

## Some Like it Hot! A new high temperature MFM comes online at the IRM

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 IRM

Applications of different magnetic imaging techniques are becoming increasingly popular in the study of rock magnetism. "Seeing magnetism" undoubtedly provides better understanding of known phenomena and discovery of important new ones. To this extent, a variety of techniques involving different underlying principles exist, allowing observations at room, low and high temperatures:

Magnetic force microscopy (MFM) allows imaging micromagnetic features that are 10's of nanometers in scale, using thin-sections or polished surfaces of natural samples. The origins of magnetic remanence can thus be investigated at the nanoscale, within single grains.

Optical techniques such as Bitter, magneto-optical Kerr effect (MOKE) and TEM-based techniques (electron holography, Lorentz microscopy) have also been used to image magnetic structures in natural materials and synthetic analogues. While each of these imaging techniques has advantages and disadvantages related to resolution, ease of use and flexibility, almost all of them are confined to room T observations due to various technical limitations of the existing systems.

Exceptions to this have been attempts to image low-T micromagnetic structures in the monoclinic phase of magnetite ( $T < 122$  K) using prototype instruments (e.g., LT MFM Moloni et al, 1996) or transmission electron microscopes equipped with cooling stages (e.g., Kasama et al. 2012). Magnetic imaging above room-T has been confined to low-resolution optical techniques either using the Bitter technique to investigate MD particles (10-50  $\mu\text{m}$ ) of magnetite and titanomagnetite up to  $T \sim 600$  K, or with MOKE to study mm-sized single crystals of magnetite up to the Curie temperature  $\sim 850$  K (e.g., Muxworthy and Williams 2006).

Here at the IRM we have had a room-temperature MFM since 1992. Its primary uses have been to investigate the internal structures of domain walls and domains, the relationship between mineral microstructures and magnetic domain configurations in natural and synthetic (titano)magnetites in varying remanent states, magnetization structures in silicate-hosted inclusions and oxidized titanomagnetites, and the effects of shock remagnetization

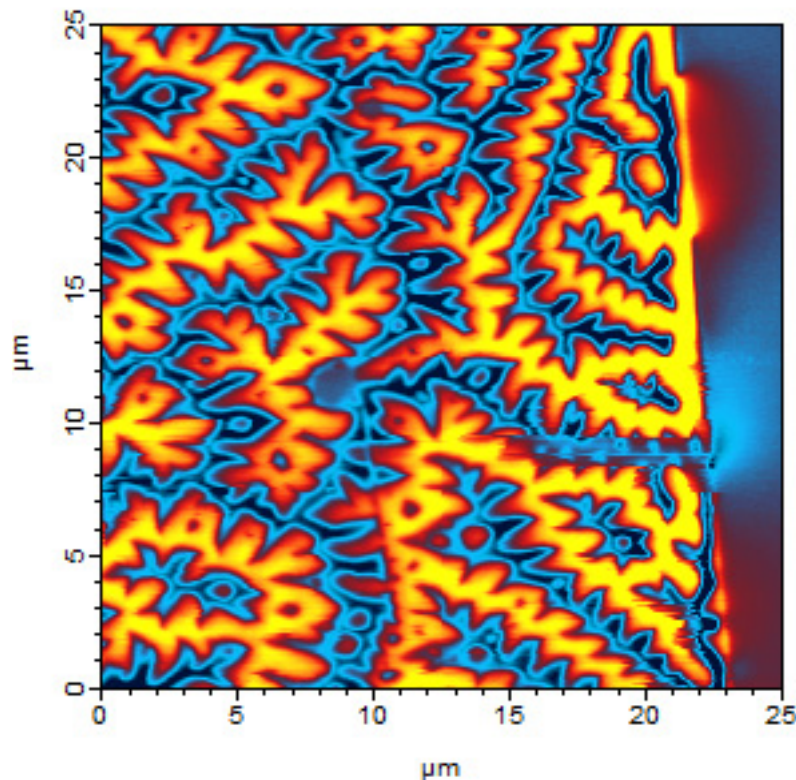


Fig 1. A magnetic image of a Mount St. Helens Ti-magnetite acquired on the new MFM

on domain states in iron sulfides.

However, many advances have been made in MFM instrumentation over the past 20 years, including the possibility of making observations over a range of controlled temperatures: the latest generation MFMs are outfitted with environmental stages for precise heating and cooling in controlled atmospheres up to  $T = 650-700$  K ( $377-427^\circ$  C) which allow a variety of important rock-magnetic processes to be studied directly in single grains. Furthermore, MFM samples are usually in a form suitable for reflected light, EM and microprobe analyses, allowing complementary chemical, compositional, and structural characterizations on the same grains imaged magnetically.

We are very proud to introduce one such new generation MFM, an Asylum Research MFP-3D MFM, to our quiver of magnetic imaging instrumentation. Together with our old MFM, and the Bitter and MOKE capabilities, the new MFM will improve our (and our guests') ability

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# Visiting Fellows' Reports

## Gegham Style: Applying Obsidian Magnetic Studies in Central Armenia

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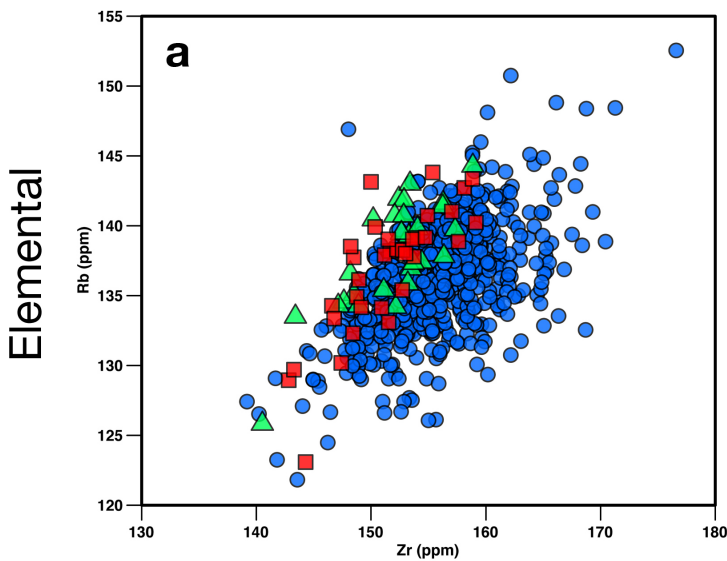
The Gegham highlands in central Armenia encompass more than a hundred Quaternary volcanoes, including a series of archaeologically significant obsidian sources. In particular, the extensive Gutansar complex was the volcanic origin of most obsidian artifacts discovered at

two sites recently excavated by the Hrazdan Gorge Palaeolithic Project, a collaboration including the University of Connecticut's Palaeolithic Studies Lab and Armenia's Institute of Archaeology and Ethnography. Using portable X-ray fluorescence, I analyzed 1700 obsidian artifacts from Lower Palaeolithic Nor Geghi-1 and Middle Palaeolithic Lusakert-1. Based on the artifacts' elemental "fingerprints," 93% were matched to the Gutansar volcanic complex.

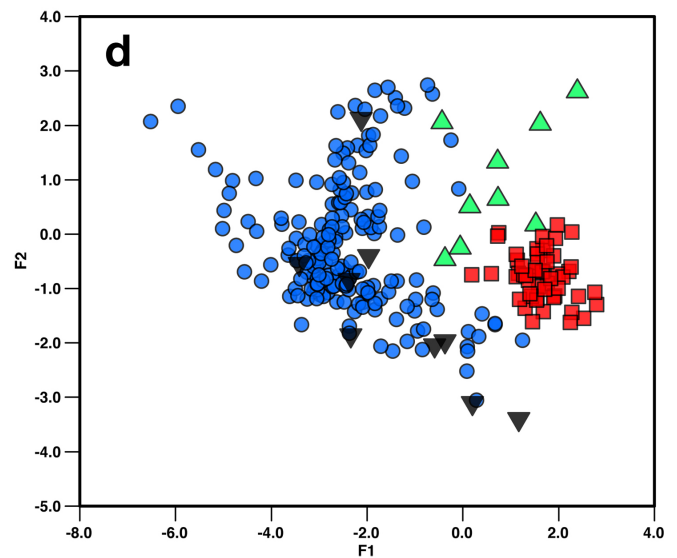
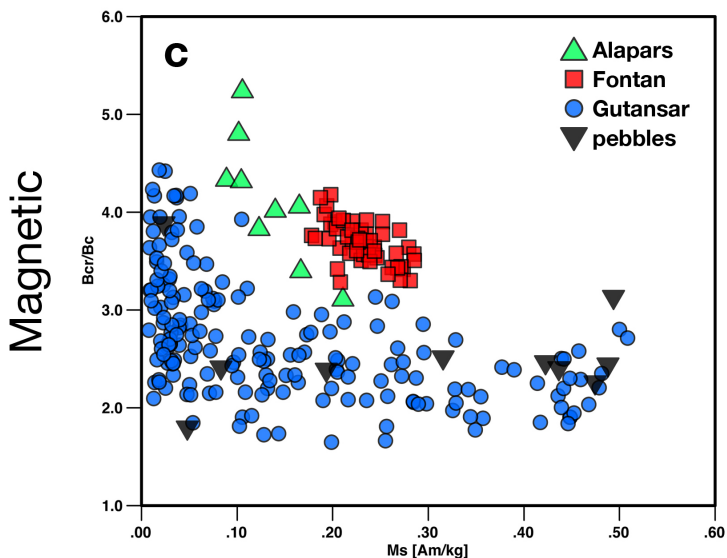
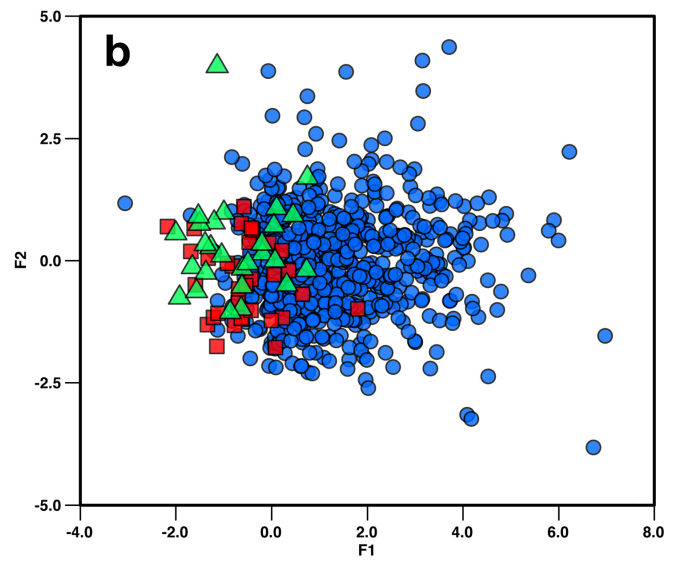
For hundreds of millennia, early humans used obsidian available from flows, domes, and intrusions scattered across more than 80 km<sup>2</sup> around Gutansar. However, any patterns that could reveal how past peoples used the different obsidian outcrops are currently invisible. Obsidian is chemically homogeneous throughout this complex. Thus, traditional chemical sourcing methods cannot discern obsidian from different areas of the Gutansar complex.

My research at the IRM uses the magnetic properties of Gutansar obsidian as a means to identify patterns in stone tools crucial to reconstruct settlement, subsistence, and other behaviors during the Palaeolithic.

### Scatterplots



### Discriminant Analysis



Without distinguishing obsidian tools from varied parts of this complex, we are left with a largely incomplete picture of early human activities due to missing information about resource acquisition in the Gegham highlands. My work seeks to remedy the gap by using magnetic properties to recognize different quarrying sites. Pilot studies in Turkey and California have established the validity of this approach (Frahm and Feinberg 2013).

The pioneering work of McDougall et al. (1983) used NRM, Mr, and  $\chi$  to match obsidian artifacts to their origins, albeit with mixed success. My preliminary work showed that hysteresis parameters (Bc, Bcr, Mr, Ms) can discern not only obsidian flows that are geochemically identical but also different quarries within a flow. The parameters reflect differences in compositions, sizes and shapes of magnetic grains in obsidian. At the IRM in July, I acquired hysteresis loops and backfield curves for 250 specimens from 25 Gutansar sampling locations.

The plots here show some of my initial results. Plots A and B demonstrate that obsidian from three volcanic centers at this complex -- Fontan, Alapars, and Gutansar proper -- cannot be discerned by elemental composition. Magnetic hysteresis parameters, however, can differentiate them. Plots C and D show that hysteresis parameters largely separate the three volcanic centers, and they reveal that obsidian pebbles found at Nor Geghi-1 match the Gutansar flow rather than the Fontan and Alapars domes. Other data show that different portions of the Fontan dome have obsidian with distinct magnetic properties, enabling quarry identification.

Because magnetic properties of obsidian vary on scales different from geochemistry, their measurement may reveal previously unnoticed patterns in lithic assemblages. The ultimate goal uncovering new insights about human behavior in the Gegham highlands.

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**The next Visiting  
Fellowship  
application deadline is  
October 30, 2013.**

## Current Articles

A list of current research articles dealing with various topics in the physics and chemistry of magnetism is a regular feature of the IRM Quarterly. Articles published in familiar geology and geophysics journals are included; special emphasis is given to current articles from physics, chemistry, and materials-science journals. Most are taken from ISI Web of Knowledge, after which they are subjected to Procrustean culling for this newsletter. An extensive reference list of articles (primarily about rock magnetism, the physics and chemistry of magnetism, and some paleomagnetism) is continually updated at the IRM. This list, with more than 10,000 references, is available free of charge. Your contributions both to the list and to the Current Articles section of the IRM Quarterly are always welcome.

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to study the fundamental physics of magnetic behavior at elevated temperatures by direct imaging of domain structures in ambient or applied fields.

The MFP-3D is a significant upgrade to our existing room temperature MFM. Unlike our current MFM that can accommodate samples up to 1 cm in diameter, the new unit can scan much larger samples, up to 10 cm in diameter and 2.5 cm in thickness. This larger sample size can easily accommodate a variety of geological samples such as 1" cores and thin section slides without the "downsizing" necessary to fit our older MFM and will allow a greater range of samples to be studied. The MFP-3D can also collect high-pixel-density images (up to 5k x 5k) with high-speed data capture up to 5 MHz, allowing imaging of large structures and small features in the same image, at rates substantially higher than those of the old instrument. In addition to room-temperature imaging, the MFP-3D also allows high-temperature imaging up to 400° C. This is accomplished with a separate high temperature environmental stage (PolyHeater) that sits on the MFP-3D XY scanning stage. Heating and cooling can be done in air or controlled atmospheres (e.g., He, Ar, CO, etc). Sample temperature can be maintained to better than 0.2° C precision with accuracy to 0.5° C and temperature overshoots less than 0.2° C. All control and measurement functions are fully programmable for custom heating-cooling imaging experiments. The heating stage can be used with standard cantilever holders for heating up to 250° C or Asylum Research PEEK™ cantilever holders for T> 250° C. The cantilevers and tips are made from silicon or silicon nitride while the MFM tips are coated with CoCr and all can be heated to much higher temperatures than 400° C and survive. However, sample size is more restricted for HT imaging with a maximum sample diameter of 20 mm and sample thickness of 2 mm but 8 mm diameter is recommended.

The T range of the new MFM reaches the Curie point of important magnetic minerals such as titanomagnetite and pyrrhotite, enabling observation of phenomena related to full thermoremanence (TRM) acquisition. Partial TRMs, involving heating to temperatures less than the Curie point, are of central importance in paleointensity work and in paleomagnetic overprinting, and the new MFM will provide entirely new ways of studying these as well, in minerals like magnetite. Magnetic mineral alteration via oxidation and reduction can also be examined at the nanometer scale using the controlled atmosphere chamber of the new MFM.

### Magnetic Force Microscopy

Magnetic force microscopy is one member of a family of imaging modes collectively called scanning probe microscopy, which also includes the scanning tunneling microscope, atomic force microscope and the scanning SQUID microscope.

Unlike the scanning SQUID microscope, which is designed to map magnetic fields from geological thin sections at ~100 μm resolution to study natural remanent

magnetization, the MFM is used to map the fine scaled (~50 nm - 100 μm) micromagnetic structures within ferromagnetic particles.

It was developed in 1987 shortly after the invention of the atomic force microscope (AFM), of which it is a special mode of operation, and is a well-developed technique for high resolution imaging with limited sample preparation and without the need for high vacuum conditions. It has been widely used since the early 1990s in the fundamental research of magnetic materials, as well as the development of magnetic recording components. An introduction to MFM operating principles appeared in the IRM Quarterly soon after our first instrument was acquired (IRMQ v2, n4, 1992). Like the AFM, a sharp tip, a few micrometers long and less than 50 nm in diameter, interacts with the specimens surface forces, therefore mapping the surface topography; the MFM tips are coated with magnetic material to enable mapping of the magnetic stray fields above a sample by scanning above the surface. The ferromagnetic tip is mounted on the free end of a flexible cantilever that is 100-200 μm in length. Minute forces on the magnetic tip are detected by changes in the resonant frequency of the cantilever induced by magnetostatic interactions between the ferromagnetic tip and the sample's stray magnetic field at typical tip-sample separations less than 100 nm. As the tip is scanned across the sample, an image of the inhomogeneous stray magnetic fields deriving from such features as domain walls, vortex structures, and surface termination of domains is produced. The MFM does not provide direct information on the near surface distribution of magnetization. Instead, it is a measure of the local stray field components of micromagnetic structures resulting from the distributions of magnetic poles where the magnetization diverges or has a component normal to the surface. In certain respects this is similar to the classical Bitter decoration technique but with spatial resolutions 1000 times better, and with the ability to distinguish polarity.

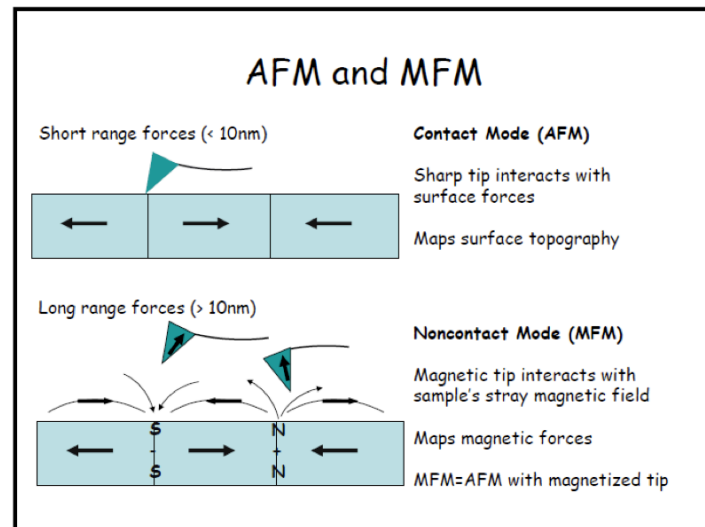


Fig 2. (from Moskowitz, 2008) Vertical gradients of stray field above domain walls exert forces on the tip, attractive or repulsive according to field polarity and tip moment orientation. Vertical force gradients alter the tip oscillation amplitude and phase.

### Basics of magnetic contrast information

The following summary is derived largely from the recent review by Ferri et al (2012). The cantilever (incorporating the magnetic tip) is essentially a spring, with a force-displacement constant  $k=F/z$ . In the absence of sample-tip interaction forces, the cantilever has a resonant frequency

$$\omega_0 = 2\pi f_0 = \sqrt{\frac{k}{m}}$$

where  $m$  is the mass. The cantilever during operation is excited to vibrate close to its resonance frequency, and a laser deflection sensor monitors the tip oscillation amplitude and phase shift with respect to the drive signal.

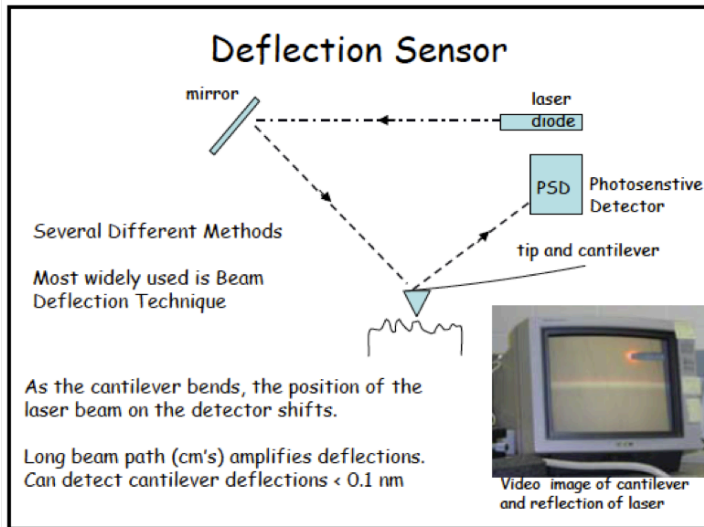


Figure 3. Schematic of an MFM deflection sensor which monitors the tip's oscillation amplitude and phase shift.

When there are probe-sample interaction forces that vary vertically (with vertical derivative  $dF_i/dz$ ), the effect is equivalent to a change in the cantilever spring constant,  $k_{\text{eff}} = k - dF_i/dz$ . Because the force exerted on a magnetic dipole (like the tip) is proportional to the gradient of the magnetic field, the MFM signal is (to a first approximation) related to the second derivative of the stray field distribution ( $dF_i/dz \sim d^2Bz/dz^2$ ).

An attractive interaction with  $dF_i/dz > 0$  will "soften" the cantilever spring, decreasing its resonance frequency:

$$f = f_0 \sqrt{1 - \frac{\partial F_i / \partial z}{k}}$$

Such a shift leads to changes of the oscillation amplitude of the probe and of its phase, which are measurable quantities that can be used to map the lateral variation of  $dF_i/dz$ . A common detection method uses the amplitude signal and is referred to as amplitude modulation (AM). The cantilever is driven slightly off resonance ( $\omega_D$ ), where the slope of the amplitude versus frequency curve is high, thus maximizing the signal obtained from a given force derivative (Fig. 4).

The force gradient  $dF_i/dz$  may originate from a range of sources, including long range electrostatic probe-

sample interactions ( $>10$  nm,  $< \sim 100$  nm), nanoscale ( $<10$  nm) van der Waals and/or Coulombic Forces, damping, or capillary forces. However, MFM relies on those forces that arise from a long-range ( $>10$  nm) magnetostatic coupling between probe and sample. This coupling is dependent on the internal magnetic structure of the probe, which greatly complicates the mechanism of contrast formation. In general a magnetized body, when brought into the stray field of a sample, will possess magnetic potential energy  $E$ :

$$E = -\mu_0 \int \vec{M}_{\text{tip}} \cdot \vec{H}_{\text{sample}} dV_{\text{tip}}$$

where  $\mu_0$  is the vacuum permeability. The force acting on a magnetic tip can be calculated by:

$$\vec{F} = -\vec{\nabla} E = \mu_0 \int \vec{\nabla} (\vec{M}_{\text{tip}} \cdot \vec{H}_{\text{sample}}) dV_{\text{tip}}$$

The integration needs to be carried out over the tip volume, or rather its magnetized part.

A limitation of MFM is that the magnetic configuration of the sensing probes is rarely known in detail. It is therefore not possible to model the measured signal from first principles and MFM may generally not be used quantitatively to measure stray fields, let alone magnetization. MFM can, however, be used to compare the experimentally detected stray field variation of a micromagnetic object to that obtained from certain model calculations, or to observe changes that result from applied fields or demagnetization treatments.

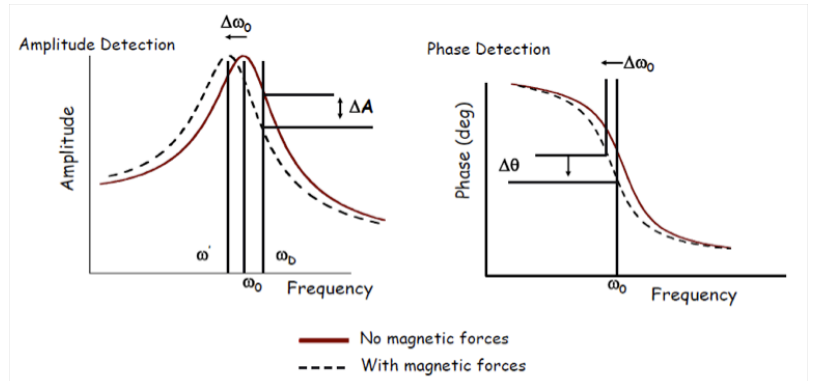


Figure 4. (from Moskowitz, 2008) Changes in the magnetic force on the tip produce changes in resonance frequency, amplitude, and phase of the cantilever oscillations.

### Requirements for MFM tips

The cantilever/tip assembly is obviously the critical element of a magnetic force microscope. The tip shape is important due to the long-range nature of magnetic forces. Originally, electrochemically etched wires of cobalt or nickel were used as cantilevers. Thanks to the widespread use of AFM, cantilevers with integrated sharp tips are now fabricated in large numbers out of silicon-based materials. These tips can be coated with a thin layer of magnetic material for the purpose of MFM observations. A lot of effort has been spent on the optimization of magnetic tips in order to obtain quantitative information

from MFM data. Coating of conventional tips generates a pattern of magnetic domains, which reduces the effective magnetic moment of the tip. The exact domain structure is unknown and can even change during MFM operation. Nevertheless, some information on the magnetization state of selected probes has been acquired using electron holography.

The spatial resolution in MFM imaging is related to the tip-sample distance, but also to the magnetized part of the tip that is actually exposed to the sample stray field. Thus in order to improve lateral resolution, it is beneficial to restrict the magnetically sensitive region to the smallest possible size. Ideally the effective volume of the probe would consist of a small single-domain ferromagnetic particle located at the probe apex. So-called supertips have been developed based on this idea. However, there is a physical lower limit for the dimensions because an ultra-small particle becomes superparamagnetic, and this limitation becomes more severe for high-temperature operation.

The demand for a strong signal, produced by a small sensitive volume, indicates the need to maximize the magnetic moment in the tip. For this reason a single domain tip will give the best results and is also easier to describe theoretically. Materials with a high saturation magnetization should be used in order to limit the required volume. The well-defined magnetic state of a tip should be stable during scanning and it should interfere as little as possible with the sample's magnetization. A high switching field of the tip can be obtained through the influence of shape anisotropy, which forces the magnetization vector field near the probe apex to align with its axis of symmetry. Ultimately, the smallest detail from which a sufficient signal-to-noise ratio can be gained is determined by the sensitivity of the deflection sensor, as well as the noise characteristics of the cantilever.

### Imaging procedure

Different modes of operations of the MFM exist, mainly subdivided under two main categories: [Constant] Contact mode and Dynamic mode.

In "Contact mode", the probe tip is dragged across the

surface of the specimen to 'feel' the topography, this results in high friction and repulsive forces (also the typical AFM mode of operation).

In "Dynamic mode", the probe oscillates flexurally above the surface, "tapping" it as the specimen is scanned. "AC mode", "tapping mode" and "intermittent mode" are trademarked terms for dynamic mode (by Quate/Asylum, Bruker and Topometrics, respectively).

"Dynamic mode" is further subdivided into Attractive mode, where the attractive forces dominate (can be either in contact or non-contact, i.e. the tip oscillates above the specimen either touching/tapping the surface or not) and Repulsive mode, where the repulsive forces dominate (is always in contact).

"Mixed mode" is also a sub-category of "Dynamic mode", which is a switching between "Attractive" and "Repulsive" modes.

Attractive mode broadens the potential well and generates a positive Phase shift, whereas Repulsive mode tightens the potential well and generates a negative Phase shift, therefore generating a sharper image. To get into Repulsive mode it is thus necessary to lower the Phase shift, which is achieved by hitting the surface harder, increasing the free amplitude with respect to the setpoint amplitude. Many other modes of operation also exist, but these are probably the most common.

Quantitative information on the sample stray field can only be derived from MFM images when topographic signal contributions are not included. This is especially important when the tip is brought very close to the sample (in order to improve resolution), since non-magnetic forces become increasingly stronger. The solution to this problem is to keep the topography influence constant by letting the tip follow the surface height profile. This constant distance mode places higher demands on instrument stability, because it is sensitive to drift.

To separate the two signal contributions the tapping-lift method is employed. This involves measuring the topography on a first scan, and the magnetic information during a second scan on the same profile. The difference in height  $\Delta h$  between the two scans, the lift-height, is specified by the user. Topography is measured in dynamic

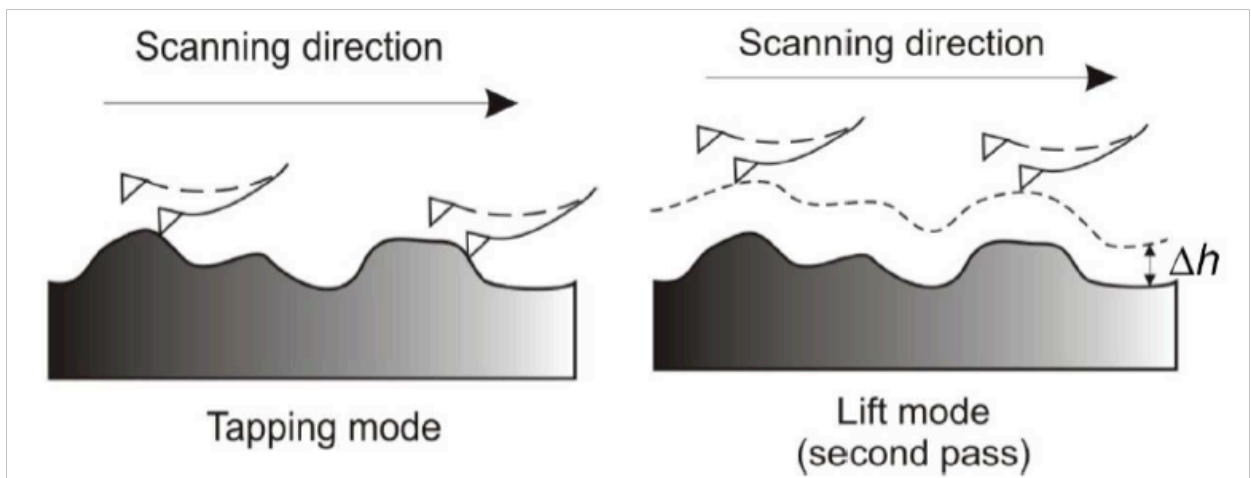
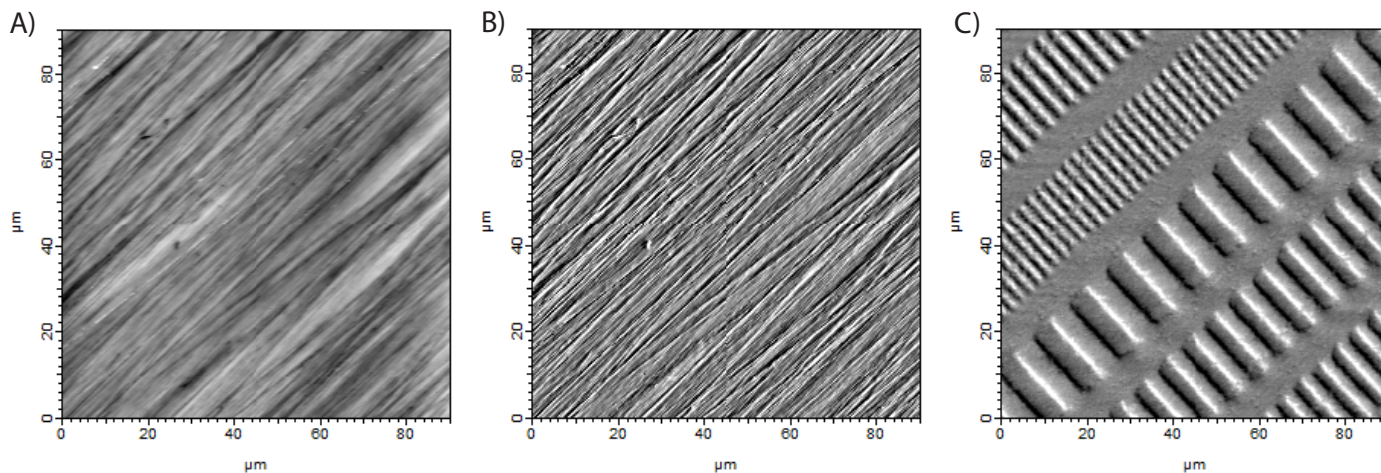


Figure 5: Outline of the lift mode principle. Magnetic information is recorded during the second pass (right panel). The constant height difference between the two scan lines is the lift height  $\Delta h$ .



**Figure 6: MFM images acquired on a hard drive. A) shows a scan of the topography (tapping mode); in B) another topography image is shown, acquired in tap and lift mode; C) magnetic image with topography subtracted (different sized bits are clearly visible).**

amplitude modulation (AM) mode and the data is recorded to one image. This height data is also used to move the tip at a constant local distance above the surface during the second scan, with feedback turned off. Theoretically, topographic contributions are thus eliminated from the second scan image.

Magnetic data can be recorded either as variations in amplitude, frequency or phase of the cantilever oscillation. It is argued that phase detection and frequency modulation give the best results, with higher signal-to-noise ratio.

#### Examples of new areas of research

The added functionality of the new MFM will greatly improve our abilities in magnetic imaging at room temperatures and will allow us for the first time to systematically study changes in magnetic structures within single and interacting grains at elevated temperatures. This will provide new insights into the acquisition, retention, and alteration of magnetization in nature and will open the way for new and exciting rock magnetic and paleomagnetic discoveries such as the origin and stability of TRM and VRM in MD grains.

Fundamental studies of magnetization at high temperatures (300-800 K) are important, because it is within this temperature range that blocking and acquisition of TRM and VRM occur in nature. Despite many fundamental advances in rock magnetism, the detailed response of micromagnetic structures with respect to time (VRM), temperature (TRM), and chemical change (CRM) and their connections to bulk behavior remain poorly understood for MD materials. For example, major efforts are now underway in retrieving high-resolution, reliable absolute paleointensity data from Precambrian rocks, extraterrestrial materials, and other igneous and archeological materials. It is therefore timely to renew investigations into the temperature and field response of micromagnetic structures most responsible for natural remanent magnetization.

The following is a brief summary of some of the research that can be accomplished with a HT-HFM.

**Thermoremanent Magnetization and Paleointensity:** The maximum temperatures available so far for MFM imaging

(400° C) allow imaging the acquisition of thermoremanent magnetization and the complementary process of thermal demagnetization in many natural samples containing common magnetic minerals such as titanomagnetite, titanohematite, and pyrrhotite. For magnetite-bearing rocks, HT-MFM would still allow to better understand the partial thermal remanent magnetization (pTRM) steps that are used in nearly all Thellier-based paleointensity techniques. One of the major causes of failed paleointensity experiments is associated with samples containing a significant fraction of pseudosingle domain grains that exhibit strong pTRM tails during the acquisition of laboratory-induced TRMs. The HT-MFM would allow the study of domain states in such samples, in order to better understand the origin of these pTRM tails.

#### Thermoremanent Magnetization and Cation Ordering:

Very recent work (Bowles et al., 2012) has shown that thermal annealing at relatively low temperatures (350-450° C) can affect the distribution of ferrous and ferric cations in the octahedral and tetrahedral sites of the inverse-spinel structure of intermediate-composition natural titanomagnetites, and thereby cause very significant changes in the Curie temperature without any change in chemical composition or crystallographic structure. Unlike the well-documented similar phenomenon in magnetoferrite (e.g. Harrison and Putnis, 1999a; b), where significant cation reordering occurs only at temperatures hundreds of degrees above the Curie point, these titanomagnetites reorder at temperatures near and even below  $T_c$ . This raises very interesting and important questions about how TRM is acquired or demagnetized in a material for which the Curie temperature is not a constant but is itself a function of thermal history.

The ability of the new HT MFM to operate in said temperature range enables to observe the domain configurations associated with different remanent states and different states of cation order, and to gain insight into the micromagnetic processes that operate as the states evolve.

**Chemical Remanent Magnetization:** CRM is one of the most common forms of remagnetization in nature. Although there have been many experimental studies

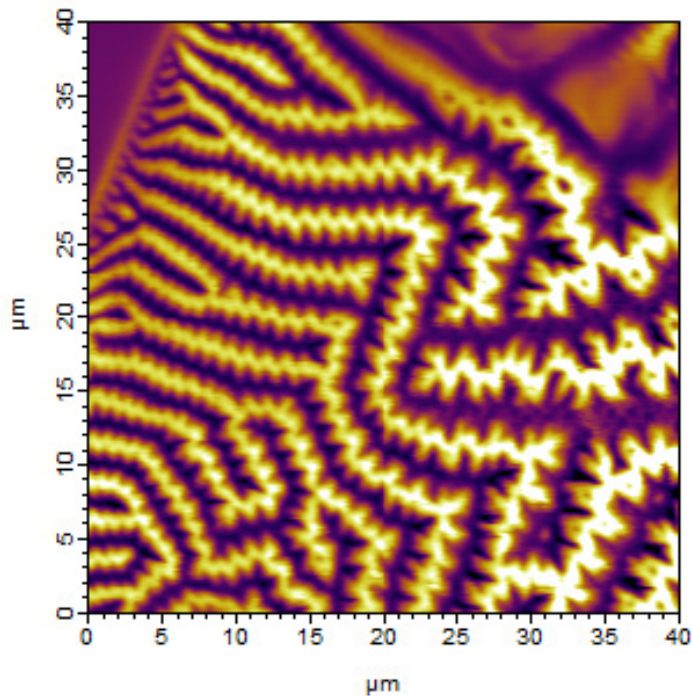


Fig 7. MFM magnetic image of a Mount St. Helens Ti-Magnetite in bright color contrasts.

on CRM in different natural chemical systems, laboratory studies of CRM under controlled conditions are still some of the most challenging experiments in rock magnetism (e.g. Dunlop and Özdemir, 1997). With the exception of the detailed work on hematite-ilmenite intergrowths and lamellar magnetism (McEnroe et al., 2001; Harrison et al., 2002; McEnroe et al., 2008), there has been little work on the direct imaging of micro-magnetic structures carrying CRM in other common magnetic minerals (Harrison et al., 2002; Krása et al., 2005) and none at high temperatures where many of these processes occur. It would be possible to use the high-temperature MFM to directly observe the magnetic consequences of incipient oxidation and reduction with controlled time-temperature imaging. These types of experiments will lead to a better understanding of alteration processes that commonly affect rocks in igneous and sedimentary environments as well as during typical laboratory heating protocols for paleointensity measurements.

#### First Images

Figures 1, 7 and 8 show some of the first results of measurements with the new instrument. All of these are room-T images of homogeneous Al- and Mg- substituted titanomagnetites (~TM30) from Mt. St. Helens (Bowles et al., 2013).

The crystallographic orientations of the imaged particles are unknown, and the high contrast indicates that magnetization is generally oriented at high angles to the surface. Wavy boundaries between the light and dark regions resemble the wall patterns observed in TM60 by Moskowitz et al. (1988).

We are still working on the techniques of high-T specimen preparation and measurement, and "hot" images will be forthcoming!

#### Acknowledgements

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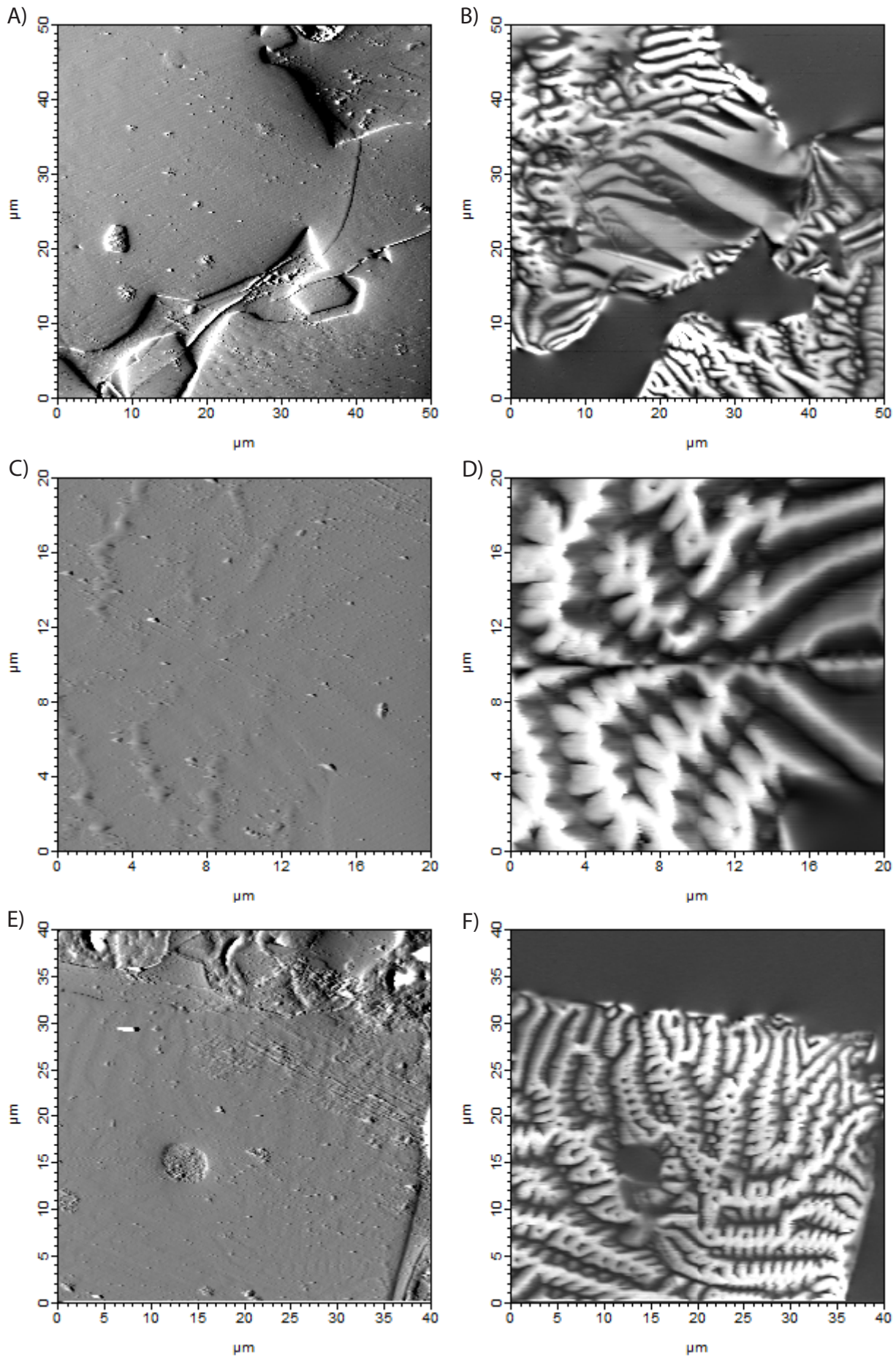


Fig 8. MFM topographic (tap and lift mode) and magnetic images of Mount St. Helens Ti-magnetite crystals.

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