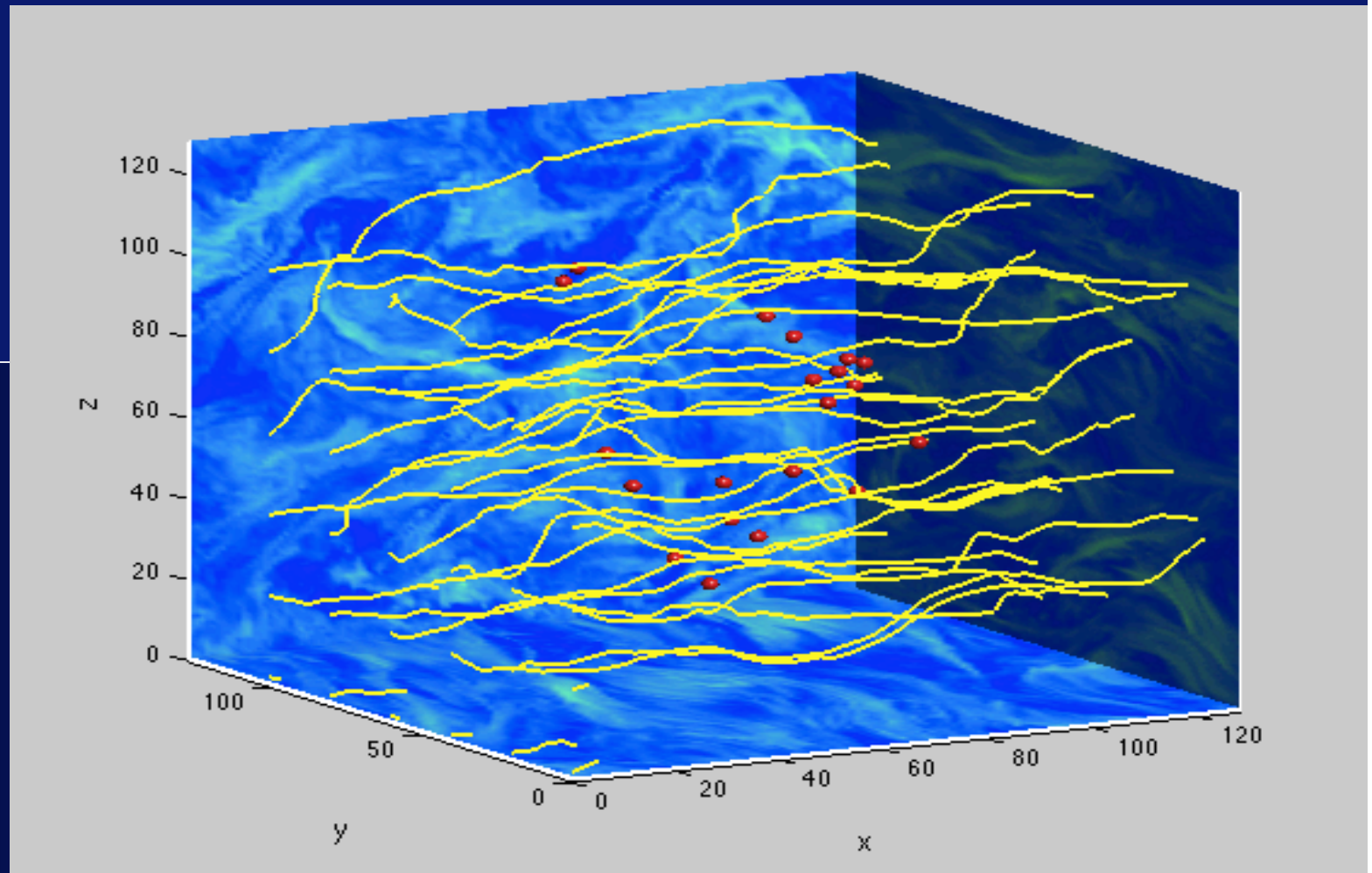


Superdiffusion and Streaming of Cosmic Rays

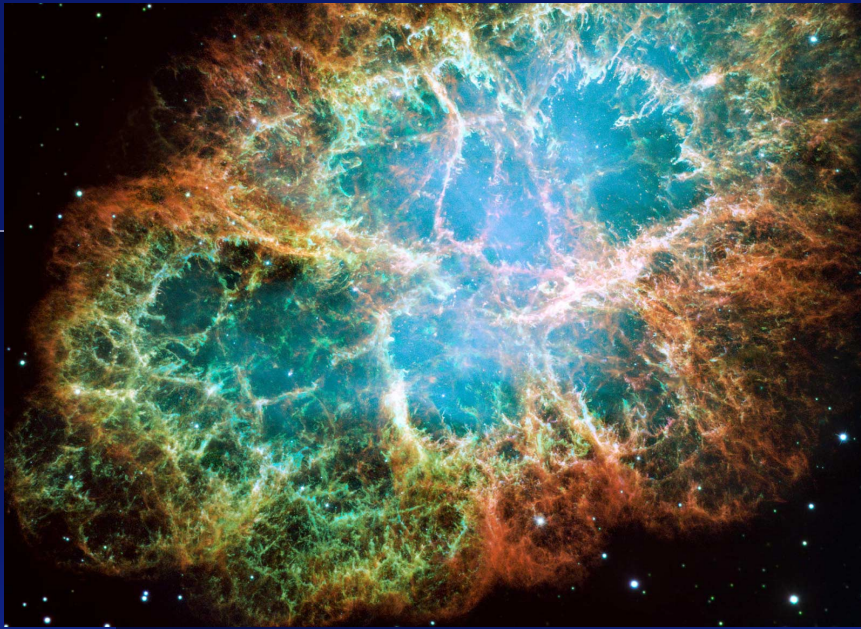


Alex Lazarian (Astronomy, Physics)

Special Thanks to G. Eyink, G. Kowal, E. Vishniac, H. Yan

Astrophysical fluids are generically turbulent

$$Re = LV/\nu = (L^2/\nu)/(L/V) = \tau_{diff}/\tau_{eddy}$$

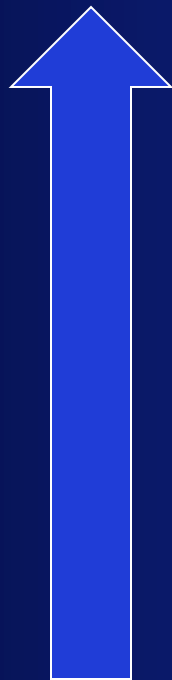


Astrophysical flows have $Re > 10^{10}$.

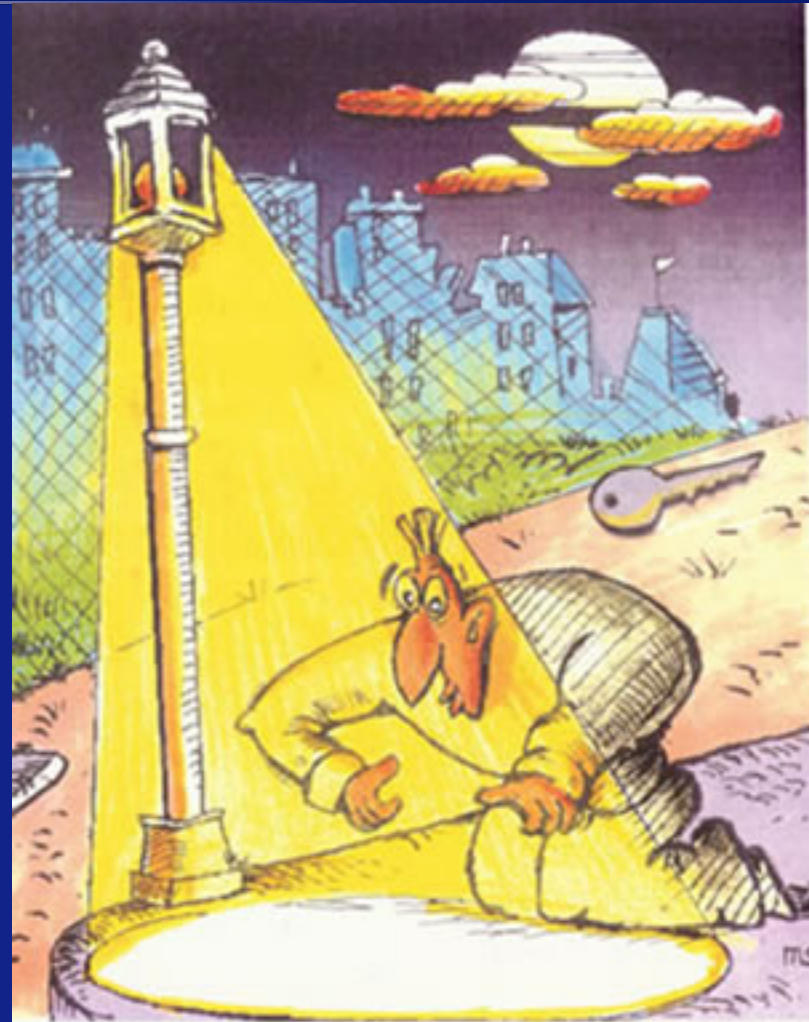


A lot of research is driven by what we can currently simulate, but simulating realistic turbulence is challenging/impossible

Real world

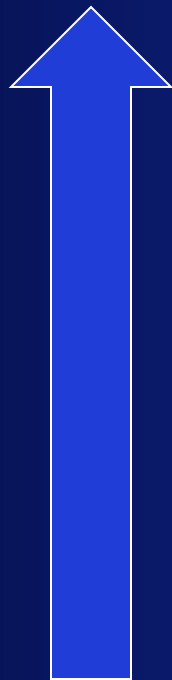


Numerical simulations



The studies extrapolate from low resolution numerical simulations to very different astrophysical regimes, while turbulence does require high resolution

Real world



Numerical simulations



Efforts scale as Re^4

Differences in Re can be more than 10^{10}

Main Points

Flux is not frozen in turbulent media

Turbulent magnetic fields present superdiffusion

MHD turbulence suppresses streaming instability, but does not create “streaming catastrophe”

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*Flux is
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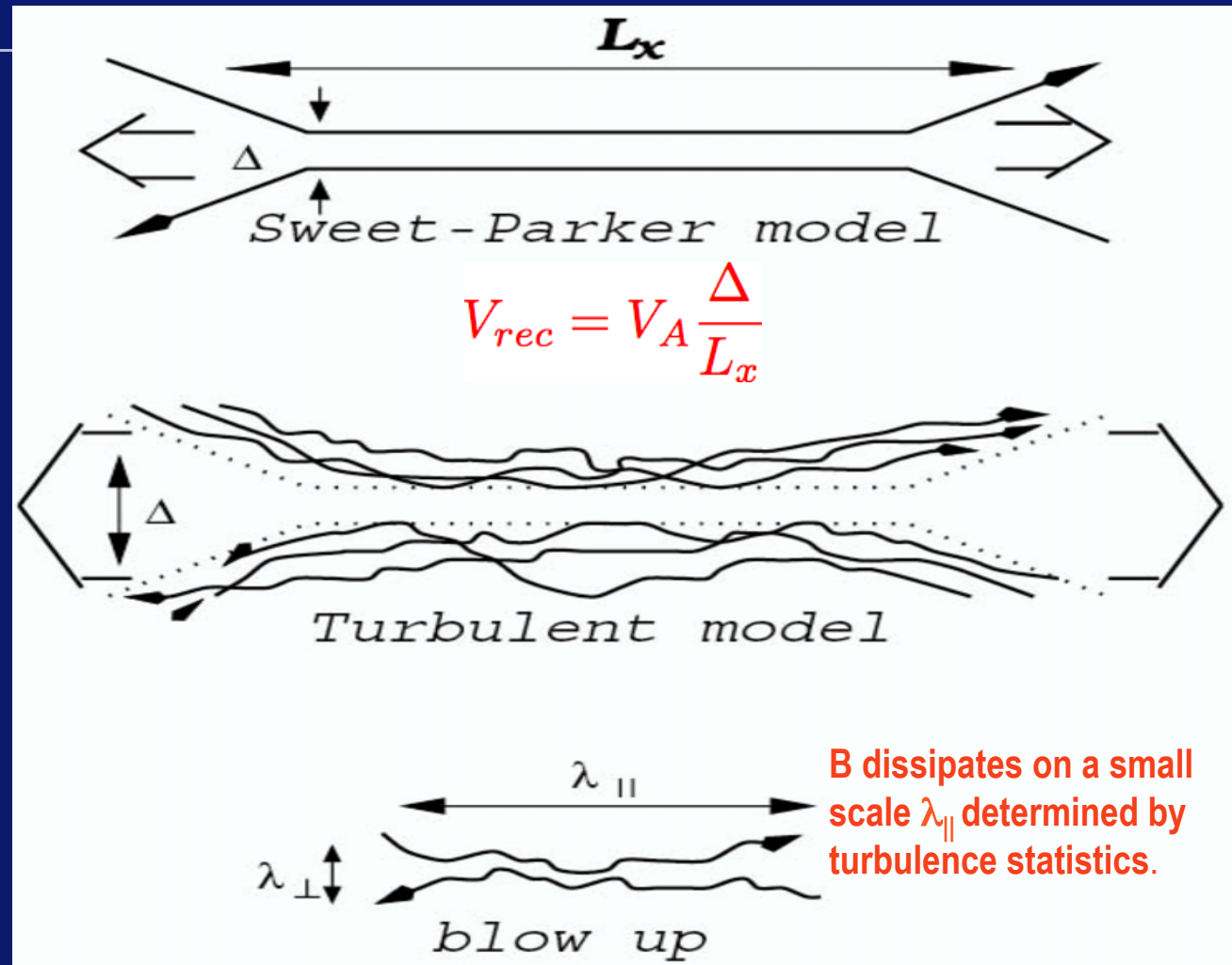
LV99 model extends Sweet-Parker model for turbulent astrophysical plasmas and makes reconnection fast

Turbulent reconnection:

Outflow is determined by field wandering.

Key element:

L/λ_{\parallel} reconnection simultaneous events



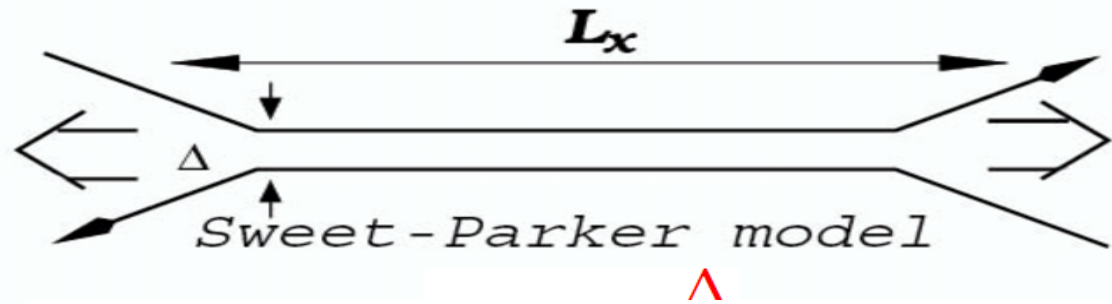
AL & Vishniac (1999)

henceforth referred to as LV99

LV99 model extends Sweet-Parker model for turbulent astrophysical plasmas and makes reconnection fast

Turbulent reconnection:

Outflow is determined by field wandering.



Without turbulence:

molecular diffusion coefficient $D \sim 10^{-5} \text{ cm}^2/\text{sec}$
(\leftarrow It's for small molecules in water.)

\rightarrow Mixing time $\sim (\text{size of the cup})^2/D \sim 10^7 \text{ sec} \sim 0.3 \text{ year} !$

L/λ_{\parallel} reconnection
simultaneous events

AL & Vishniac (1999)

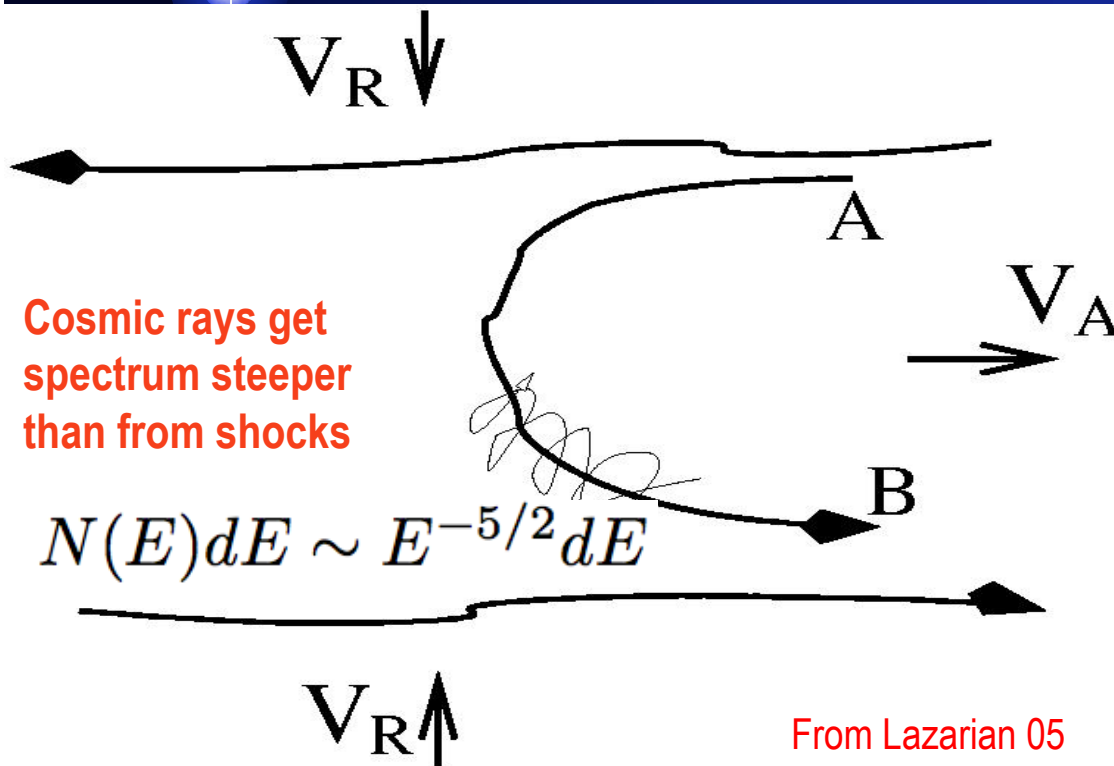
henceforth referred to as LV99

Example: Good correspondence between LV99 turbulent reconnection and solar wind data



Lalescu et al. 2015

In LV99 model energetic particles get accelerated by First Order Fermi mechanism (as discussed by Gianfranco)



(cp. Drake 2006).

Published in De Gouveia Dal Pino & Lazarian 2003

Applications to pulsars, microquasars, solar flare acceleration (De Gouveia Dal Pino & Lazarian 00, 03, 05, Lazarian 05).

Big Implication: LV99 predicts that magnetic field in *turbulent fluids* is not frozen in



Hannes Alfvén

Instead of flux freezing condition one should consider flux diffusion by turbulent flow. This has dramatic consequences for many areas of astrophysics including star formation!

Violation of magnetic field frozen in condition in turbulent fluids proven in Eyink (2011). The equivalence of this and LV99 approach was demonstrated in Eyink, Lazarian & Vishniac 2011.

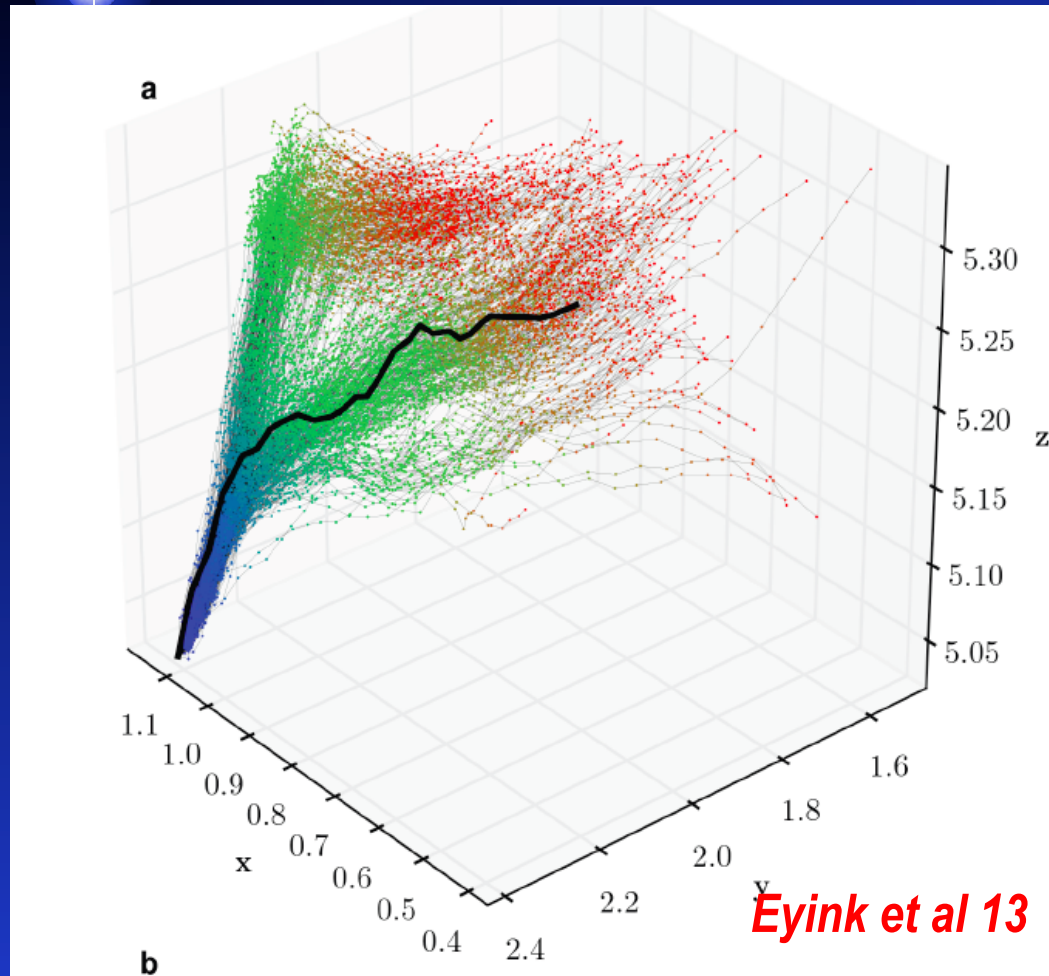
Eyink, AL & Vishniac 2011 related LV99 to the concept of Richardson diffusion



$$\langle |\mathbf{x}_1(t) - \mathbf{x}_2(t)|^2 \rangle \sim t^3$$

Richardson's law

Eyink, AL & Vishniac 2011 related LV99 to the well-known concept of Richardson diffusion

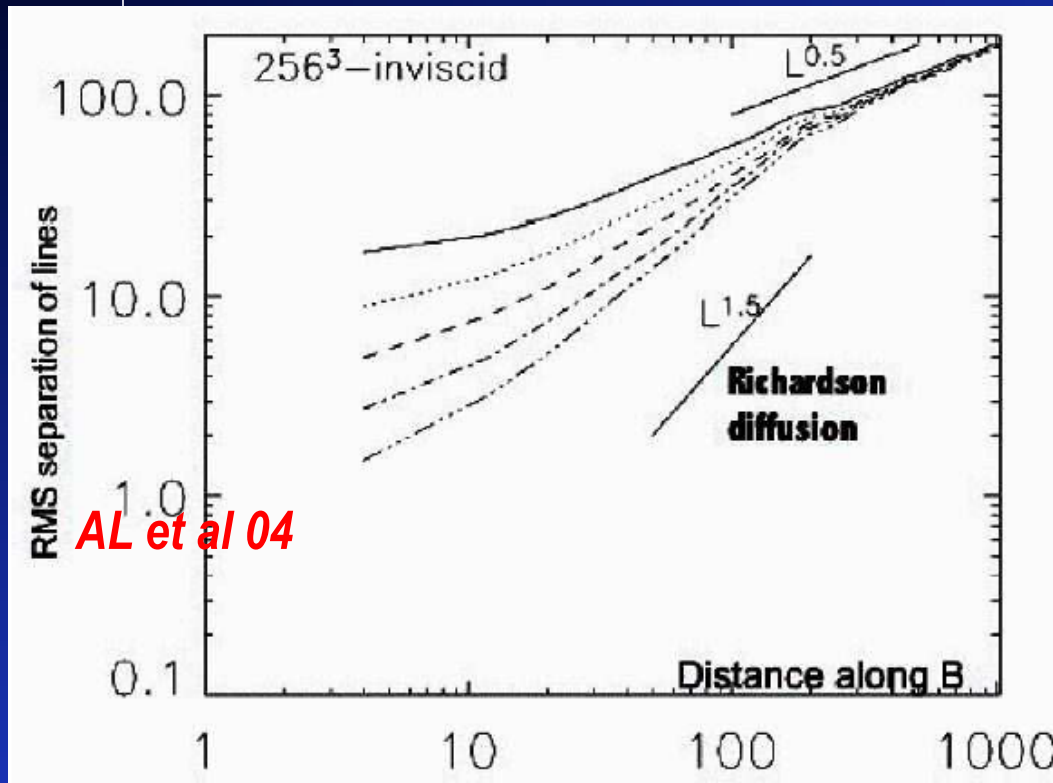


$$\langle |\mathbf{x}_1(t) - \mathbf{x}_2(t)|^2 \rangle \sim t^3$$

Magnetic diffusion in time

If one traces magnetic field lines in the presence of Richardson diffusion than one gets the LV99 result for field wandering

Richardson diffusion measured in MHD



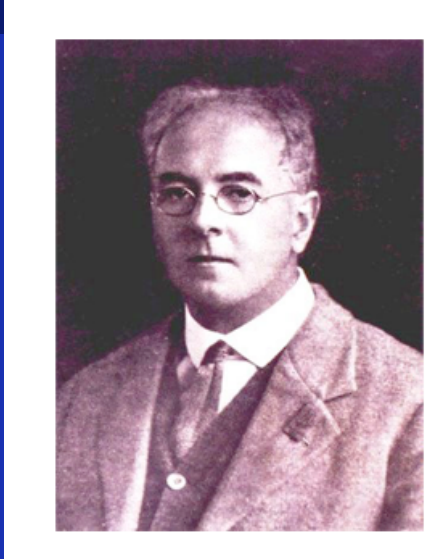
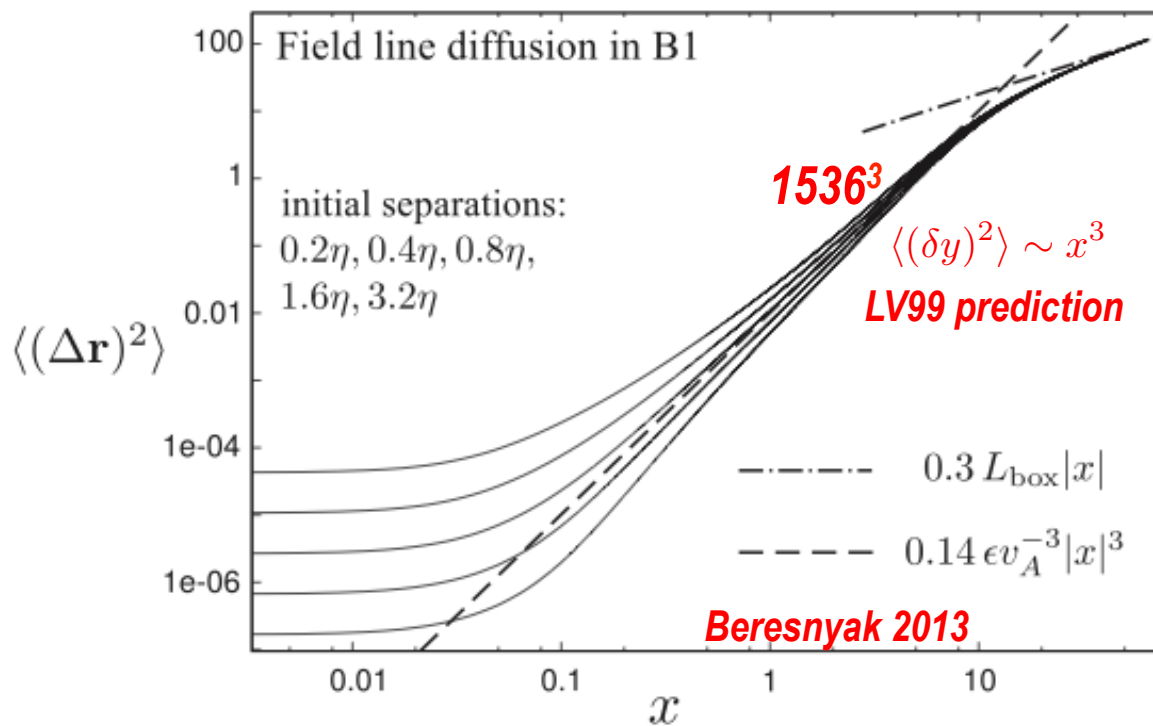
$$\langle (\delta y)^2 \rangle \sim x^3$$

We decided to keep the term
Richardson diffusion

Magnetic diffusion in space: field wandering

If one traces magnetic field lines in the presence of Richardson diffusion than one gets the LV99 result for field wandering

Richardson diffusion measured in MHD



$$\langle(\delta y)^2\rangle \sim x^3$$

We decided to keep the term Richardson diffusion

Magnetic diffusion in space: field wandering

Main Points

Flux is not frozen in turbulent media

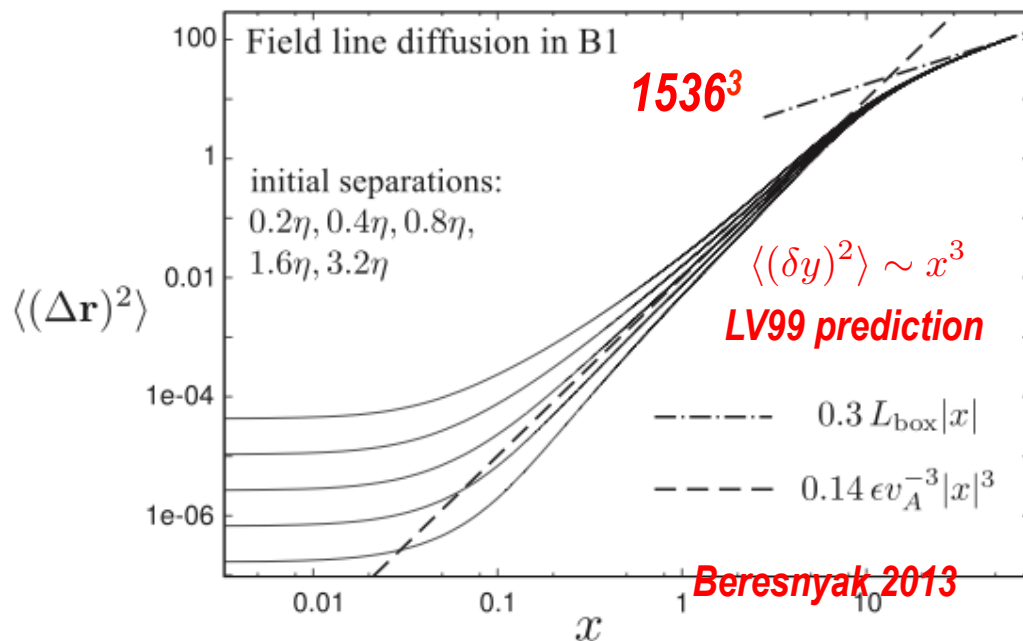
Turbulent magnetic fields present superdiffusion

MHD turbulence suppresses streaming instability, but does not create “streaming catastrophe”

Richardson diffusion in space means superdiffusion (superballistic behavior) for CRs following magnetic field

$$\langle (\delta y)^2 \rangle \sim x^3$$

Superdiffusion acts on scales x less than the injection scale of MHD turbulence



Injection scale of turbulence in the Galaxy is about 100 pc

Diffusion perpendicular to mean magnetic field direction is determined by magnetic field line wandering

Realized by Jokipii & Parker 69, Jokipii 73 but turbulence model was not right

In fact, this motivated my work in turbulent magnetic reconnection

The study with modern understanding of MHD turbulence is in AL& Vishniac 99

Strong subAlfvenic turbulence at scales $s < l_{\text{trans}}$ results in superdiffusion:

$$\frac{d}{ds} \ell_{\perp}^2 \sim D_{\perp}^B(\ell) \sim (\delta u_{\ell}/v_A)^2 \ell_{\parallel} \sim L \left(\frac{\ell_{\perp}}{L} \right)^{4/3} M_A^{4/3},$$

$$\ell_{\perp}^2 \sim \frac{s^3}{27L} M_A^4,$$

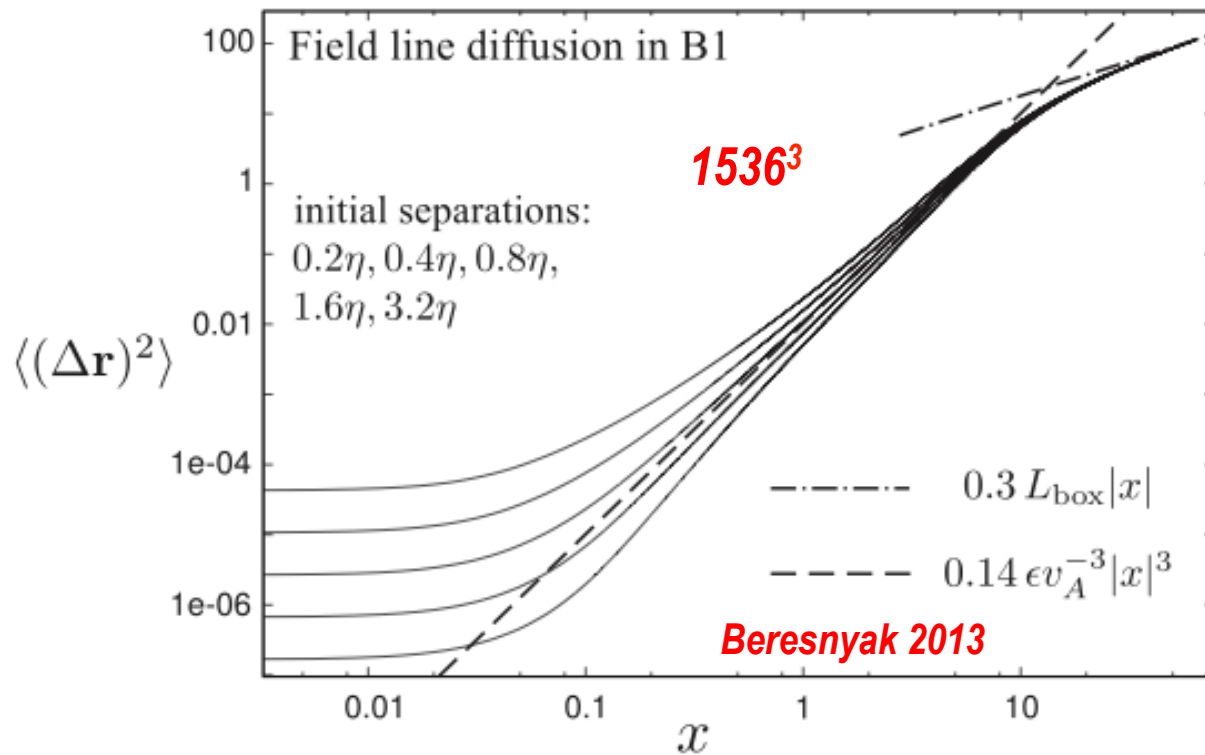
At scales $s > l_{\text{trans}}$ results in ordinary diffusion:

$$\ell_{\perp}^2 \sim sLM_A^4.$$

Diffusion is not applicable on scales less than the turbulence injection scale L

On scales $\leq L$ and $s \leq \text{mfp}$, CRs trace magnetic field divergence

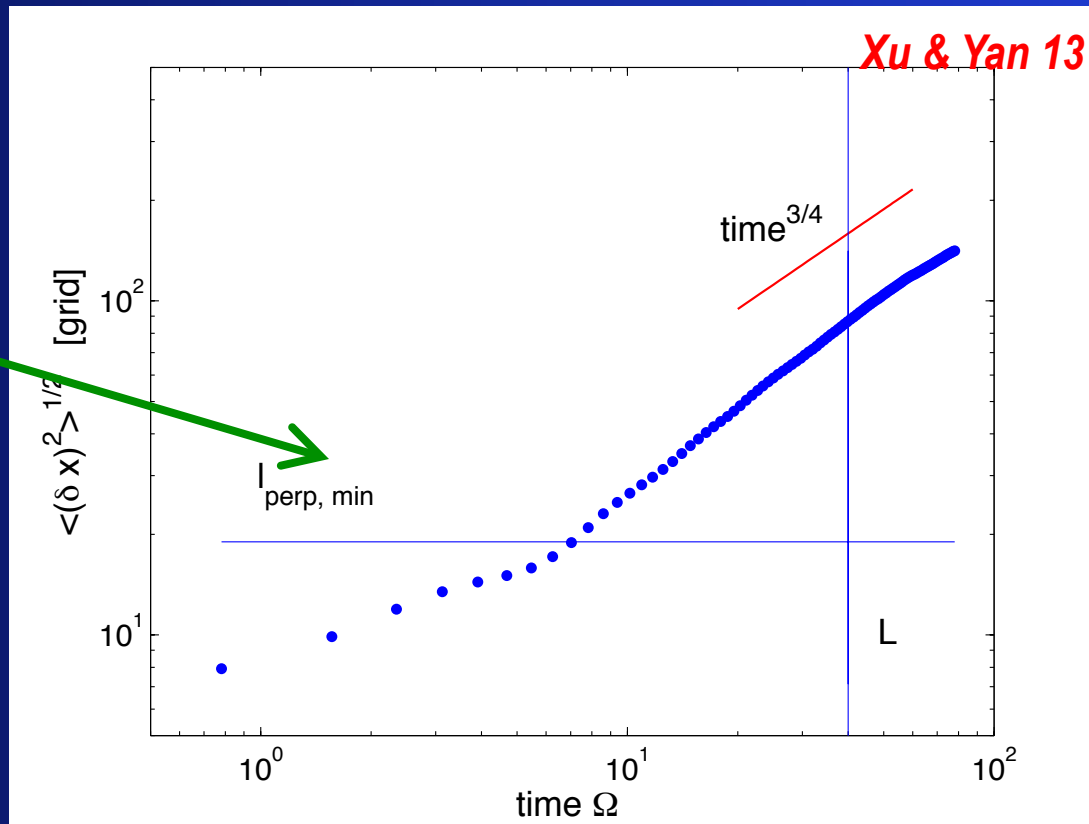
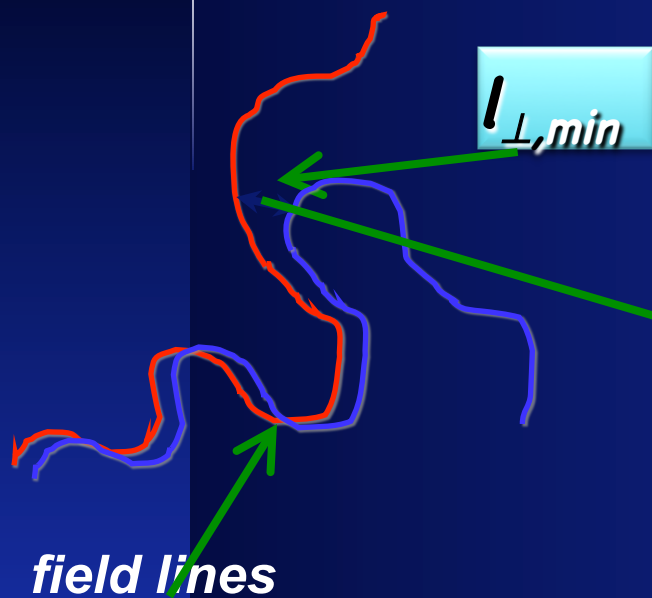
$$\ell_{\perp}^2 \sim \frac{s^3}{27L} M_A^4,$$



For scales $\leq L$

Explosive separation of magnetic field lines is described analytically in AL & Vishniac 1999. Separation $\sim X^{3/2}$

If CRs diffuse along the magnetic field lines the perpendicular motion is still superdiffusive



To compare with

$$l_{\perp, \text{CR}}^2 \sim \frac{(D_{\parallel} \delta t)^{3/2}}{27L} M_A^4, \quad M_A < 1,$$

The dependence on forth power of Alfven Mach number is also confirmed

On scales $s > L$ and $s \gg mfp$ the ordinary diffusion is present (AL06, Yan & AL08)

$$D_{\perp, \text{global}} \approx D_{\parallel} M_A^4,$$

On scales $< L$ and $s < mfp$, CRs trace magnetic field divergence

$$l_{\perp}^2 \sim \frac{s^3}{27L} M_A^4,$$

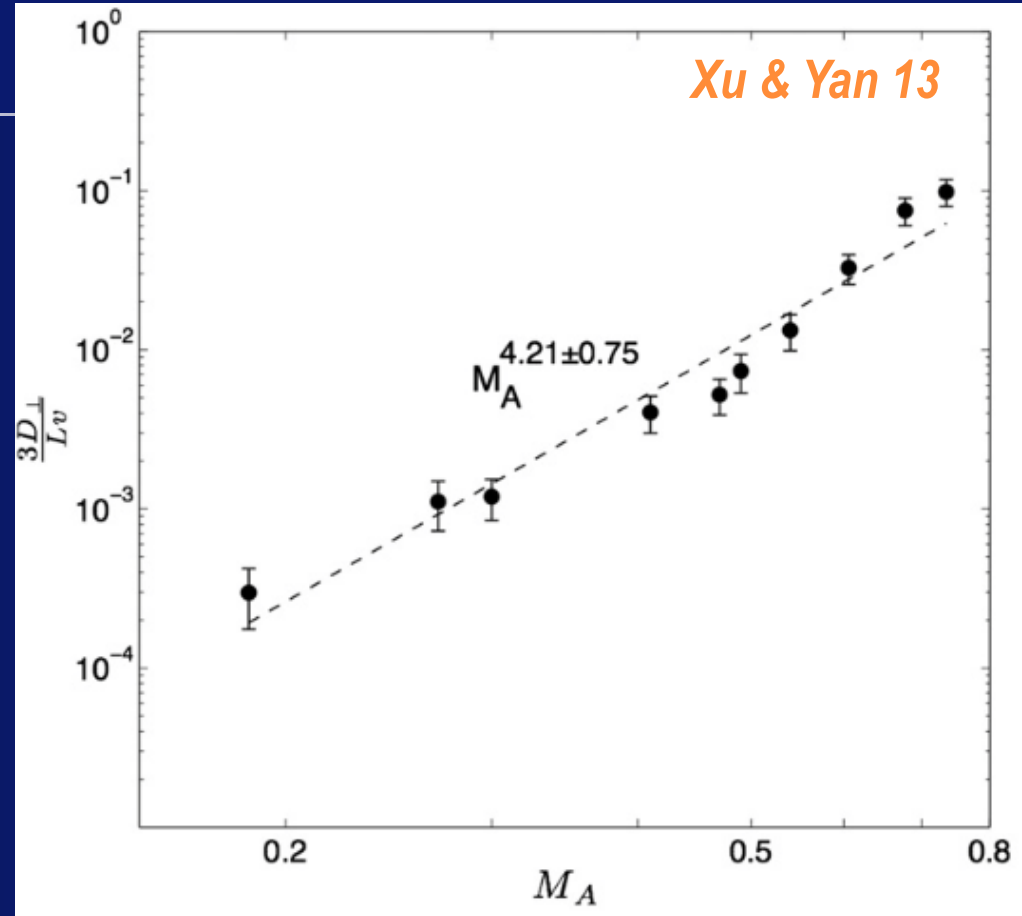
On scales $< L$ and $s > mfp$, CRs trace magnetic field divergence, s is covered in diffusion process

$$l_{\perp, \text{CR}}^2 \sim \frac{(D_{\parallel} \delta t)^{3/2}}{27L} M_A^4, \quad M_A < 1,$$

Differs from the textbook (see Jokipii & Parker 69) M_A^2 dependence

The dependence on forth power of Alfvén Mach number is also confirmed

Diffusive regime



To compare with

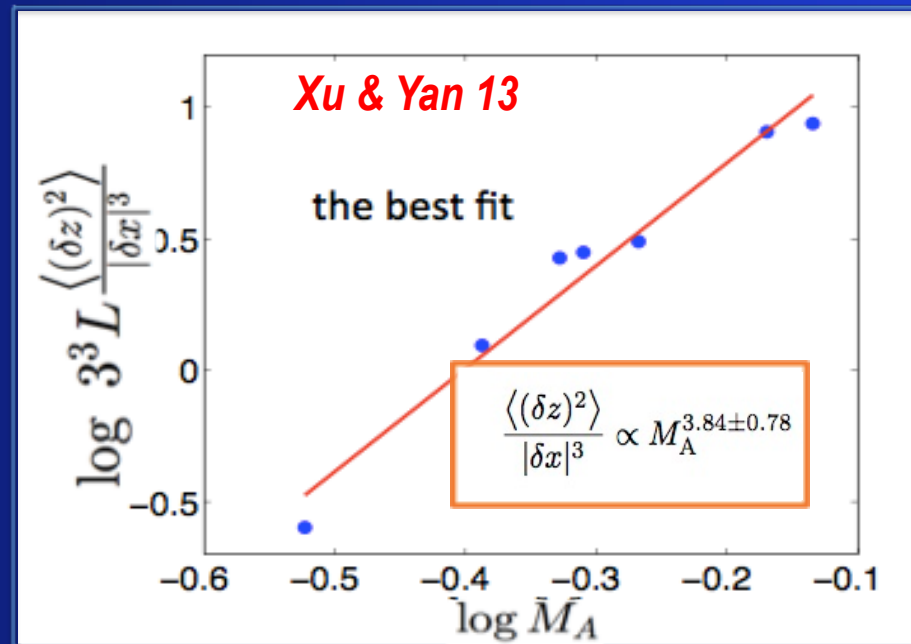
$$D_{\perp, \text{global}} \approx D_{\parallel} M_A^4,$$

in AL06, Yan & AL08

The dependence on forth power of Alfvén Mach number is also confirmed

Superdiffusive regime

$$\langle (\delta z)^2 \rangle = \frac{|\delta x|^3}{3^3 L} M_A^4$$



To compare with

$$\ell_{\perp}^2 \sim \frac{s^3}{27L} M_A^4,$$

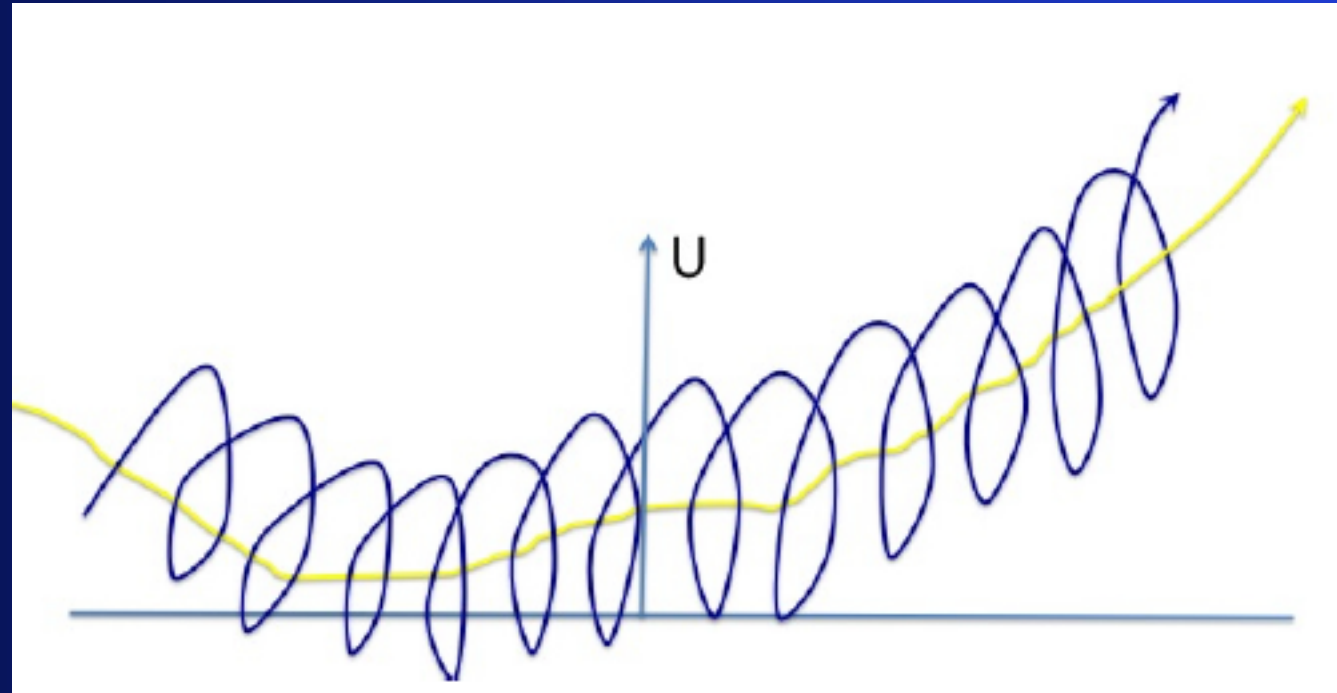
AL & Vishniac 1999;
Yan & AL 2008

Superdiffusion changes the accepted formalism for parallel and perpendicular shock acceleration

$$\frac{\kappa_{\perp}}{\kappa_{\parallel}} = \frac{1}{1 + (\lambda_{CR}/r_L)^2}$$

Accepted expression

In reality



Lazarian & Yan 2013

Main Points

Flux is not frozen in turbulent media

Turbulent magnetic fields present superdiffusion

MHD turbulence suppresses streaming instability, but does not create “streaming catastrophe”

Streaming instability damping by Alfvénic turbulence is suggested by Yan & AL 02 and quantified by Farmer & Goldreich 04

$$\Gamma_{cr} \approx \Omega_B \frac{n_{cr}(> \gamma)}{n_i} \left(\frac{v_{stream}}{V_A} - 1 \right),$$

Streaming instability growth rate

FG04 considered only damping by strong turbulence and assumed that the energy is injected with Alfvén velocity. They also came to the conclusion that streaming instability should be suppressed in the Galaxy, which presented a problem with explaining cosmic ray isotropy measured.

The result of the pioneering FG04 study were overinterpreted in later works:

- 1. That this is the model of damping of Alfvén waves in turbulence in general sense.*
- 2. That it covers both subAlfvénic and superAlfvénic damping of streaming.*
- 3. That indeed “streaming catastrophe” exists for realistic Mikly Way model.*

Results in FG04, however, do not cover all important regimes of turbulence

$$M_A \equiv \frac{V_L}{V_A} = \frac{\delta B}{B}$$

AL16

Type of MHD turbulence	Injection velocity	Range of scales	Spectrum $E(k)$	Instability damping rate and r_L range
Weak	$V_L < V_A$	$[l_{trans}, L]$	k_{\perp}^{-2}	$\frac{V_A M_A^{8/3}}{r_L^{2/3} L^{1/3}}, \quad LM_A^4 < r_L < LM_A$
Strong subAlfvenic	$V_L < V_A$	$[l_{min}, l_{trans}]$	$k_{\perp}^{-5/3}$	$\frac{V_A M_A^2}{r_L^{1/2} L^{1/2}}, \quad \frac{l_{min}^{4/3}}{L^{1/3}} < r_L < LM_A^4$
Hydro-like superAlfvenic	$V_L > V_A$	$[l_A, L]$	$k^{-5/3}$	$\frac{V_A M_A}{r_L^{2/3} L^{1/3}}, \quad l_A < r_L < L$
Strong superAlfvenic	$V_L > V_A$	$[l_{min}, l_A]$	$k_{\perp}^{-5/3}$	$\frac{V_A M_A^{3/2}}{r_L^{1/2} L^{1/2}}, \quad \frac{l_{min}^{4/3}}{L^{1/3}} M_A < r_L < l_A$

For subAlfvenic turbulence the range for the strong turbulence damping is limited to $< LM_A^4$, It extends from LM_A^4 to LM_A for weak turbulence

As for the suppression of streaming in the Galaxy FG04 uses a number of assumptions

1. The level of turbulent dissipation is estimated assuming that all the heating of halo gas is due to turbulent damping

This is not true as in the presence of streaming additional heating comes from CRs.

2. Streaming instability damping is induced by strong turbulence in the halo

Even for the turbulent dissipation rate that FG04 assumes (excessive) the damping of instability for high energy CRs is due not to strong turbulence, but to weak turbulence and therefore is reduced. Indeed $r_L < LM_A^4$ and their assumption implies $M_A = 0.2$

3. A simple relation between the streaming velocity and degree anisotropy is assumed

This ignores superdiffusion of magnetic field lines that we discussed.

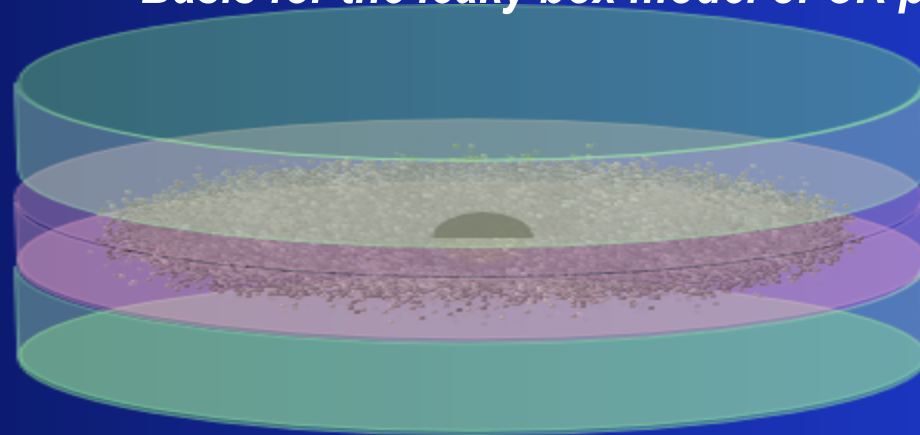
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$$\Gamma_{cr} \approx \Omega_B \frac{n_{cr}(> \gamma)}{n_i} \left(\frac{v_{stream}}{V_A} - 1 \right),$$

Streaming instability

Basis for the leaky box model of CR propagation

Disk



Halo

Quantified in Farmer & Goldreich 04 for strong transAlfvénic turbulence

Farmer & Goldreich 2004 challenged the "leaky box" for CR confinement and isotropization claiming that streaming instability cannot exist in the presence of turbulence

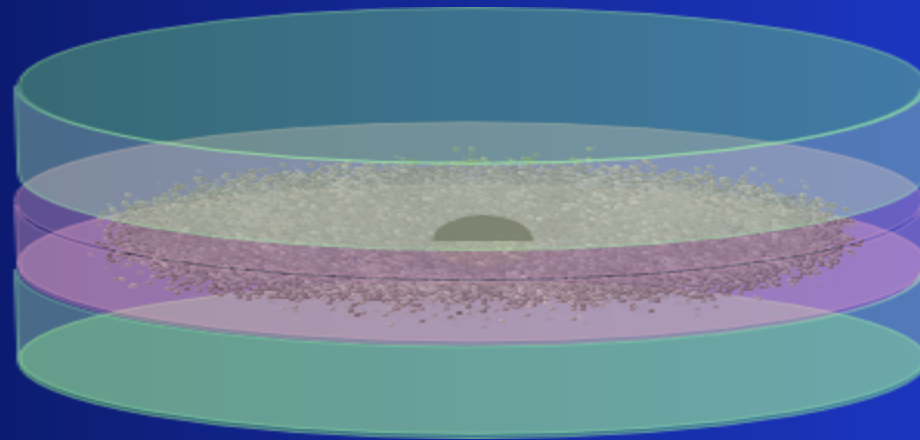
The model of damping by strong turbulence is used

The turbulence level was estimated using

$$\epsilon_{\text{turb.dissipation}} = \text{radiation cooling}$$

Problem: cosmic rays stream and do not isotropize

Disk



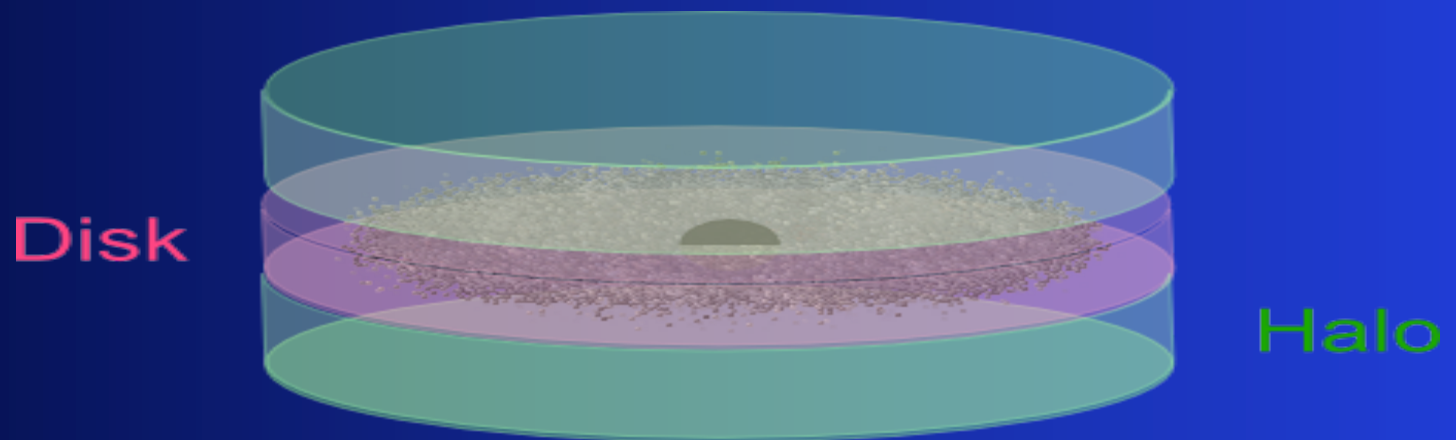
Halo

The work meant the crisis of the existing models of CR isotropization!

Detailed calculations in AL16 show that "leaky box" model is valid if it is accounted that scattering is by weak turbulence and the level of turbulence in Halo is small

The model of damping by weak turbulence is used

The turbulence level was estimated using $\epsilon_{\text{turb.dissipation}} < \text{radiation cooling}$



CRs stream in the disk where turbulence is transAlfvenic and randomize by streaming instability in the halo. Streaming CR and not turbulence dissipation is the source of halo healing.

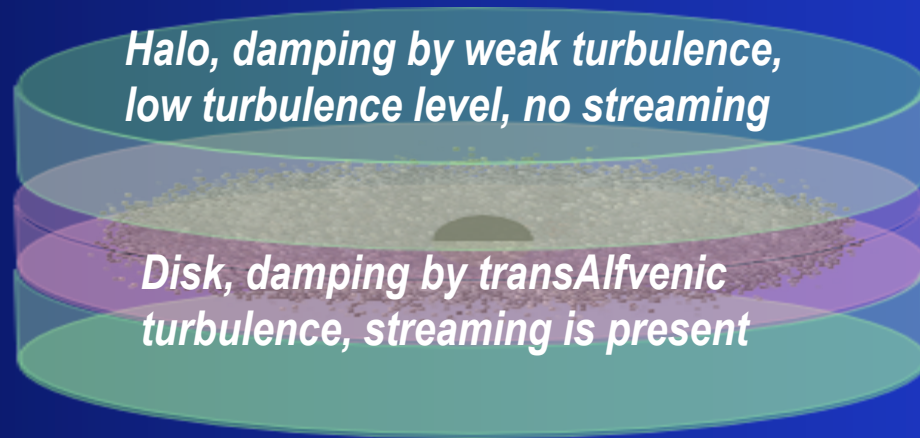
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New understanding:

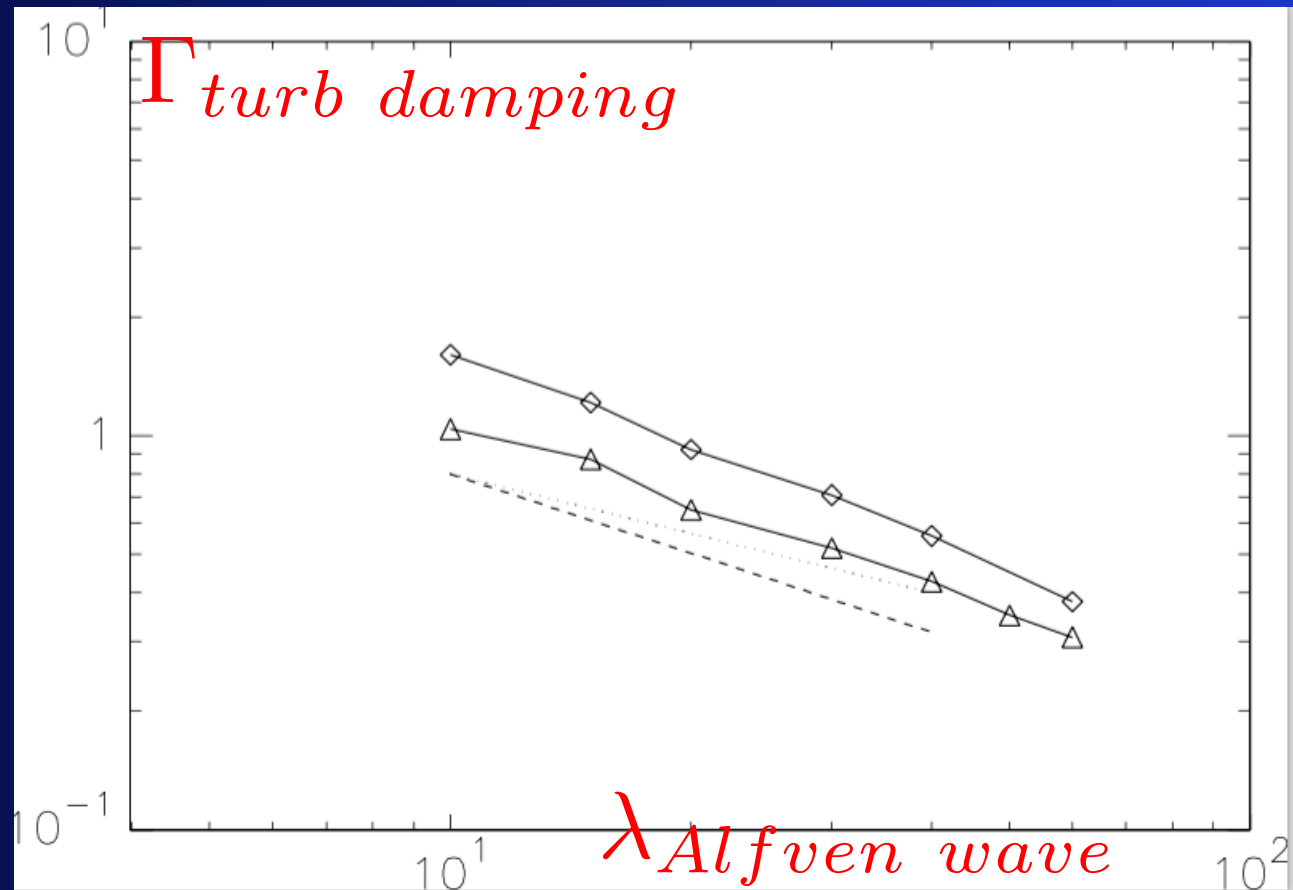
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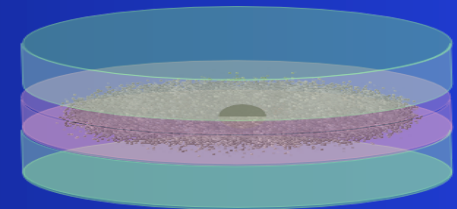
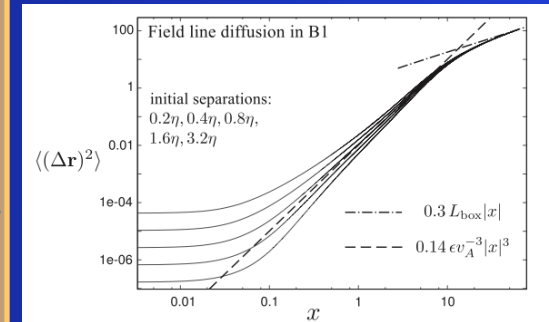
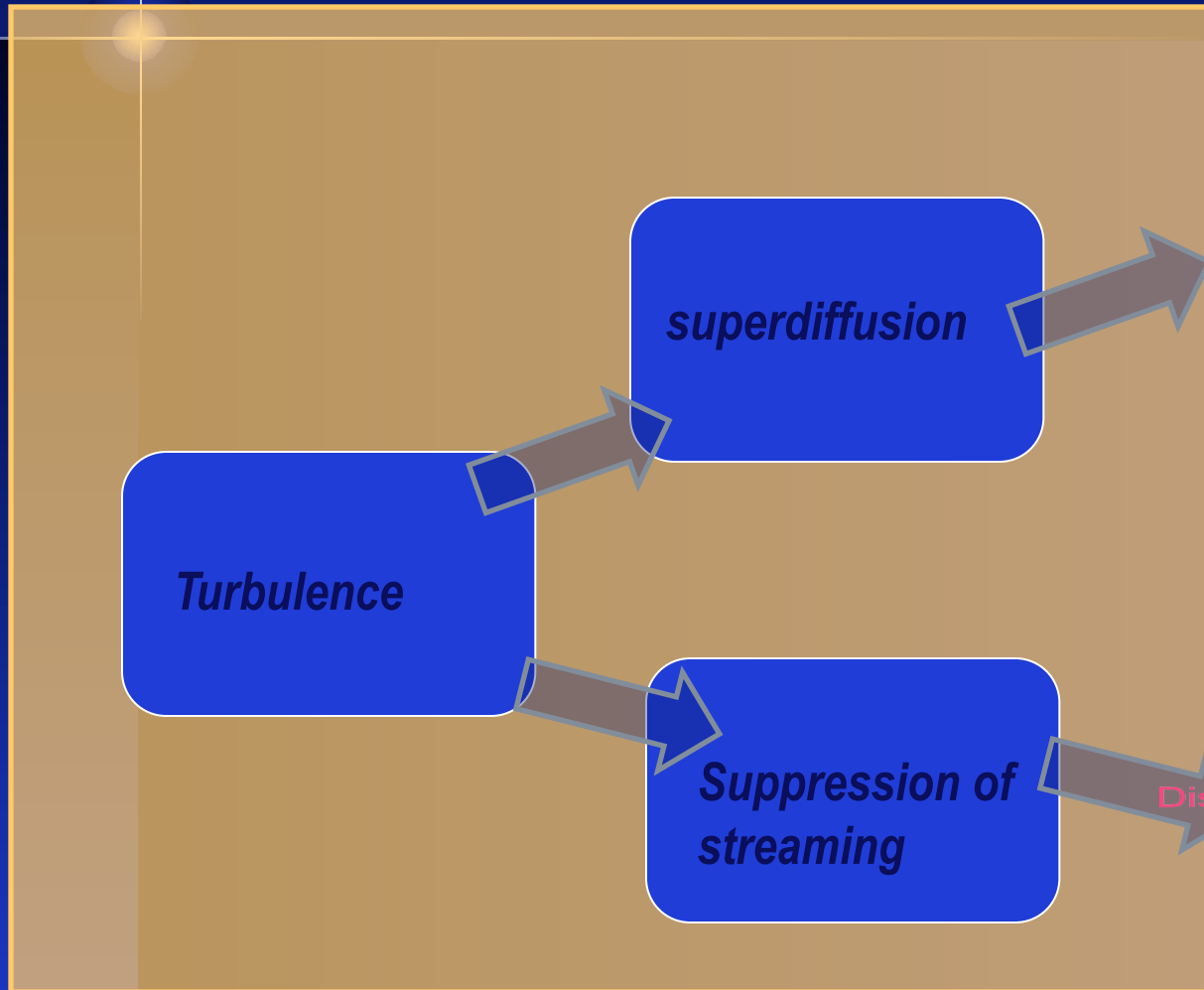
Halo

CRs stream in the disk where turbulence is transAlfvenic and randomize by streaming instability in the halo. Streaming CR and not turbulence dissipation is the source of halo healing.

Numerical simulations confirm the predicted scaling of the Alfvén wave damping



Turbulence induces superdiffusion and damping of Alfvén waves



Hal...