

Re-acceleration Model for Radio Relics in Galaxy Clusters

+ turbulent acceleration in the postshock region

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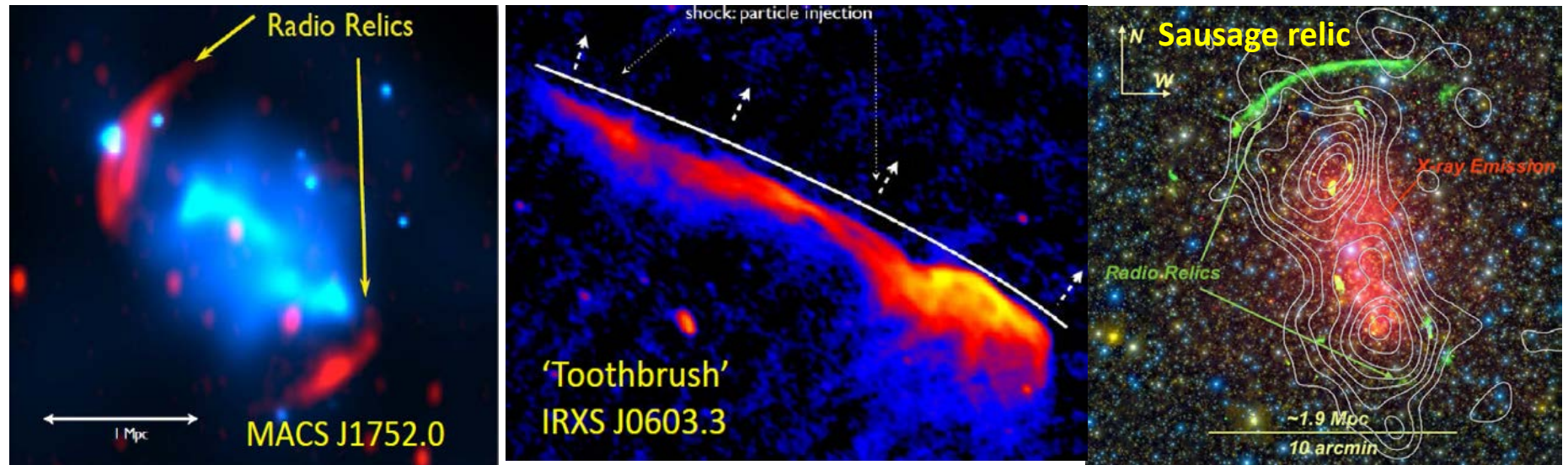
Dongsu Ryu

UNIST (Ulsan National Institute of Science & Technology), Korea

T. W. Jones

University of Minnesota

Radio relics: diffuse radio sources found mainly in merging clusters

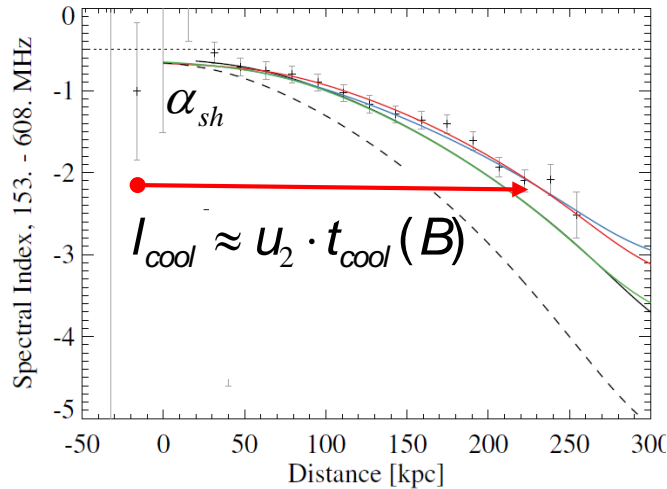


- elongated morphology
- spectral aging behind the shock (due to radiative cooling)
- power-law like integrated radio spectrum
- high polarization up to 50 % (B field compression)

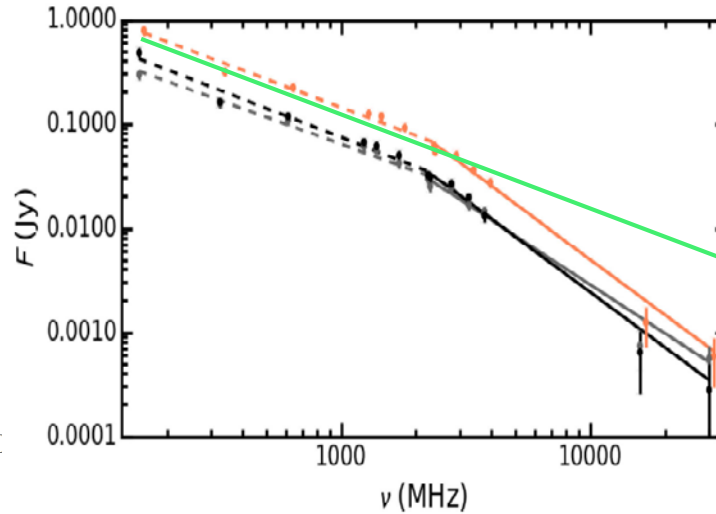
→ synchrotron radiation emitted by \sim GeV electrons accelerated at structure formation shocks

Radio vs. X-ray Observations of Radio Relics

Spectral index profile



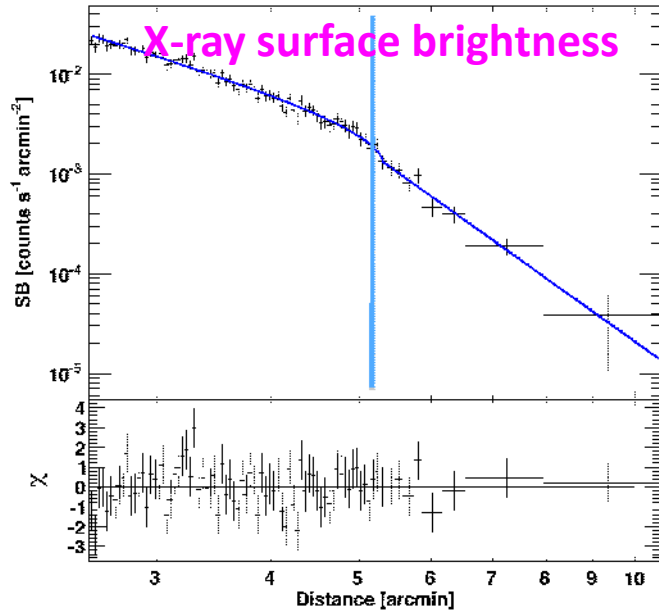
Integrated radio spectrum



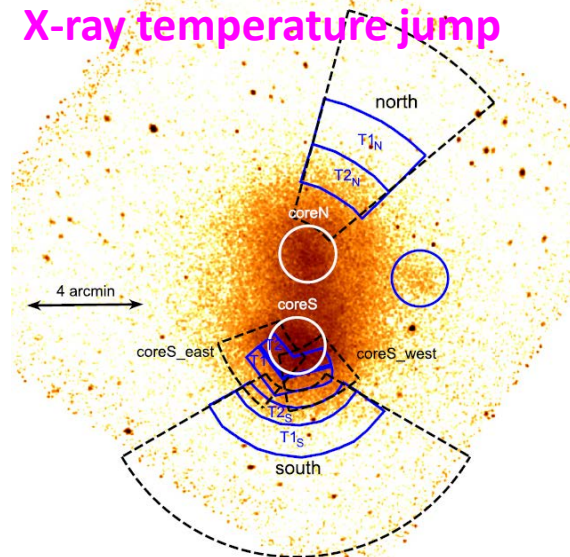
$$M_{radio}^2 = \frac{(3 + 2\alpha_{sh})}{(2\alpha_{sh} - 1)}$$

in situ injection model

$$\alpha_{integ} \approx \alpha_{sh} + 0.5$$



X-ray temperature jump

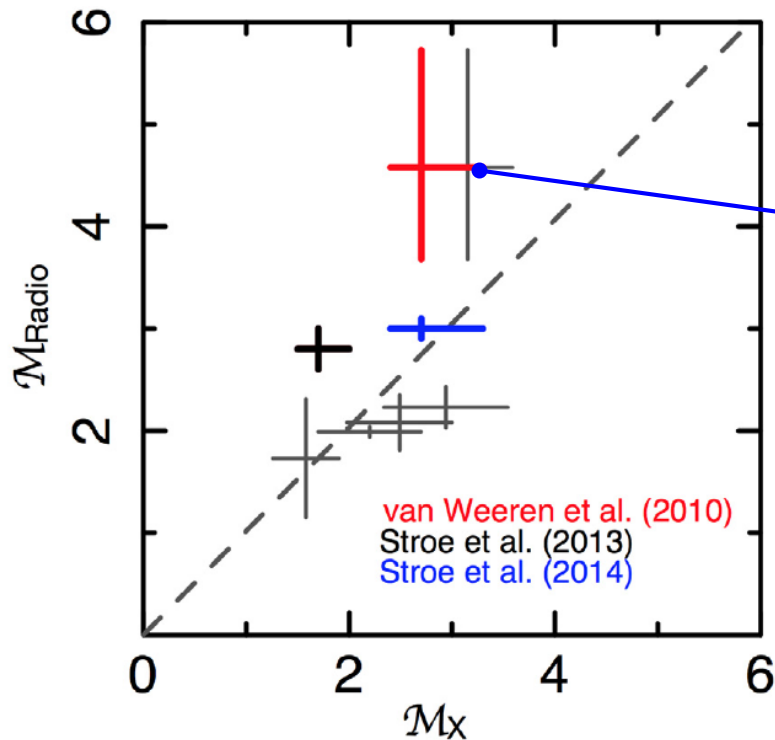


$$\frac{T_2}{T_1} = \frac{(M_X^2 + 3)(5M_X^2 - 1)}{16M_X^2}$$

$c_{s,1}$ = sound speed
 V_s = shock speed

* Projection effects

Some puzzles in DSA model with *in situ* injection only



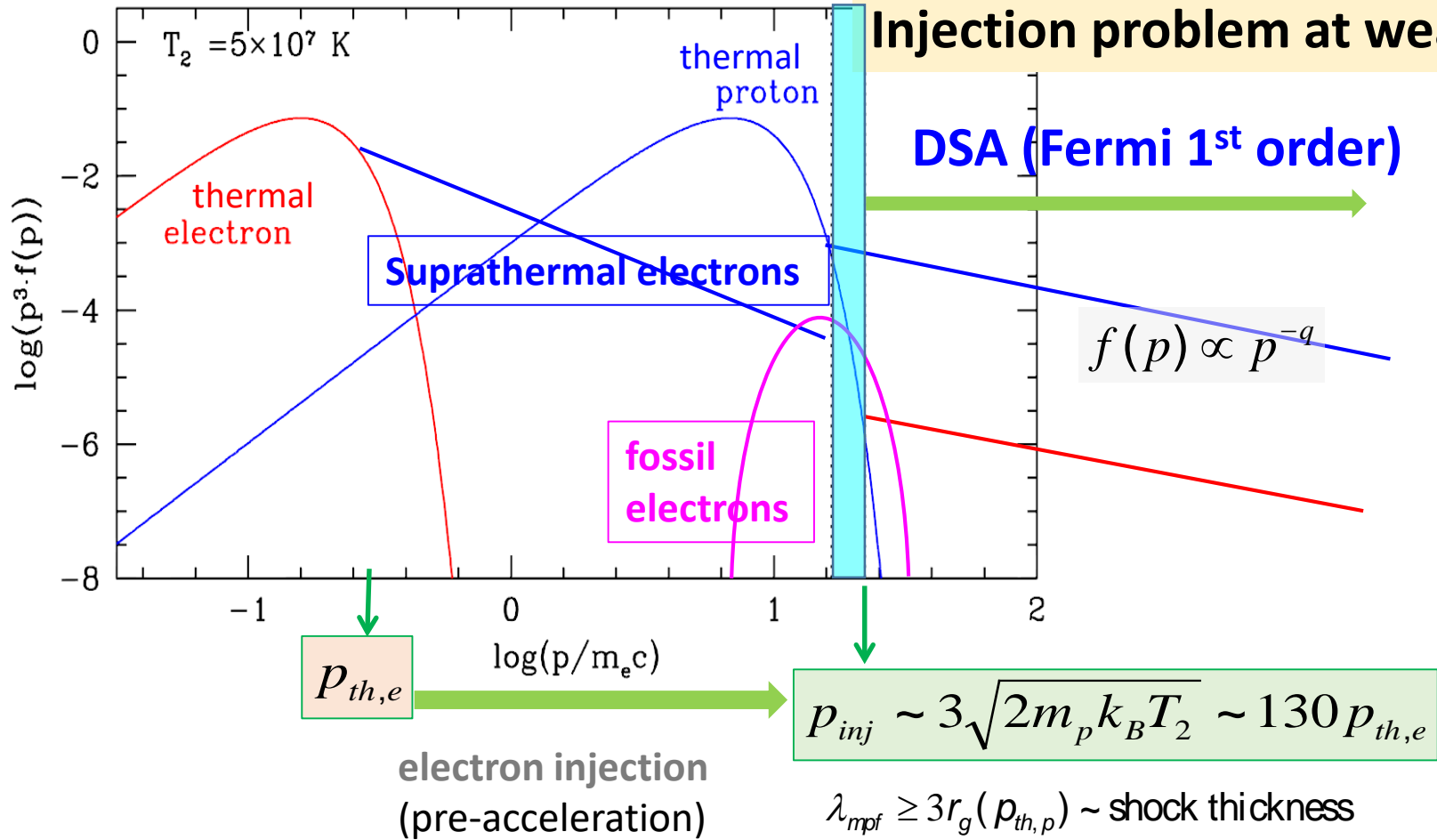
(1) For some radio relics, $M_{radio} > M_X$

Sausage relic: $M_{radio} \approx 4.6$, $M_X \approx 3$

(2) Low injection/acceleration efficiency expected for weak shocks that form in ICM ($M < 3$)

Akamatsu et al. 2015

Injection problem at weak shocks



Electrons need to be pre-accelerated to p_{inj} in order to get injected into DSA process

In-situ injection at the shock from **thermal electrons** is **very inefficient at weak shocks**.

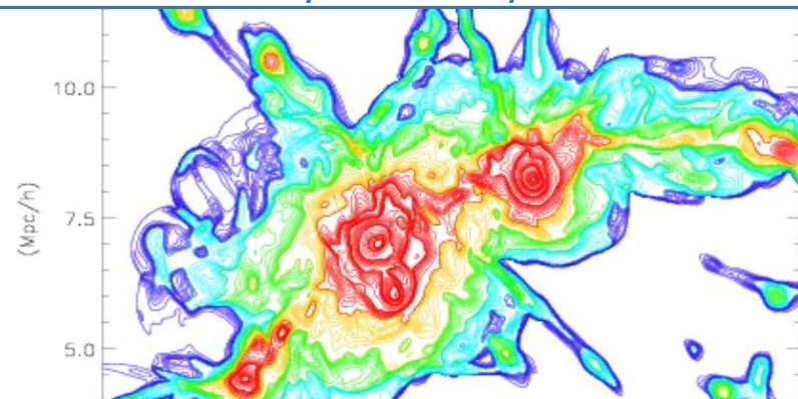
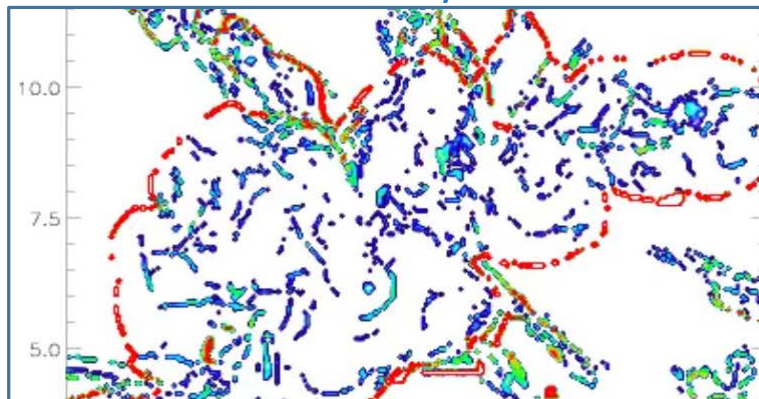
If there are **suprathermal/fossil** electrons in the preshock region
 → **enhance injection into DSA**

Ubiquitous presence of structure formation shocks in ICM

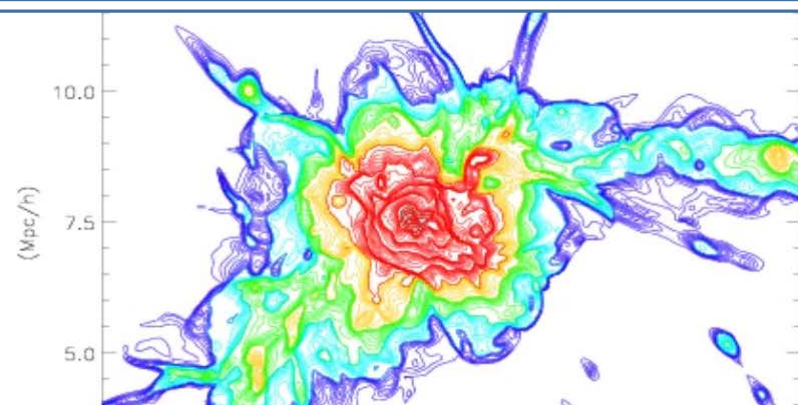
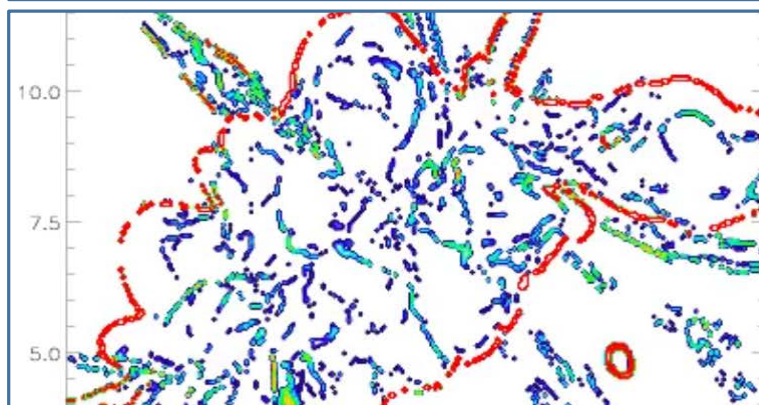
Shock distribution by Mach number

X-ray emissivity

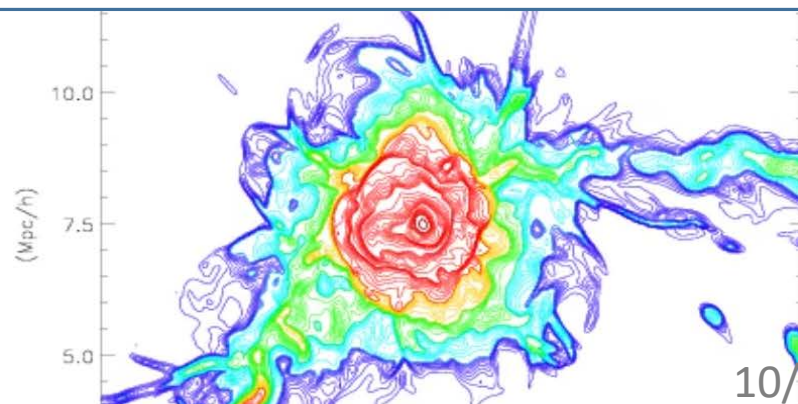
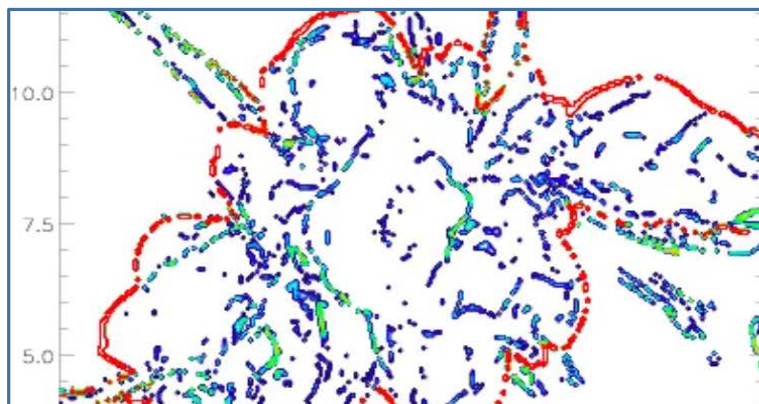
Z=0.65



Z=0.30

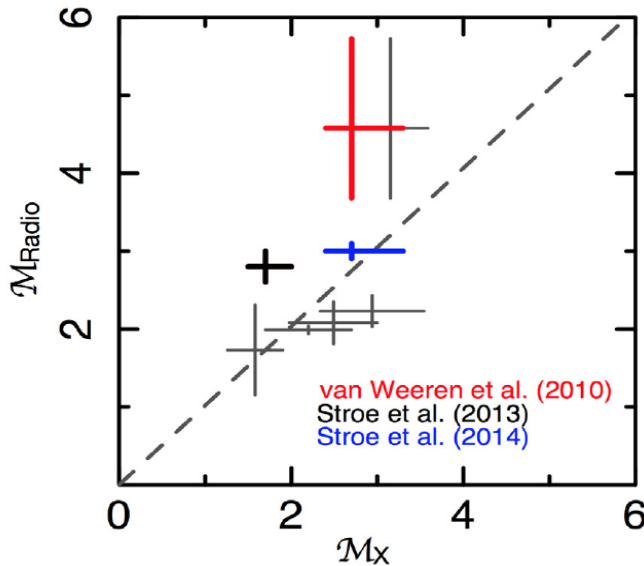


Z=0.22



Ha, Ryu,
& Kang
2016

Some puzzles with *in situ* injection model



(1) For some radio relics, $M_{\text{radio}} > M_X$

(2) Low acceleration efficiency of weak shocks
($M < 3$)

(3) Only ~10 % of merging clusters host radio relics,
while numerous shocks are expected to form
in ICM.

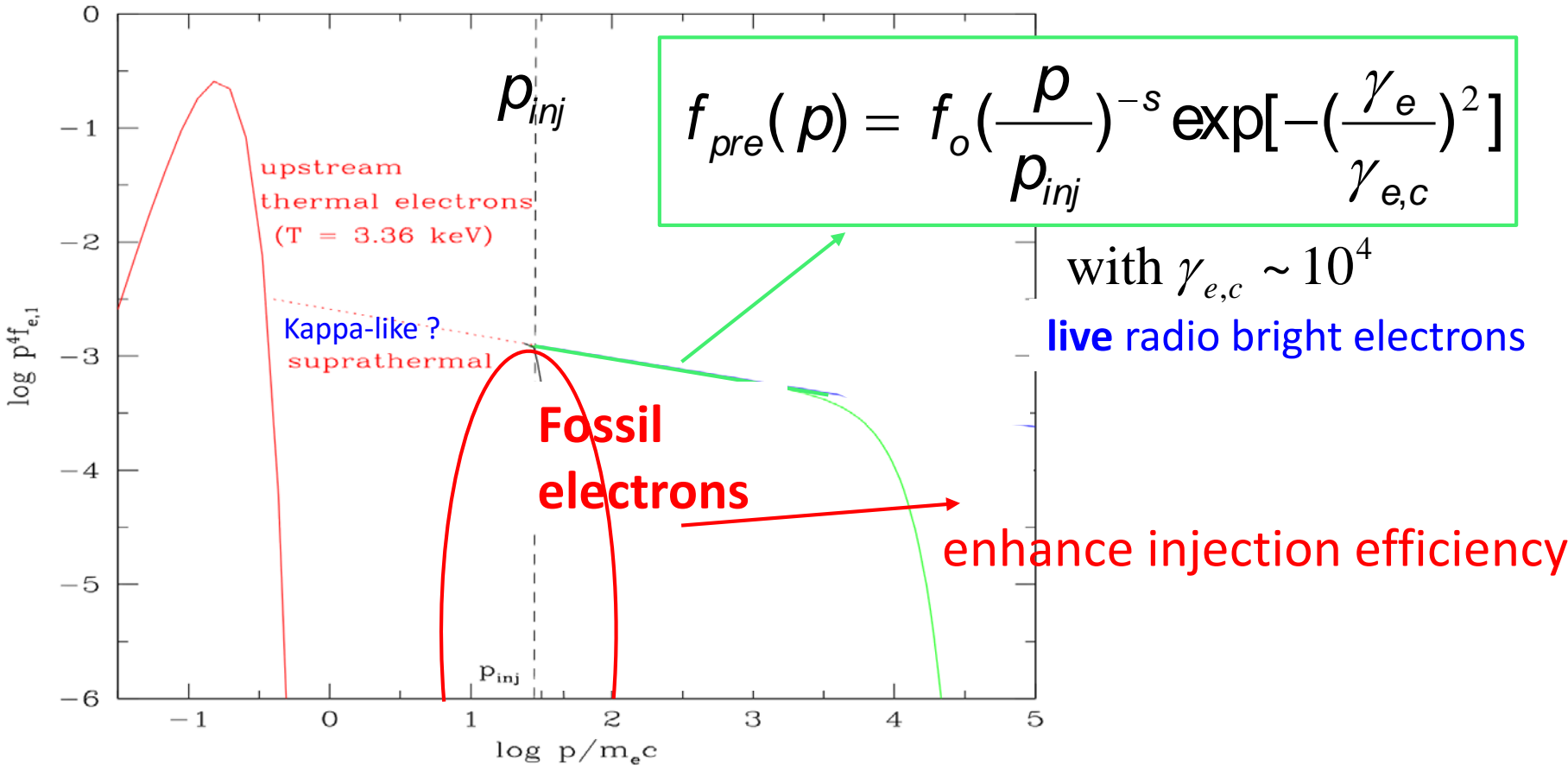
(4) Some X-ray shocks without associated radio relics

(3) & (4) → Not all shock can accelerate electrons and become radio relics.

Reacceleration model can solve these puzzles:

a radio relic forms when a weak shock encounters the ICM plasma with **pre-existing** (live/fossil) electrons.

Pre-existing electrons: fossil or live (radio-bright)



radio spectral index α is determined by M_s , s , & $\gamma_{e,c}$

$\Rightarrow M_{radio} \neq M_s$ or M_x (not necessarily)

Diffusion-Convection Equation for electron distribution function

$$\frac{\partial f}{\partial t} + u \frac{\partial f}{\partial x} = \frac{p}{3} \frac{\partial u}{\partial x} \frac{\partial f}{\partial p} + \frac{\partial}{\partial x} \left(\kappa \frac{\partial f}{\partial x} \right) + \frac{1}{p^2} \frac{\partial}{\partial p} \left(p^2 D_{pp} \frac{\partial f}{\partial p} \right) + \frac{1}{p^2} \frac{\partial}{\partial p} (p^2 \dot{p} f)$$

spatial diffusion=
Fermi 1st order acc.

Momentum diffusion =
stochastic acceleration

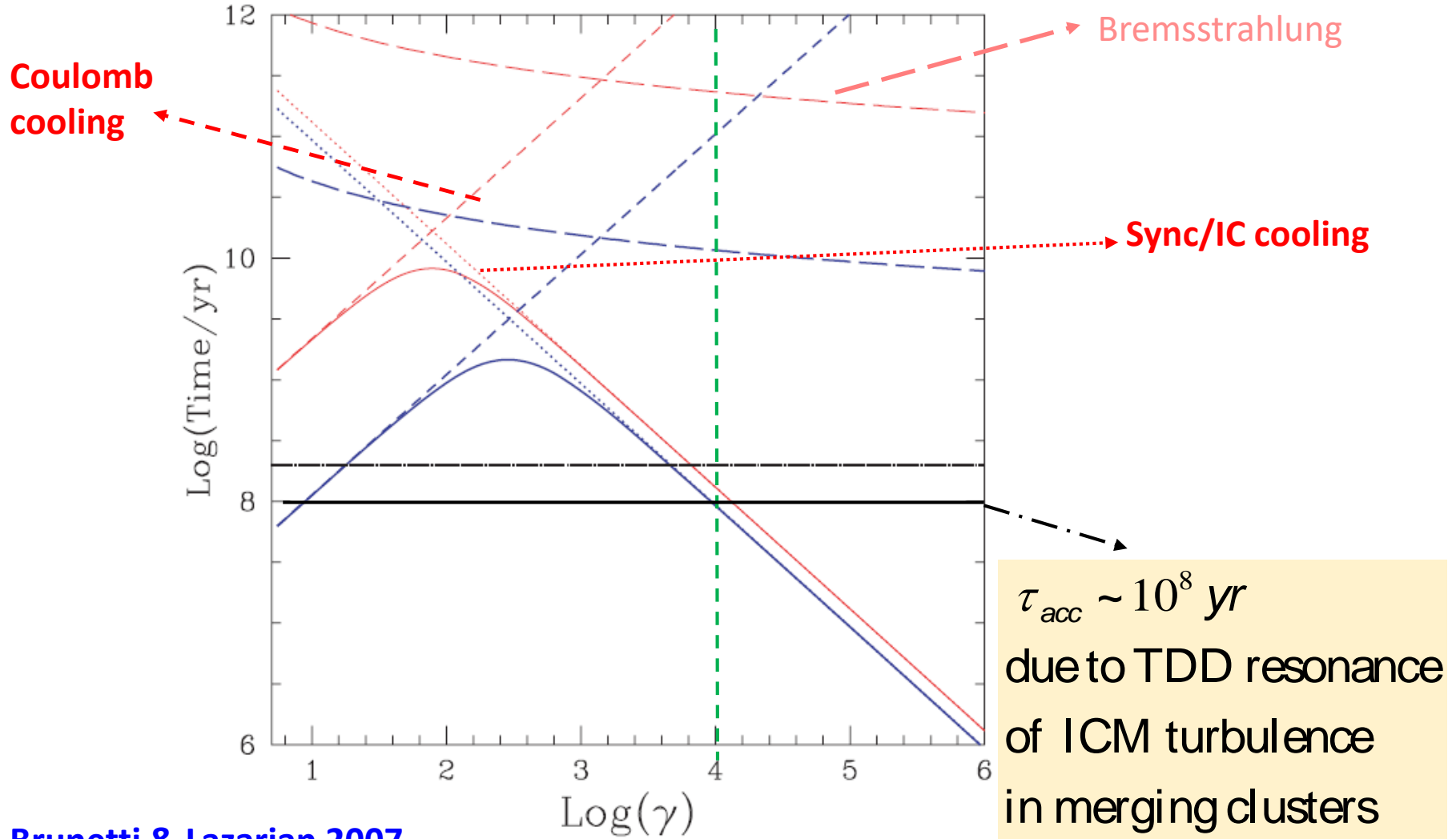
Coulomb/
Synchrotron/iC
losses

momentum diffusion due to TTD: $D_{pp} \approx \frac{p^2}{4\tau_{acc}} \cdot \exp\left(-\frac{R}{R_{decay}}\right)$

where $\tau_{acc} \sim 10^8$ yr : systematic acceleration time **Brunetti & Lazarian 2007, 2011**

simple model for decay of turbulence with $R_{decay} \sim 100$ kpc

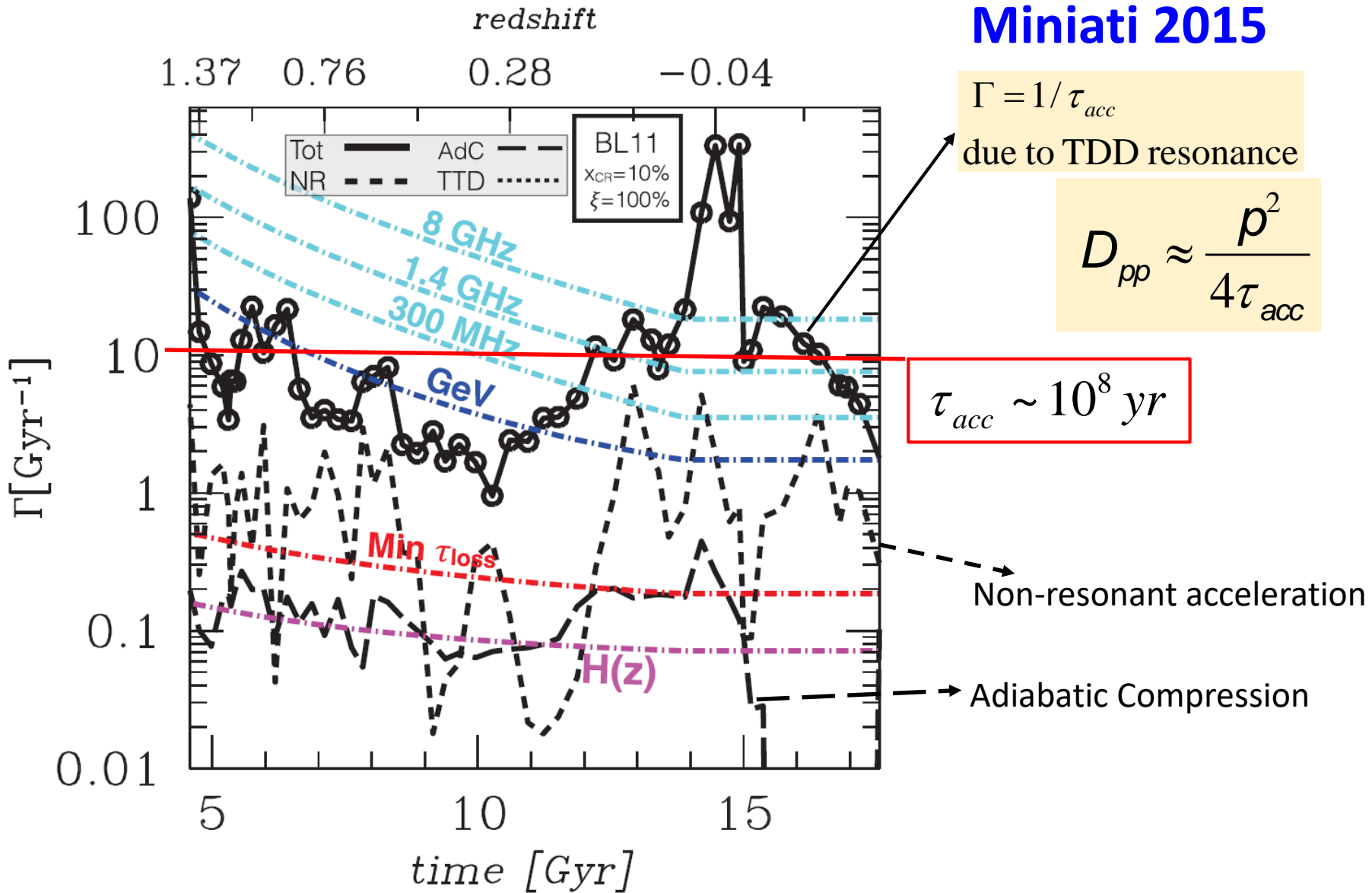
Electron lifetime in the ICM : $B \approx 0.5 \mu\text{G}$, $n_{\text{H}} \approx 10^{-4} \text{ cm}^{-3}$



Brunetti & Lazarian 2007

Acceleration Rate due to ICM Turbulence based on Brunetti & Lazarian 2011

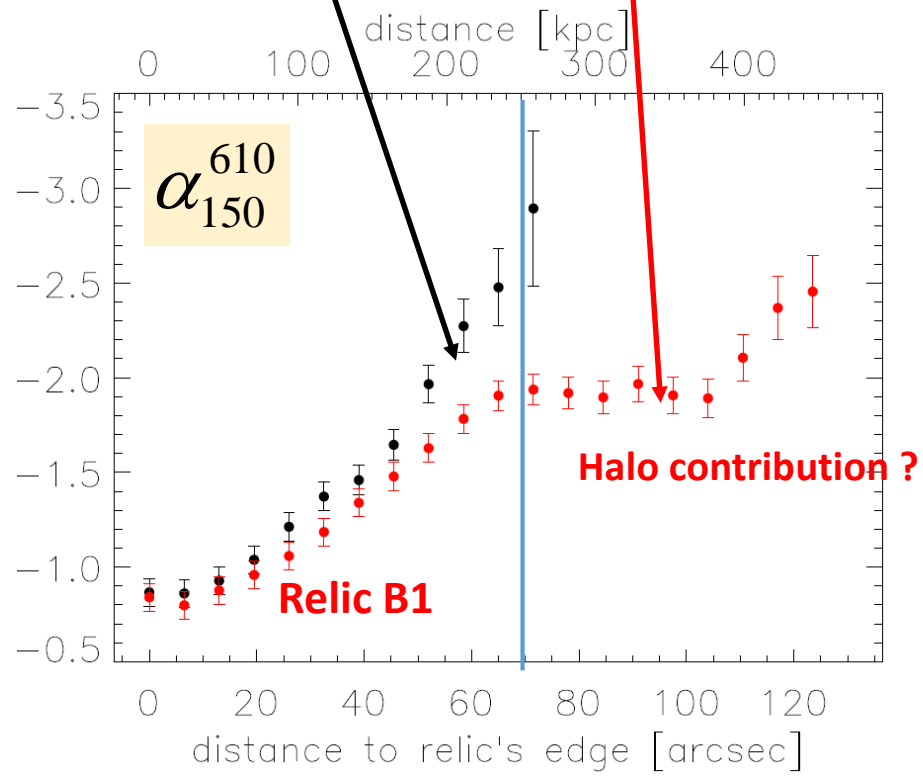
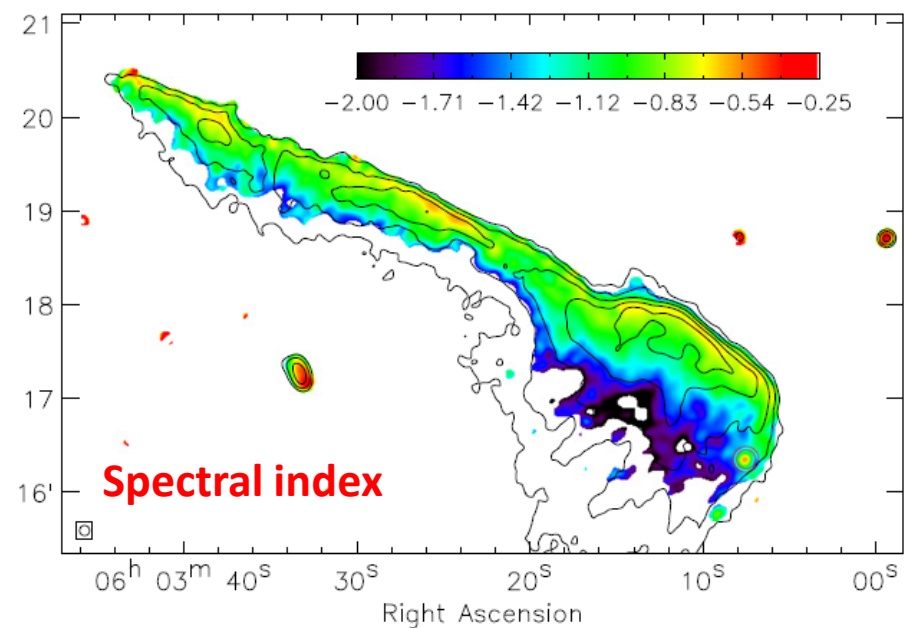
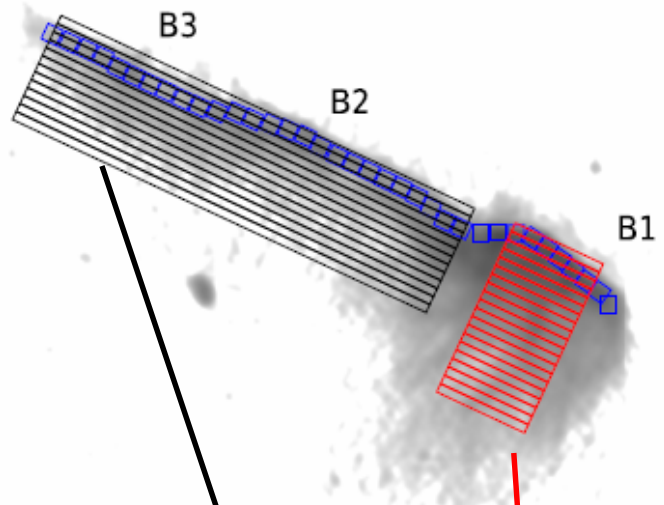
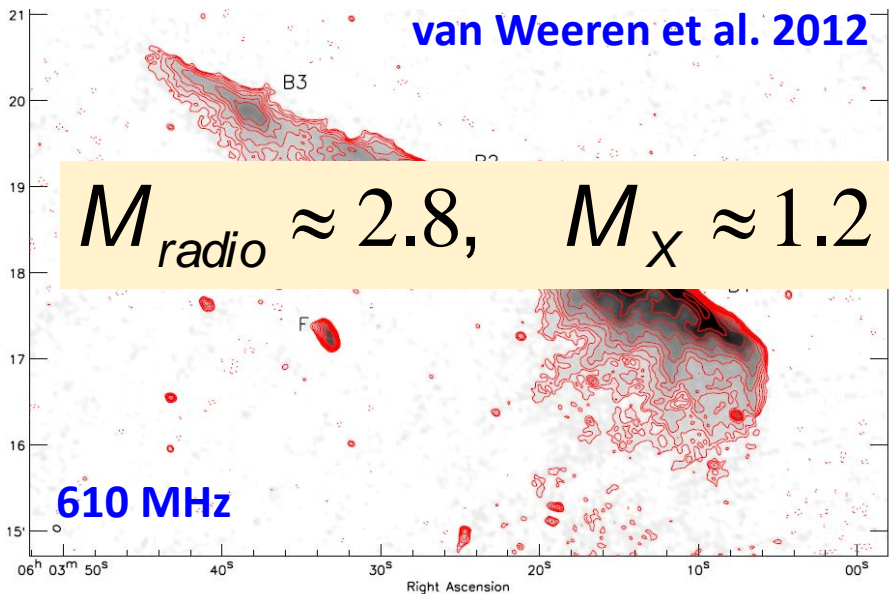
Miniati 2015

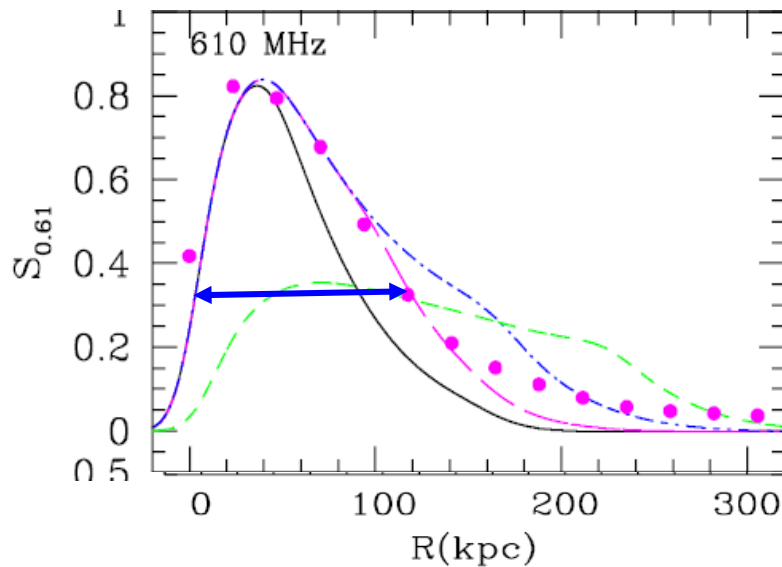


Structure formation simulations → Turbulence in ICM → Turbulence Acceleration rate

Toothbrush relic in RX J0603.3

van Weeren et al. 2016





Shock Parameters

$$M_s \approx 1.4, \quad kT_1 = 6.9 \text{ keV},$$

$$u_s \approx 1.8 \times 10^3 \text{ km/s}$$

radio flux profile at 610 MHz,

$$\Delta l \sim 100 \text{ kpc}$$

Relic Width at a given frequency \sim cooling length of electrons

$$l_{cool} \approx u_2 \cdot t_{cool}(B, z) \approx 100 \text{ kpc} \cdot W_h \cdot u_{2,3} \cdot Q(B, z) \cdot \left[\frac{v_{obs}(1+z)}{0.63 \text{ GHz}} \right]^{-1/2}$$

depends on the postshock flow speed and magnetic field strength

$$u_{2,3} = u_2 / 1000 \text{ km/s}, \quad W_h \sim 1.3$$

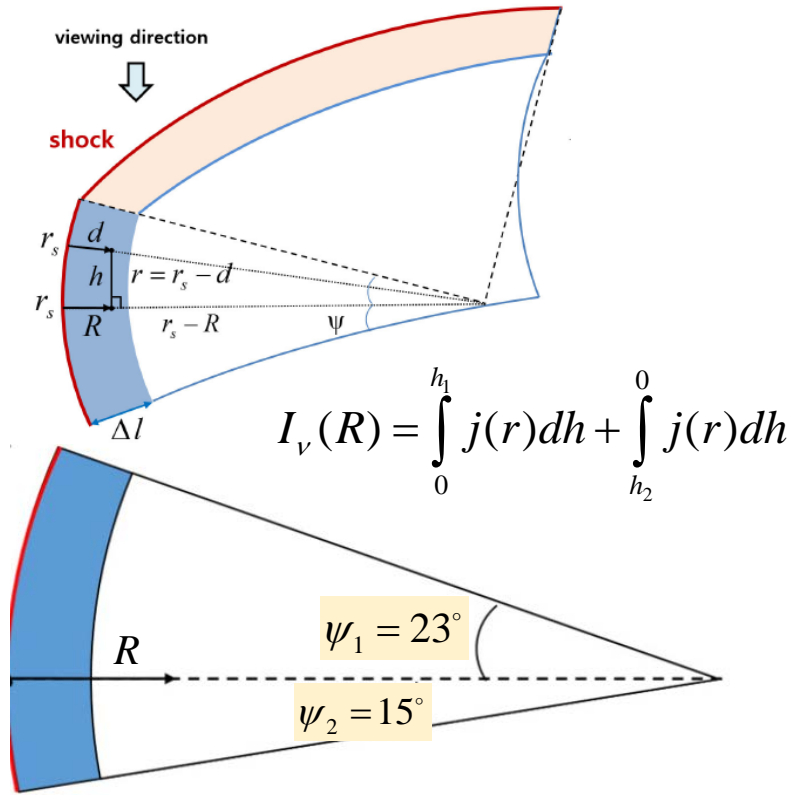
$$Q_{max}(z, B) \approx 0.6 \text{ for } B_2 \approx 2.5 \mu\text{G}$$

Assume pre-existing electrons

$$f_{pre}(p) = f_o \left(\frac{p}{p_{inj}} \right)^{-s} \exp \left[- \left(\frac{\gamma_e}{\gamma_{e,c}} \right)^2 \right]$$

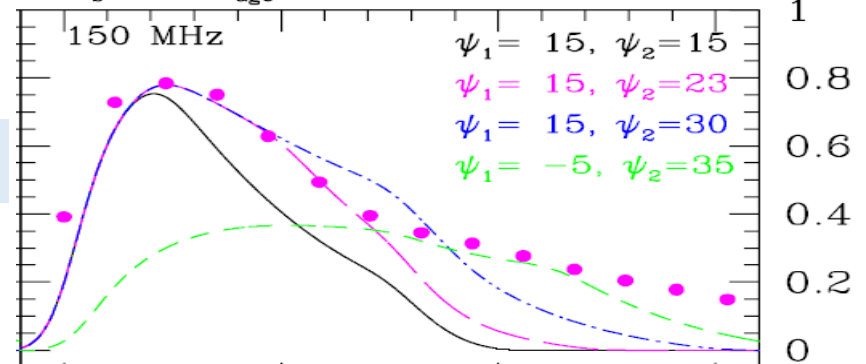
$$s = 4.6, \quad \gamma_{e,c} = 7 \times 10^4$$

Reacceleration model for the Toothbrush Relic (Kang 2016, DSA + cooling only)

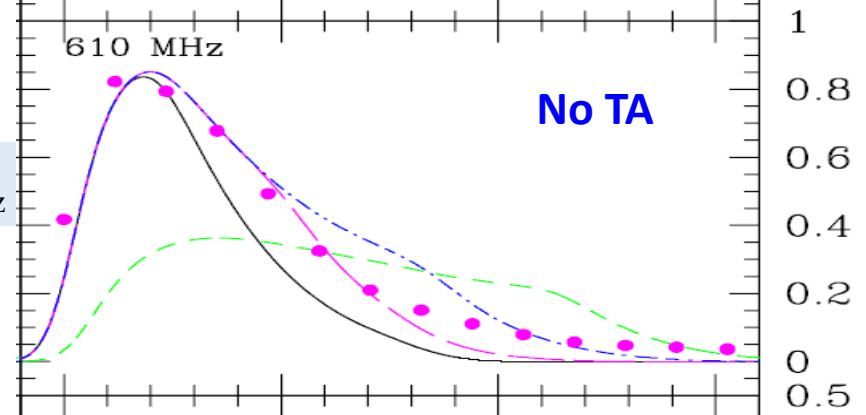


$M_s = 1.4$, $t_{\text{age}} = 110 \text{ Myr}$

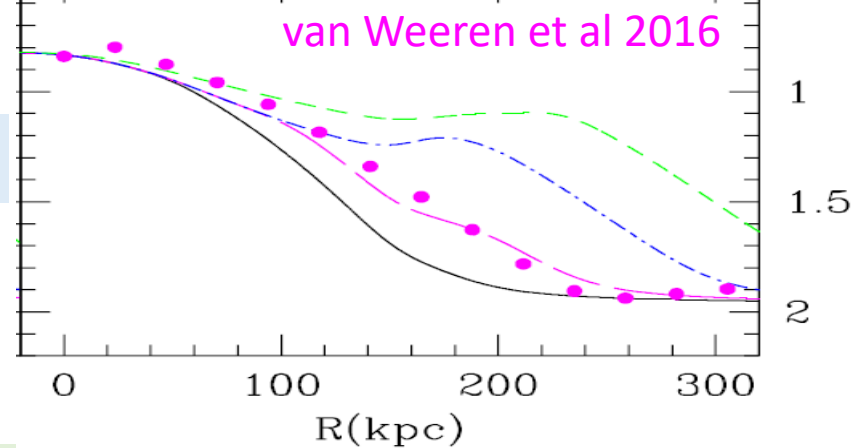
$S_{150\text{MHz}}$



$S_{610\text{MHz}}$



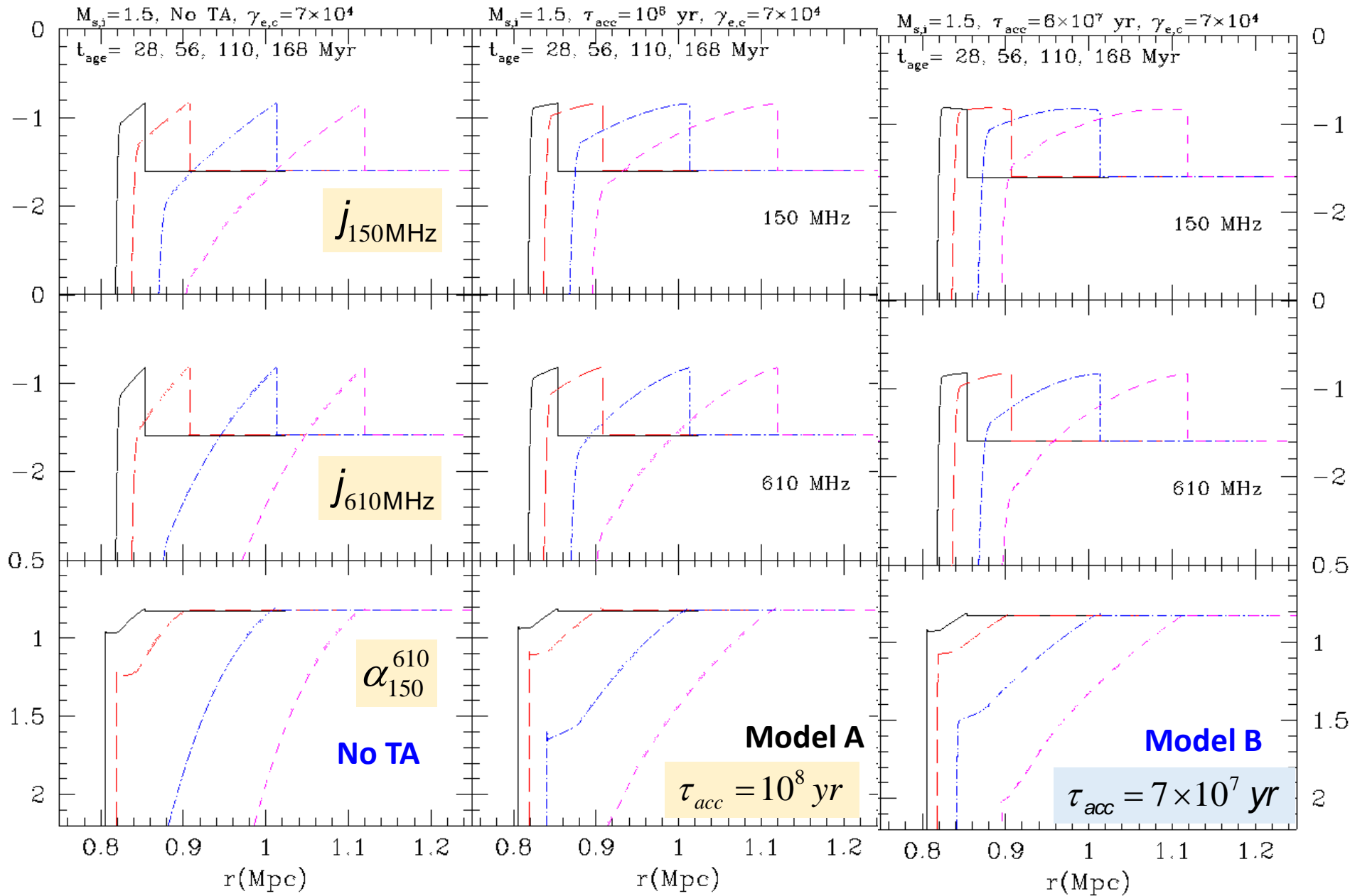
α_{150}^{610}



$M_s \approx 1.4$, $kT_1 = 6.9 \text{ keV}$, $kT_2 = 8.9 \text{ keV}$, $u_s \approx 1.8 \times 10^3 \text{ km/s}$, $B_1 = 2 \mu\text{G}$, $B_2 \approx 2.6 \mu\text{G}$

Reacceleration model for the Toothbrush Relic: Test for Turbulence Acceleration

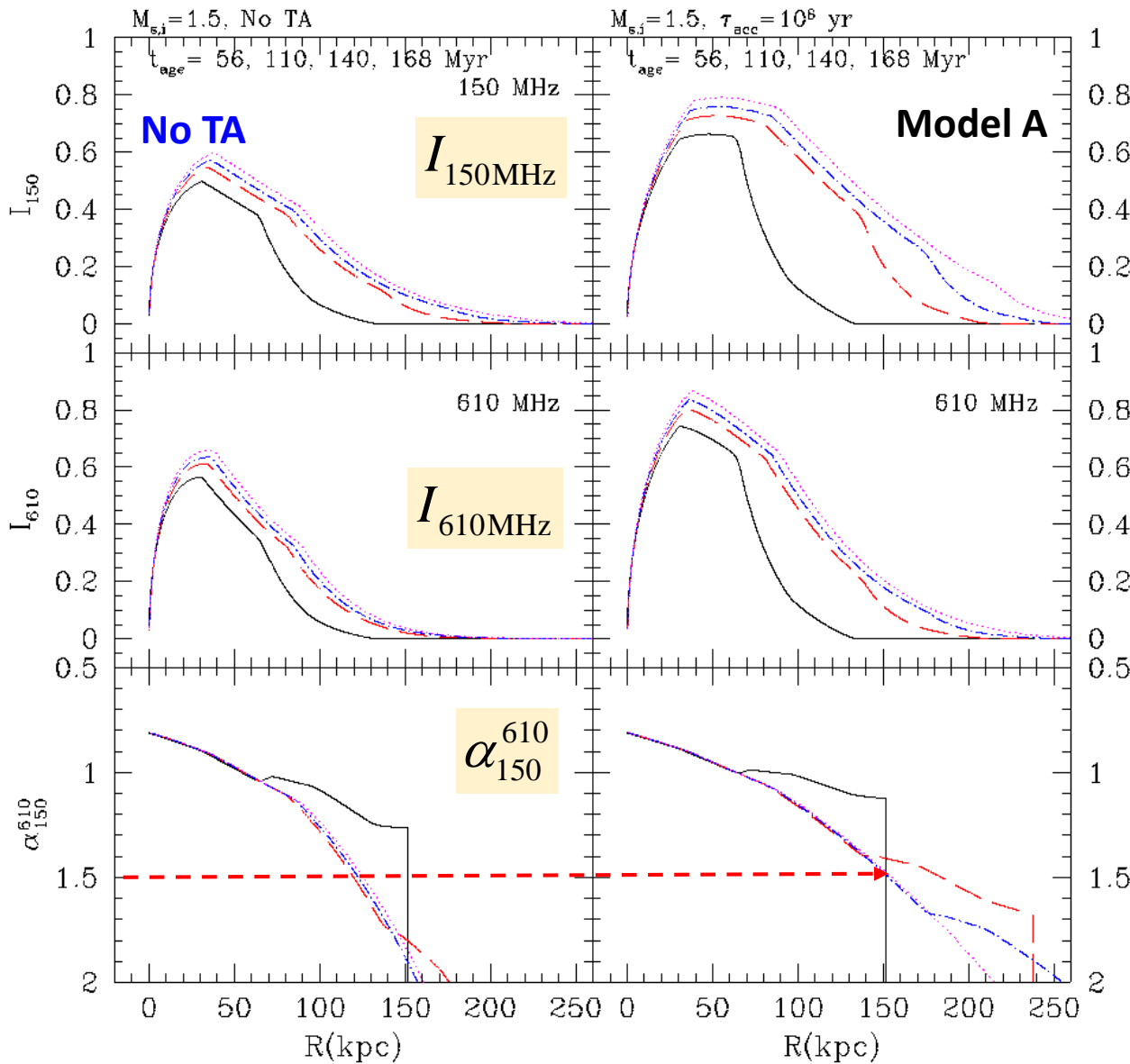
$$D_{pp} \approx \frac{p^2}{4\tau_{acc}} \cdot \exp\left(-\frac{R}{R_{decay}}\right)$$



Coulomb/Sync/iC cooling

With Turbulent Acceleration

Toothbrush Relic Model



**DSA +
Coulomb/Sync/iC cooling
+ Turbulence Acceleration**

Model A

$$D_{pp} \approx \frac{p^2}{4\tau_{acc}} \cdot \exp\left(-\frac{R}{R_{decay}}\right)$$

$$\tau_{acc} = 10^8 \text{ yr}$$

$$R_{decay} = 100 \text{ kpc}$$

surface brightness

$$I_\nu(R) = \int_0^{h_1} j(r) dh + \int_{h_2}^0 j(r) dh$$

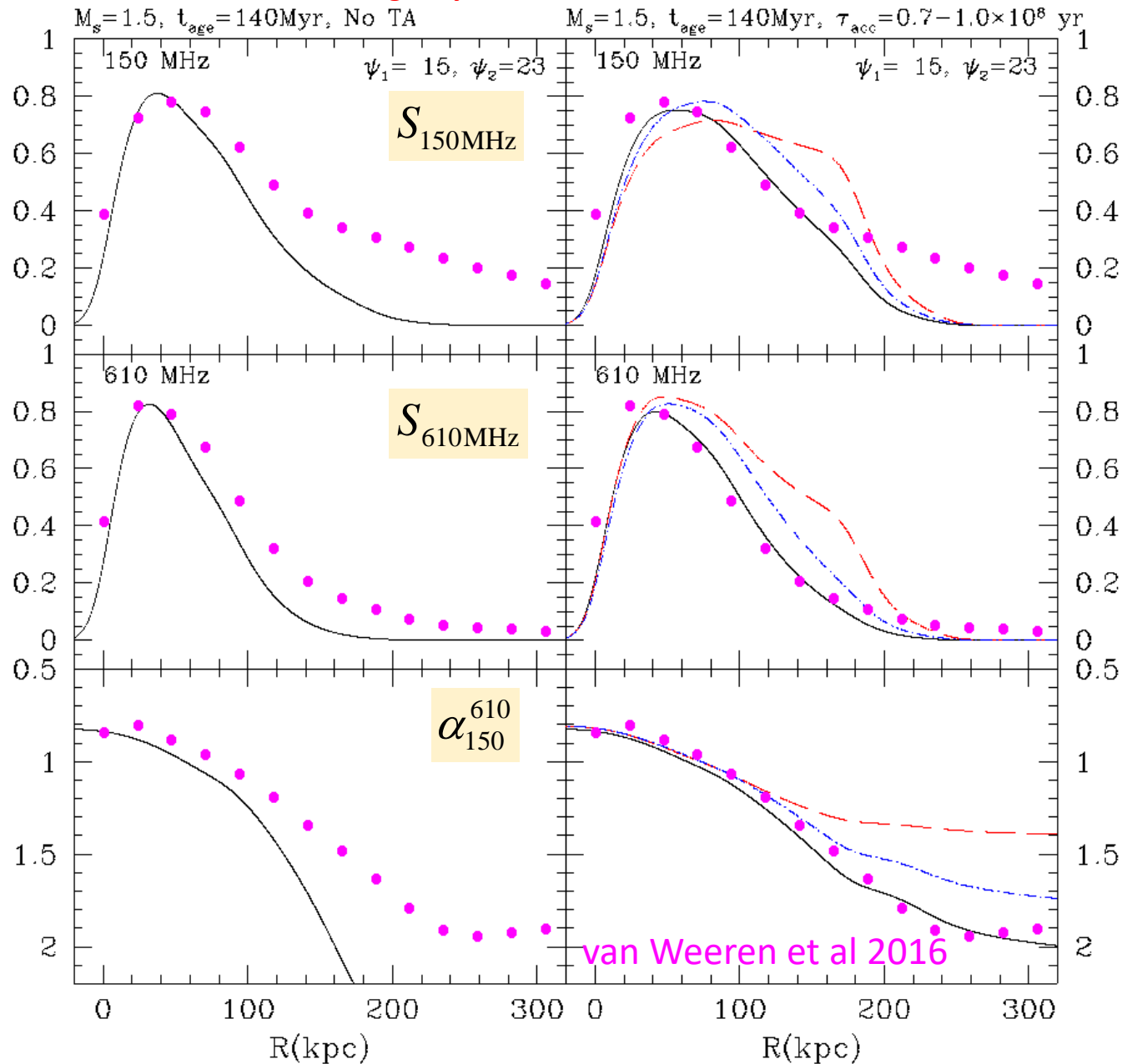
$M_s \approx 1.4$, $kT_1 = 6.9 \text{ keV}$, $kT_2 = 8.9 \text{ keV}$, $u_s \approx 1.8 \times 10^3 \text{ km/s}$, $B_1 = 2 \mu\text{G}$, $B_2 \approx 2.6 \mu\text{G}$

(Kang, Ryu, & Jones 2016)

Beam convolved radio flux profile of the Toothbrush Relic

Cooling only

With Turbulent Acceleration



Model A:

$$D_{pp} \approx \frac{p^2}{4\tau_{\text{acc}}} \cdot \exp\left(-\frac{R}{R_{\text{decay}}}\right)$$

$$\tau_{\text{acc}} = 10^8 \text{ yr}$$

Model B:

$$D_{pp} \approx \frac{p^2}{4\tau_{\text{acc}}} \cdot \exp\left(-\frac{R}{R_{\text{decay}}}\right)$$

$$\tau_{\text{acc}} = 7 \times 10^7 \text{ yr}$$

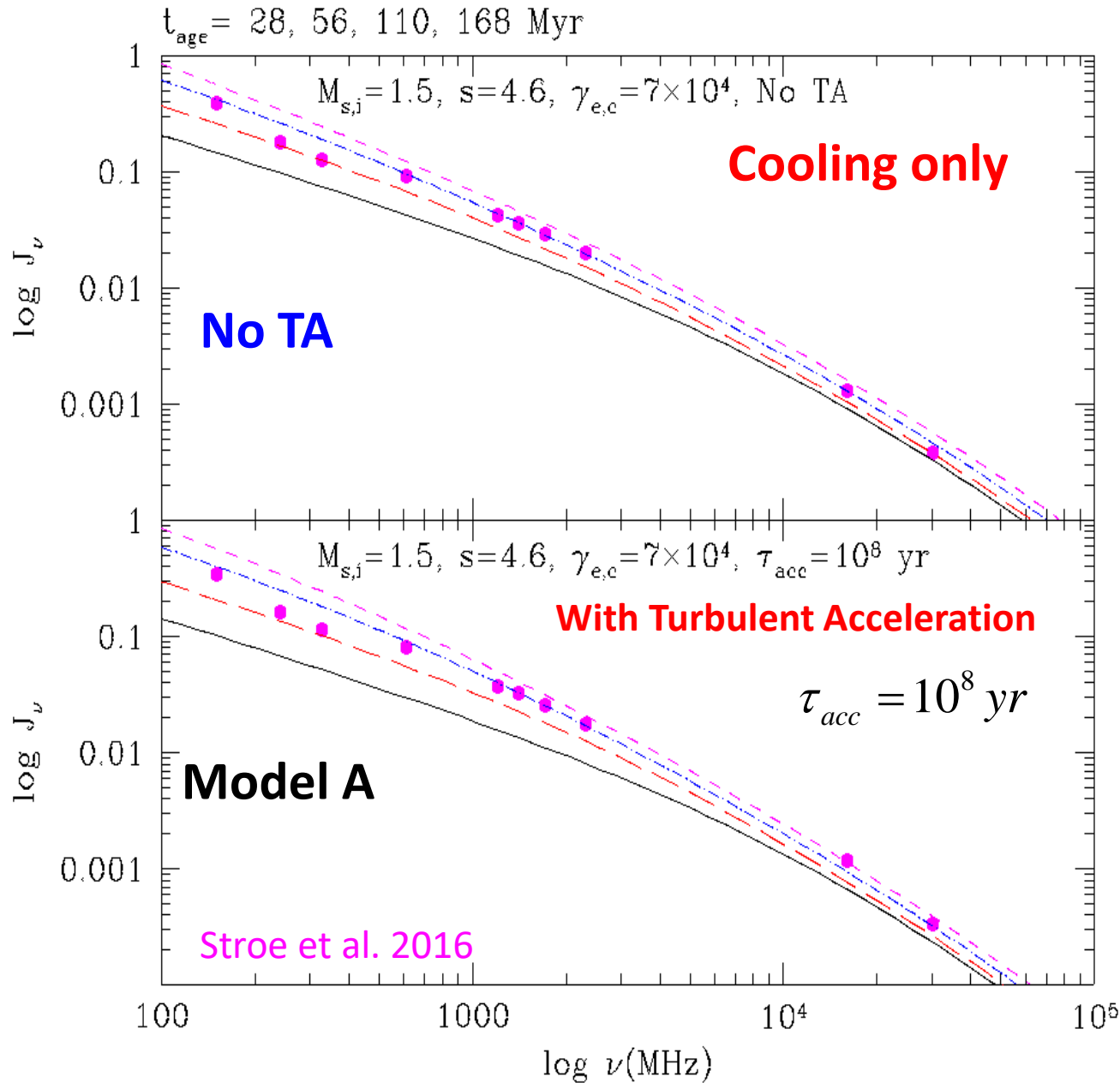
Model C:

$$D_{pp} \approx \frac{p^2}{4\tau_{\text{acc}}}$$

$$\tau_{\text{acc}} = 10^8 \text{ yr}$$

Gaussian Smoothing with
23.5 kpc width

Integrated Radio Spectrum of the Toothbrush relic



$$D_{pp} \approx \frac{p^2}{4\tau_{acc}} \cdot \exp\left(-\frac{R}{R_{decay}}\right)$$

$$\tau_{acc} = 10^8 \text{ yr}$$

$$R_{decay} = 100 \text{ kpc}$$

Summary

- **Re-acceleration model** explains observed features of radio relics

(1) relative rareness, compared to the expected number of shocks in the ICM

(2) radio spectral index α is determined by M_s , s , & $\gamma_{e,c}$ so $M_{radio} \neq M_x$

(3) solves the low acceleration efficiency problem for weak shocks with $M < 3$

- **Turbulent acceleration by ICM turbulence**

due to TTD resonance with fast mode waves (Brunetti & Lazarian 2007, 2011)

$$D_{pp} \approx \frac{p^2}{4\tau_{acc}} \cdot \exp\left(-\frac{R}{R_{decay}}\right) \quad \text{with } \tau_{acc} \sim 10^8 \text{ yrs, } R_{decay} \approx 100 \text{ kpc}$$

The test results are consistent with the observational data of the Toothbrush and the Sausage relics.