^{t+T} u(x, s) ds.$$

Then, assuming (2.2), we see that (2.4)–(2.7) fail, i.e.,

$$(2.12) \quad \overline{u'(x, t)} \neq 0, \quad \overline{\overline{u(x, t)v(x, t)}} \neq \overline{u(x, t)}\overline{v(x, t)}, \quad \overline{\overline{u(x, t)v'(x, t)}} \neq 0,$$

$$(2.13) \quad \overline{\overline{u(x, t)v(x, t)}} \neq \overline{u(x, t)}\overline{v(x, t)} + \overline{u'(x, t)v'(x, t)},$$

However, (2.8) holds with some adjustments:

$$(2.14) \quad \overline{\frac{\partial}{\partial x} u(x, t)} = \frac{\partial}{\partial x} \overline{u(x, t)}, \quad \overline{\frac{\partial}{\partial x} u'(x, t)} = \frac{\partial}{\partial x} \overline{u'(x, t)} \neq 0$$

and we see that (2.11) repairs the problem in (2.9) since now

$$(2.15) \quad \overline{\frac{\partial}{\partial t} u(x, t)} = \frac{1}{2T} (u(x, t+T) - u(x, t-T)) = \frac{\partial}{\partial t} \overline{u(x, t)}.$$

3. TURBULENT KINETIC ENERGY

In the Boussinesq Approximation for incompressible media we have

$$(3.1) \quad \frac{Du}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x} + fv + F_{rx}$$

$$(3.2) \quad \frac{Dv}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial y} - fu + F_{ry}$$

$$(3.3) \quad \frac{Dw}{Dt} = -\frac{1}{\rho_0} \frac{\partial p}{\partial z} + g \frac{\theta}{\theta_0} + F_{rz}$$

$$(3.4) \quad \frac{D\theta}{Dt} = -w \frac{d\theta_0}{dz}$$

$$(3.5) \quad \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0,$$

where u, v, w are the velocity components, p is pressure, θ is potential temperature and $\rho_0(z), \theta_0(z)$ are given functions. Note that by (3.5)

$$(3.6) \quad \begin{aligned} \frac{Du}{Dt} &\equiv \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \\ &= \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} + u \underbrace{\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right)}_0 \\ &= \frac{\partial u}{\partial t} + \frac{\partial uu}{\partial x} + \frac{\partial uv}{\partial y} + \frac{\partial uw}{\partial z}. \end{aligned}$$

Also note that by (2.8) and (3.5) we have

$$(3.7) \quad \frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} + \frac{\partial \bar{w}}{\partial z} = 0.$$

We set (cf. [1])

$$(3.8) \quad \frac{\bar{D}}{Dt} = \frac{\partial}{\partial t} + \bar{u} \frac{\partial}{\partial x} + \bar{v} \frac{\partial}{\partial y} + \bar{w} \frac{\partial}{\partial z}.$$

Definition 3.1. Instantaneous turbulent kinetic energy ITKE (x, y, z, t) is a function defined by

$$(3.9) \quad \text{ITKE}(x, y, z, t) := \frac{1}{2} ((u')^2 + (v')^2 + (w')^2).$$

A standard notion of turbulent kinetic energy

$$(3.10) \quad TKE(x, y, z) \equiv \frac{1}{2} \overline{((u')^2 + (v')^2 + (w')^2)}$$

and various balance equations are available for TKE¹ Note that

$$(3.12) \quad TKE(x, y, z) = \overline{\text{ITKE}(x, y, z, t)}$$

¹The balance equation

$$(3.11) \quad \frac{\bar{D}}{Dt}(TKE) = MP + BPL + TR - \varepsilon$$

is known but differs from our balance.

Theorem 3.1. *Let u, v, w, p, θ be the functions of (x, y, z, t) satisfying the Boussinesq approximation system of partial differential equations (3.1)–(3.5). Then*

$$(3.13) \quad \overline{\frac{D}{Dt}(ITKE)} = MP + BPL + TR - \varepsilon,$$

where the mechanical production

$$(3.14) \quad \begin{aligned} MP \equiv & -\overline{u'u'} \frac{\partial \bar{u}}{\partial x} - \overline{u'v'} \frac{\partial \bar{u}}{\partial y} - \overline{u'w'} \frac{\partial \bar{u}}{\partial z} \\ & - \overline{v'u'} \frac{\partial \bar{v}}{\partial x} - \overline{v'v'} \frac{\partial \bar{v}}{\partial y} - \overline{v'w'} \frac{\partial \bar{v}}{\partial z} \\ & - \overline{w'u'} \frac{\partial \bar{w}}{\partial x} - \overline{w'v'} \frac{\partial \bar{w}}{\partial y} - \overline{w'w'} \frac{\partial \bar{w}}{\partial z} \\ & - \overline{u'u'} \frac{\partial u'}{\partial x} - \overline{u'v'} \frac{\partial u'}{\partial y} - \overline{u'w'} \frac{\partial u'}{\partial z} \\ & - \overline{v'u'} \frac{\partial v'}{\partial x} - \overline{v'v'} \frac{\partial v'}{\partial y} - \overline{v'w'} \frac{\partial v'}{\partial z} \\ & - \overline{w'u'} \frac{\partial w'}{\partial x} - \overline{w'v'} \frac{\partial w'}{\partial y} - \overline{w'w'} \frac{\partial w'}{\partial z}, \end{aligned}$$

the buoyancy production or loss

$$(3.15) \quad BPL \equiv \overline{w'\theta'} \frac{g}{\theta_0},$$

the transport redistribution

$$(3.16) \quad TR \equiv -\frac{1}{\rho_0} \overline{\left(u' \frac{\partial p'}{\partial x} + v' \frac{\partial p'}{\partial y} + w' \frac{\partial p'}{\partial z} \right)} = -\frac{1}{\rho_0} \overline{\mathbf{u}' \cdot \nabla p'},$$

and the frictional dissipation

$$(3.17) \quad \varepsilon \equiv -\overline{(u'F_{rx}' + v'F_{ry}' + w'F_{rz}')}) = -\overline{\mathbf{u}' \cdot \mathbf{F}_r'}.$$

Proof. We average (3.6) with use of (2.6) and (3.7) to obtain

$$(3.18) \quad \begin{aligned} \overline{\frac{Du}{Dt}} &= \overline{\frac{\partial u}{\partial t}} + \frac{\partial}{\partial x}(\overline{u\bar{u}} + \overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u\bar{v}} + \overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u\bar{w}} + \overline{u'w'}) \\ &= \overline{\frac{\partial u}{\partial t}} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} + \bar{w} \frac{\partial \bar{u}}{\partial z} + \frac{\partial}{\partial x}(\overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u'w'}) \\ &= \overline{\frac{\partial u'}{\partial t}} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} + \bar{w} \frac{\partial \bar{u}}{\partial z} + \frac{\partial}{\partial x}(\overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u'w'}). \end{aligned}$$

Then the mean equations of (3.1)–(3.5) will be (3.7) and

$$(3.19) \quad \frac{\overline{\partial u}}{\partial t} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} + \bar{w} \frac{\partial \bar{u}}{\partial z} = -\frac{1}{\rho_0} \frac{\partial \bar{p}}{\partial x} + f\bar{v} - \left[\frac{\partial}{\partial x}(\overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u'w'}) \right] + \overline{F_{rx}},$$

$$(3.20) \quad \frac{\overline{\partial v}}{\partial t} + \bar{u} \frac{\partial \bar{v}}{\partial x} + \bar{v} \frac{\partial \bar{v}}{\partial y} + \bar{w} \frac{\partial \bar{v}}{\partial z} = -\frac{1}{\rho_0} \frac{\partial \bar{p}}{\partial y} - f\bar{u} - \left[\frac{\partial}{\partial x}(\overline{v'u'}) + \frac{\partial}{\partial y}(\overline{v'v'}) + \frac{\partial}{\partial z}(\overline{v'w'}) \right] + \overline{F_{ry}},$$

$$(3.21) \quad \frac{\overline{\partial w}}{\partial t} + \bar{u} \frac{\partial \bar{w}}{\partial x} + \bar{v} \frac{\partial \bar{w}}{\partial y} + \bar{w} \frac{\partial \bar{w}}{\partial z} = -\frac{1}{\rho_0} \frac{\partial \bar{p}}{\partial z} + g \frac{\bar{\theta}}{\theta_0} - \left[\frac{\partial}{\partial x}(\overline{w'u'}) + \frac{\partial}{\partial y}(\overline{w'v'}) + \frac{\partial}{\partial z}(\overline{w'w'}) \right] + \overline{F_{rz}},$$

$$(3.22) \quad \frac{\overline{\partial \theta}}{\partial t} + \bar{u} \frac{\partial \bar{\theta}}{\partial x} + \bar{v} \frac{\partial \bar{\theta}}{\partial y} + \bar{w} \frac{\partial \bar{\theta}}{\partial z} = -\bar{w} \frac{d\theta_0}{dz} - \left[\frac{\partial}{\partial x}(\overline{\theta'u'}) + \frac{\partial}{\partial y}(\overline{\theta'v'}) + \frac{\partial}{\partial z}(\overline{\theta'w'}) \right].$$

Now we first subtract the left-hand side (3.19) from the left-hand side (3.1) and multiply the result by u' , i.e.,

$$(3.23) \quad \begin{aligned} & u' \left(\frac{\partial u}{\partial t} - \frac{\overline{\partial u}}{\partial t} + u' \frac{\partial u}{\partial x} + \bar{u} \frac{\partial u'}{\partial x} + v' \frac{\partial u}{\partial y} + \bar{v} \frac{\partial u'}{\partial y} + w' \frac{\partial u}{\partial z} + \bar{w} \frac{\partial u'}{\partial z} \right) \\ &= u' \left(\frac{\partial u'}{\partial t} - \frac{\overline{\partial u}}{\partial t} + u' \frac{\partial u}{\partial x} + \bar{u} \frac{\partial u'}{\partial x} + v' \frac{\partial u}{\partial y} + \bar{v} \frac{\partial u'}{\partial y} + w' \frac{\partial u}{\partial z} + \bar{w} \frac{\partial u'}{\partial z} \right). \end{aligned}$$

Next, we subtract the left-hand side (3.20) from the left-hand side (3.2) and multiply the result by v' , i.e.,

$$(3.24) \quad \begin{aligned} & v' \left(\frac{\partial v}{\partial t} - \frac{\overline{\partial v}}{\partial t} + u' \frac{\partial v}{\partial x} + \bar{u} \frac{\partial v'}{\partial x} + v' \frac{\partial v}{\partial y} + \bar{v} \frac{\partial v'}{\partial y} + w' \frac{\partial v}{\partial z} + \bar{w} \frac{\partial v'}{\partial z} \right) \\ &= v' \left(\frac{\partial v'}{\partial t} - \frac{\overline{\partial v}}{\partial t} + u' \frac{\partial v}{\partial x} + \bar{u} \frac{\partial v'}{\partial x} + v' \frac{\partial v}{\partial y} + \bar{v} \frac{\partial v'}{\partial y} + w' \frac{\partial v}{\partial z} + \bar{w} \frac{\partial v'}{\partial z} \right). \end{aligned}$$

Finally, we subtract the left-hand side (3.21) from the left-hand side (3.3) and multiply the result by w' , i.e.,

$$(3.25) \quad \begin{aligned} & w' \left(\frac{\partial w}{\partial t} - \frac{\overline{\partial w}}{\partial t} + u' \frac{\partial w}{\partial x} + \bar{u} \frac{\partial w'}{\partial x} + v' \frac{\partial w}{\partial y} + \bar{v} \frac{\partial w'}{\partial y} + w' \frac{\partial w}{\partial z} + \bar{w} \frac{\partial w'}{\partial z} \right) \\ &= w' \left(\frac{\partial w'}{\partial t} - \frac{\overline{\partial w}}{\partial t} + u' \frac{\partial w}{\partial x} + \bar{u} \frac{\partial w'}{\partial x} + v' \frac{\partial w}{\partial y} + \bar{v} \frac{\partial w'}{\partial y} + w' \frac{\partial w}{\partial z} + \bar{w} \frac{\partial w'}{\partial z} \right). \end{aligned}$$

We subtract the right-hand side (3.19) from the right-hand side (3.1) and multiply the result by u' , i.e.,

$$(3.26) \quad u' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial x} + f v' + \frac{\partial}{\partial x}(\overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u'w'}) + F_{rx}' \right)$$

Then we subtract the right-hand side (3.20) from the right-hand side (3.2) and multiply the result by v' , i.e.,

$$(3.27) \quad v' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial y} - fu' + \frac{\partial}{\partial x}(\overline{v'u'}) + \frac{\partial}{\partial y}(\overline{v'v'}) + \frac{\partial}{\partial z}(\overline{v'w'}) + F_{ry}' \right)$$

Finally, we subtract the right-hand side (3.21) from the right-hand side (3.3) and multiply the result by w' , i.e.,

$$(3.28) \quad w' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + g \frac{\theta'}{\theta_0} + \frac{\partial}{\partial x}(\overline{w'u'}) + \frac{\partial}{\partial y}(\overline{w'v'}) + \frac{\partial}{\partial z}(\overline{w'w'}) + F_{rz}' \right).$$

The sum of (3.23)–(3.25) equals the sum of (3.26)–(3.28), so

$$(3.29) \quad \begin{aligned} & u' \left(\frac{\partial u'}{\partial t} - \frac{\overline{\partial u}}{\partial t} + u' \frac{\partial u}{\partial x} + \overline{u} \frac{\partial u'}{\partial x} + v' \frac{\partial u}{\partial y} + \overline{v} \frac{\partial u'}{\partial y} + w' \frac{\partial u}{\partial z} + \overline{w} \frac{\partial u'}{\partial z} \right) \\ & + v' \left(\frac{\partial v'}{\partial t} - \frac{\overline{\partial v}}{\partial t} + u' \frac{\partial v}{\partial x} + \overline{u} \frac{\partial v'}{\partial x} + v' \frac{\partial v}{\partial y} + \overline{v} \frac{\partial v'}{\partial y} + w' \frac{\partial v}{\partial z} + \overline{w} \frac{\partial v'}{\partial z} \right) \\ & + w' \left(\frac{\partial w'}{\partial t} - \frac{\overline{\partial w}}{\partial t} + u' \frac{\partial w}{\partial x} + \overline{u} \frac{\partial w'}{\partial x} + v' \frac{\partial w}{\partial y} + \overline{v} \frac{\partial w'}{\partial y} + w' \frac{\partial w}{\partial z} + \overline{w} \frac{\partial w'}{\partial z} \right) \\ & = u' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial x} + fv' + \frac{\partial}{\partial x}(\overline{u'u'}) + \frac{\partial}{\partial y}(\overline{u'v'}) + \frac{\partial}{\partial z}(\overline{u'w'}) + F_{rx}' \right) \\ & + v' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial y} - fu' + \frac{\partial}{\partial x}(\overline{v'u'}) + \frac{\partial}{\partial y}(\overline{v'v'}) + \frac{\partial}{\partial z}(\overline{v'w'}) + F_{ry}' \right) \\ & + w' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + g \frac{\theta'}{\theta_0} + \frac{\partial}{\partial x}(\overline{w'u'}) + \frac{\partial}{\partial y}(\overline{w'v'}) + \frac{\partial}{\partial z}(\overline{w'w'}) + F_{rz}' \right) \end{aligned}$$

or

$$\begin{aligned}
& \frac{\partial}{\partial t} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right] - \left[u' \frac{\partial \bar{u}}{\partial t} + v' \frac{\partial \bar{v}}{\partial t} + w' \frac{\partial \bar{w}}{\partial t} \right] \\
& + u' u' \frac{\partial u}{\partial x} + u' \bar{u} \frac{\partial u'}{\partial x} + u' v' \frac{\partial u}{\partial y} + u' \bar{v} \frac{\partial u'}{\partial y} + u' w' \frac{\partial u}{\partial z} + u' \bar{w} \frac{\partial u'}{\partial z} \\
& + v' u' \frac{\partial v}{\partial x} + v' \bar{u} \frac{\partial v'}{\partial x} + v' v' \frac{\partial v}{\partial y} + v' \bar{v} \frac{\partial v'}{\partial y} + v' w' \frac{\partial v}{\partial z} + v' \bar{w} \frac{\partial v'}{\partial z} \\
(3.30) \quad & + w' u' \frac{\partial w}{\partial x} + w' \bar{u} \frac{\partial w'}{\partial x} + w' v' \frac{\partial w}{\partial y} + w' \bar{v} \frac{\partial w'}{\partial y} + w' w' \frac{\partial w}{\partial z} + w' \bar{w} \frac{\partial w'}{\partial z} \\
& = u' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial x} + f v' + \frac{\partial}{\partial x} (\overline{u' u'}) + \frac{\partial}{\partial y} (\overline{u' v'}) + \frac{\partial}{\partial z} (\overline{u' w'}) + F_{rx}' \right) \\
& + v' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial y} - f u' + \frac{\partial}{\partial x} (\overline{v' u'}) + \frac{\partial}{\partial y} (\overline{v' v'}) + \frac{\partial}{\partial z} (\overline{v' w'}) + F_{ry}' \right) \\
& + w' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + g \frac{\theta'}{\theta_0} + \frac{\partial}{\partial x} (\overline{w' u'}) + \frac{\partial}{\partial y} (\overline{w' v'}) + \frac{\partial}{\partial z} (\overline{w' w'}) + F_{rz}' \right)
\end{aligned}$$

We can rewrite (3.30) in the following way:

$$\begin{aligned}
& \frac{\partial}{\partial t} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right] + \bar{u} \frac{\partial}{\partial x} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right] \\
& + \bar{v} \frac{\partial}{\partial y} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right] + \bar{w} \frac{\partial}{\partial z} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right] \\
& - \left[u' \frac{\partial \bar{u}}{\partial t} + v' \frac{\partial \bar{v}}{\partial t} + w' \frac{\partial \bar{w}}{\partial t} \right] \\
(3.31) \quad & + u' u' \frac{\partial u}{\partial x} + u' v' \frac{\partial u}{\partial y} + u' w' \frac{\partial u}{\partial z} + v' u' \frac{\partial v}{\partial x} + v' v' \frac{\partial v}{\partial y} + v' w' \frac{\partial v}{\partial z} \\
& + w' u' \frac{\partial w}{\partial x} + w' v' \frac{\partial w}{\partial y} + w' w' \frac{\partial w}{\partial z} \\
& = u' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial x} + f v' + \frac{\partial}{\partial x} (\overline{u' u'}) + \frac{\partial}{\partial y} (\overline{u' v'}) + \frac{\partial}{\partial z} (\overline{u' w'}) + F_{rx}' \right) \\
& + v' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial y} - f u' + \frac{\partial}{\partial x} (\overline{v' u'}) + \frac{\partial}{\partial y} (\overline{v' v'}) + \frac{\partial}{\partial z} (\overline{v' w'}) + F_{ry}' \right) \\
& + w' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + g \frac{\theta'}{\theta_0} + \frac{\partial}{\partial x} (\overline{w' u'}) + \frac{\partial}{\partial y} (\overline{w' v'}) + \frac{\partial}{\partial z} (\overline{w' w'}) + F_{rz}' \right)
\end{aligned}$$

or

$$\begin{aligned}
& \overline{\frac{D}{Dt} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right]} \\
& - \left[u' \frac{\partial \bar{u}}{\partial t} + v' \frac{\partial \bar{v}}{\partial t} + w' \frac{\partial \bar{w}}{\partial t} \right] \\
& + u'u' \frac{\partial u}{\partial x} + u'v' \frac{\partial u}{\partial y} + u'w' \frac{\partial u}{\partial z} + v'u' \frac{\partial v}{\partial x} + v'v' \frac{\partial v}{\partial y} + v'w' \frac{\partial v}{\partial z} \\
(3.32) \quad & + w'u' \frac{\partial w}{\partial x} + w'v' \frac{\partial w}{\partial y} + w'w' \frac{\partial w}{\partial z} \\
& = u' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial x} + f v' + \frac{\partial}{\partial x} (\overline{u'u'}) + \frac{\partial}{\partial y} (\overline{u'v'}) + \frac{\partial}{\partial z} (\overline{u'w'}) + F_{rx}' \right) \\
& + v' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial y} - f u' + \frac{\partial}{\partial x} (\overline{v'u'}) + \frac{\partial}{\partial y} (\overline{v'v'}) + \frac{\partial}{\partial z} (\overline{v'w'}) + F_{ry}' \right) \\
& + w' \left(-\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + g \frac{\theta'}{\theta_0} + \frac{\partial}{\partial x} (\overline{w'u'}) + \frac{\partial}{\partial y} (\overline{w'v'}) + \frac{\partial}{\partial z} (\overline{w'w'}) + F_{rz}' \right)
\end{aligned}$$

We average (3.32) and obtain using (2.4)

$$\begin{aligned}
(3.33) \quad & \overline{\frac{D}{Dt} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right]} + \overline{u'u' \frac{\partial \bar{u}}{\partial x} + u'v' \frac{\partial \bar{u}}{\partial y} + u'w' \frac{\partial \bar{u}}{\partial z}} \\
& + \overline{v'u' \frac{\partial \bar{v}}{\partial x} + v'v' \frac{\partial \bar{v}}{\partial y} + v'w' \frac{\partial \bar{v}}{\partial z}} + \overline{w'u' \frac{\partial \bar{w}}{\partial x} + w'v' \frac{\partial \bar{w}}{\partial y} + w'w' \frac{\partial \bar{w}}{\partial z}} \\
& = -\frac{1}{\rho_0} \overline{\left(u' \frac{\partial p'}{\partial x} + v' \frac{\partial p'}{\partial y} + w' \frac{\partial p'}{\partial z} \right)} + \overline{w'\theta' \frac{g}{\theta_0}} + \overline{u'F_{rx}' + v'F_{ry}' + w'F_{rz}'}
\end{aligned}$$

or using also (2.4), (2.7), (2.8)

$$\begin{aligned}
(3.34) \quad & \overline{\frac{D}{Dt} \left[\frac{1}{2} ((u')^2 + (v')^2 + (w')^2) \right]} + \overline{u'u' \frac{\partial \bar{u}}{\partial x} + u'v' \frac{\partial \bar{u}}{\partial y} + u'w' \frac{\partial \bar{u}}{\partial z}} \\
& + \overline{v'u' \frac{\partial \bar{v}}{\partial x} + v'v' \frac{\partial \bar{v}}{\partial y} + v'w' \frac{\partial \bar{v}}{\partial z}} + \overline{w'u' \frac{\partial \bar{w}}{\partial x} + w'v' \frac{\partial \bar{w}}{\partial y} + w'w' \frac{\partial \bar{w}}{\partial z}} \\
& + \overline{u'u' \frac{\partial u'}{\partial x} + u'v' \frac{\partial u'}{\partial y} + u'w' \frac{\partial u'}{\partial z} + v'u' \frac{\partial v'}{\partial x} + v'v' \frac{\partial v'}{\partial y} + v'w' \frac{\partial v'}{\partial z}} \\
& + \overline{w'u' \frac{\partial w'}{\partial x} + w'v' \frac{\partial w'}{\partial y} + w'w' \frac{\partial w'}{\partial z}} \\
& = -\frac{1}{\rho_0} \overline{\left(u' \frac{\partial p'}{\partial x} + v' \frac{\partial p'}{\partial y} + w' \frac{\partial p'}{\partial z} \right)} + \overline{w'\theta' \frac{g}{\theta_0}} + \overline{u'F_{rx}' + v'F_{ry}' + w'F_{rz}'}
\end{aligned}$$

□

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