
Large n multi-vortices and multi-monopoles

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Introduction

Topological Solitons in a Gauge-Higgs theory:

$$\mathcal{L} = -FF - |D\phi|^2 - V(|\phi|)$$

Vortex

$$\pi_1(U(1)) = \mathbb{Z}$$

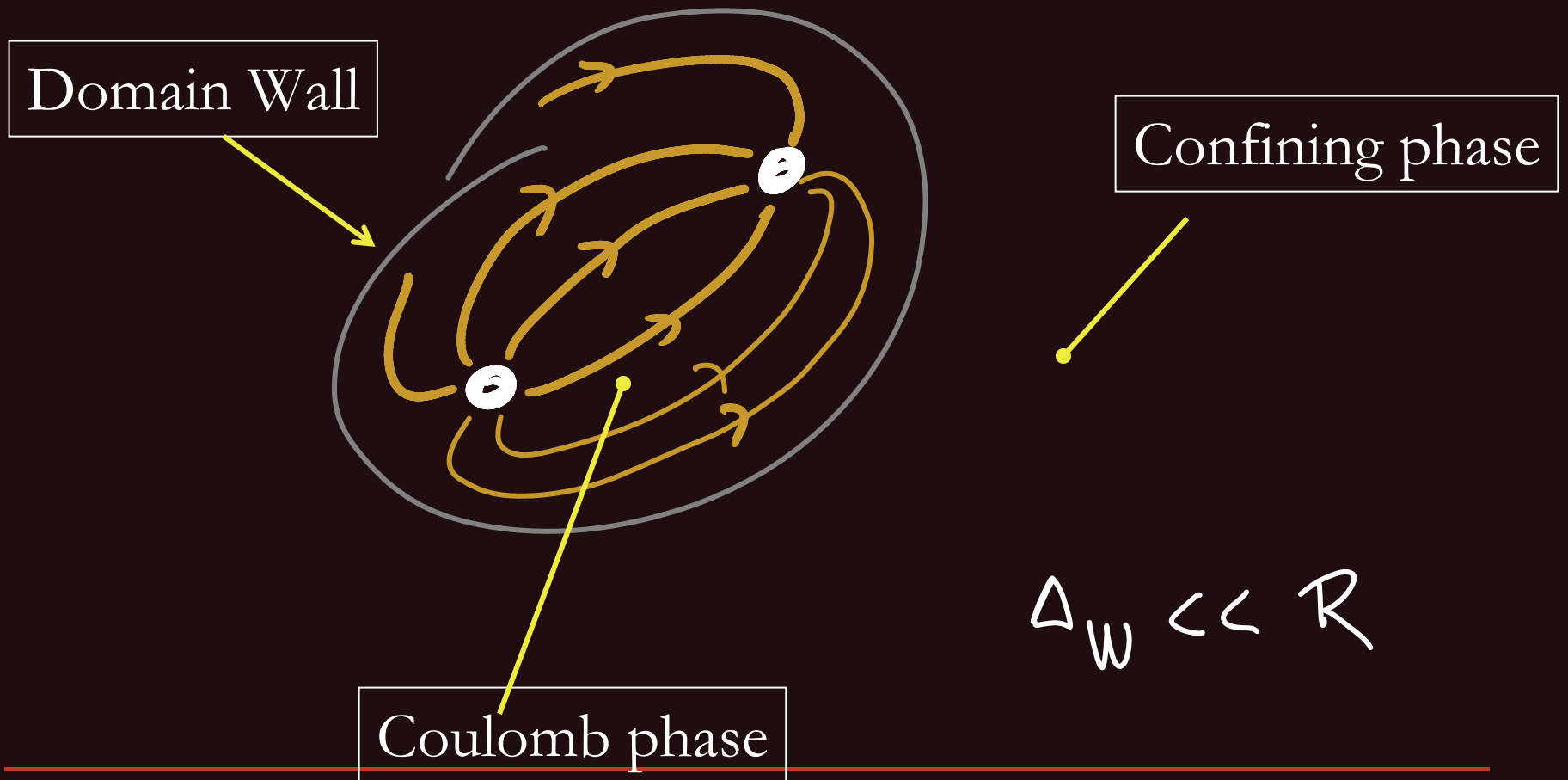
Monopole

$$\pi_2\left(\frac{SU(2)}{U(1)}\right) = \mathbb{Z}$$

Large n



Bag Models of Hadrons



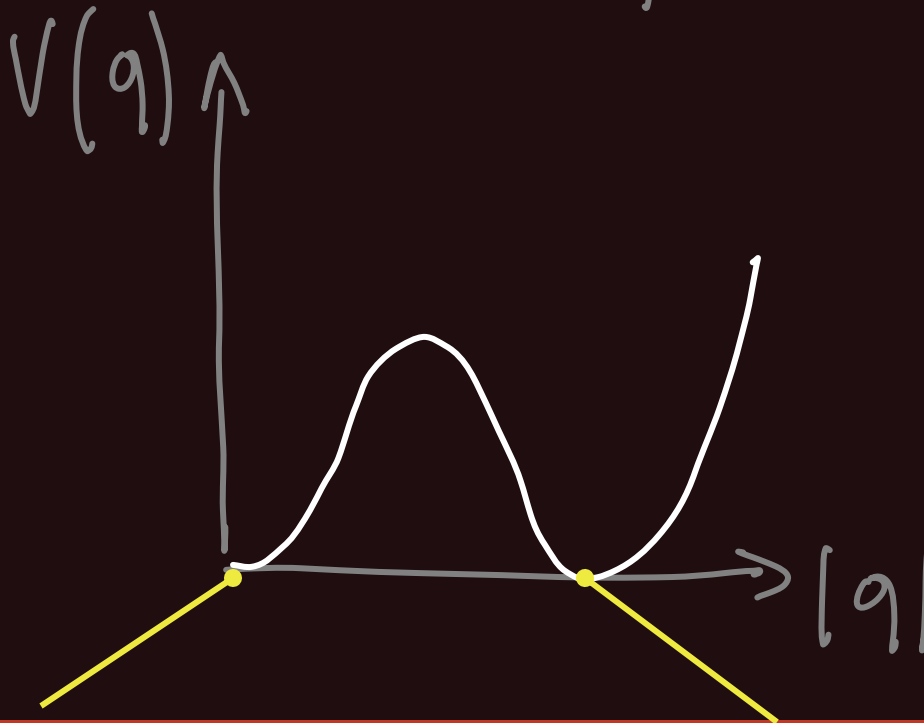
Genesis of the idea

$$N=2 \xrightarrow{W(\Phi)} N=1$$

Many discrete vacua survive

$$U(N) \rightarrow U(1)^N$$

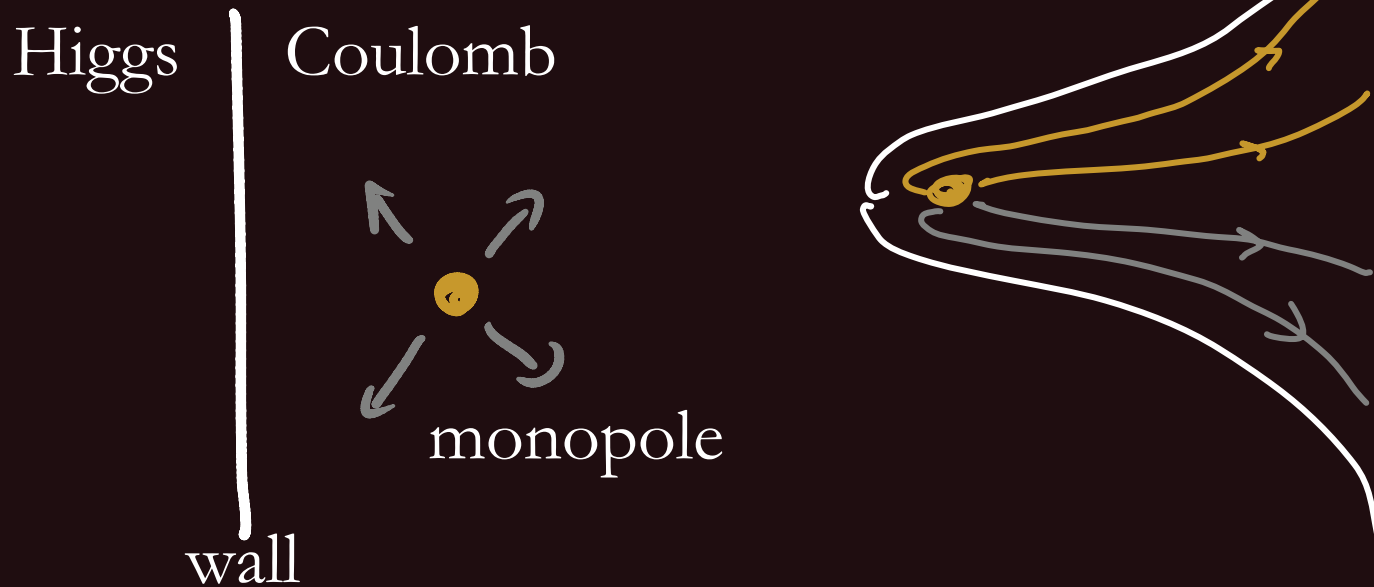
Toy model



Coulomb phase

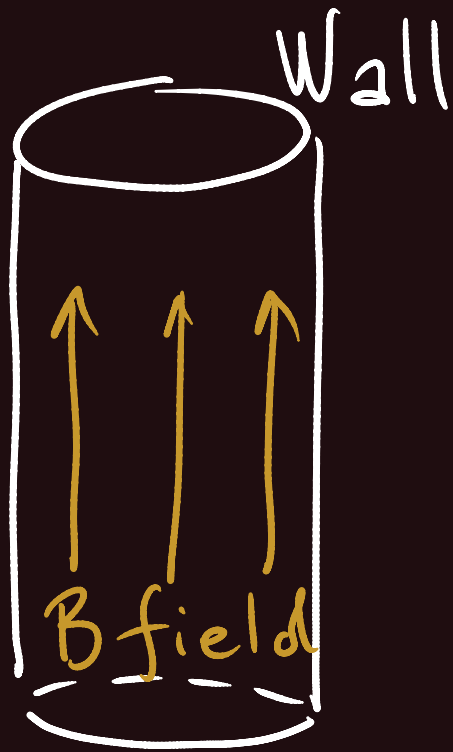
Higgs phase

Vortex and Wall are “made of the same stuff”



Question:

Is there a limit in which this
equivalence is exact?



$$T(R) = \frac{\Phi_B^2}{2\pi R^2} + T_W 2\pi R$$

Repulsion

Attraction

$$R_V \sim \frac{n^{2/3}}{T_W^{1/3}}$$

$$T_V \sim T_W^{2/3} n^{2/3}$$

Softly Broken N=2

$$W = \tilde{q} \bar{\Phi} Q - m \tilde{q} Q + \mu \bar{\Phi}^2$$

Color-Flavor locked phase

$$\phi = m$$

$$\tilde{q} q = \mu m$$

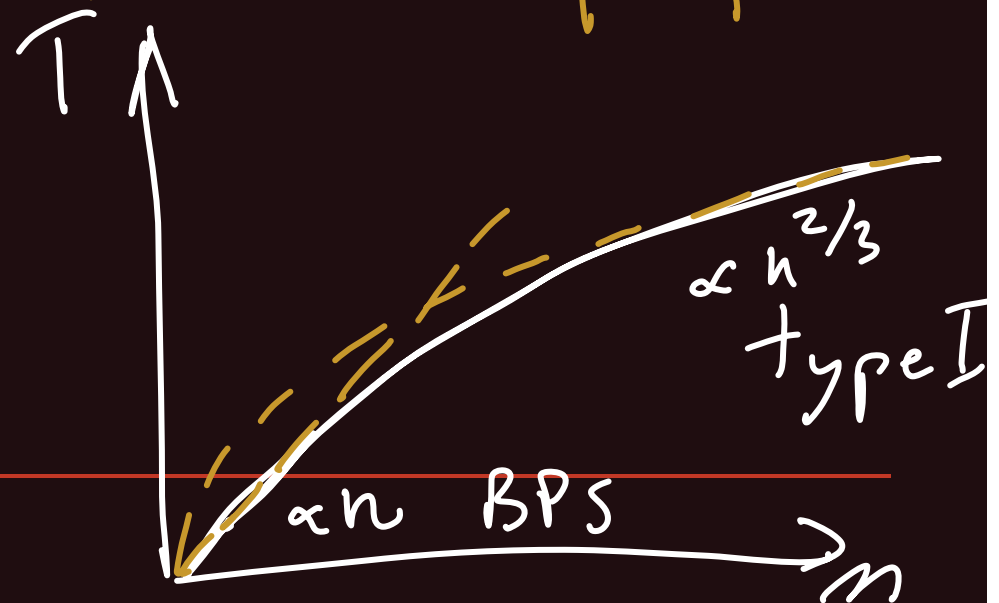
Coulomb phase

$$\phi = 0$$

$$q = \tilde{q} = 0$$

$$\frac{e^2 \mu}{m} \ll 1$$

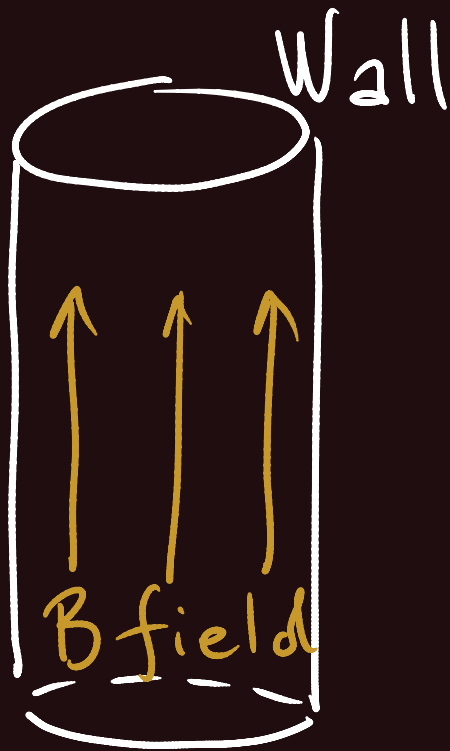
Almost BPS



And now the conjecture...

BPS potential

$$V(\varrho) = \frac{e^2}{2} \left(|\varrho|^2 - \nu \right)$$



$$T(R) = \frac{2\pi n^2}{e^2 R^2} + \frac{e^2 \nu^2}{2} \pi R^2$$

$$T_V = 2\pi n \nu$$

It saturates exactly the BPS bound!

Multi-vortex ansatz

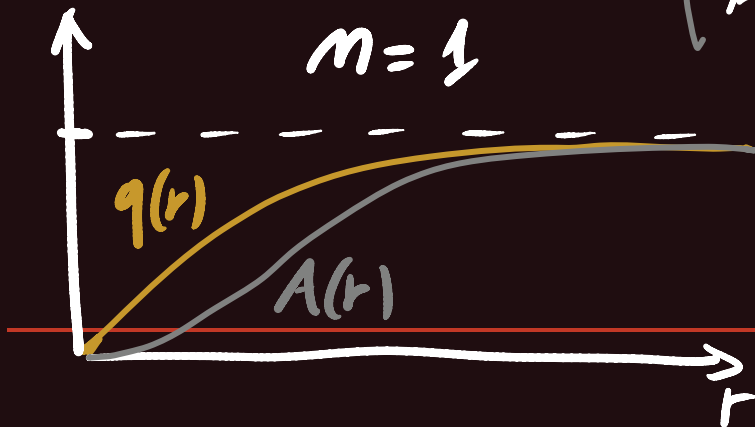
$$\varphi = e^{in\vartheta} q(r)$$

$$A_\vartheta = \frac{n}{r} A(r)$$

Profile functions

Differential
Equations:

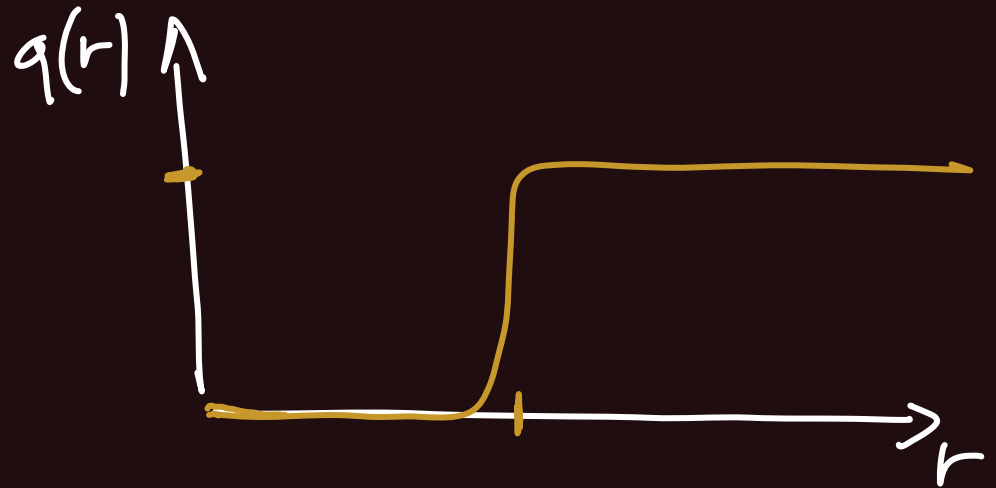
$$\begin{cases} q'' + \frac{1}{r} q' - n^2 \frac{(1-A)^2}{r^2} q - \frac{\delta V}{\delta q} = 0 \\ A'' - \frac{1}{r} A' + e^2 (1-A) q^2 = 0 \end{cases}$$



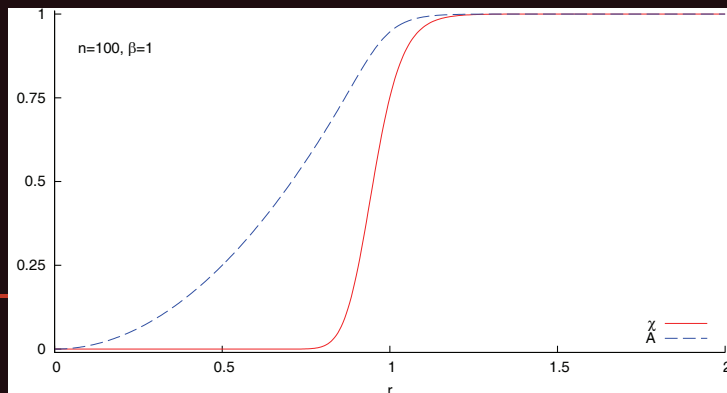
Numerical Proof

Length Rescaling

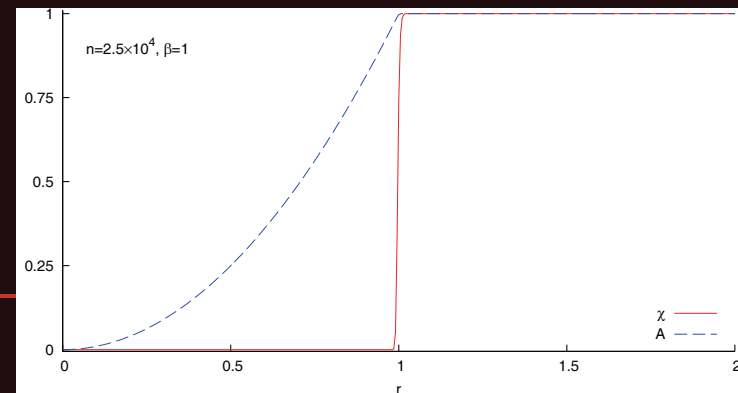
$$r \rightarrow \frac{r}{n}$$



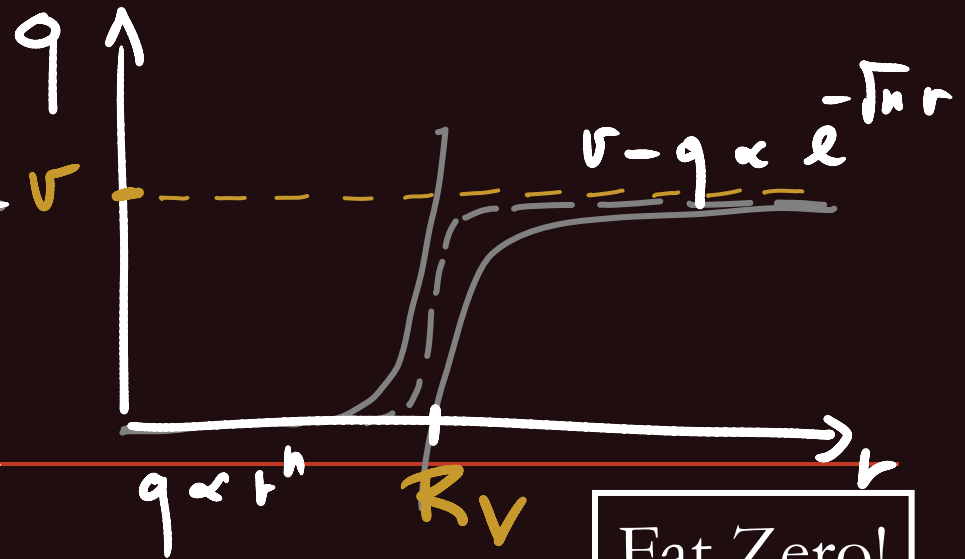
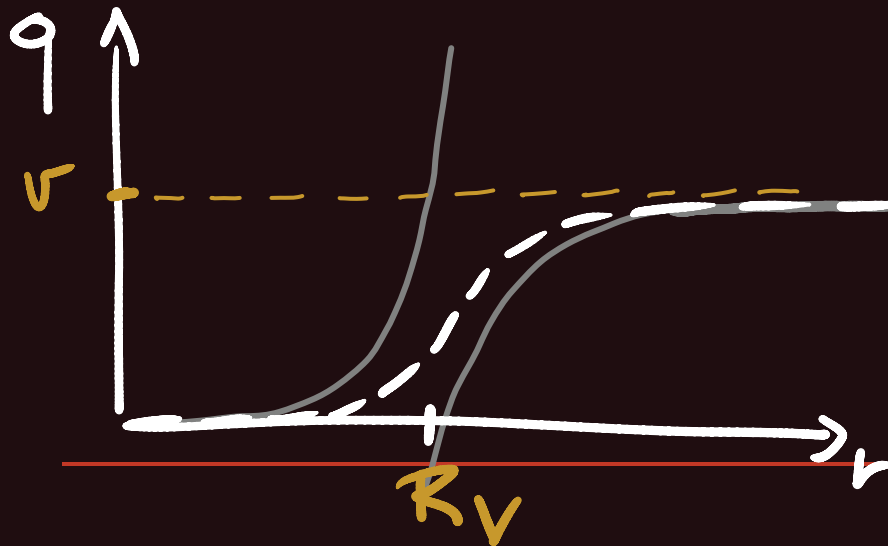
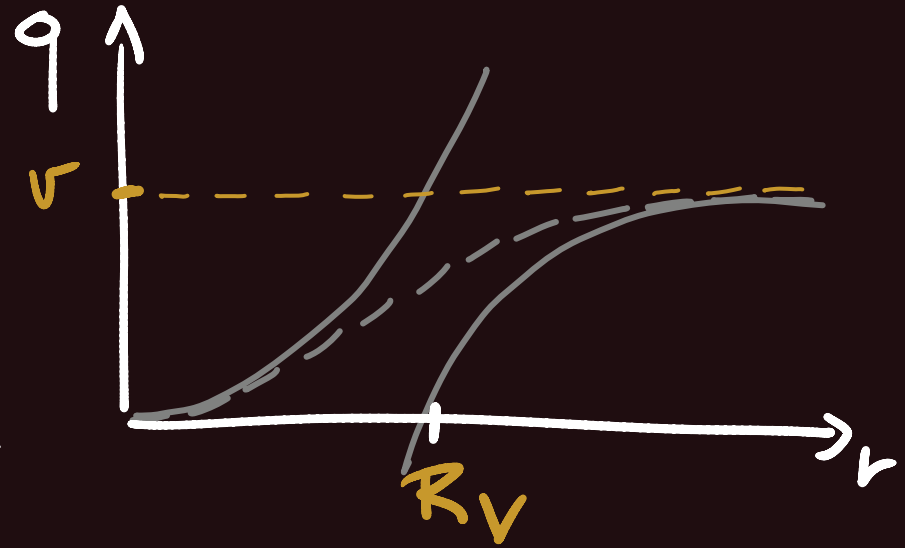
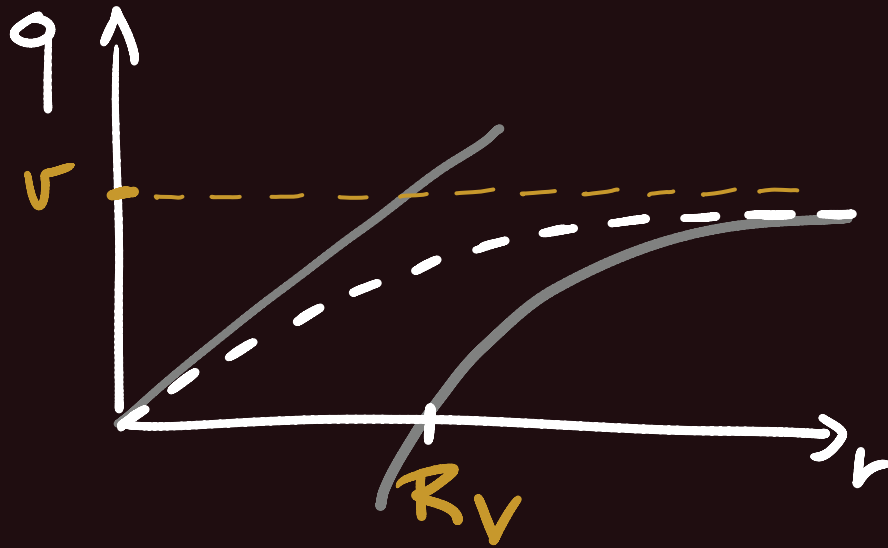
BPS $n=100$



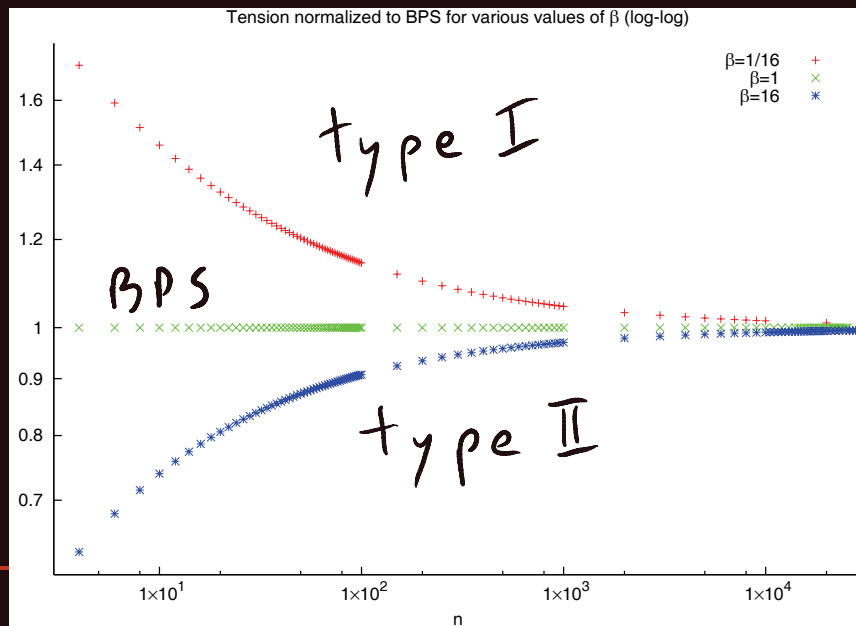
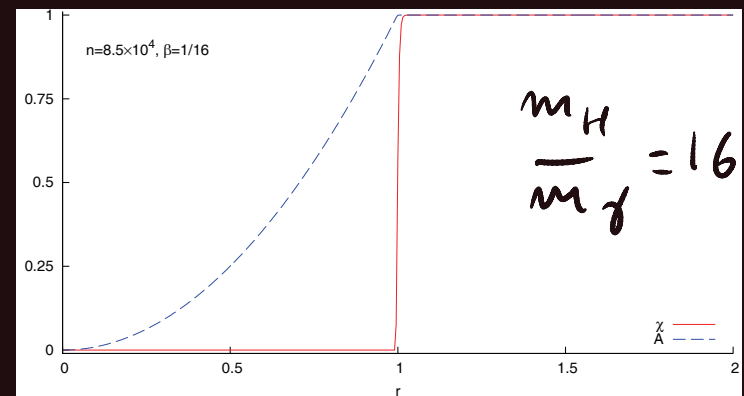
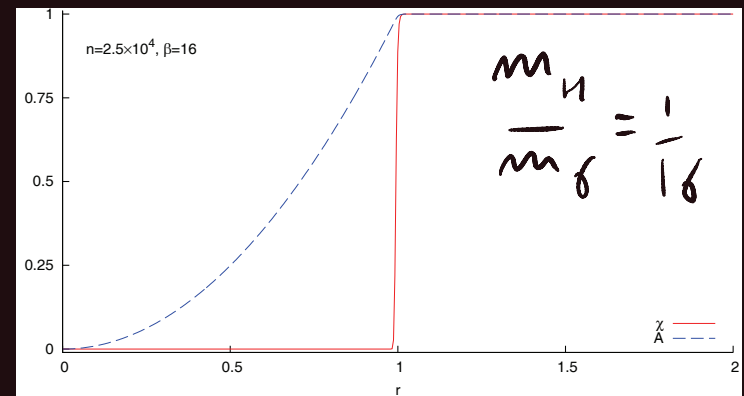
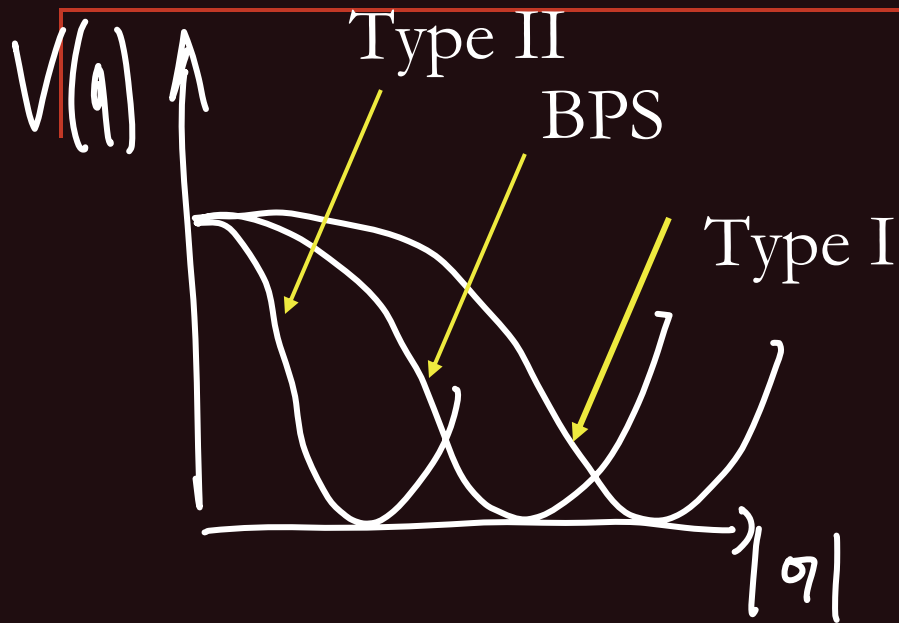
BPS $n=25,000$



Why a step function?



Fat Zero!



It works for every potential!

with S.B.Gudnason

Another example...



Confining Strings

$SU(N)$



$T(\kappa, N)$

Tension of the k_string

$$R(\kappa, N) = \frac{T(\kappa, N)}{T(1, N)}$$

Casimir formula

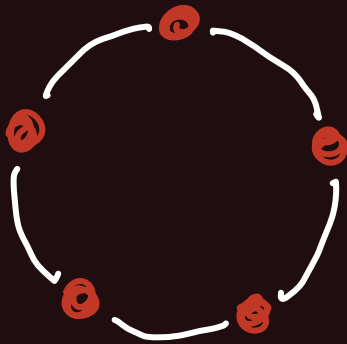
Sine formula

Assumption: n strings are vortices of a “dual” $U(1)$

Center of $SU(N)$ $\cong \mathbb{Z}_N$

$$\begin{pmatrix} e^{i \frac{2\pi k}{N}} & & \\ & \ddots & \\ & & e^{i \frac{2\pi k}{N}} \end{pmatrix}$$

Interpolating $U(1)$

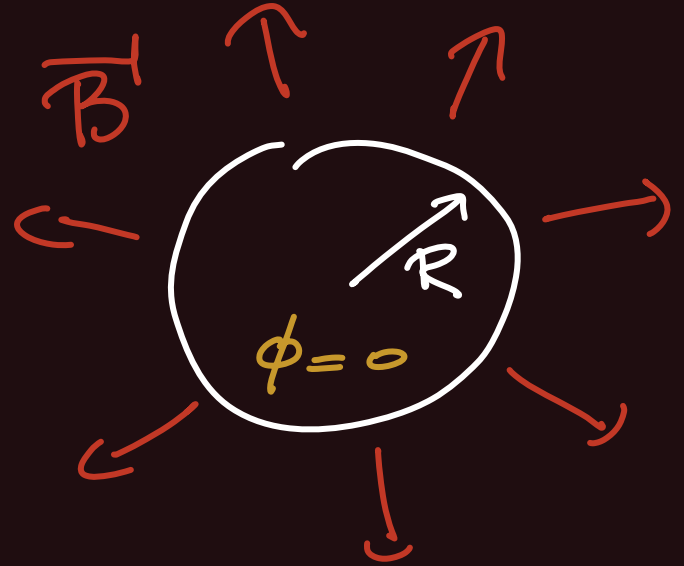


$$\begin{pmatrix} 1 & & \\ & \ddots & \\ & & 1 \end{pmatrix} - (N-1)$$

Unique up to gauge invariance

Magnetic Bag

$$E(B) = \int_R^\infty \frac{B^2}{2} = \frac{2\pi n^2}{R}$$



$$E(\phi) = \int_R^\infty \frac{|\nabla\phi|^2}{2} = 2\pi v^2 R$$

$$M(R) = \frac{2\pi n^2}{R} + 2\pi v^2 R$$

The minimization gives:

$$\left\{ \begin{array}{l} R_m = \frac{n}{v} \\ M_m = 4\pi n v \end{array} \right.$$

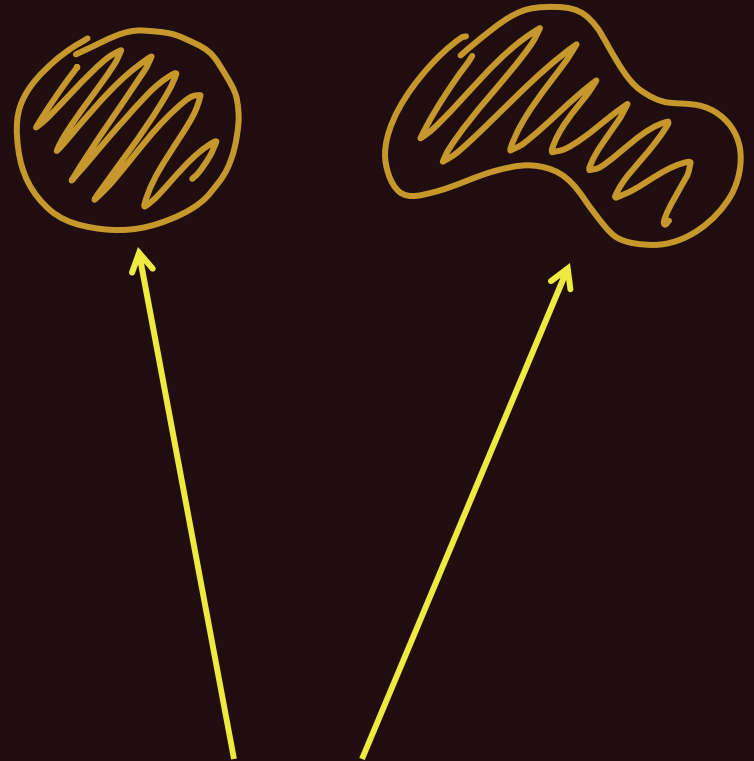
It saturates exactly the BPS bound

Moduli Space

Multi-vortices

$$T(R) = \frac{\phi_B^2}{2\pi R^2} + \epsilon_0 \pi R^2$$

$$T(A) = \frac{\phi_B^2}{2A} + \epsilon_0 A$$

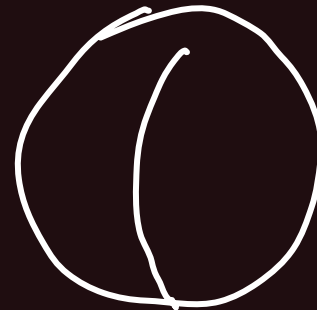


The moduli space is the closed surfaces with constant area!

Axial Symmetric multi-monopole

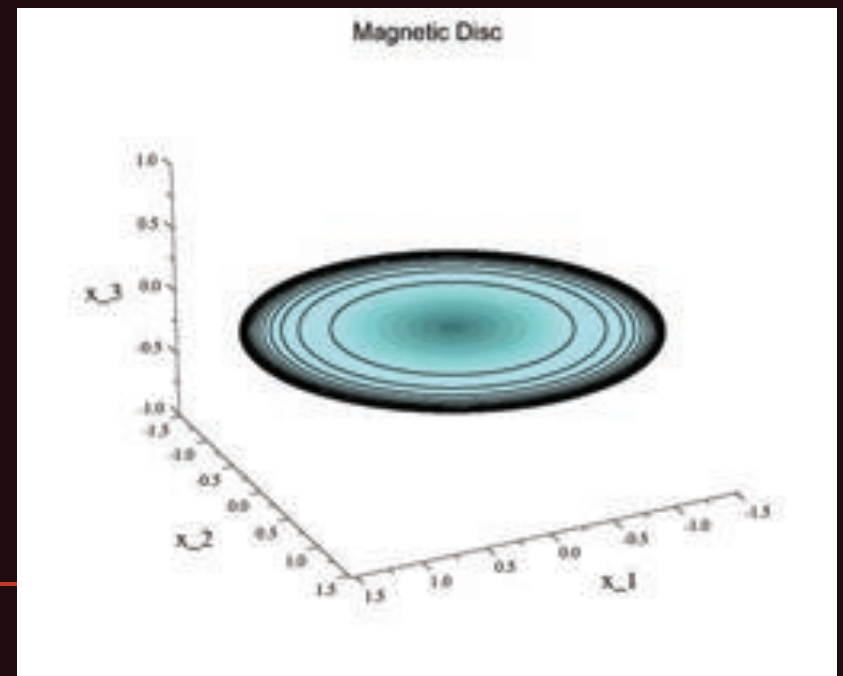


$$S^2 \quad |r| \rightarrow \infty$$



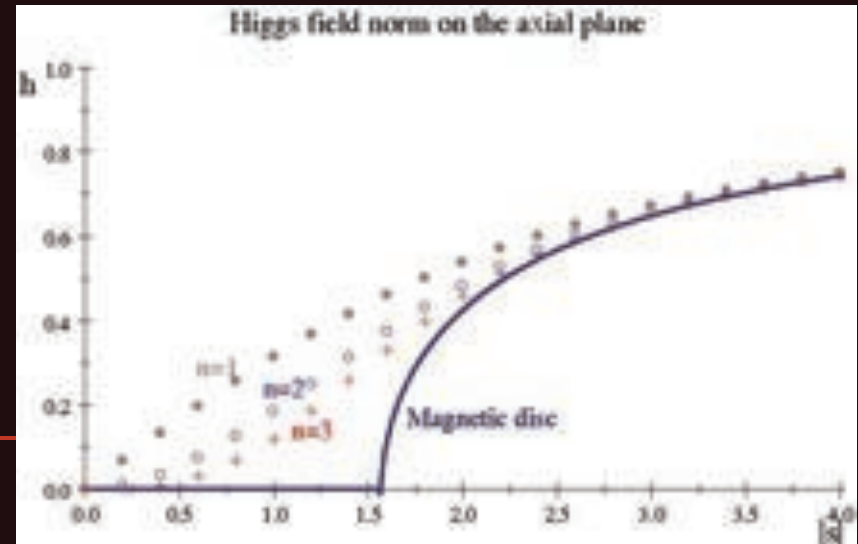
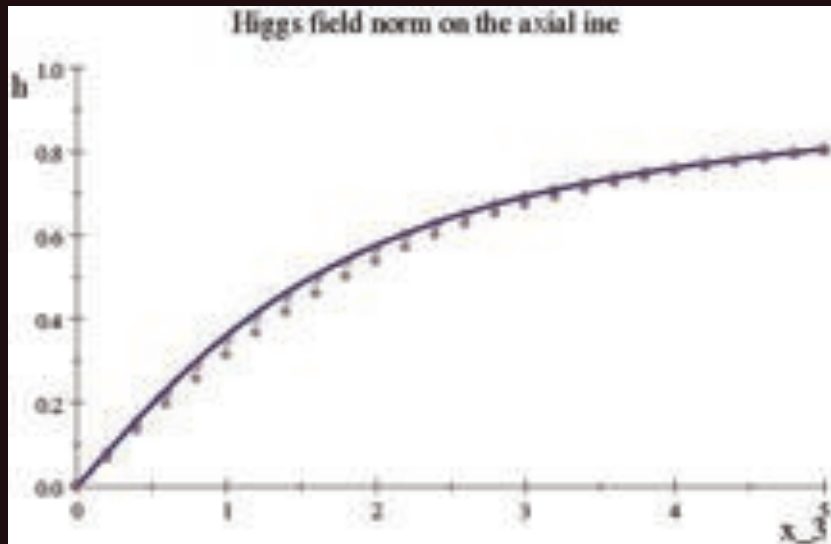
$$S^2 \quad |\phi| = v$$

$$\left\{ \begin{array}{l} \phi_3 \sim z \\ \phi_1 + i\phi_2 \sim (x + iy)^n \end{array} \right.$$

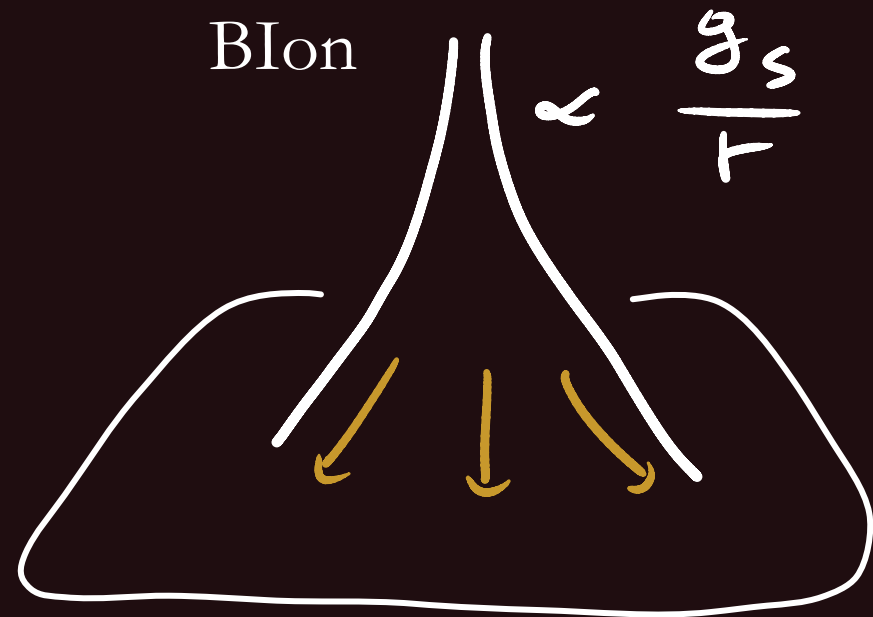
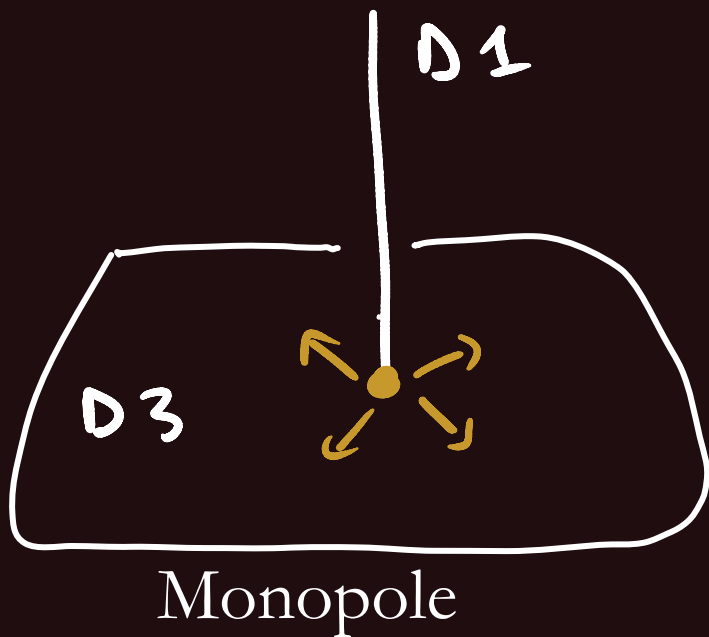


Ward-Prasad-Rossi VS Magnetic Disc

$$\phi(r, \theta) = \begin{cases} 1 - \frac{1}{r} \sum_{\ell} \frac{(-1)^{\ell}}{2\ell+1} \left(\frac{\pi}{2r}\right)^{2\ell} P_{2\ell}(\cos\theta) & r \geq \frac{\sqrt{u}}{2} \\ \frac{2}{\sqrt{u}} \sum_{\ell} \frac{(-1)^{\ell}}{2\ell+1} \left(\frac{2r}{\sqrt{u}}\right)^{2\ell+1} P_{2\ell+1}(\cos\theta) & r \leq \frac{\sqrt{u}}{2} \end{cases}$$

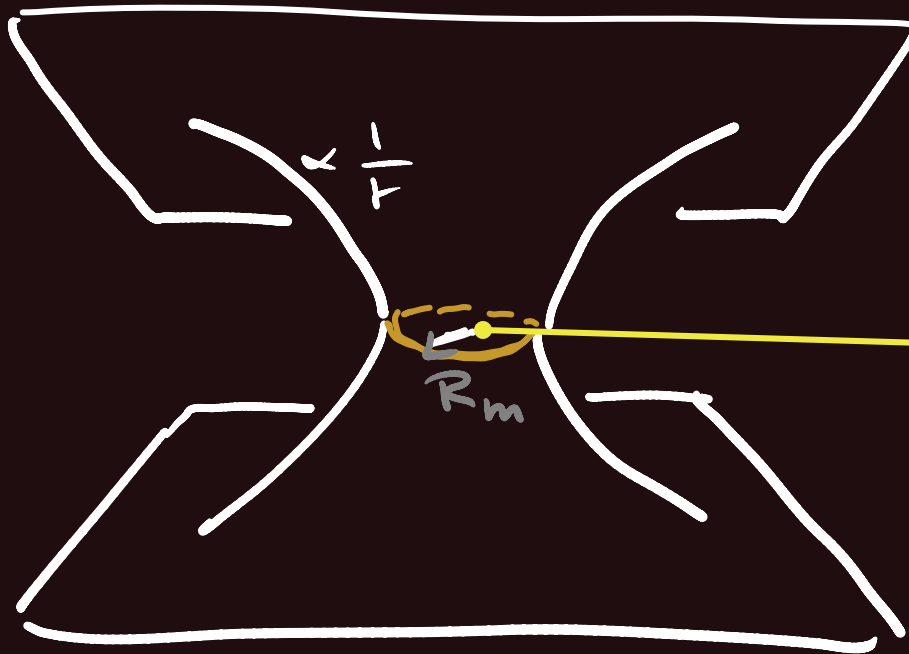


String Theory



The BIon is a solution of the abelian DBI action

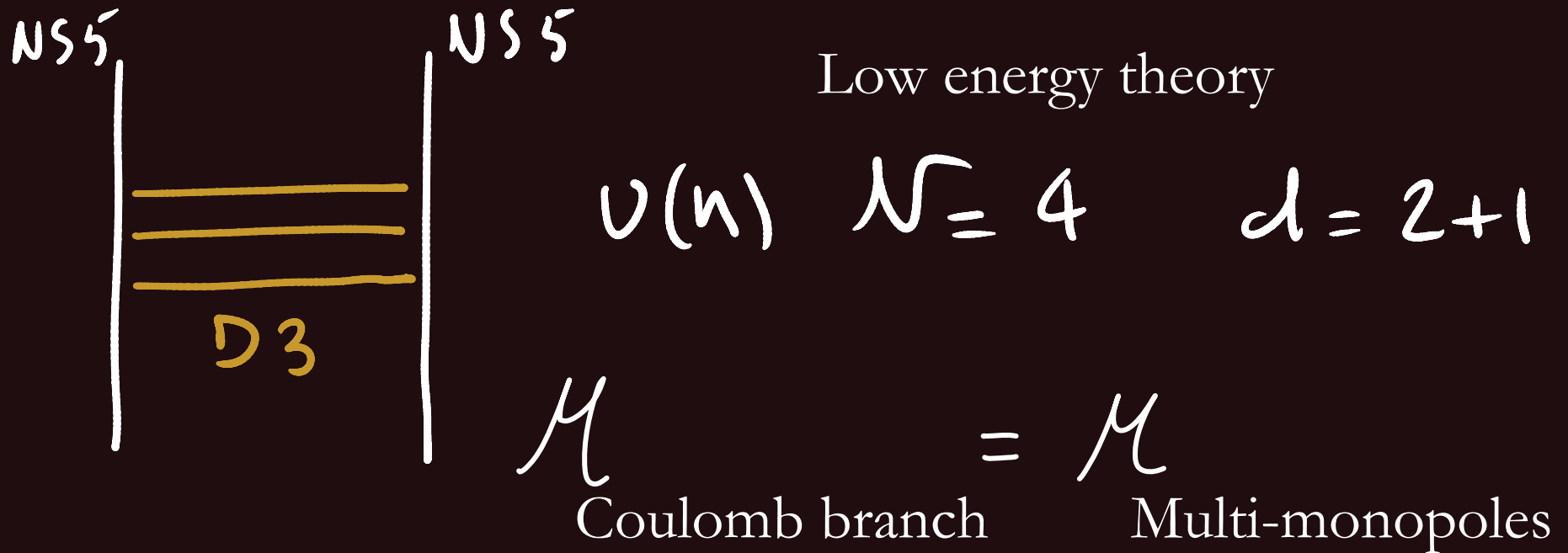
Magnetic bag in string theory



Inside the bag the D3 branes are overlapped

It is like two “cutted” abelian Bions

't Hooft and multi-monopole large n

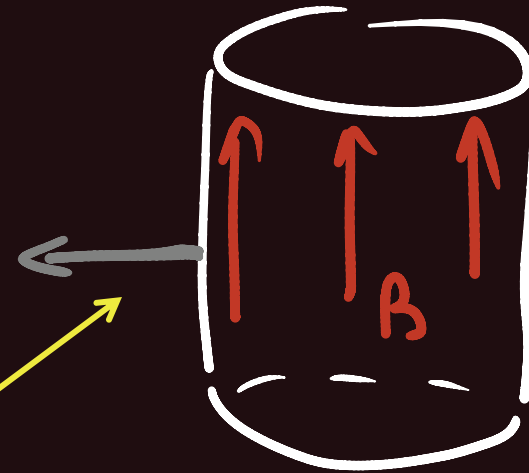


$$\frac{1}{g_{YM}^2} = \frac{5}{g_s}$$

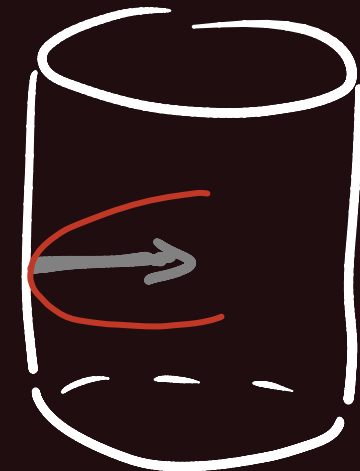
$$R_m = \frac{5}{5}$$

't Hooft rescaling = Length rescaling

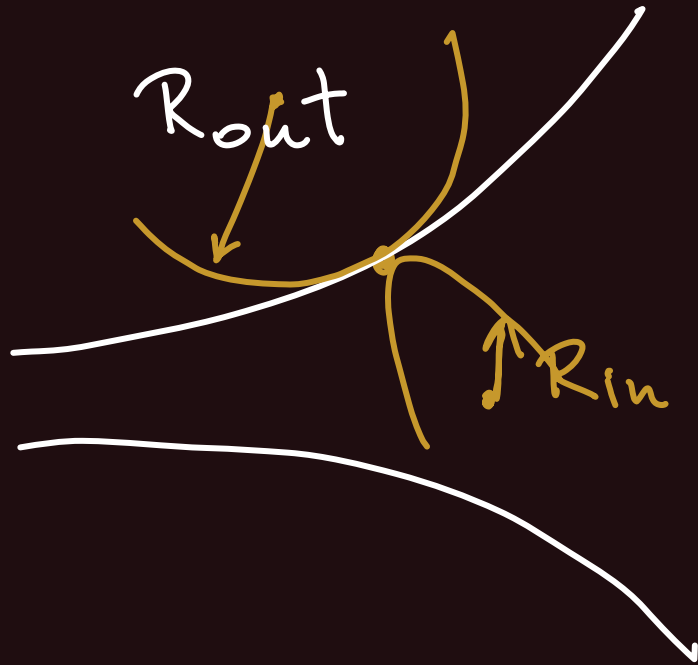
Soliton Junctions



$$\frac{T'}{2\pi R} = -\frac{B^2}{2} + \frac{T_w}{R}$$



Master Equation



$$\vec{B} = \vec{\nabla} \varphi$$

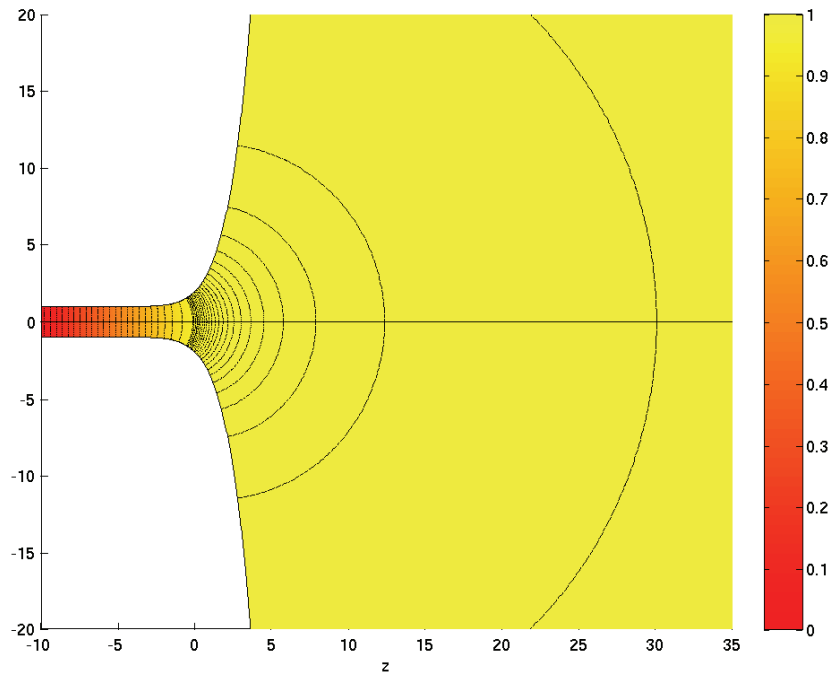
$$\nabla^2 \varphi = 0$$

+ Neuman

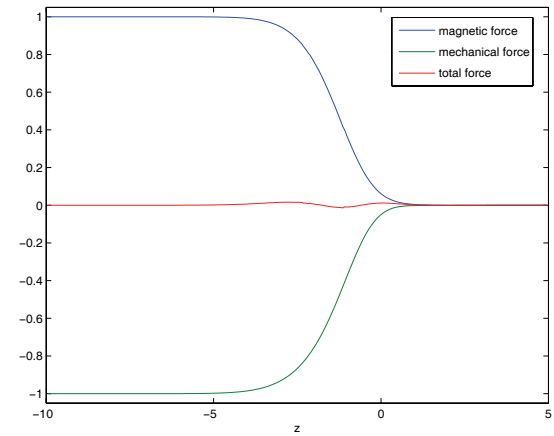
$$-\frac{B^2}{2} + T_W \left(\frac{1}{R_{in}} - \frac{1}{R_{out}} \right) = 0$$

Numerical Analysis

The Junction



Forces



Conclusions

- Large n limit of topological solitons
- Bag models description
- Simple and General

- Chern-Simons vortices (FQHE)
- Large n limit of textures (Skyrmions)
- ...