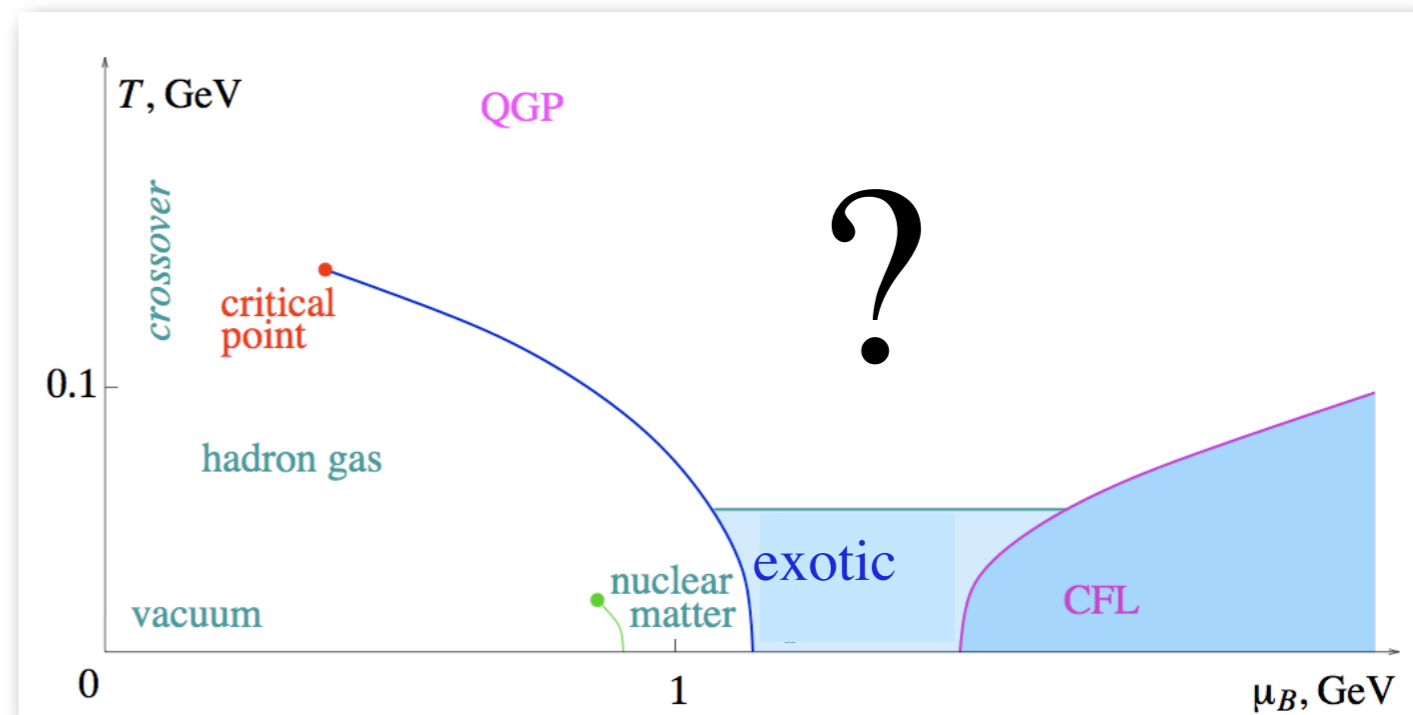


NOISE

A major challenge for physics: computing properties of strongly interacting dense matter from theory.

The obstacle: sign problems & noise in simulations.

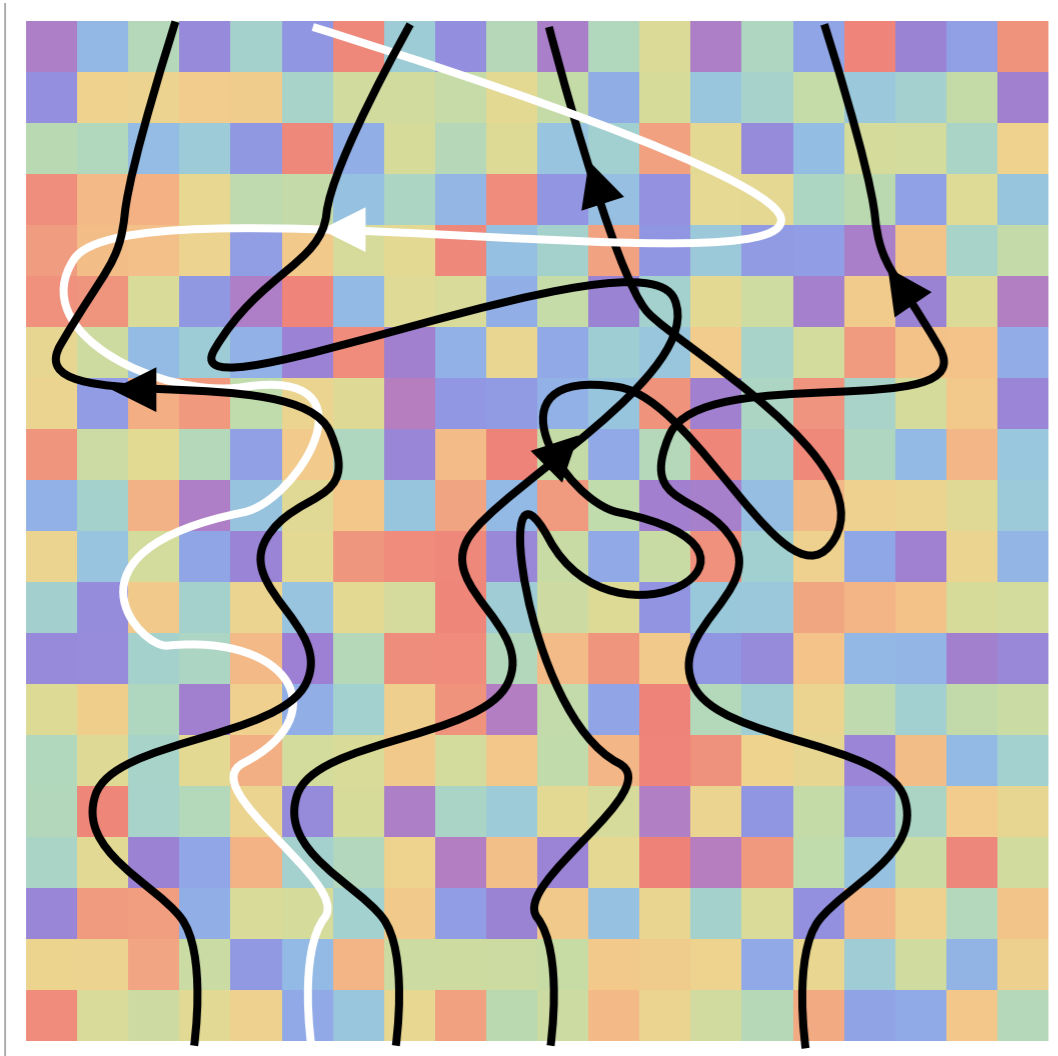
In QCD would like to map out phase diagram, but hit sign problem at nonzero μ_B ...



...or noise problems in canonical simulations

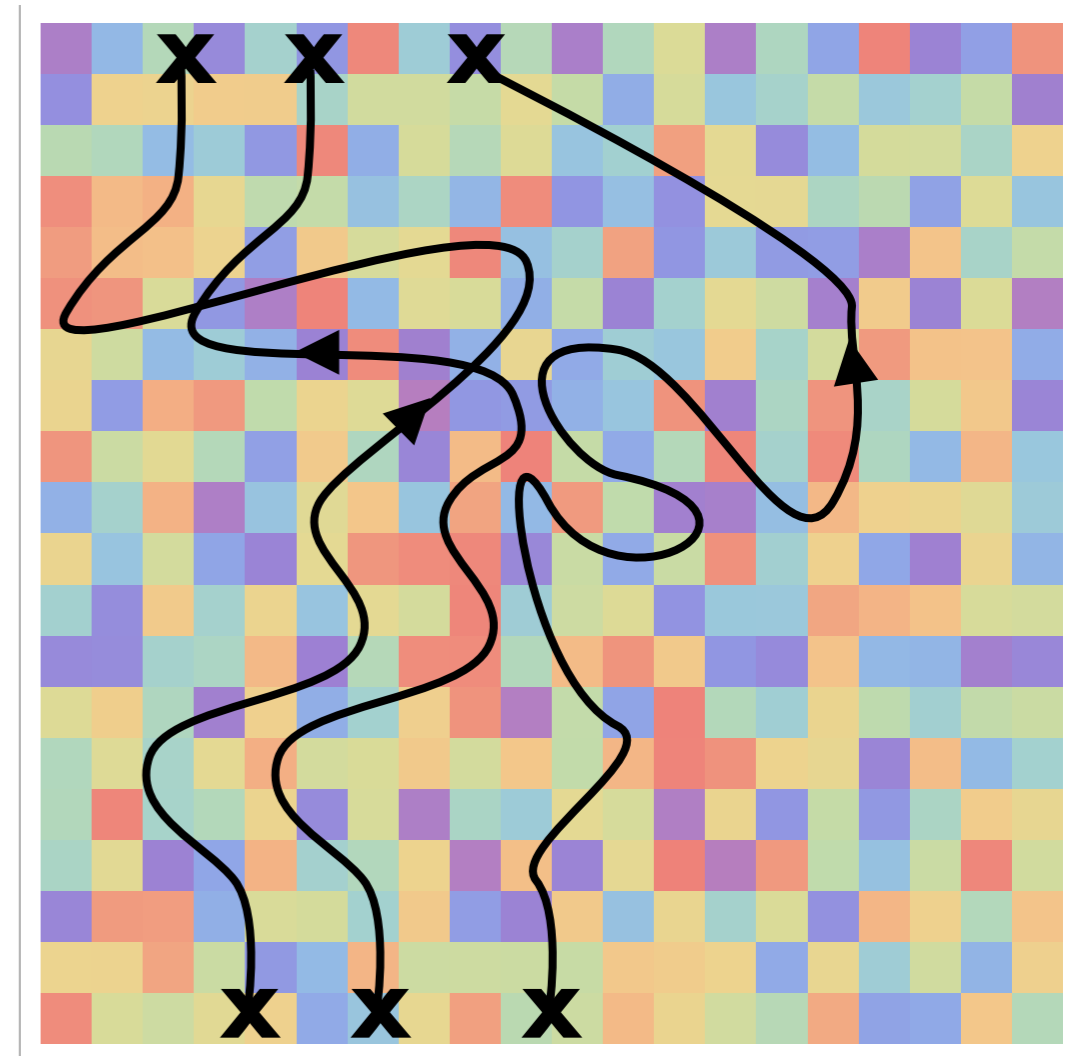
What is the sign/noise problem?

Sign Problem



Grand canonical: exponentially difficult in volume V

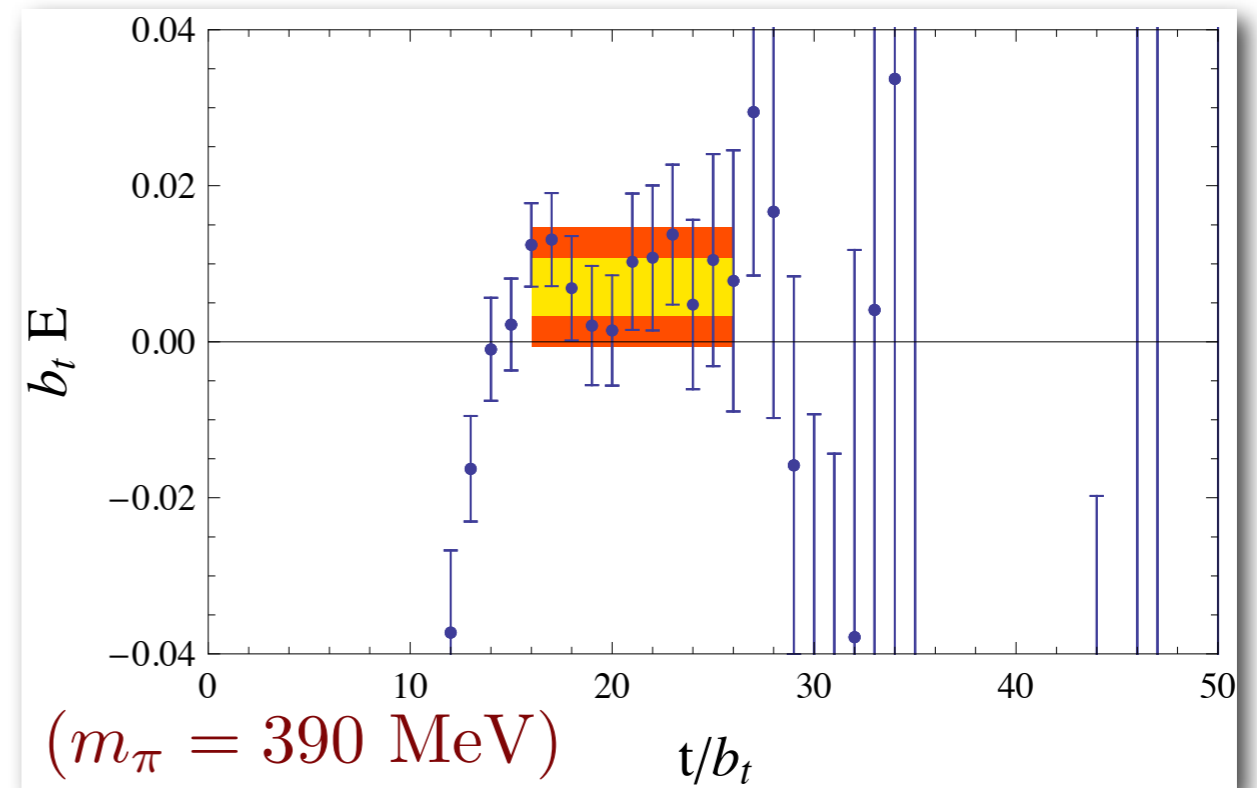
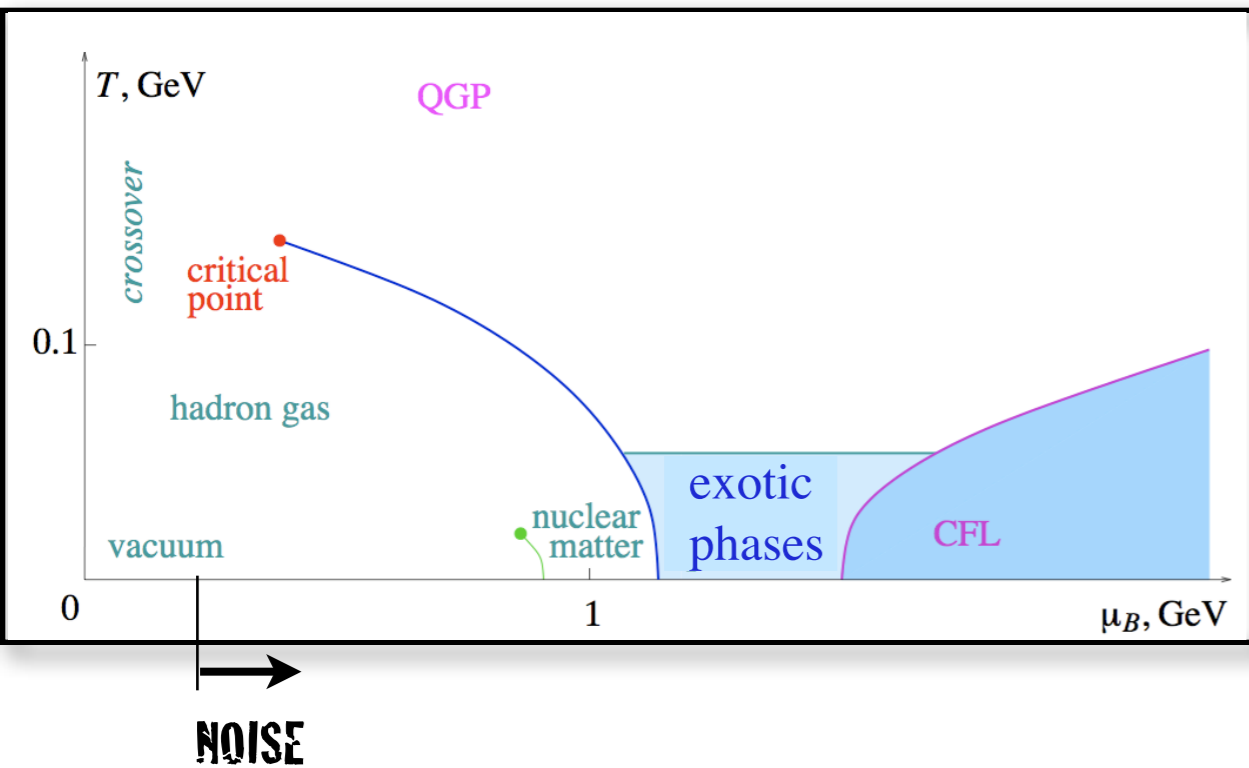
Signal/Noise Problem



Canonical: exponentially bad S/N in Euclidian time t

Triton B.E.

S. R. Beane et al. (NPLQCD),
Phys. Rev. D 80, 074501 (2009)



The obstacle:

Grand canonical - "fermion sign problem" hits you for $\mu_B > 3m_\pi/2$ ($T=0$)

Canonical (fixed B) - noise in nucleon correlator set by scale $3m_\pi/2$

This talk:

- What is the nature of noise in a noisy simulation?
- Could there be more effective ways to deal with it?
- Connections between noise, renormalization group and properties of disordered media

with:

MICHAEL ENDRES

JONG-WAN LEE

AMY NICHOLSON

S/N \sim sign problem:

In background gauge field, quarks don't know about each other.

Is the quark in a pion?

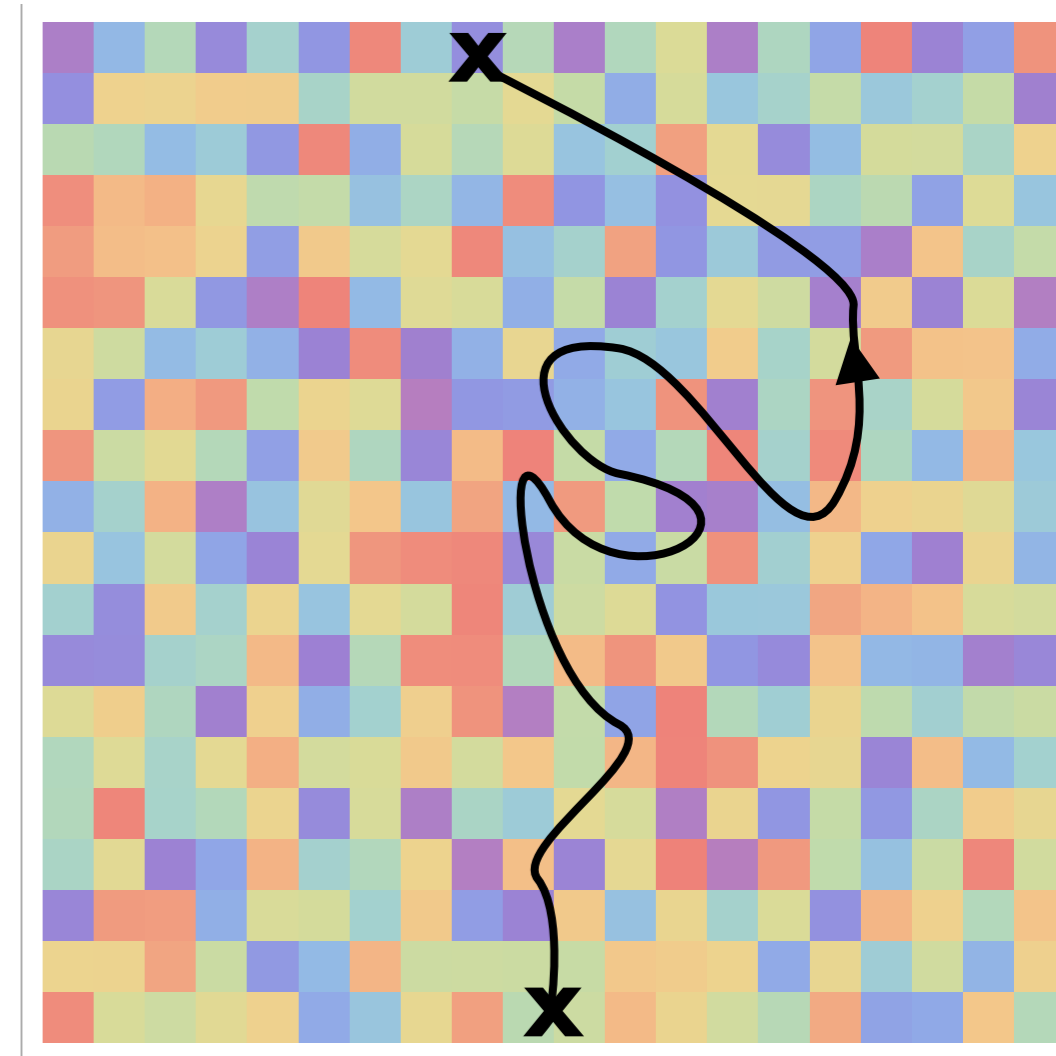
\Rightarrow big correlation after long time

Is the quark in a baryon?

\Rightarrow exponentially smaller correlation

(quarks "weigh more" in a baryon)

What is a quark to do??



S/N \sim sign problem:

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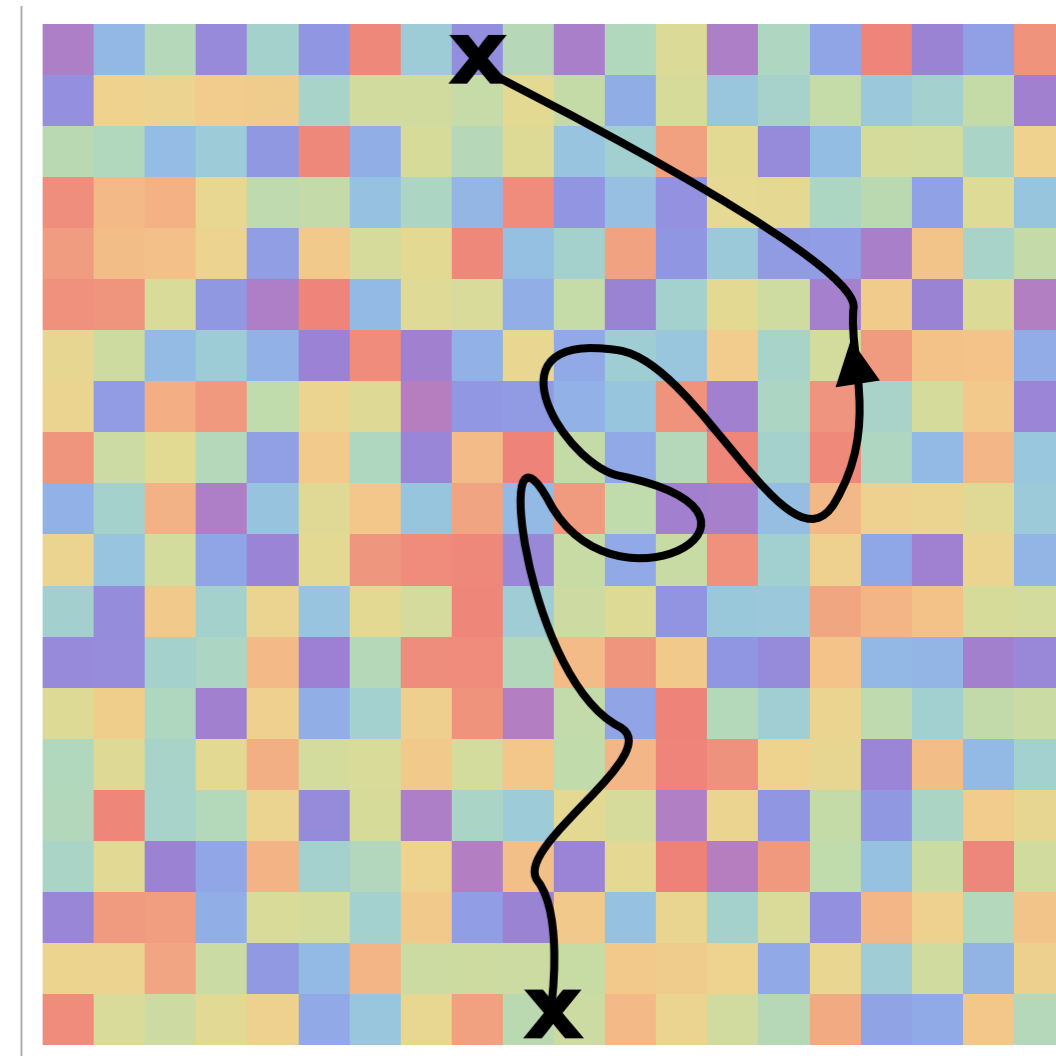
What is a quark to do??

How lattice QCD solves this problem:

- Every quark has long correlation in a background gauge field

$$“M_q” \sim M_\pi/2$$

- If the quark is in a baryon, EXPONENTIAL CANCELLATIONS when averaging gauge fields (since $M_B > 3 M_\pi/2$)

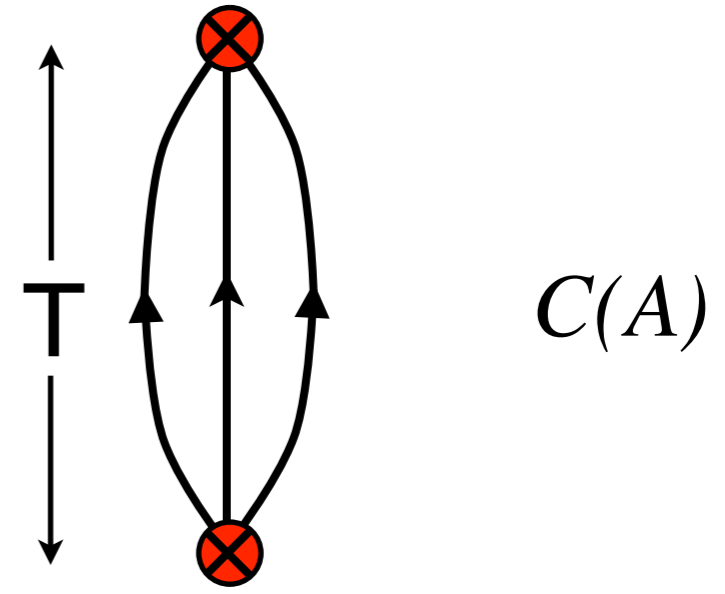


Lepage argument for the signal/noise problem

e.g: measuring the nucleon mass in LQCD:

$$\langle C \rangle = \frac{1}{N} \sum_{\{A\}} C(A) \propto e^{-M T} + \dots$$

nucleon:
lightest 3q state

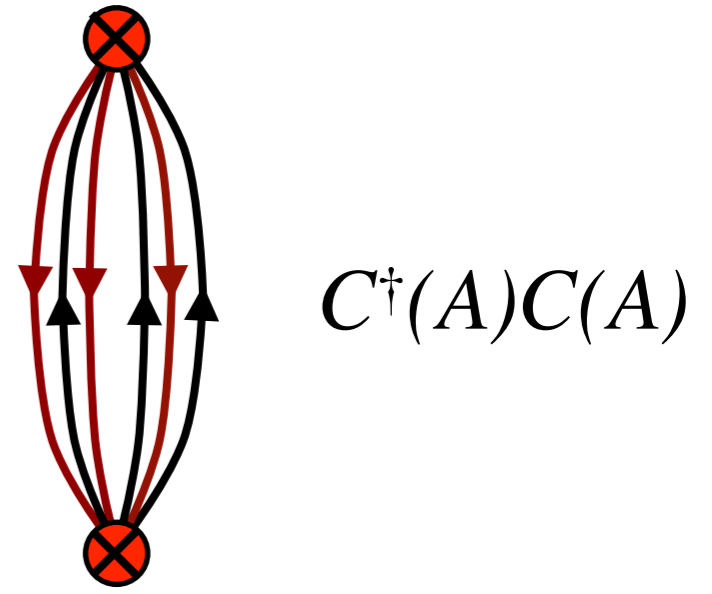


Dispersion in measurement:

$$\sigma^2 = \langle C^\dagger C \rangle = \frac{1}{N} \sum_{\{A\}} C^\dagger(A) C(A) \propto e^{-3m_\pi T} + \dots$$

3 anti-quarks
3 quarks

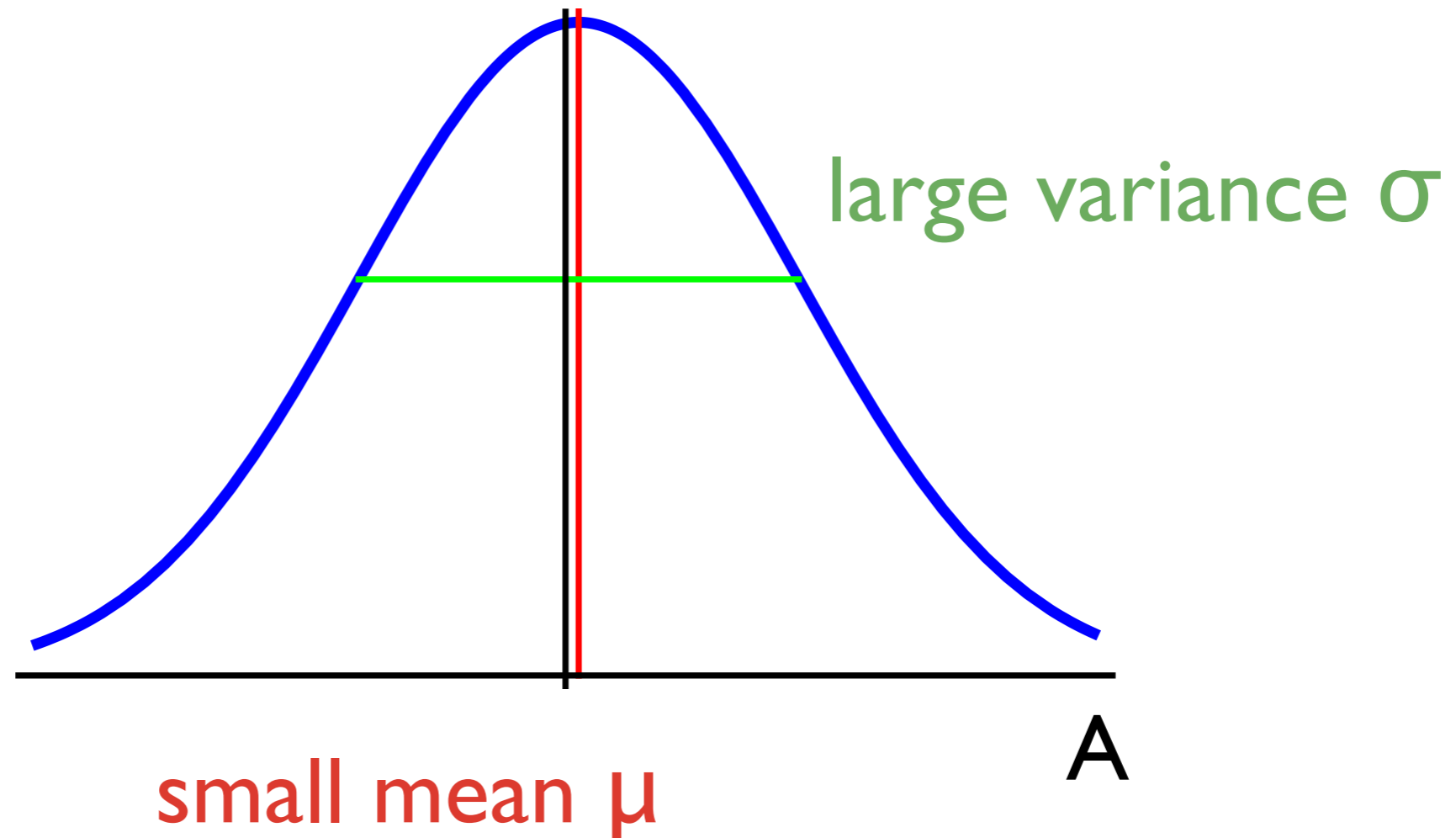
3π : lightest $3q + 3q^*$ state



$$\frac{\text{signal}}{\text{noise}} \propto \frac{1}{\sqrt{N}} e^{-(M - \frac{3}{2}m_\pi)T}$$

Basically same picture, with caveats:
"N"? Gaussian distribution?

Lepage picture for distribution of correlator for large baryon number, $C_B(t, A)$



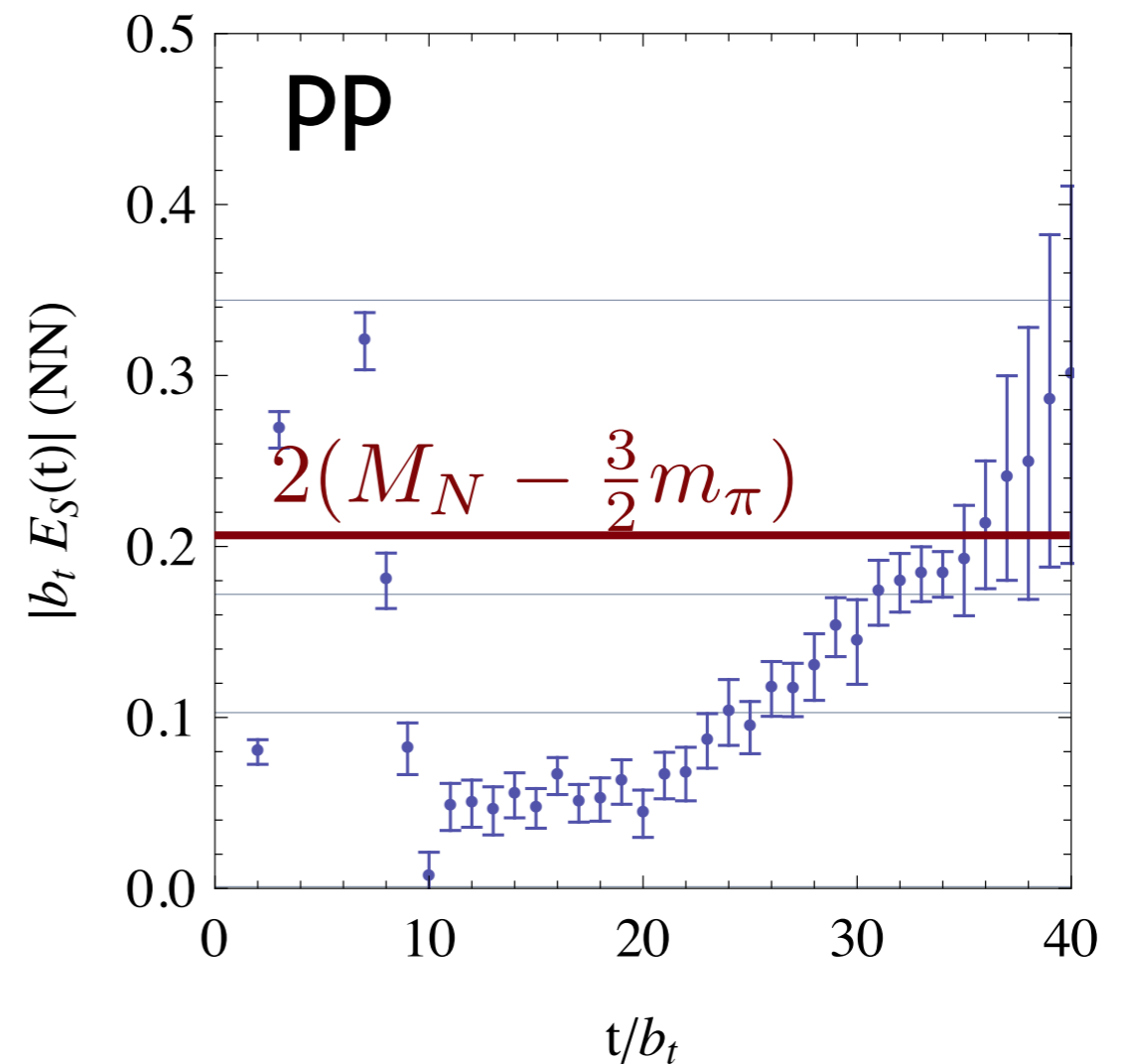
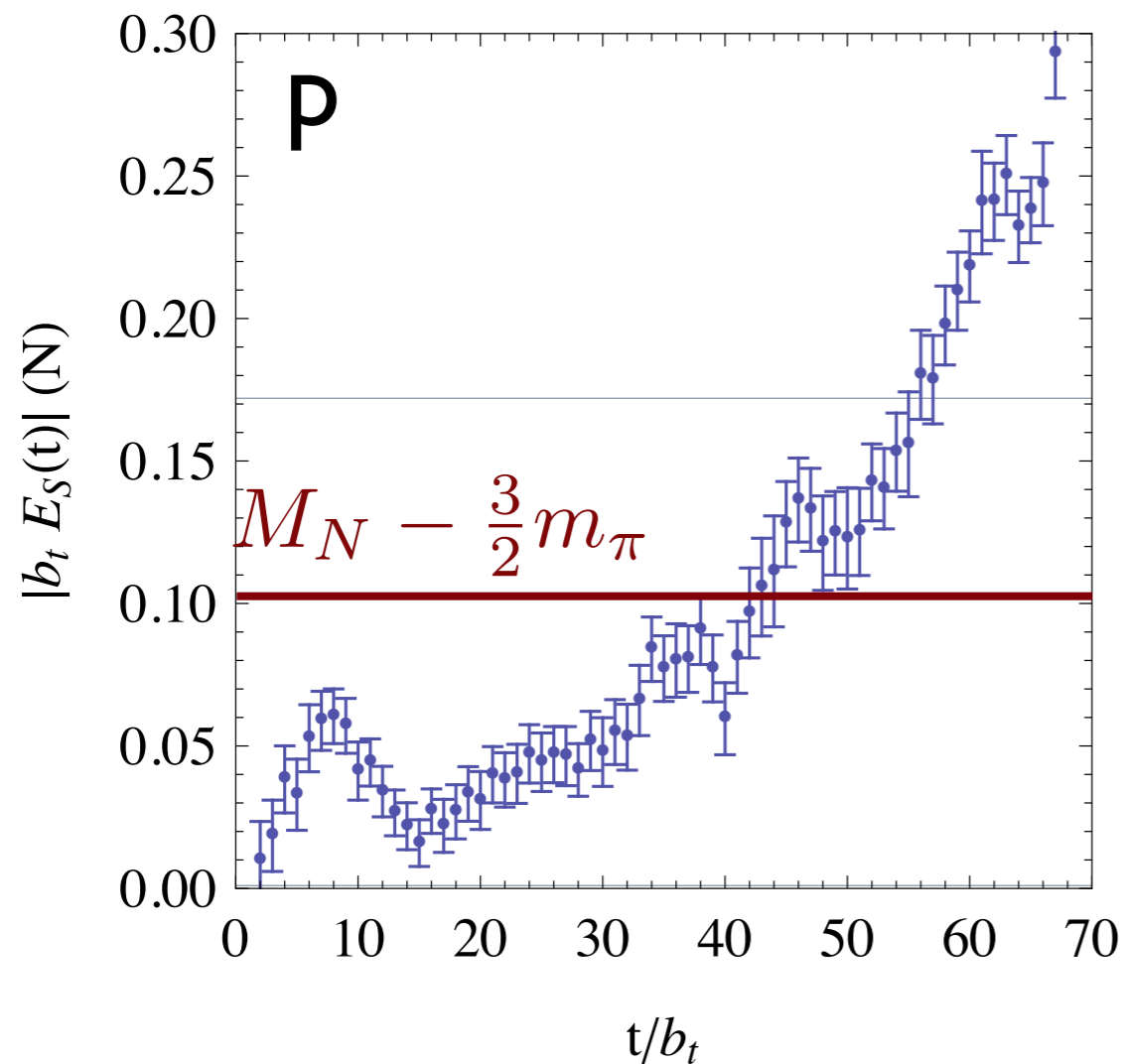
Eg, we're screwed

Actual data

S. R. Beane et al. (NPLQCD), Prog.Part.Nucl.Phys. 66 (2011) 1-40

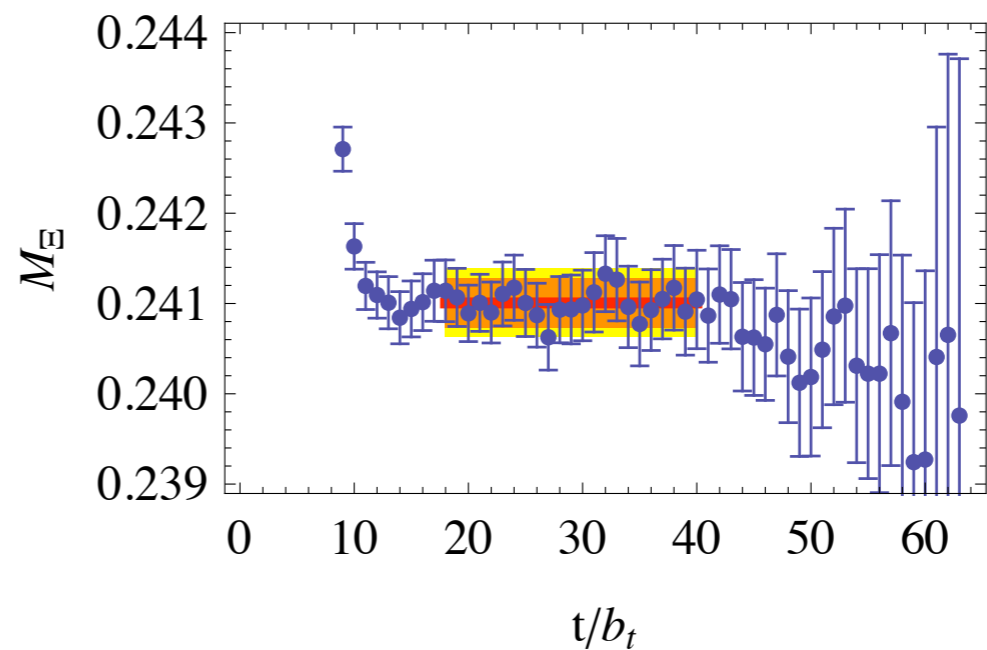
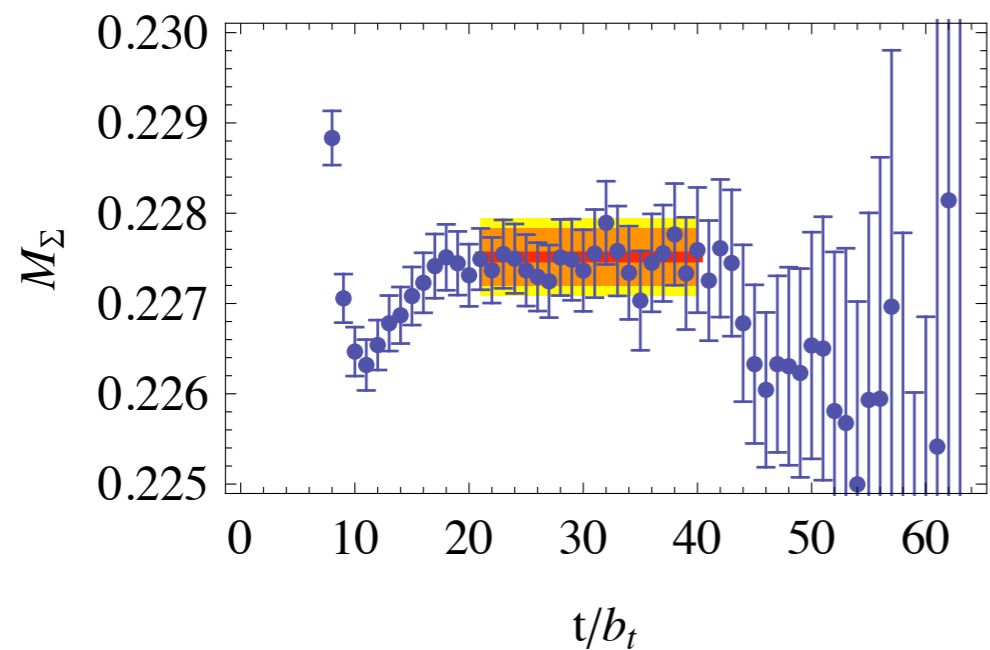
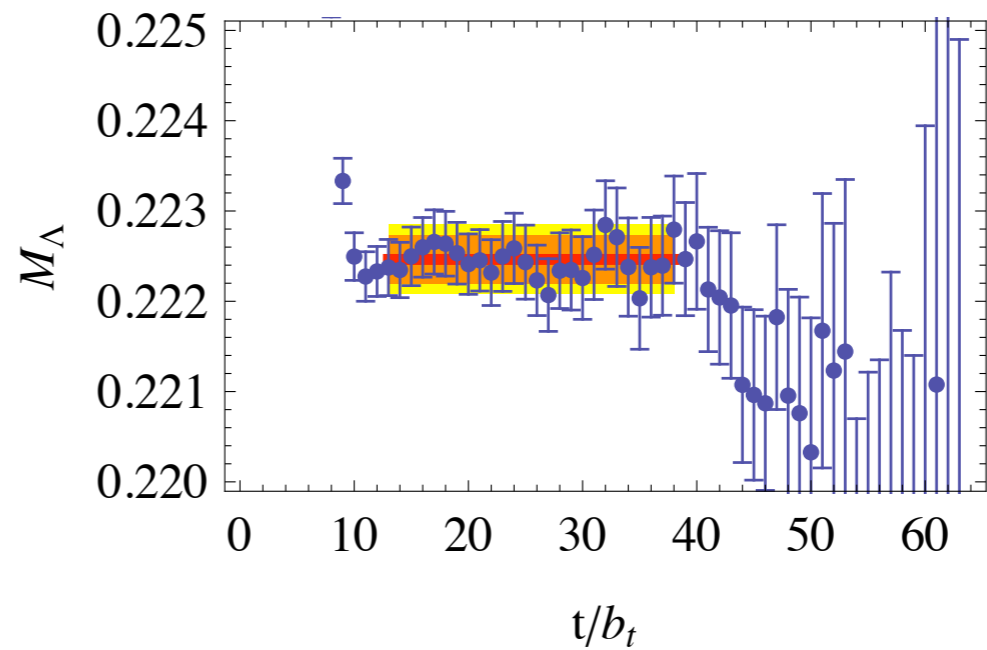
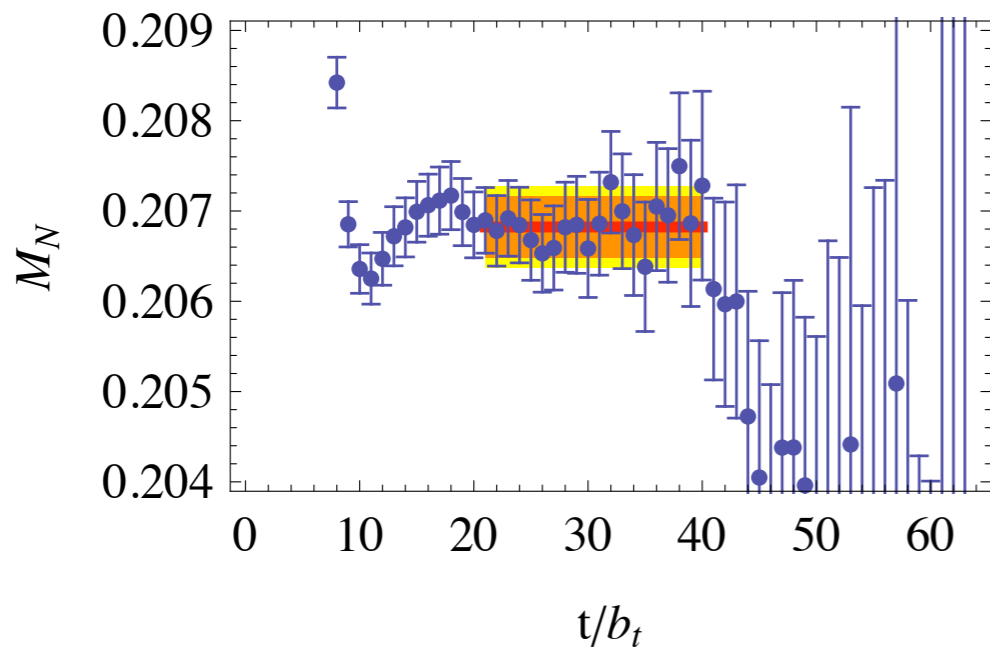
Noise/signal: $\mathcal{S}(t) = \frac{\sigma(t)}{\bar{x}(t)}$

$$E_s = -\frac{1}{t} \ln \mathcal{S}(t)$$



Lepage argument give *rough* idea of noise problem

David B. Kaplan ~ CAQCD ~ May 13, 2011



Example of high precision baryon octet effective mass plots

S. R. Beane et al., (NPLQCD) Phys. Rev. D 79, 114502 (2009)

David B. Kaplan ~ CAQCD ~ May 13, 2011

Simpler system to analyze: “unitary fermions”

(Endres, Lee, Nicholson, DK, to appear)

What are unitary fermions?

Nonrelativistic 2-particle scattering:

$$\mathcal{A} = \frac{4\pi}{M} \frac{1}{p \cot \delta - ip}$$

phase
shift



$$p = \sqrt{ME}$$

scattering
length

$$p \cot \delta = -\frac{1}{a} + \frac{1}{2} r_0 p^2 + O(p^4)$$



effective
range

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effective
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“unitary” fermions: $p \cot \delta = 0$

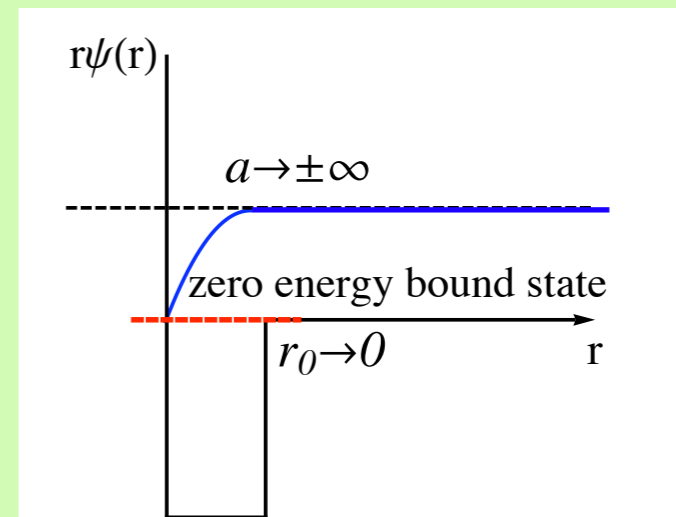
$$\delta = 90^\circ$$

$$a = \infty$$

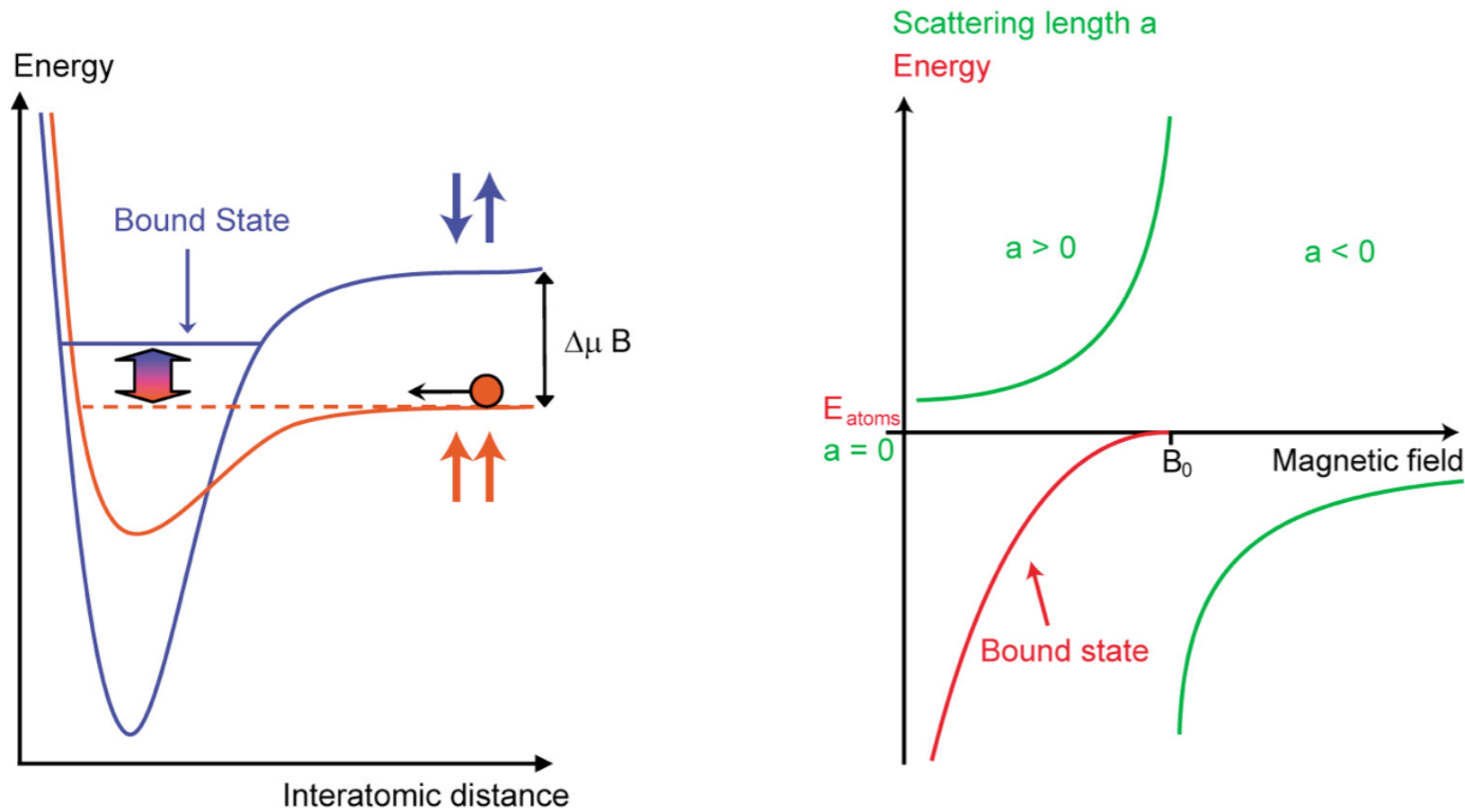
$$r_0 = 0 \dots$$

A strongly-coupled conformal system

Zero-range potential
Zero-energy bound state

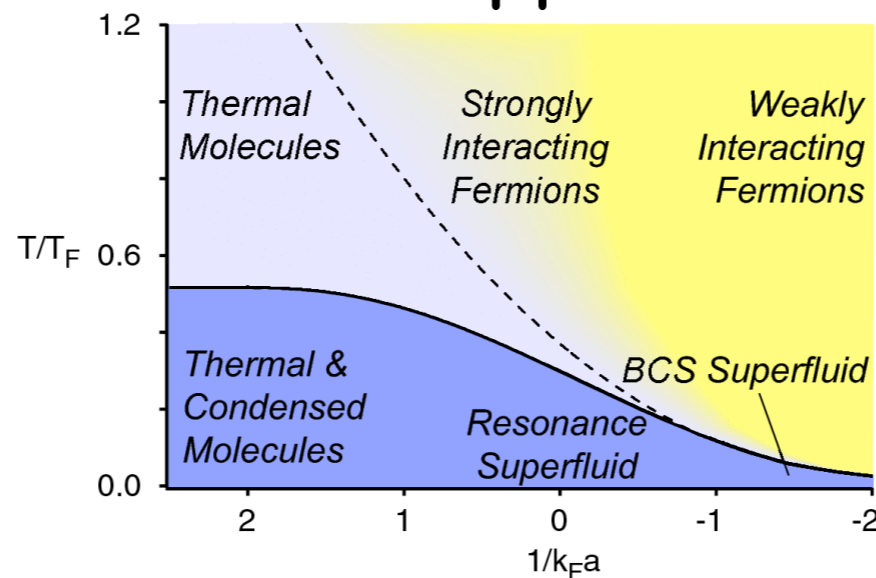


Can study universal properties of unitary fermions experimentally with trapped atoms (JILA, MIT, Innsbruck)

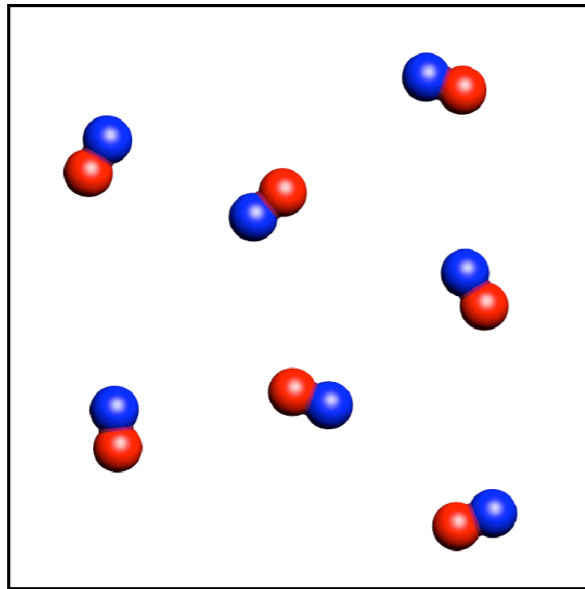


From:
Ketterle Lab web page

Feshbach resonance with trapped atoms: tune to unitarity

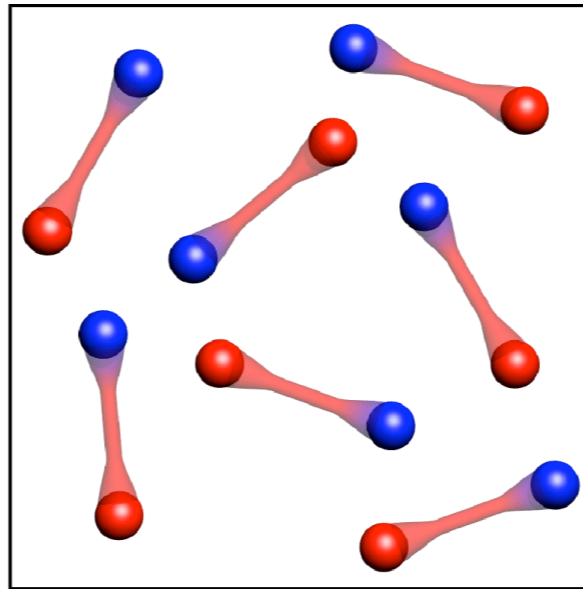


Unitary fermions



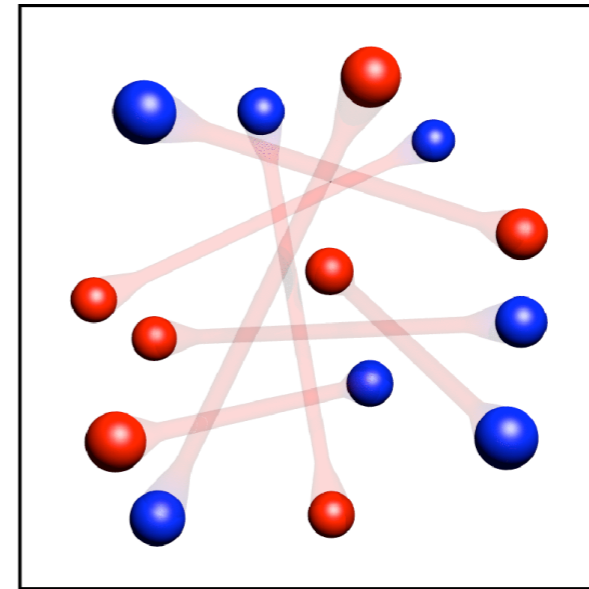
BEC of Molecules

$$\frac{1}{a} \rightarrow +\infty$$



Crossover Superfluid

$$\frac{1}{a} = 0$$



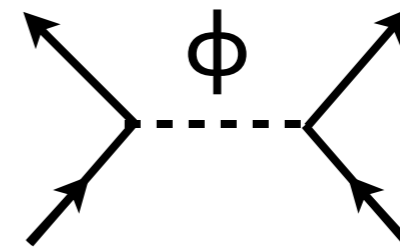
BCS state

$$\frac{1}{a} \rightarrow -\infty$$

- Less severe sign problem; no gauge symmetry; nonrelativistic (quenched)
- Have simulated up to 66 fermions on $14^3 \times 64$ lattice
- 1-2% accuracy in energies

Simulation:

2-particle interaction induced by auxiliary scalar ϕ exchange, tuned to unitarity



ϕ lives on time links
(couples to $\psi^*\psi$)

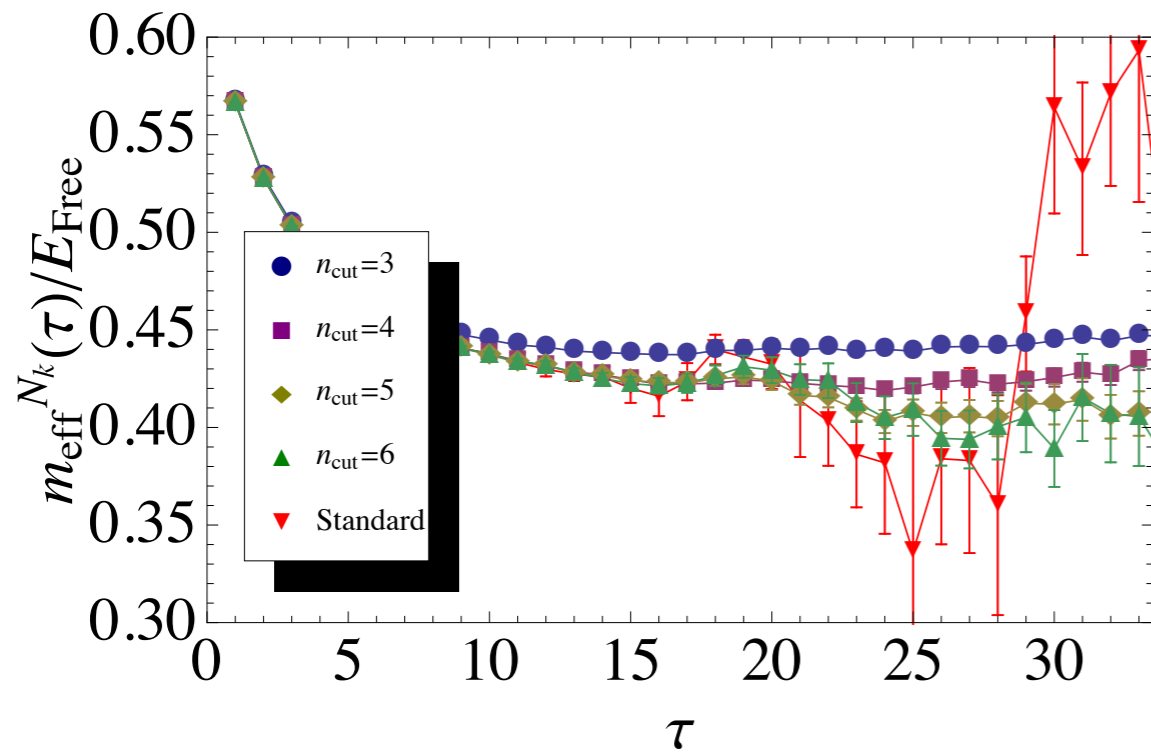
Generate an ensemble of random ϕ fields, compute average of an N-particle correlator $C_N(T; \phi)$ from $t=0$ to $t=T$

Extract ground state energy:

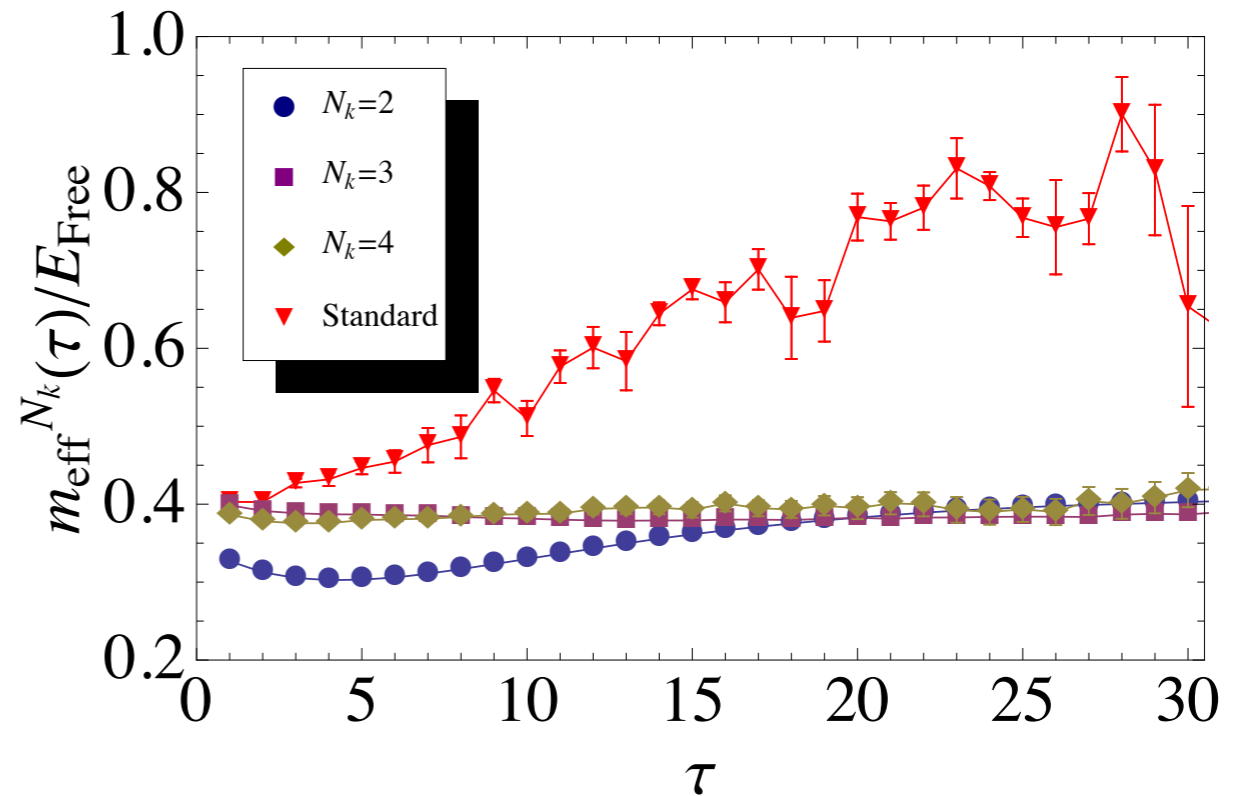
$$E_N = \lim_{T \rightarrow \infty} \left[-\frac{1}{T} \ln \langle C_N(T, \phi) \rangle_\phi \right]$$

Plot $-\frac{1}{T} \ln \langle C_N(T, \phi) \rangle_\phi$
vs T , look for "plateau"

Result of conventional calculation in **RED**



$N=16$ fermions

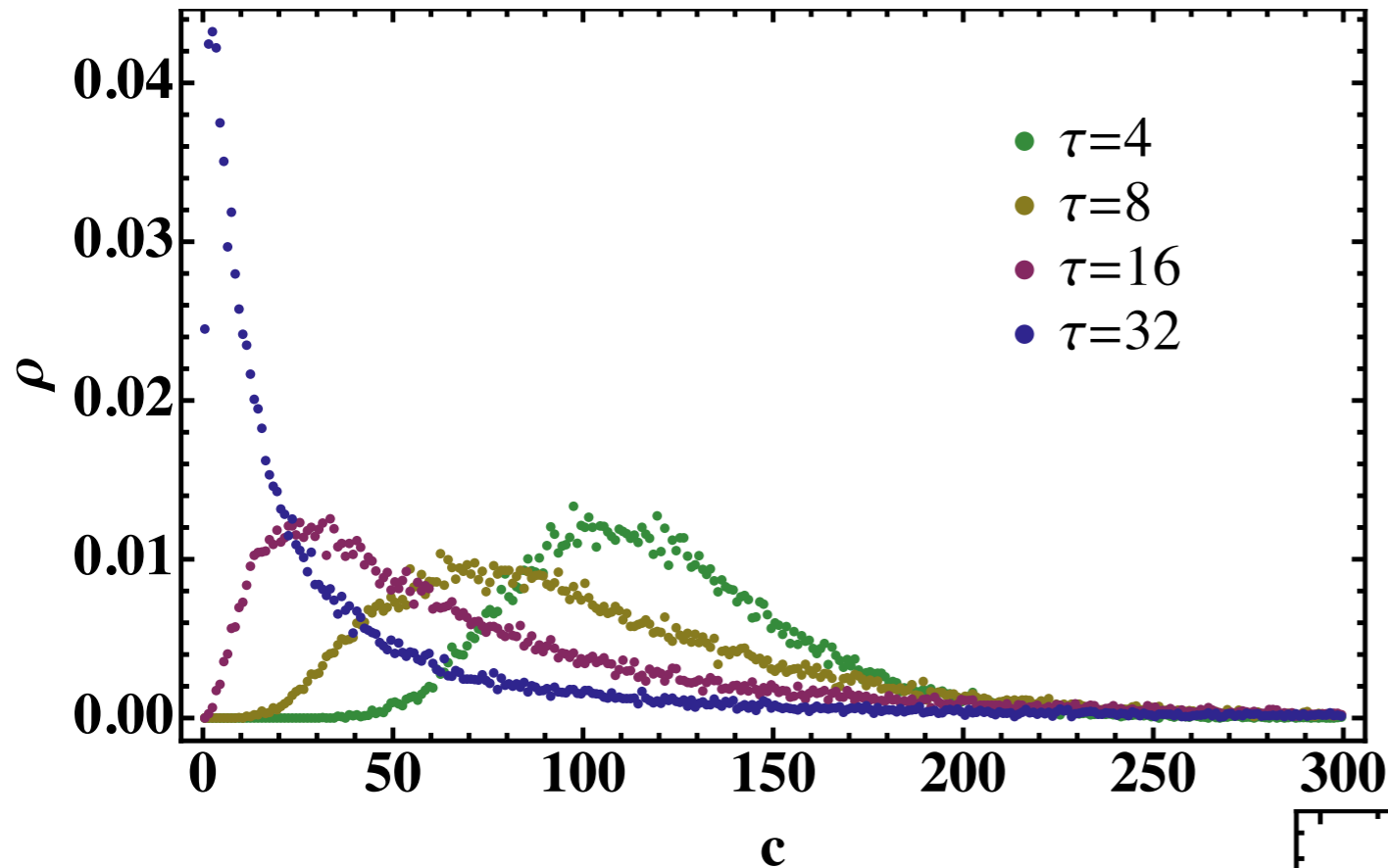


$N=46$ fermions

Conventional calculation shows:

- noise
- drift
- worse for large N

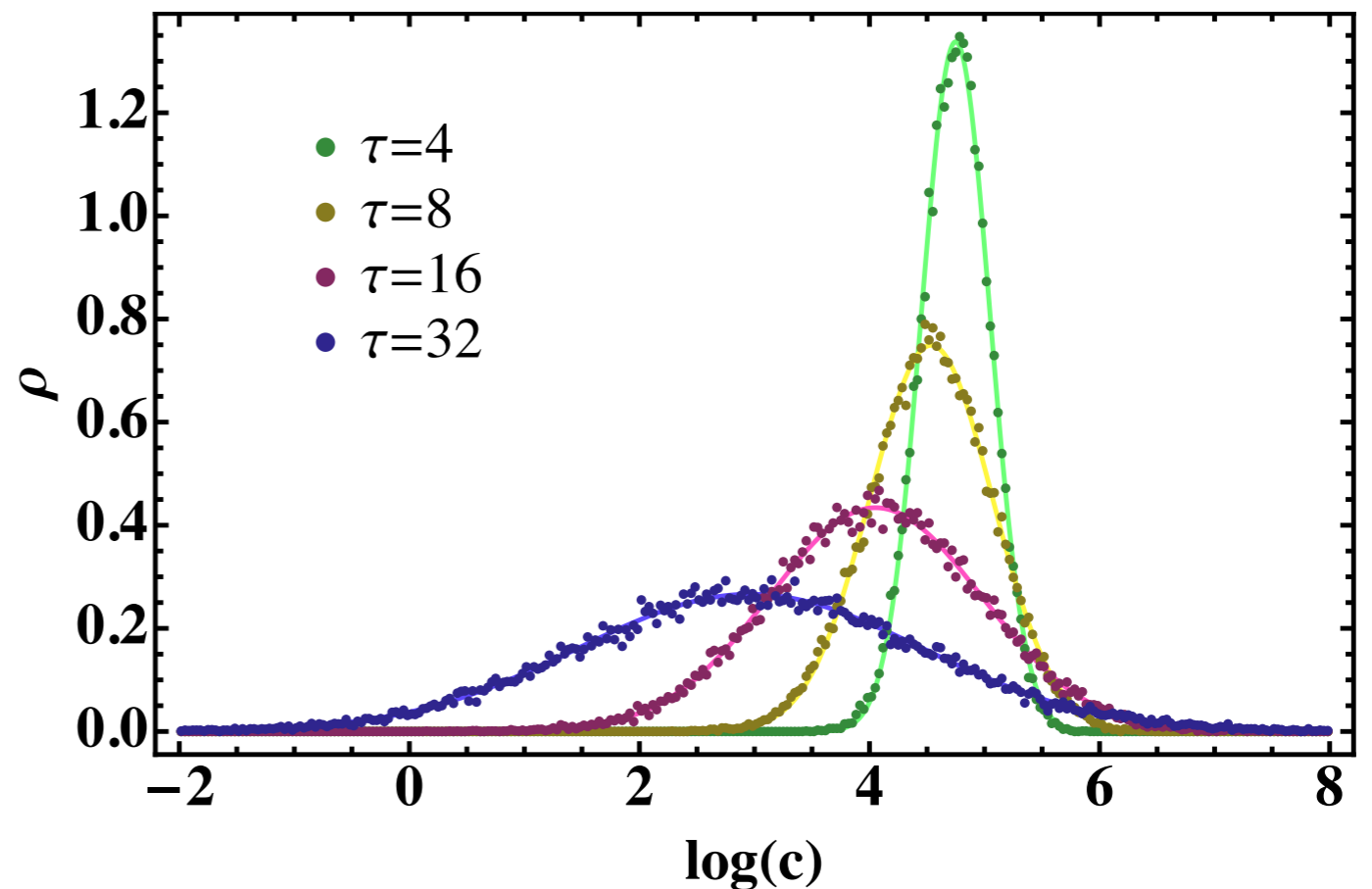
Look at raw correlator distributions: Log-normal?!



N=4 correlator
distribution at different
times (τ)



Distribution of
Log[correlator]
with Gaussian fit



Distributions close to Log-normal are ubiquitous:
arise from products of positive random numbers

$$y = \prod_{i=1}^T \left(1 + \frac{1}{T} x_i\right) \xrightarrow{T \rightarrow \infty} \exp \left[\frac{1}{T} \sum_{i=1}^T x_i \right]$$

random Normal distributed
(CLT)

$P(y)$ is Log-Normal distribution as $T \rightarrow \infty$

Aside: flow of statistics to “attractor distributions”

Example: central limit theorem

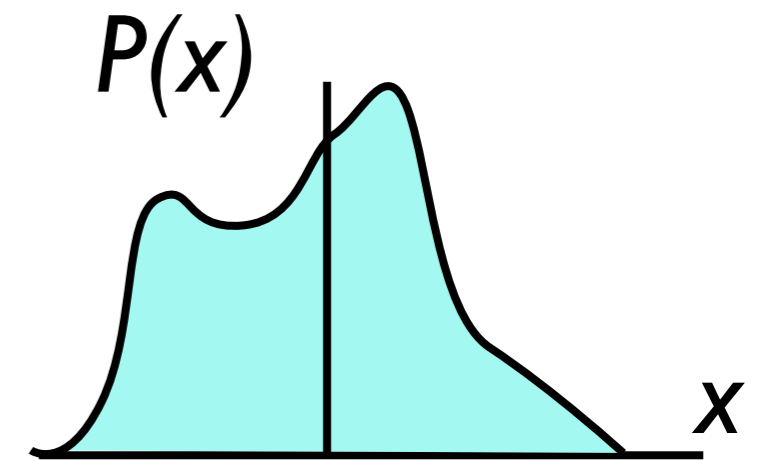
Suppose x has some distribution $P(x)$; specify by cumulants after shift and rescale so that $\mu = \langle x \rangle = 0$, $\sigma = 1$

$$P(x) \implies P(0, 1, \kappa_3, \kappa_4, \dots)$$

$\kappa_n = n^{\text{th}}$ cumulant

$$\kappa_1 = \mu = 0$$

$$\kappa_2 = \sigma^2 = 1$$



Average: $x \rightarrow \frac{x_1 + x_2}{\sqrt{2}}$

$$\kappa_n \rightarrow 2^{1-n/2} \kappa_n$$

$\kappa_1 = \mu = \text{relevant}$

$\kappa_2 = \sigma^2 = \text{marginal}$

$\kappa_{3\dots} = \text{irrelevant}$

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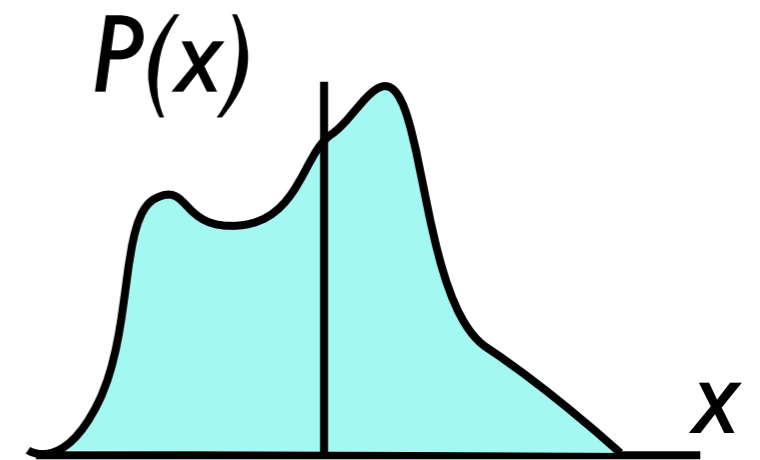
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Repeat averaging: $P(0, 1, \kappa_3, \kappa_4, \dots) \rightarrow P(0, 1, 0, 0, \dots)$
Normal distribution

Aside: flow of statistics to “attractor distributions”

Example: central limit theorem

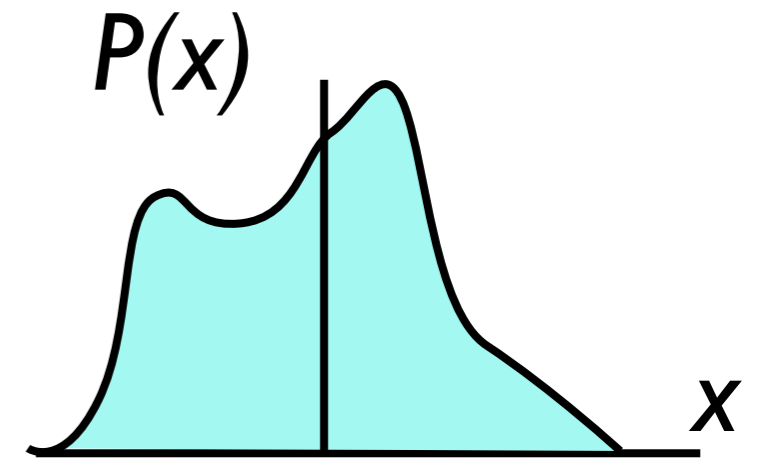
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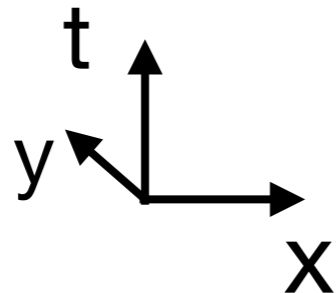
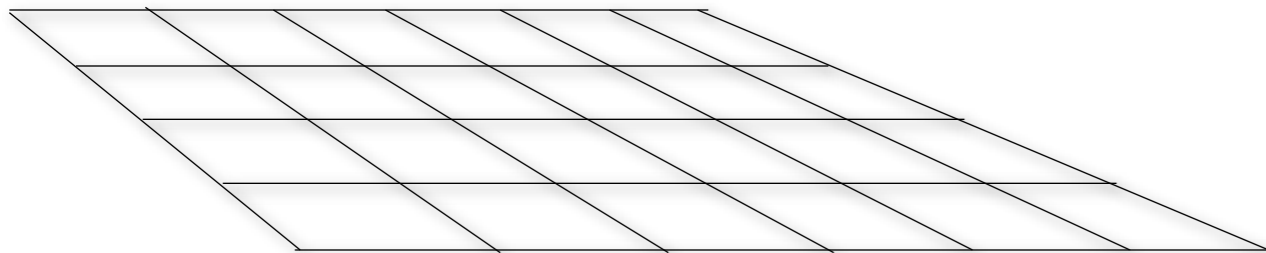
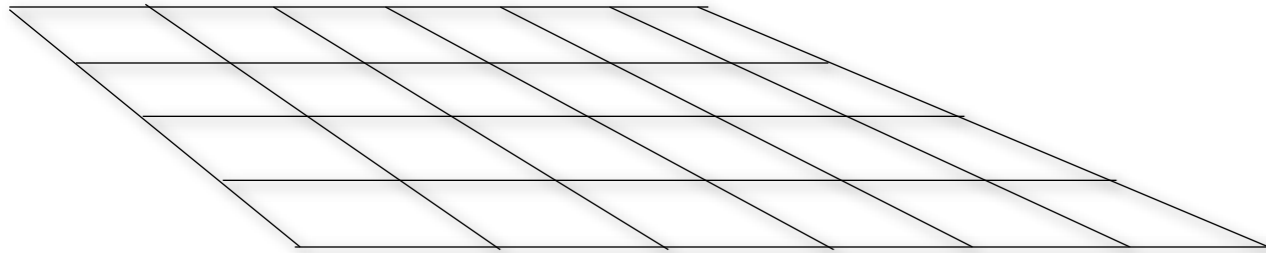
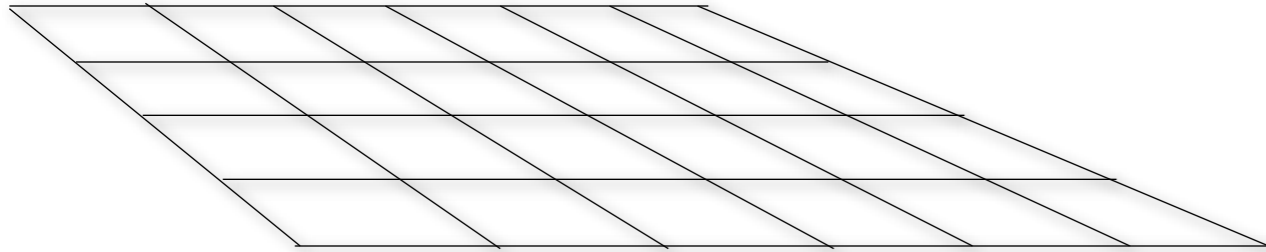
$\kappa_2 = \sigma^2 = \text{marginal}$

$\kappa_{3\dots} = \text{irrelevant}$

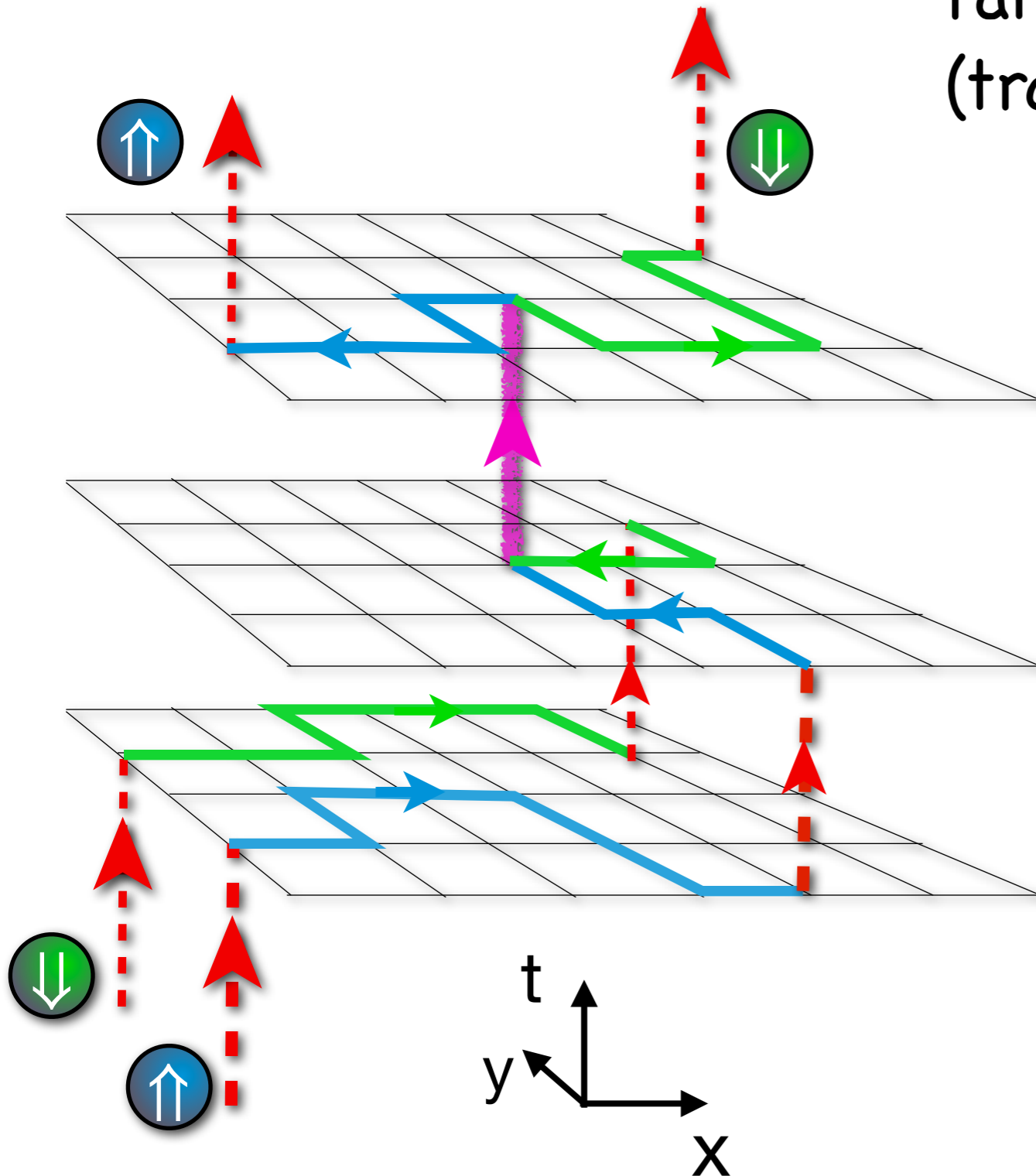
Repeat averaging: $P(0, 1, \kappa_3, \kappa_4, \dots) \rightarrow P(0, 1, 0, 0, \dots)$
Normal distribution

Just like renormalization group; κ_n like dim $2n$ operators in φ QFT

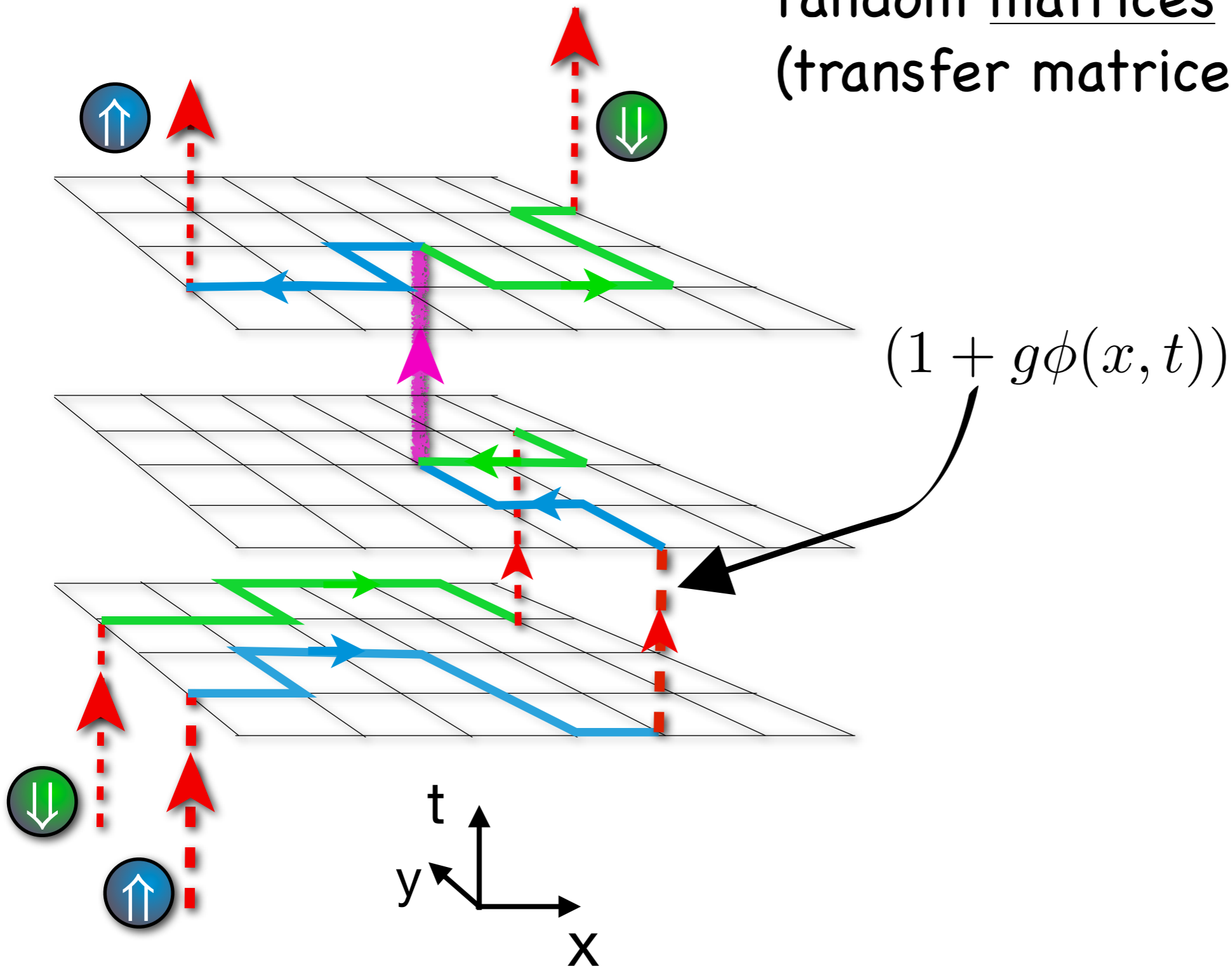
Fermion correlators
are products of
random matrices
(transfer matrices)



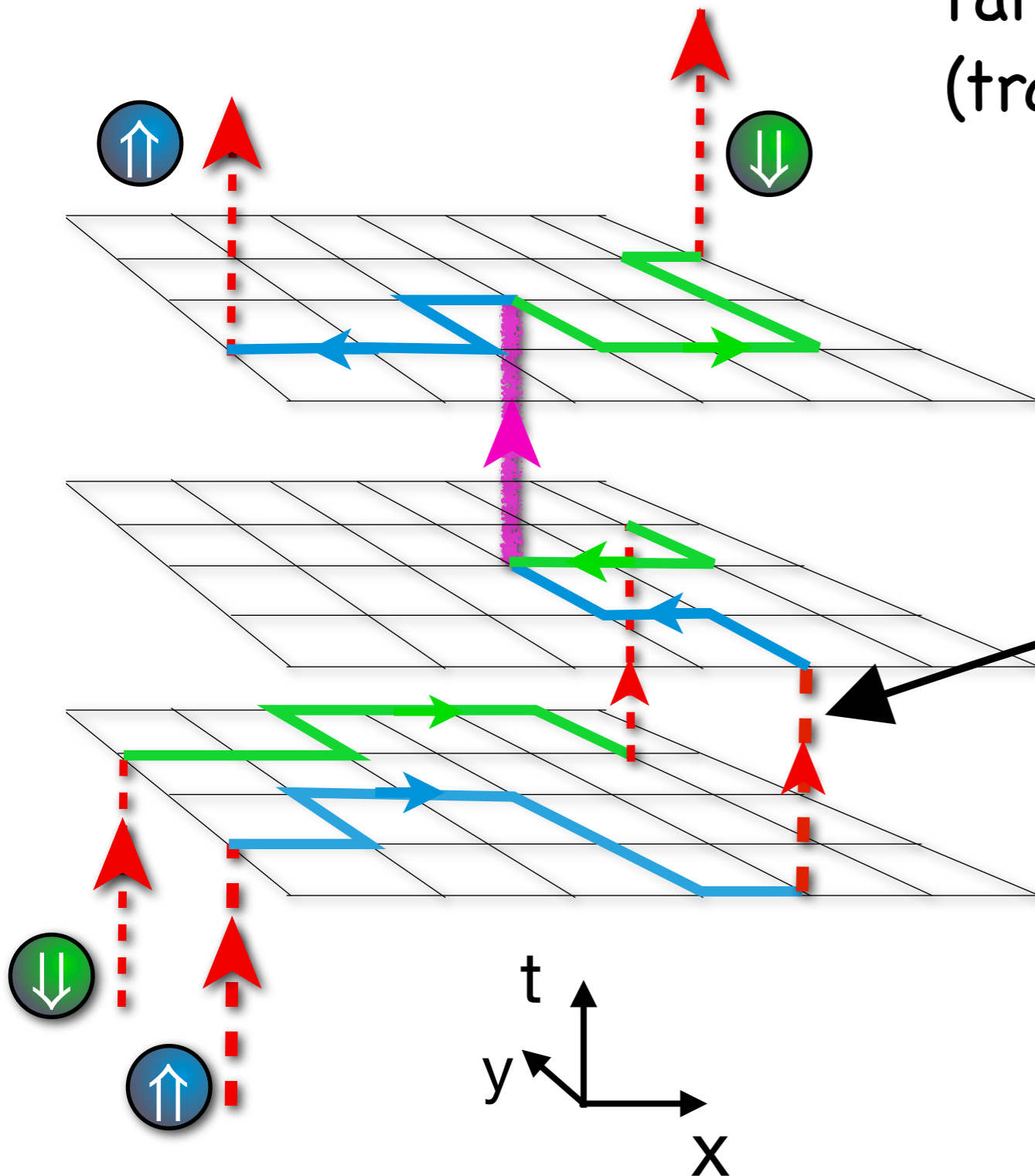
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$$(1 + g\phi(x, t))$$

Toy model:

$$C(T) = \prod_{i=1}^T (1 + g\phi_i)$$

$$\phi_i \in [-1, 1], \quad g = \frac{1}{2}$$

Toy model:

$$C(T) = \prod_{i=1}^T (1 + g\phi_i)$$

$\phi_i \in [-1, 1]$ uniform dist.

Exact answer for the "energy":

$$E(T) \equiv -\frac{1}{T} \ln \langle C(T) \rangle_{\phi} = 0$$

Compare with simulation (finite sample size), $g=1/2$

Two strategies:

- Conventional: $E \rightarrow -\frac{1}{T} \ln \left[\frac{1}{N} \sum_{i=1}^N C(T, \phi_i) \right]$

- New: use identity

$$\ln \langle C \rangle = \sum_n \frac{\kappa_n}{n!} \quad \kappa_n = \text{cumulants of } \ln C(T, \phi)$$

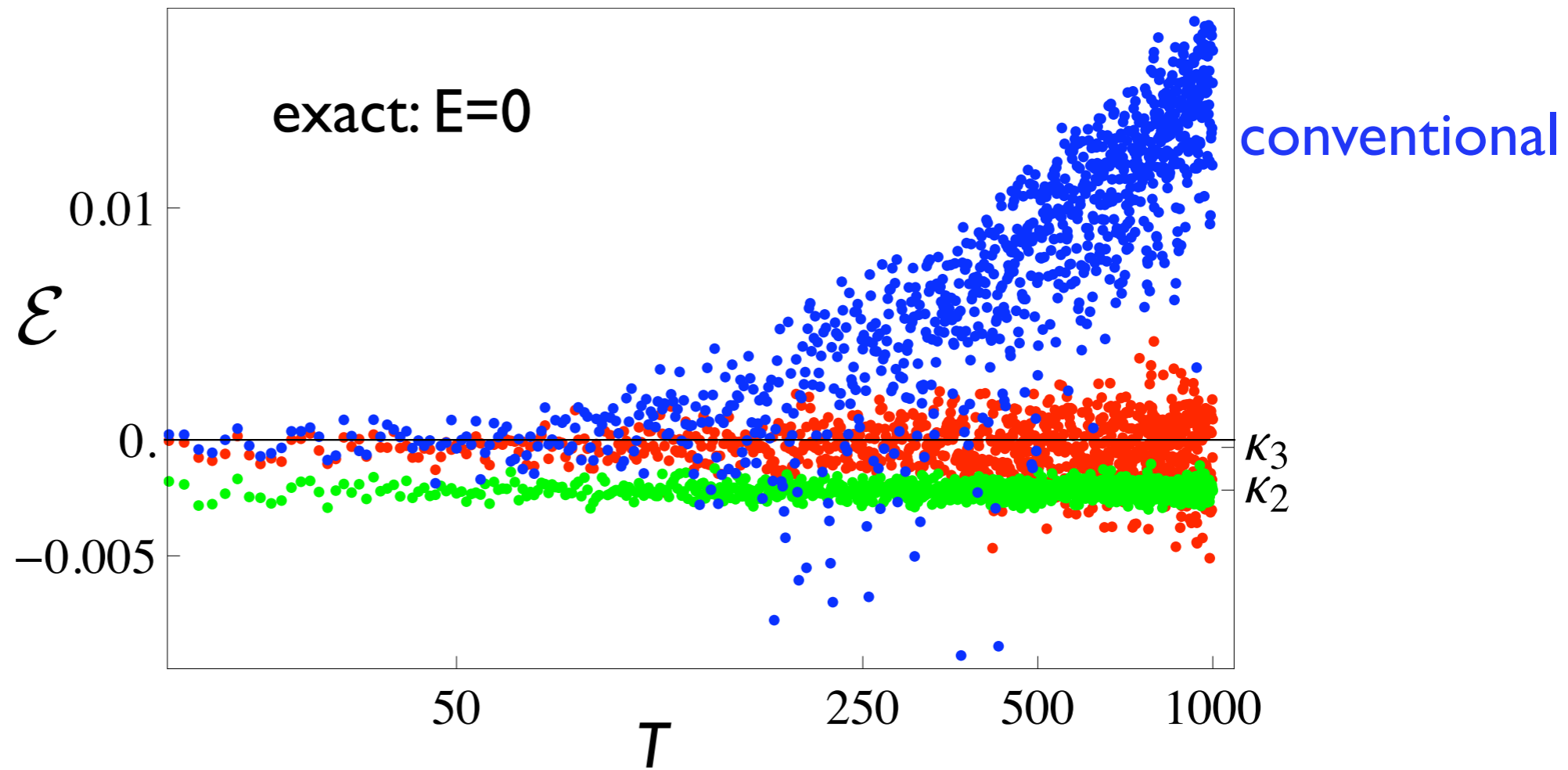
and estimate κ_n from sample for low n .

Why?

- If C distribution Log-Normal, $\kappa_{n>2}=0$;
- κ_n expansion like EFT expansion

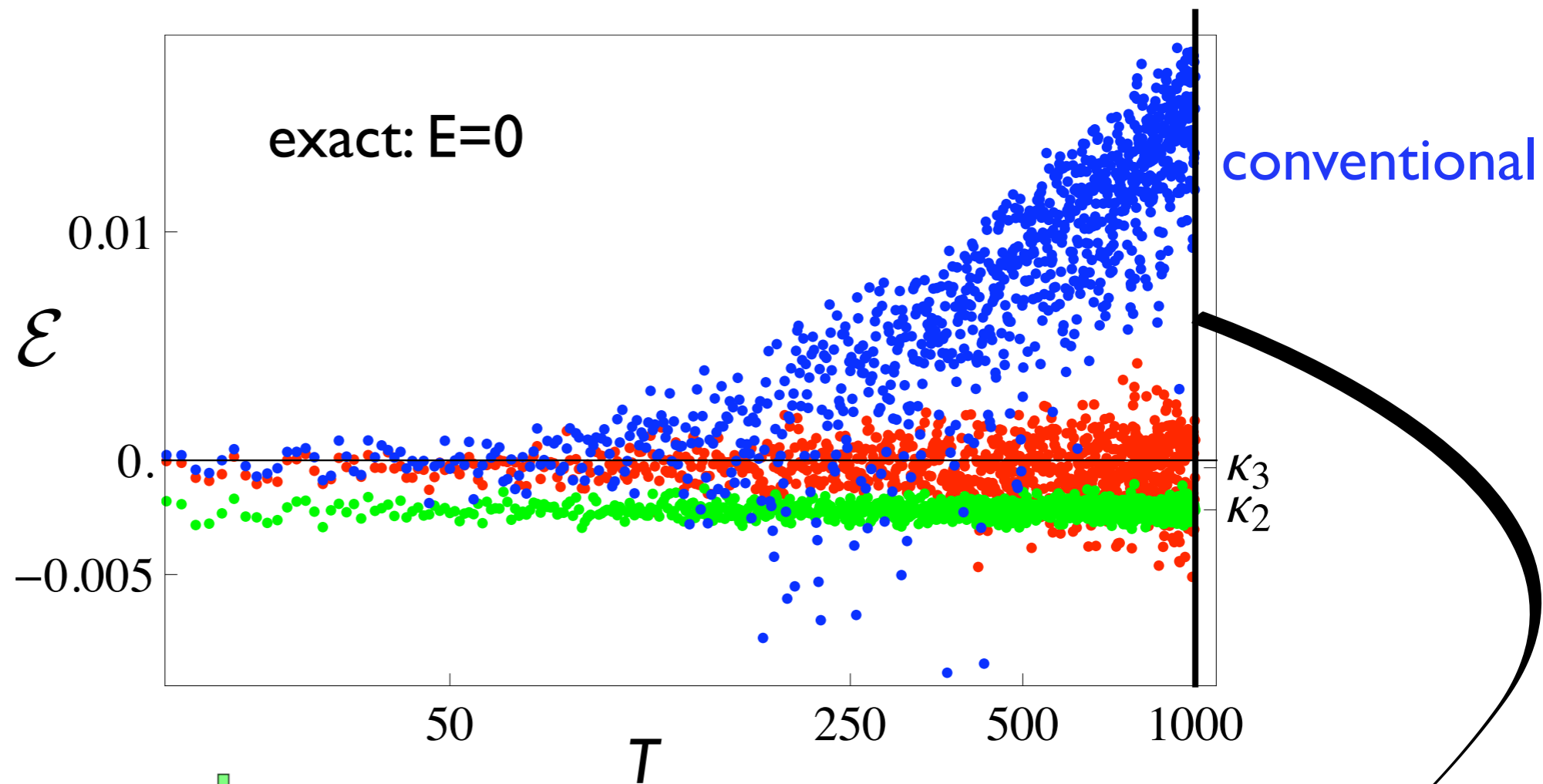
Toy model simulation

Sample size
N=50,000

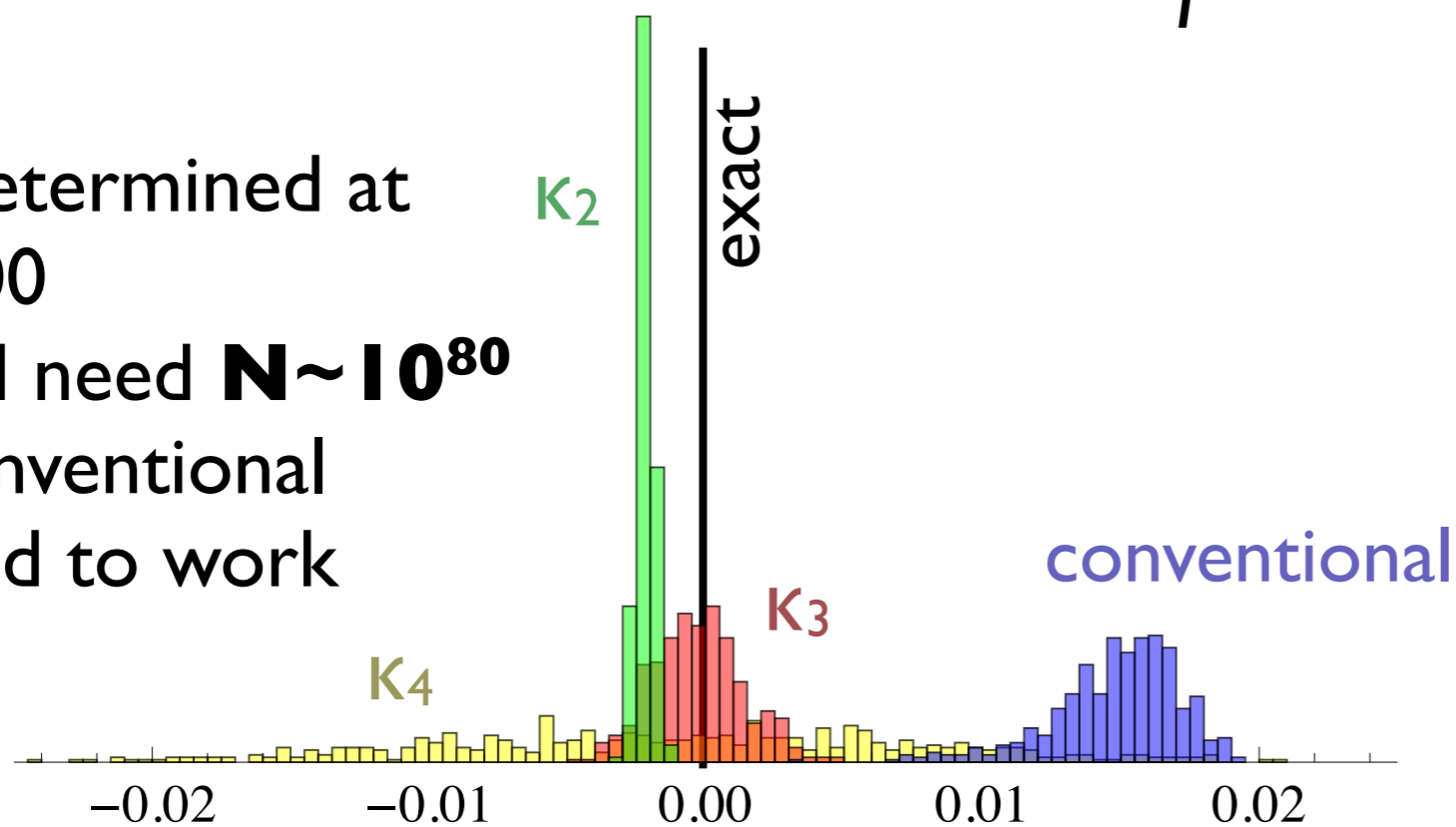


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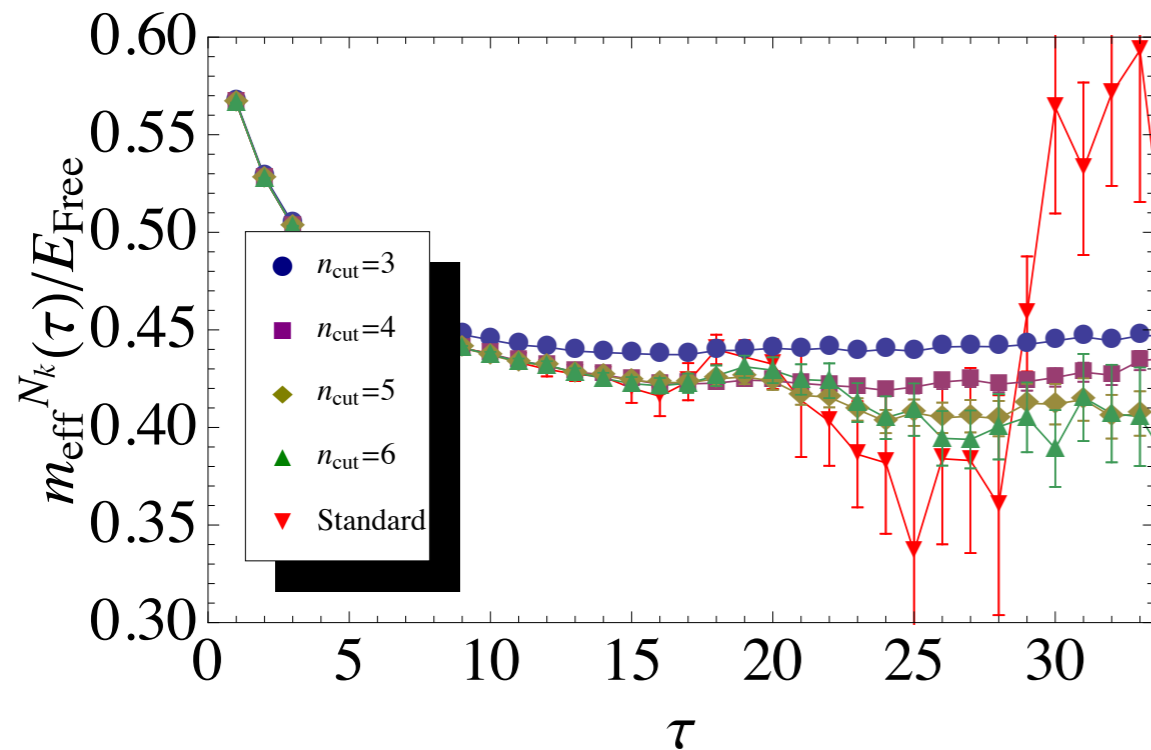


- E as determined at $T=1000$
- Would need $N \sim 10^{80}$ for conventional method to work

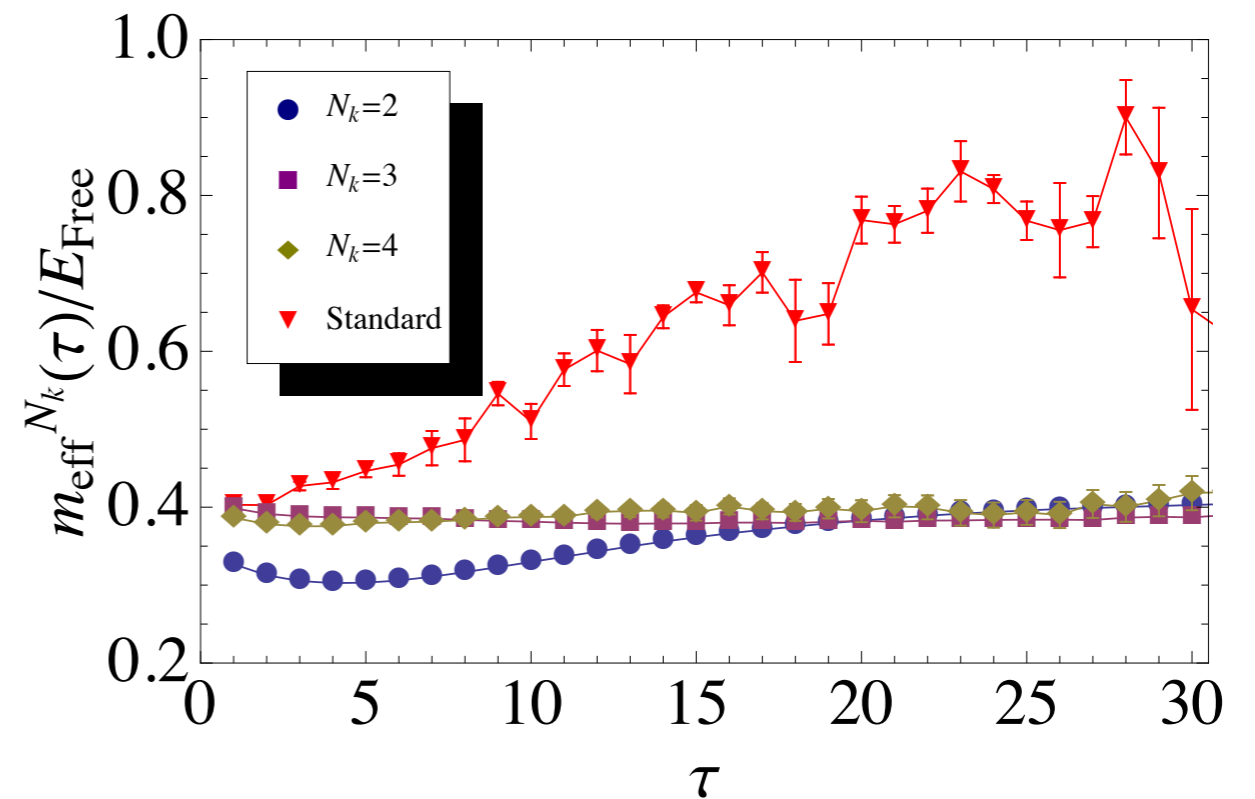


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Back to data from unitary fermion simulations



$N=16$ fermions

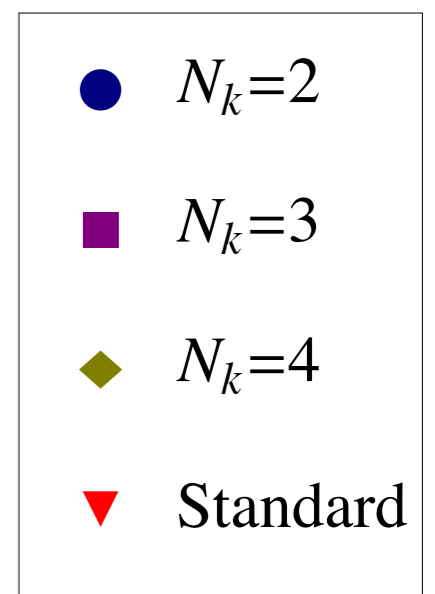


$N=46$ fermions

Conventional calculation: poor results

Cumulant expansion works much better

-- especially for large N



Toy model too simplistic;

Determining distributions of random matrix products too hard...

...but mean field analysis used for disordered media may provide connection between microphysics and correlator distributions in lattice gauge theory

Mean field argument: distribution for $y = \text{Log}[C_N(\phi, T)]$

$$\begin{aligned} P(y) &= \int [D\phi] e^{-\int d^4x \frac{1}{2} m^2 \phi^2} \delta(\ln C_N[\phi] - y) \\ &= \int \frac{ds}{2\pi} \int [D\phi] e^{is(\ln C_N - y) - \int d^4x \frac{1}{2} m^2 \phi^2} \\ &\equiv \int \frac{ds}{2\pi} \int [D\phi] e^{-S} \end{aligned}$$

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Perform semiclassical expansion of S for large N ; find $\phi_0 = \text{const.}$
Derivatives of $\text{Log}[C_N]$ w.r.t. ϕ are current correlators in N -particle Fermi gas.

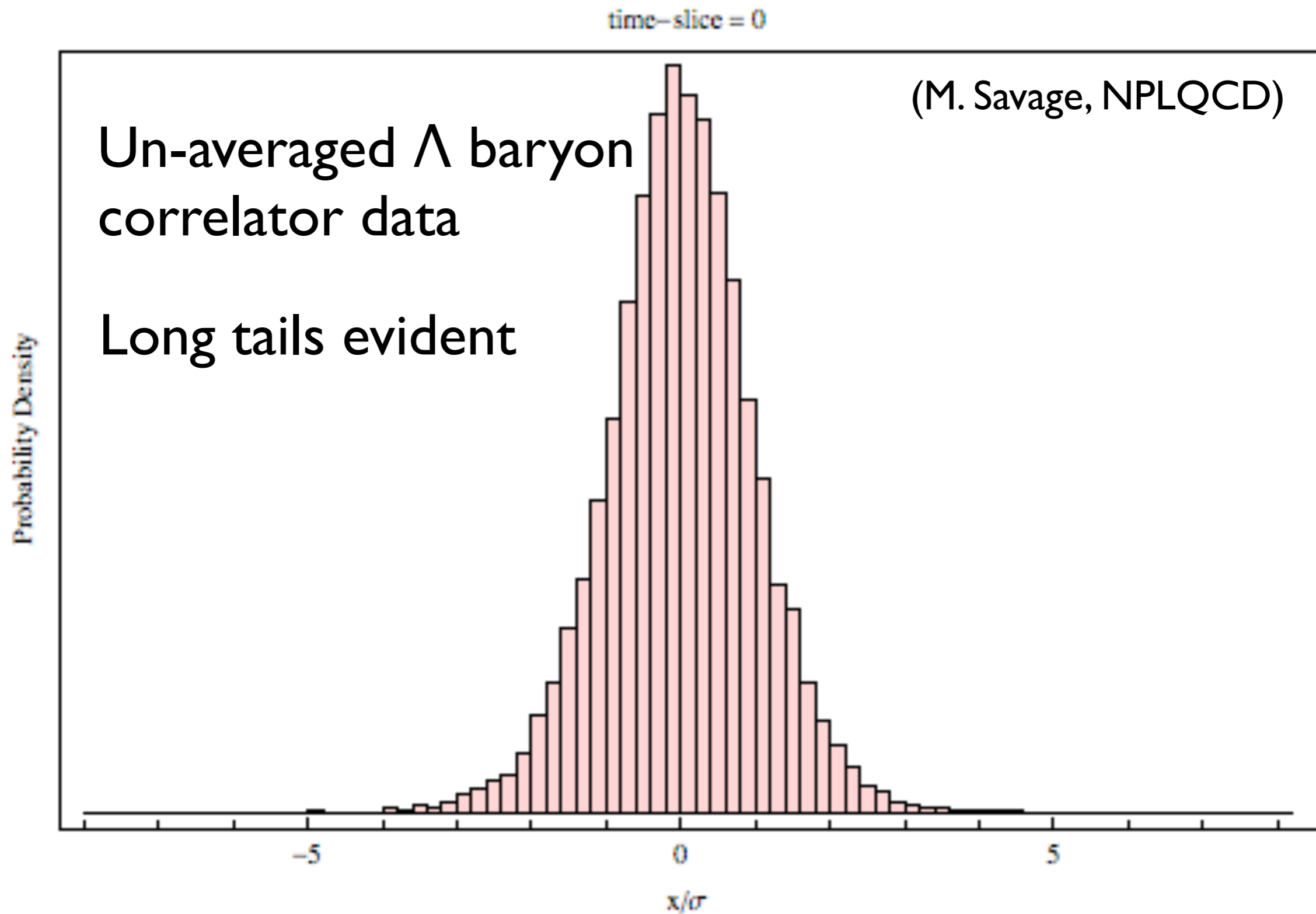
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Derivatives of $\text{Log}[C_N]$ w.r.t. ϕ are current correlators in N -particle Fermi gas.

Result: Log Normal distribution (with corrections);
 μ, σ^2 scale with N and T as seen in data

Do heavy-tailed non-gaussian distributions occur in lattice QCD? Probably, especially large baryon number



David B. Kaplan ~ CAQCD ~ May 13, 2011

Conclusions:

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- Noise distribution in multi-fermion simulations has structure
- Possibly amenable to analysis similar to EFT -- characterized by hierarchy of correlations
- By better understanding the noise, we might better hear the underlying message
- Stay optimistic and don't pay too much attention to no-go theorems!