
Orbital-like Order in Several Families of Cuprates



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Do we have a broken symmetry below T^* ?

Magnetic order in the pseudogap state of high- T_c cuprates
(examples in 4 different families: YBCO, Hg1201, LSCO, Bi2212)

Using polarized neutron diffraction over the last 6 years:
LLB-Saclay, ILL-Grenoble

associated magnetic excitations in Hg1201 (Yuan Li talk)

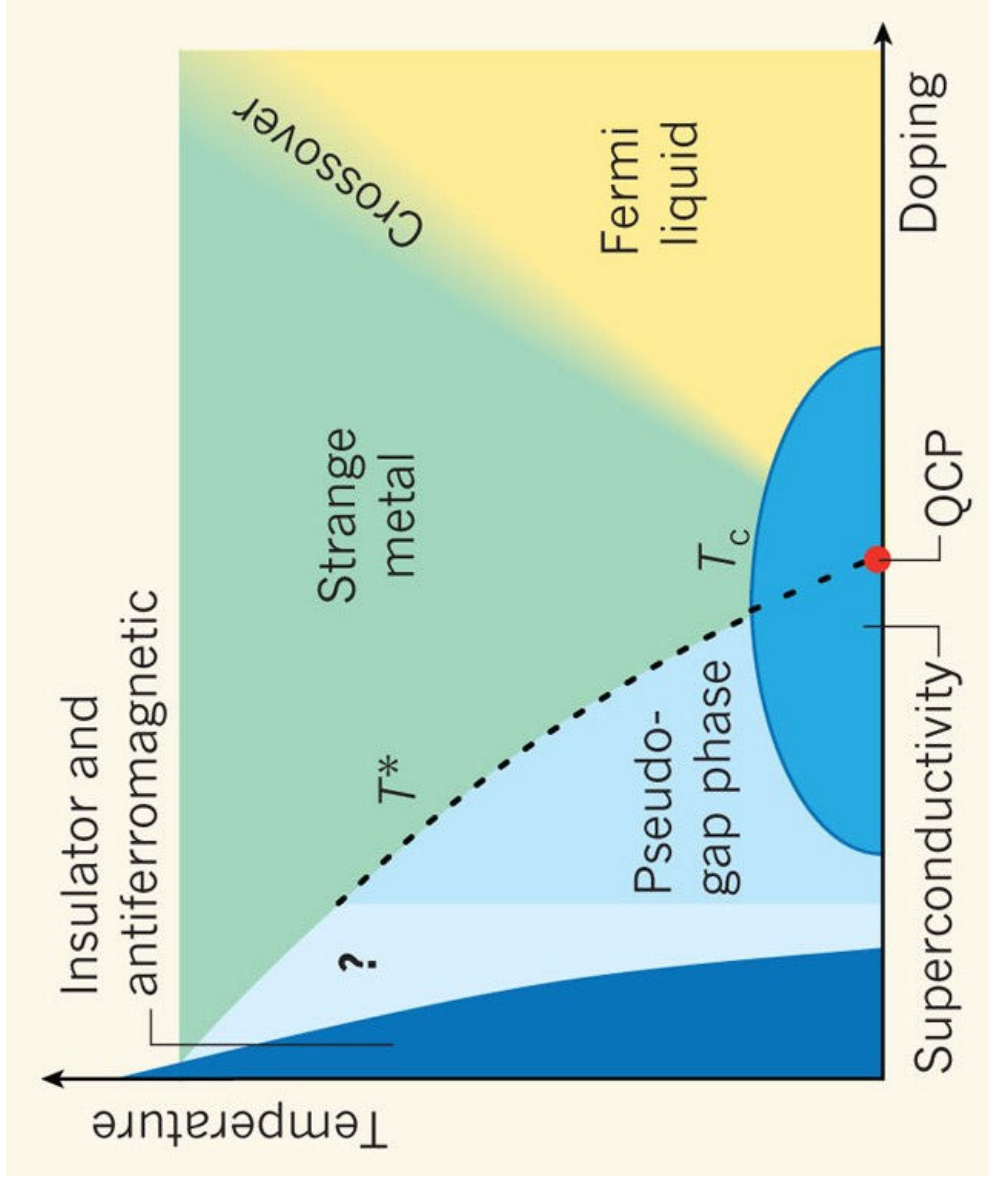
Collaborators :

- V. Balédent (PHD), B. Fauqué, S. Pailhès, Y. Sidis
(Laboratoire Léon Brillouin - Saclay) $(Y,Ca)Ba_2Cu_3O_{6+x}$
- D. Haug, C.T. Lin, V.Hinkov (MPI Stuttgart)

- X Chaud (CRETA, Grenoble), LP Regnault (CEA Grenoble)
• H.A. Mook (Oak Ridge, USA) $YBa_2Cu_3O_{6+x}$
- Yuan Li (MPI), G. Yu, M. Chan (University Minnesota)
• M. Greven (University Minnesota, USA)
• P. Steffens (ILL-Grenoble) $HgBa_2CuO_{4+x}$

- K. Conder, E. Pomjakushina, N. Christensen (Riso),
J. Mesot (PSI, Switzerland) $La_{2-x}Sr_xCuO_4$


Outline:

- 1) Motivation: Loop-current order
- 2) Neutron scattering
- 3) Long range magnetic $Q=0$ AF order in YBCO and Hg1201
- 4) Preliminary result in Bi2212
- 5) Short range magnetic order in LSCO
- 6) Updated phase diagram in YBCO and Implications

Pseudo gap: Hidden Ordered State



Order parameter?

Broken symmetry?

Quantum critical point?

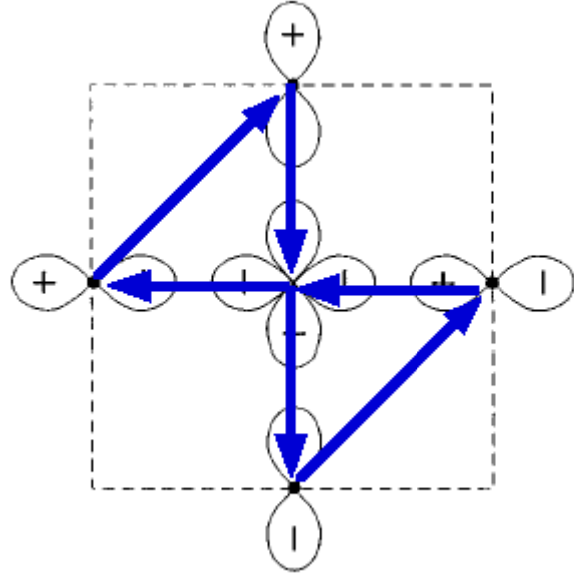
Role of the related fluctuations in the SC pairing mechanism?

Hidden order parameter: Spin Density Wave, Charge DW, more Exotic:

Circulating current,
D-Density Wave,...
Electronic Liquid Crystal

Circulating Currents: C.M. Varma, PRB 73, 155113 (2006)

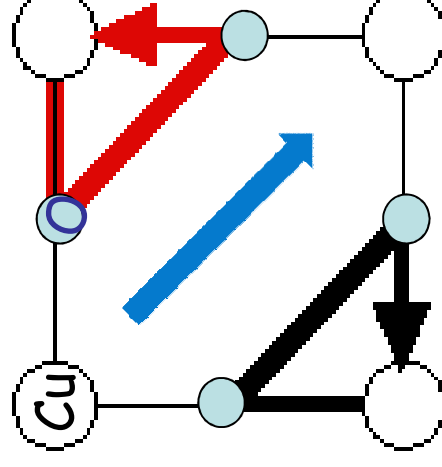
- 3-band Hubbard model: $\text{Cu}(dx^2-y^2)$, $\text{O}(2px)$, $\text{O}(2py)$



In each
unit cell:

Toroidal moment

4 possible
domains



2 current loops circulating clockwise and anti-clockwise => producing Orbital moments

Circulating currents: phase Θ_{II}

- Respects Translation Symmetry ($q = 0$ transition)

- Breaks Time-Reversal symmetry

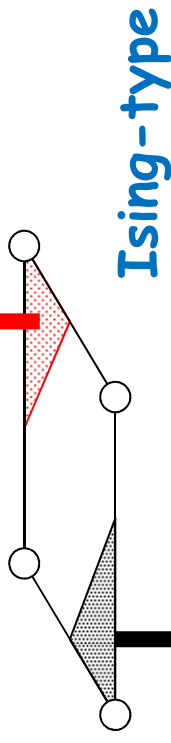
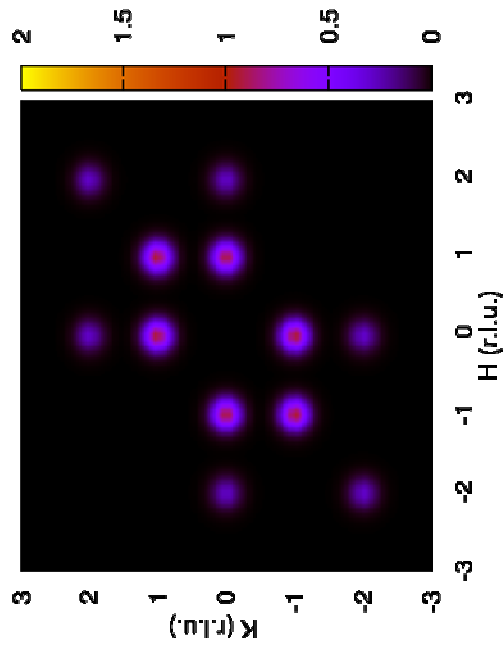
→ Dichroic signal in ARPES in BiSrCaCuO
Kaminski et al. Nature 416, 610 (2002).

- Yields a sizeable magnetic moment of $0.1 \mu\text{B}$

Circulating currents C.M. Varma, PRB (1997)(2006)

In k-space: $Q=(H,K)$

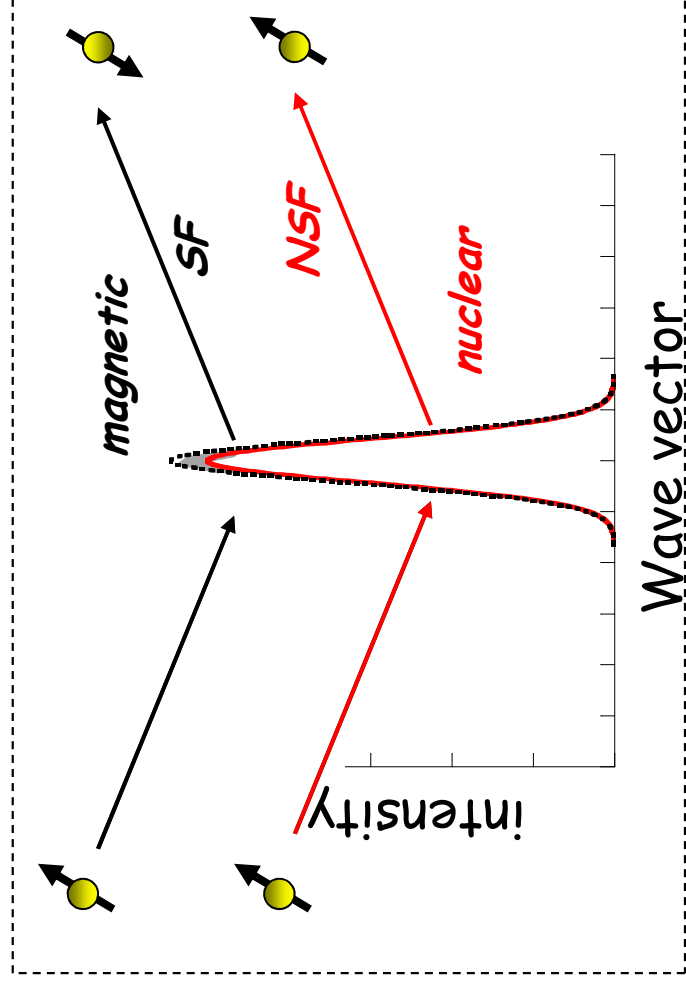
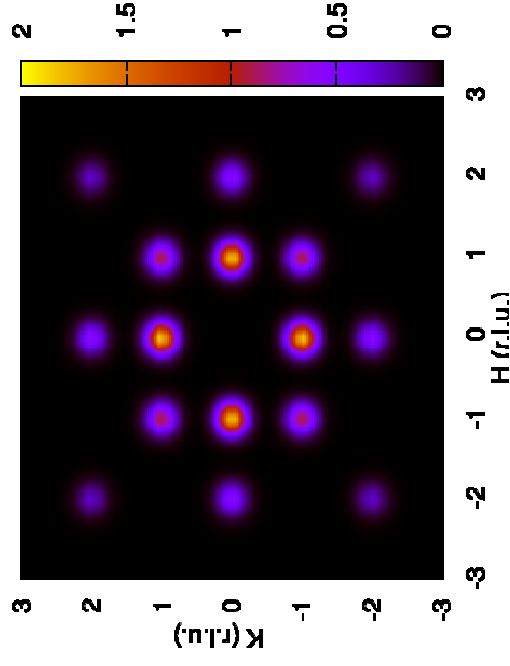
$Q=0$ Orbital AF



Ising-type

Orbital moments (\perp CuO₂ plane)
measurable with neutron diffraction

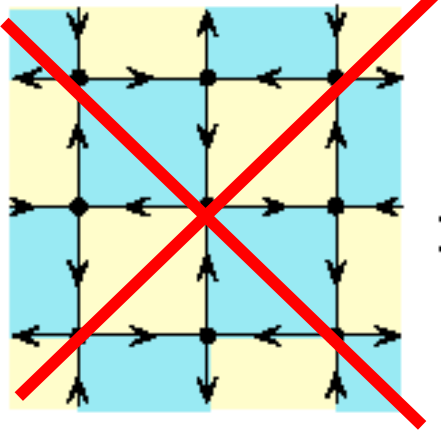
2 type of Domains at 90°



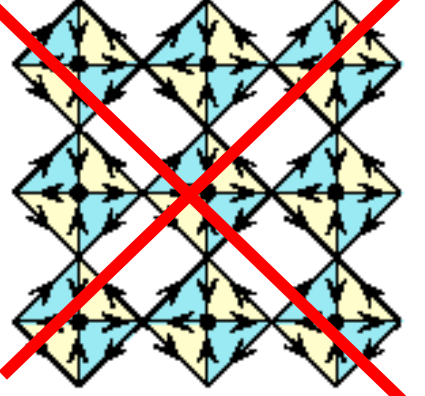
Charge currents: DDW and Circulating currents

S. Chakravarty PRB (2001)

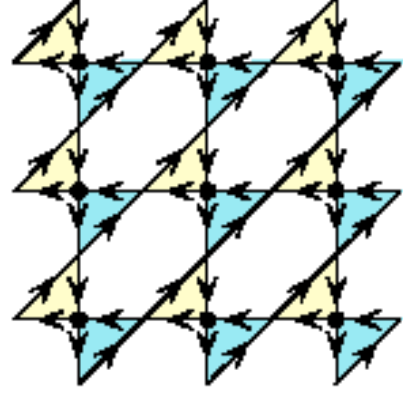
C.M. Varma, PRB (1997)(2006)



ddw



cc- θ_I

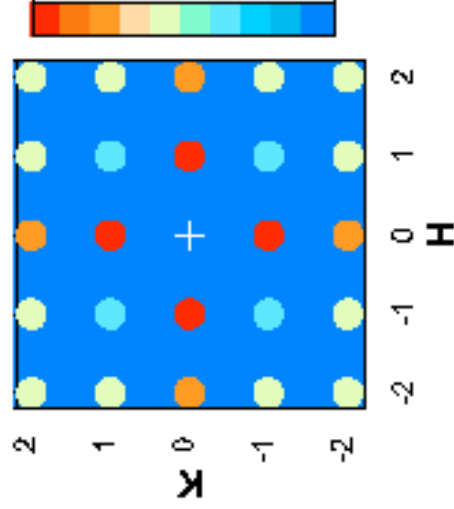
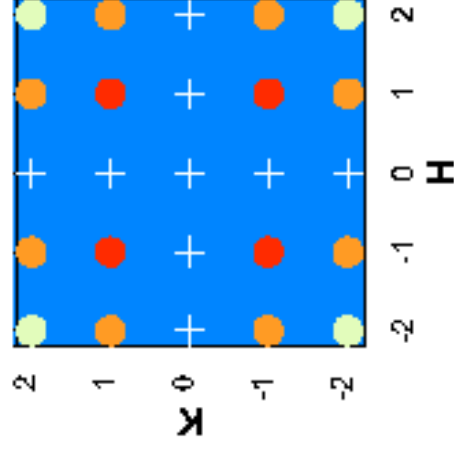
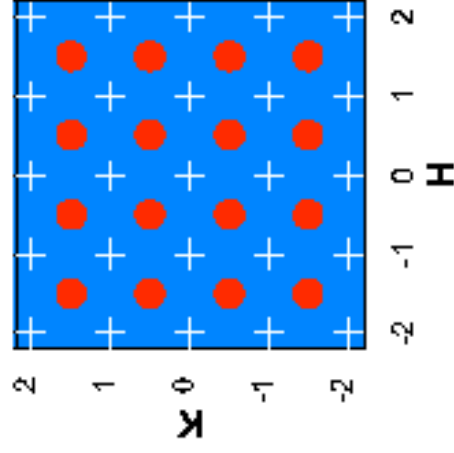


cc- θ_{II}

CC-phases
NMR
Silent Probe

Orbital moments (2D Ising) \rightarrow measurable with neutrons

In k-space: $Q=(H,K)$



Polarized neutron cross sections

- Potential due to magnetic field: $-\mu_N \cdot \mathbf{B}$

Magnetic field:

Spin Orbital

$$\mathbf{B} = \mathbf{B}_S + \mathbf{B}_L = \frac{\mu_0}{4\pi} \left\{ 2\mu_B \text{curl} \left(\frac{\mathbf{s} \wedge \hat{\mathbf{R}}}{R^2} \right) + I \frac{d\mathbf{l} \wedge \hat{\mathbf{R}}}{R^2} \right\}$$

- Elastic cross section for a given moment \vec{M} :

$$\frac{d\sigma}{d\Omega} = r_0^2 \left| \langle \pm | \vec{\sigma} \cdot \vec{M}_\perp | \pm \rangle \right|^2$$

$\vec{\sigma}$: Pauli matrices and $\vec{M}_\perp = \sum_r \exp^{-i\vec{Q}\vec{r}} (\vec{Q} \wedge \vec{M} \wedge \vec{Q})$

Neutron spin $\frac{1}{2}$: $|+\rangle$ and $|-\rangle$

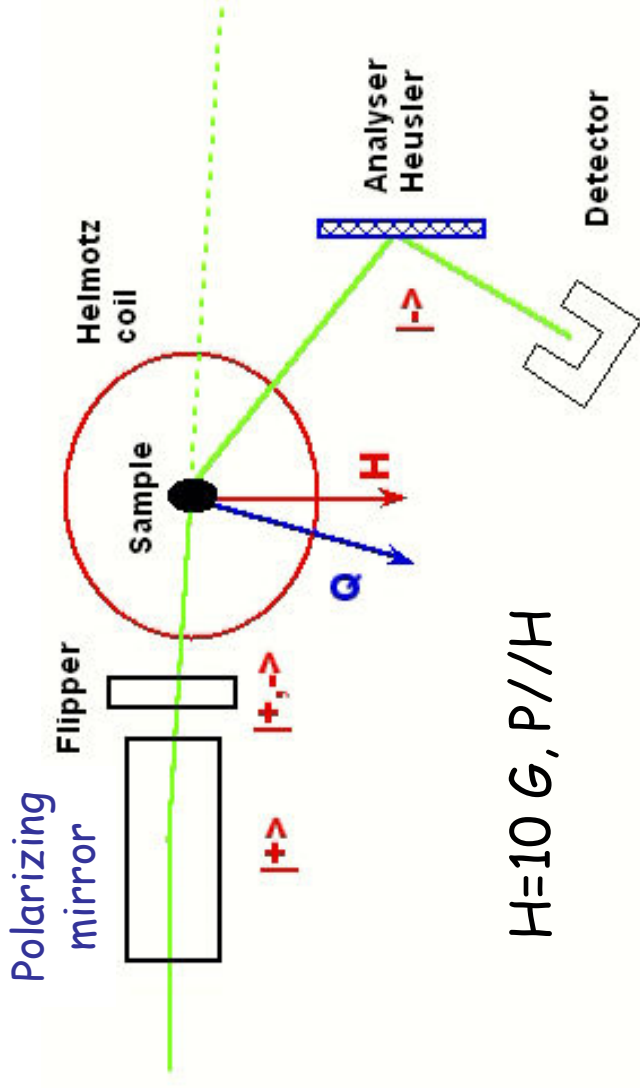
- For a polarization \vec{P} ($\parallel \vec{H}$) along a given direction z :

Non Spin-Flip: $\langle + | \vec{\sigma} \cdot \vec{M}_\perp | + \rangle = M_{\perp,z}$ and $\langle - | \vec{\sigma} \cdot \vec{M}_\perp | - \rangle = -M_{\perp,z}$

Spin-Flip: $\langle + | \vec{\sigma} \cdot \vec{M}_\perp | - \rangle = M_{\perp,x} + iM_{\perp,y}$

$\langle - | \vec{\sigma} \cdot \vec{M}_\perp | + \rangle = M_{\perp,x} - iM_{\perp,y}$

Polarized monochromatic neutron beam : 4F1-LLB



$H=10\text{ G}, P//H$

Flipping ratio:
 $R=NSF/SF=I^-/I^+ \quad (R \sim 50)$

- Nuclear Scattering
Non spin flip: $\leftarrow -|F_N| \rightarrow$
- * Magnetic Scattering

$$F_M = \langle \pm | \vec{\sigma} \cdot \vec{M}_\perp | - \rangle$$

$$\vec{M}_\perp = \vec{Q} \wedge \vec{M}_Q \wedge \vec{Q}$$

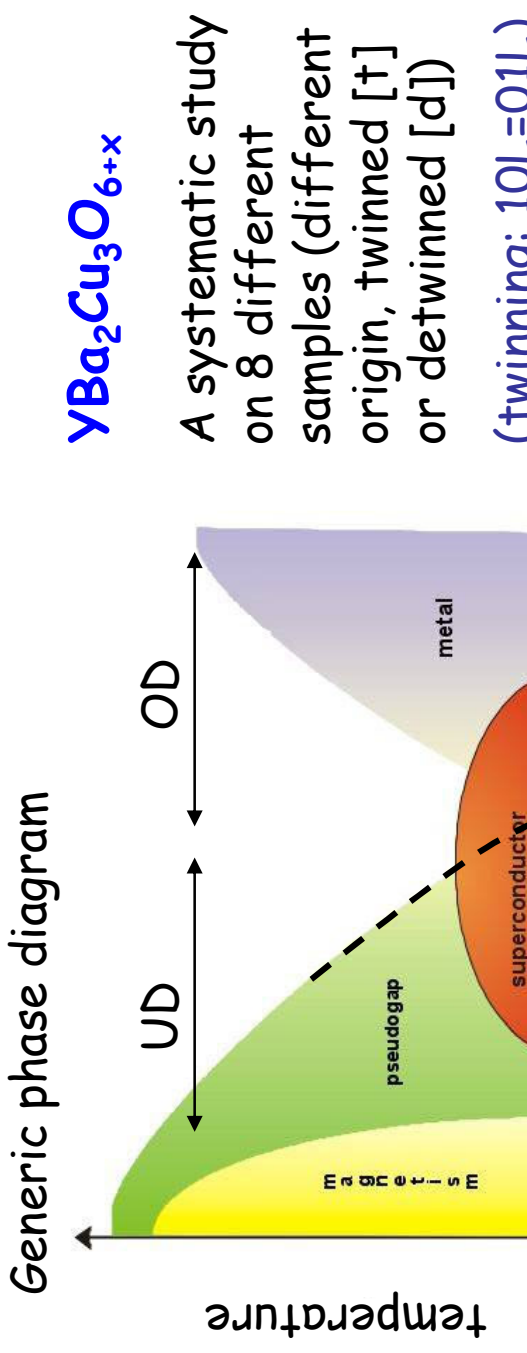
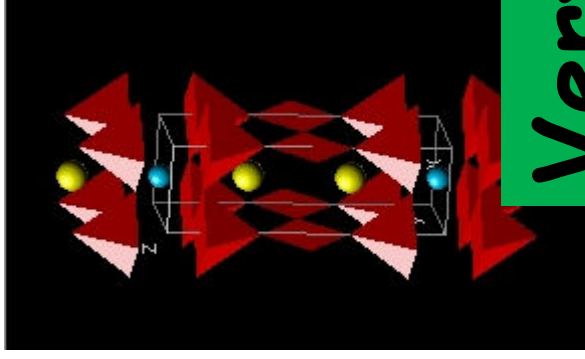
$$\vec{M}_Q = \sum \vec{M} \exp^{-i\vec{Q}\vec{r}}$$

Magnetic components $\perp \vec{Q}$
 Spin-flip components $\perp \vec{P}$

Neutron polarization:
 $p=(I^- - I^+) / (I^- + I^+) \sim 96\%$

Polarized neutron scattering experiments

on $(Y, Ca)Ba_2Cu_3O_{6+x}$



Very consistent picture

UD35 (x=0.45, $T_c=35$ K) [d]

UD54 (x=0.5, $T_c=54$ K) [t]

UD 61(x=0.6, $T_c=61$ K) [t]

UD64 (x=0.6, $T_c=64$ K) [d]

UD68 (x=0.75, $T_c=75$ K) [t]

UD89 (x=0.85, $T_c=89$ K) [t]

$Y_{0.85}Ca_{0.15}Ba_2Cu_3O_{6+x}$

OD75 (x=1, $T_c=75$ K) [t]

Underdoped $YBa_2Cu_3O_{6.6}$

H. Mook's sample ($T_c=63$ K) [t]

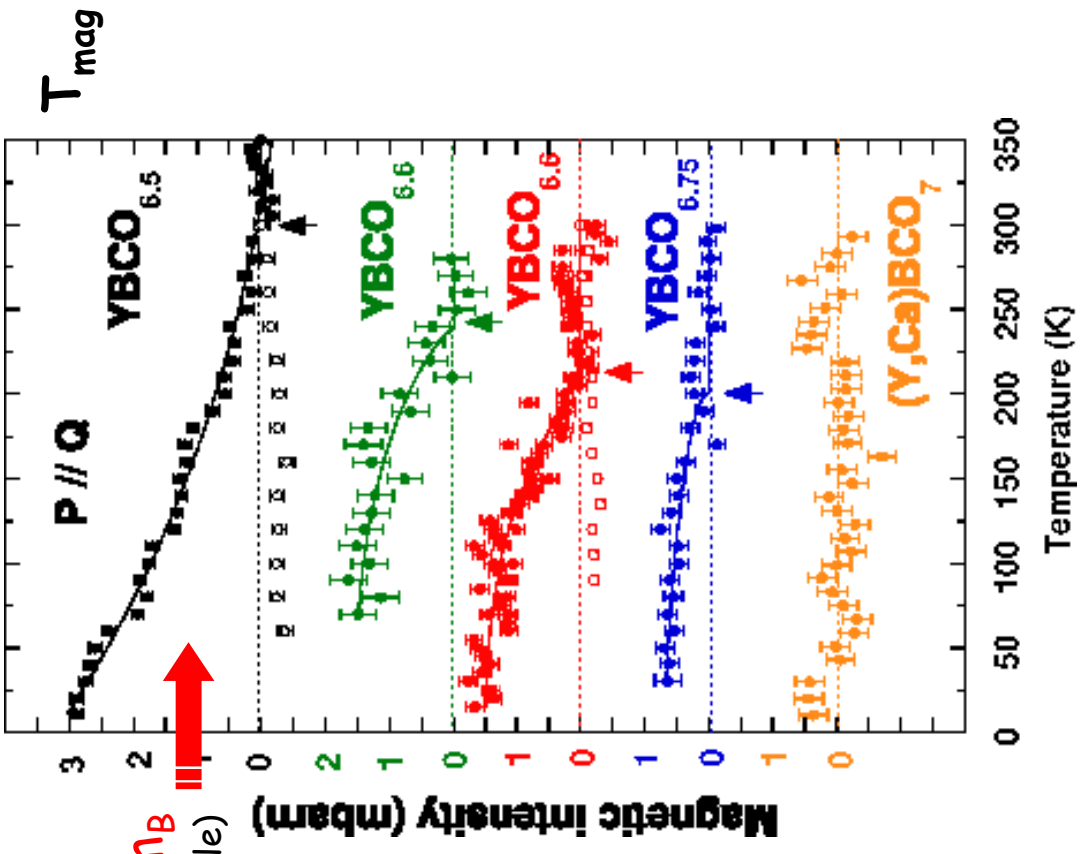
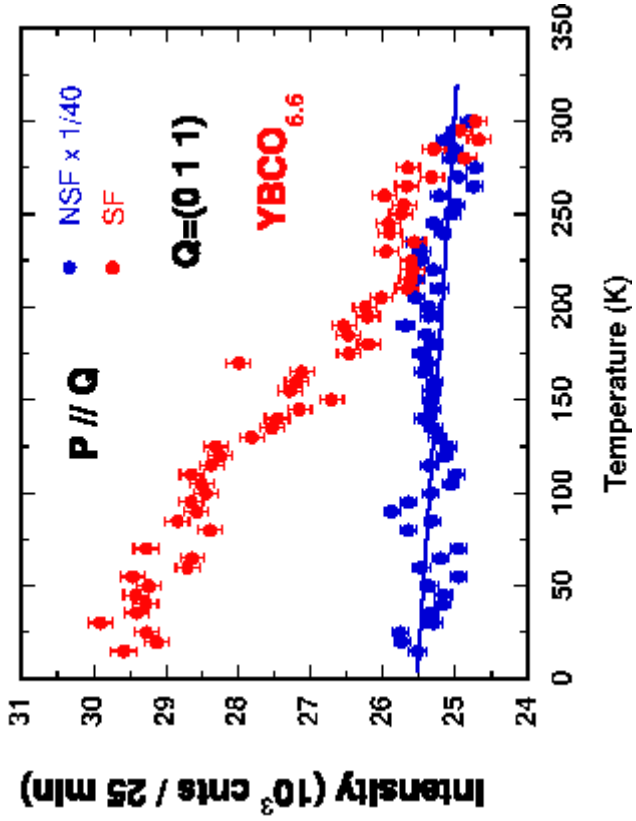
Polarized neutron results in YBCO: P//Q, Q=(0,1,1)
 B. Fauqué et al, PRL 96 197001 (2006)

$$\text{NSF: } \frac{d\sigma}{d\Omega} = |F_N|^2$$

$$\text{SF: } \frac{d\sigma}{d\Omega} = |F_M|^2 + |F_N|^2/R$$

R: flipping ratio

$m \sim 0.1 m_B$
 (per triangle)

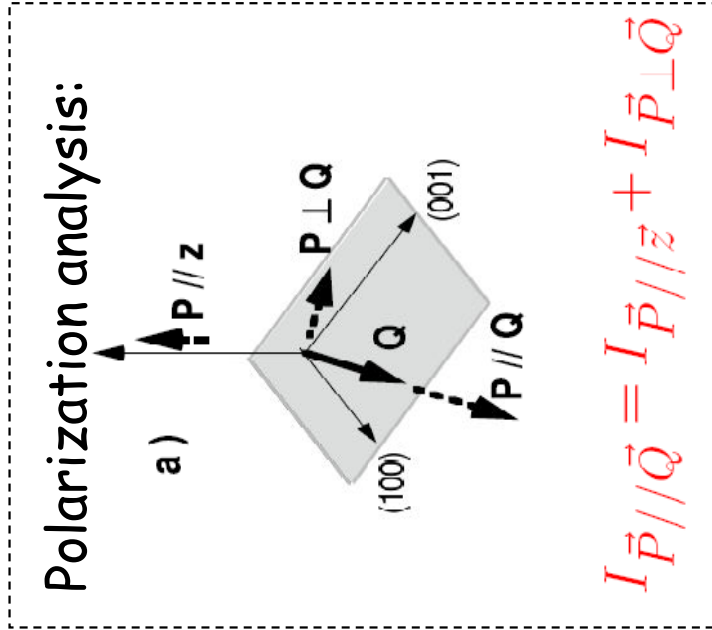
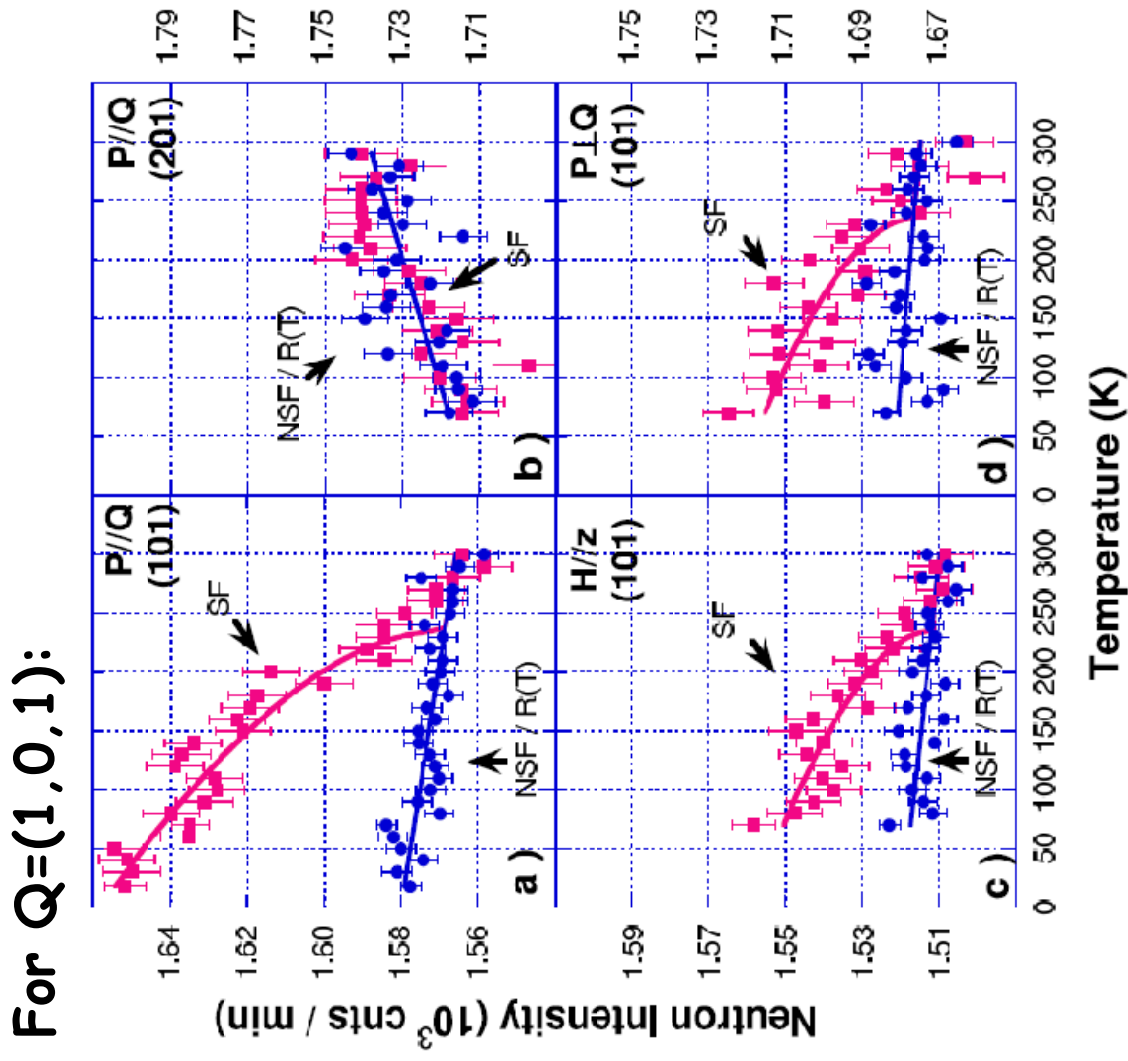


For Q=(0,1,1): $|F_N|^2 / |F_M|^2 = 400$

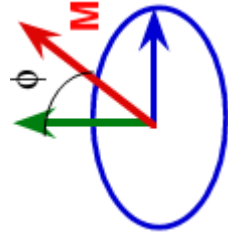
N.B. L=1, weak nuclear Bragg peak

No effect on Q=(002)
 (open symbols)

YBCO_{6.6} : H.A. Mook et al, PRB 020506(R) (2008).

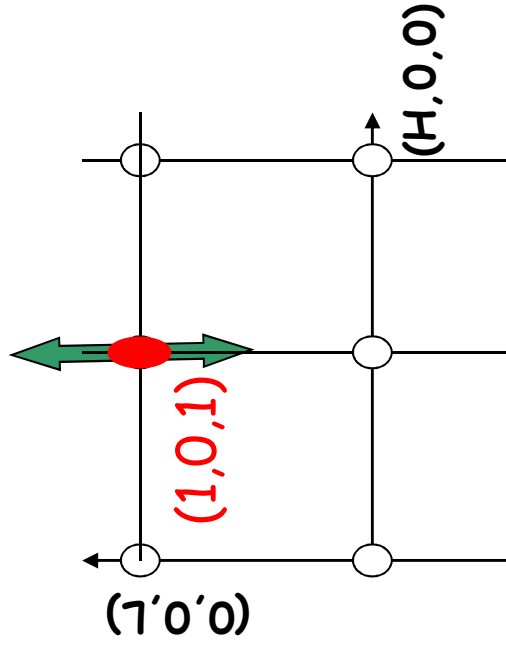


Angle (M, c^*) ~ 45 deg



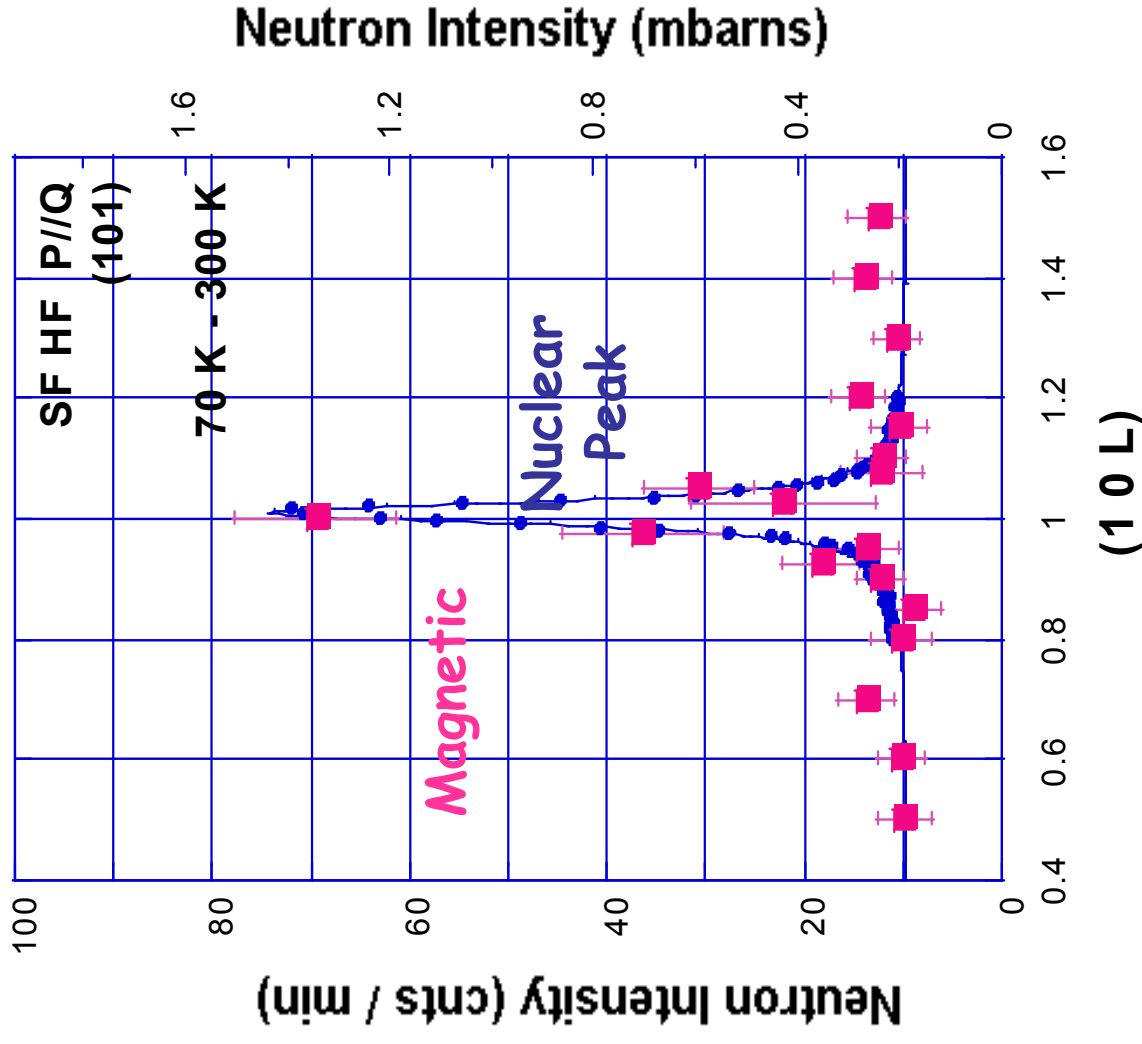
Weak or zero structure factor for $Q=(2,0,1)$

YBCO: Long range magnetic order



- L-scan

Long range order
($\xi > 50 \text{ \AA}$)

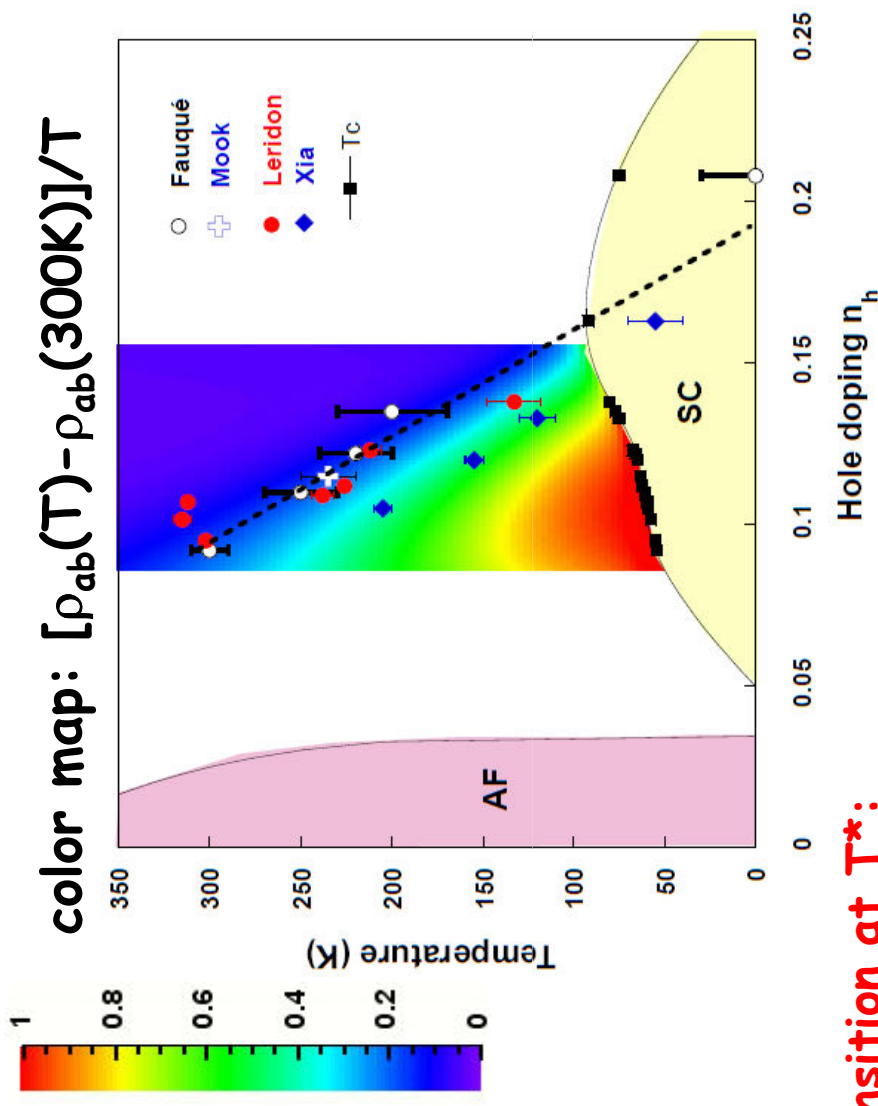
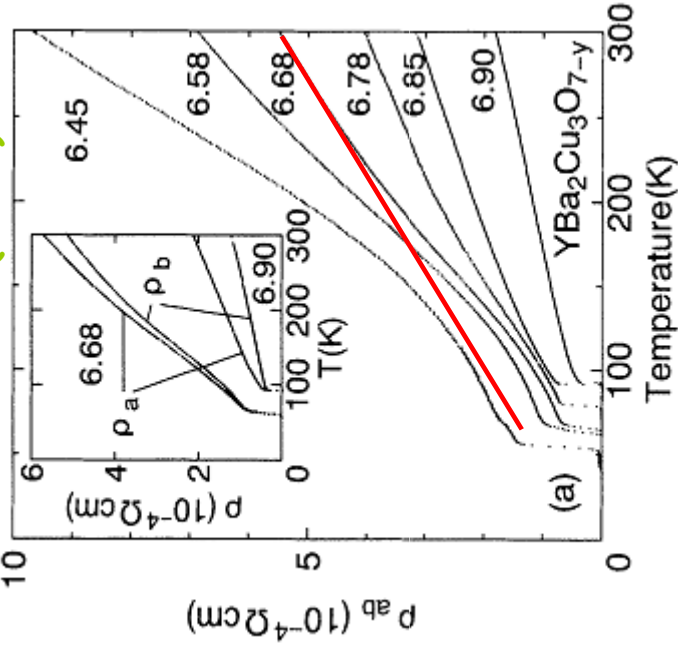


H.A. Mook et al, PRB 020506(R) (2008).

Phase diagram: Magnetic order at $T_{\text{mag}} \sim T^*$

Resistivity in YBCO

T. Ito et al PRL 70 (1993) 3995



Other reports of a phase transition at T^* :

- Polar Kerr effect

J. Xia, et al, PRL, 100, 127002 (2008)

- Uniform magnetic susceptibility

B. Leridon et al EPL, 87, 17011, (2009)

- Resonant ultrasound spectroscopy

Arkady Shekhter

B. Fauqué et al, PRL

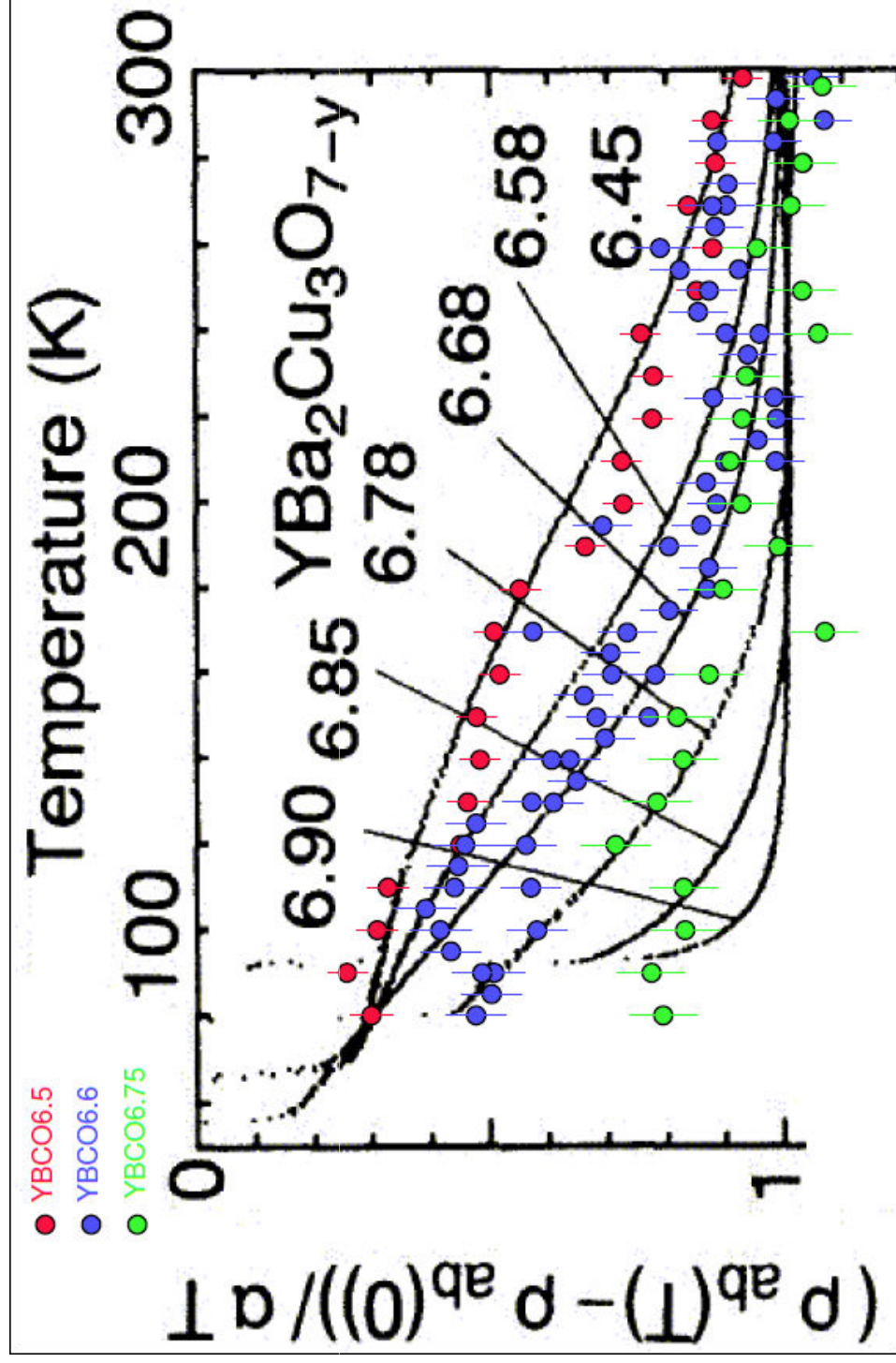
96 197001 (2006)

H.A. Mook et al, PRB 78

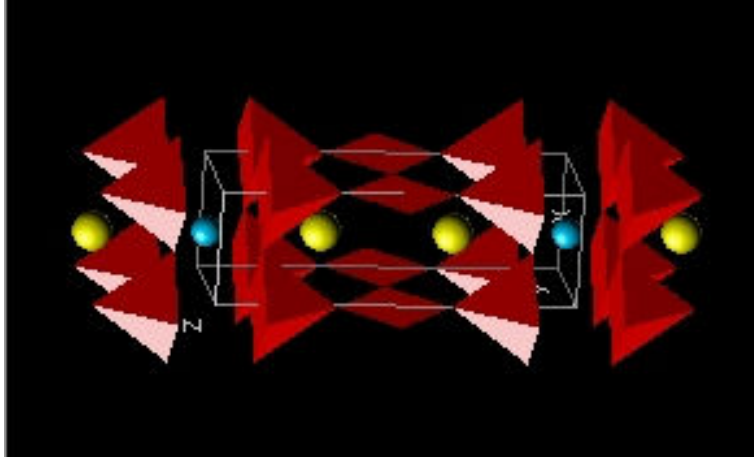
020506 (2008)

Comparison with resistivity

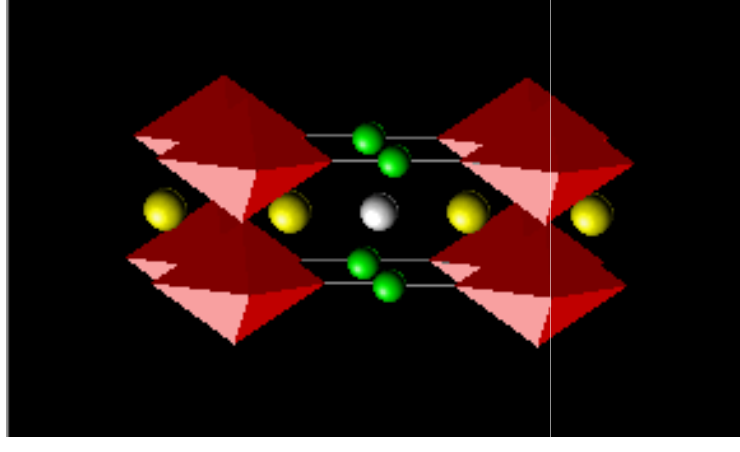
A single scaling factor with $|M|^2$ for all doping



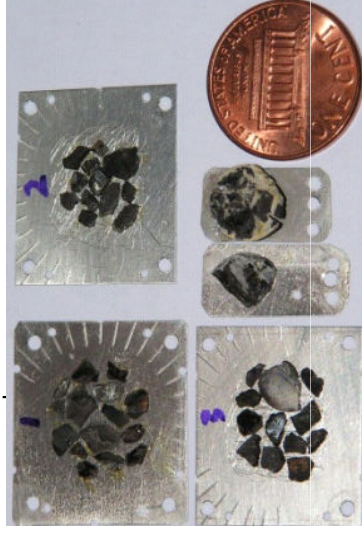
From $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ to $\text{HgBa}_2\text{CuO}_{4+x}$



- Orthorhombic
- 2 CuO₂ Planes
- distorted CuO₂ planes
- Cu-O pyramids



- Tetragonal
- 1 CuO₂ Plane
- Cu-O octahedra



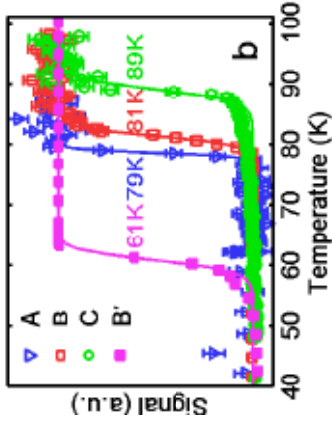
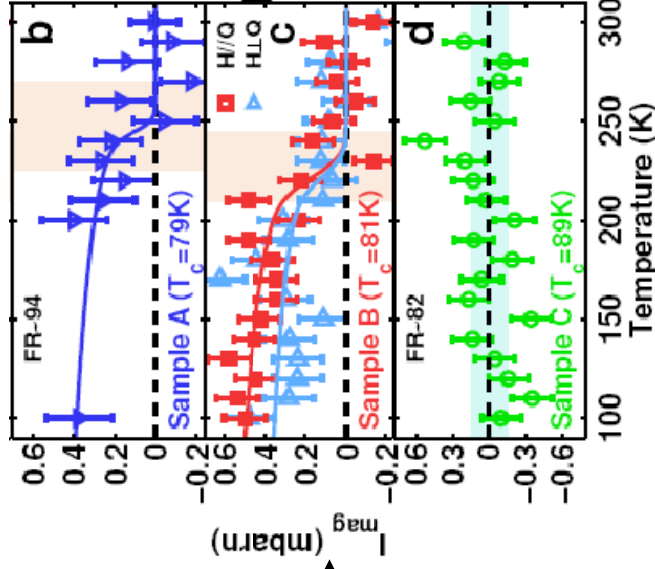
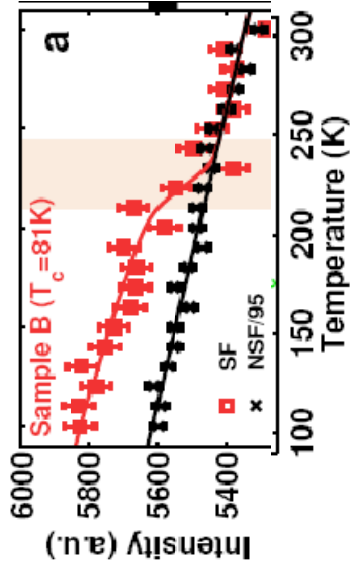
Breakthrough in single crystal synthesis in Hg1201:
 mass ~ 1 g, M. Greven (Stanford/Univ Minnesota)
 Zhao *et al.*, *Adv. Materials*, **18**, 3243 (2006).

Polarized neutron results in Hg1201

Yuan Li et al, Nature 455,372 (2008).

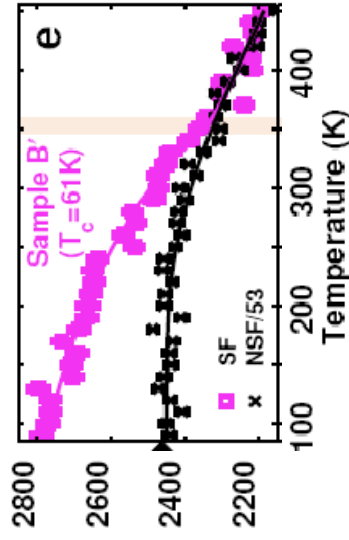
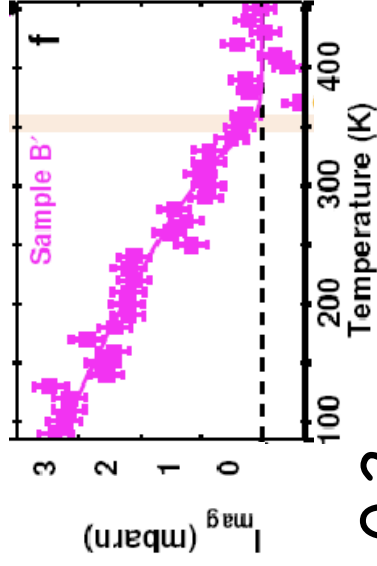
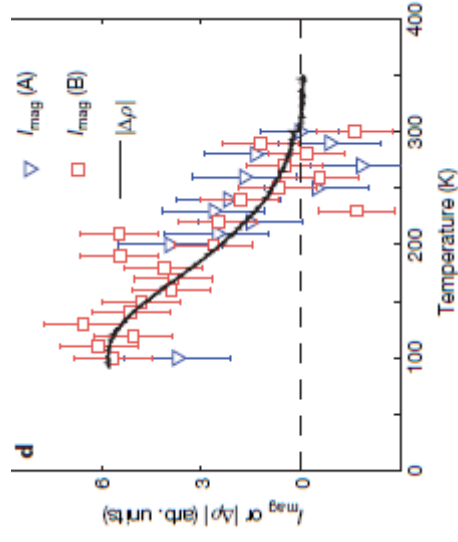
$Q=(1,0,1)$
P//Q

R > 90



Comparison Resistivity

$\Delta\rho$



$m \sim 0.2 \mu_B$

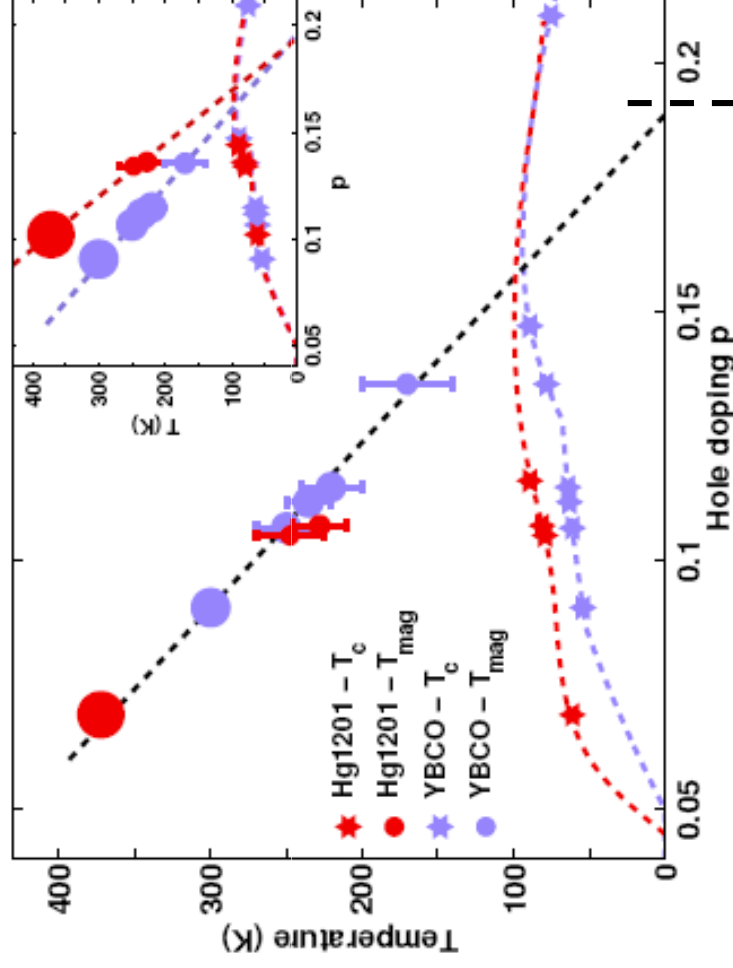
"Universal" phase diagram: from $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ to $\text{HgBa}_2\text{CuO}_{4+x}$

B. Fauqué et al, PRL 96 197001 (2006)

H.A. Mook et al, PRB 78 020506 (2008)

Y. Li et al, Nature 455 372 (2008)

$Q=0$ AF Magnetic order
at $T_{\text{mag}} \sim T^*$



p	T_c (K)	T_{mag} (K)	I_{mag} (mb)
Y 0.091	54	300 ± 10	2.8 ± 0.3
Hg 0.069 (0.102)	61	372 ± 13	4.2 ± 0.3
Y 0.107	61	250 ± 20	1.7 ± 0.2
Y 0.112	63	235 ± 15	1.6 ± 0.2
Y 0.115	64	220 ± 20	1.5 ± 0.2
Y 0.135	78	170 ± 30	0.6 ± 0.1
Hg 0.105 (0.134)	79	248 ± 23	0.5 ± 0.1
Hg 0.107 (0.136)	81	228 ± 18	0.6 ± 0.1
Hg 0.116 (0.144)	89	N/A	< 0.15
Y 0.147	89	N/A	< 0.2
Y 0.209	75	N/A	< 0.1

$T_{\text{mag}} \rightarrow 0$ at $p=0.19$, expected end point
of the pseudo-gap phase

NMR, Entropy, resistivity, ARPES....

Tallon and Loram, Physica C (2001,2004)

\rightarrow Quantum Critical Point

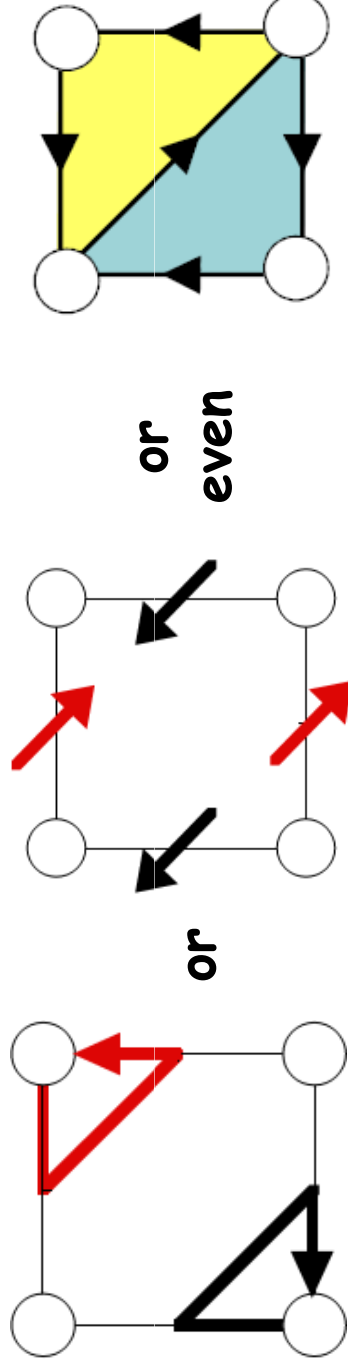
$Q=0$ AFM, What are the implications ?

- Intra unit cell magnetism (2 antiparallel moments)
- ➔ electronic nematicity by STM in Bi2212 M.J. Lawler et al Nature 2010
- Cu spins alone cannot account for this magnetic phase.
- One needs to reconsider the role of oxygens

Seamus Davis
talk

-Effective 1 band model ➔ 3 bands Hubbard model

➔ Orbital or spin magnetic moments ?



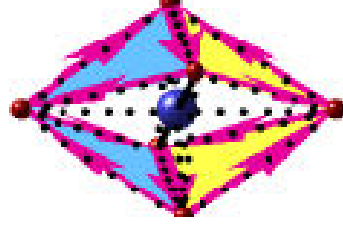
Stanescu, Philipps
PRB 69, 245104
(2002).

Chandra Varma
talk Spin moment on O sites (should have been detected by NMR.....)

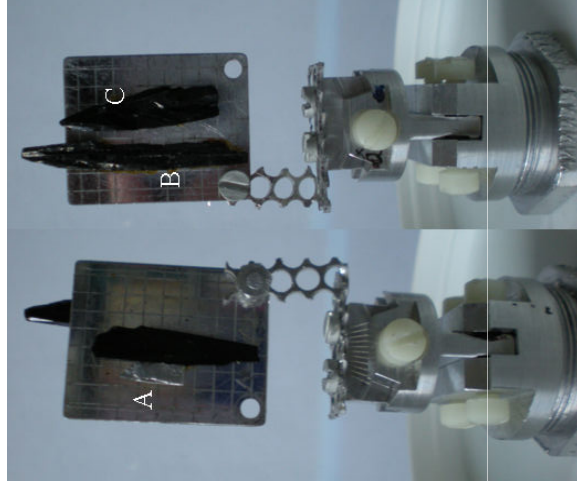
Quantum Monte-Carlo calculations

C. Weber et al, PRL 102, 017005 (2009).

➔ Angle (M, c^*) $\sim 45^\circ$ Origin ?

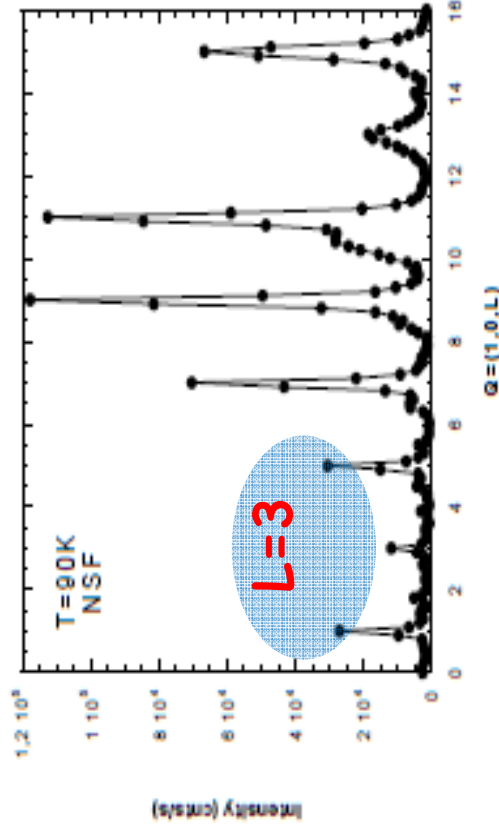
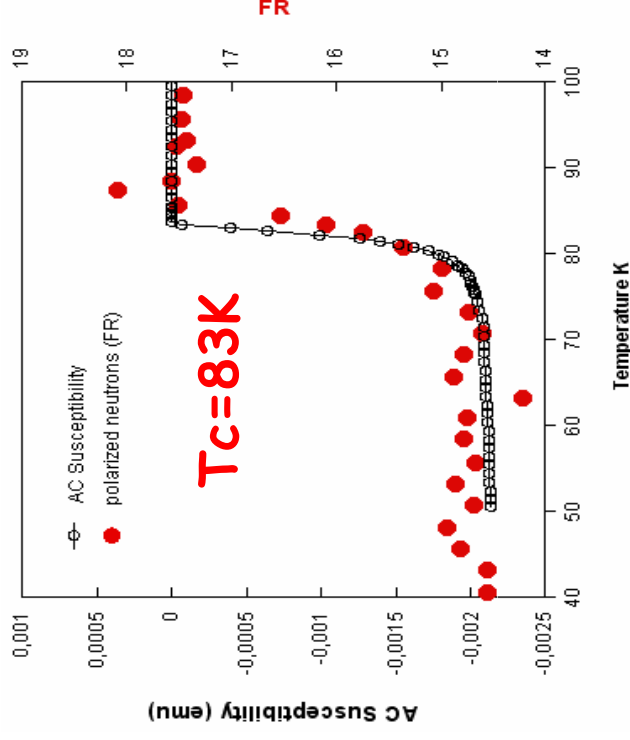


Preliminary results in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$



Large single
crystal growth

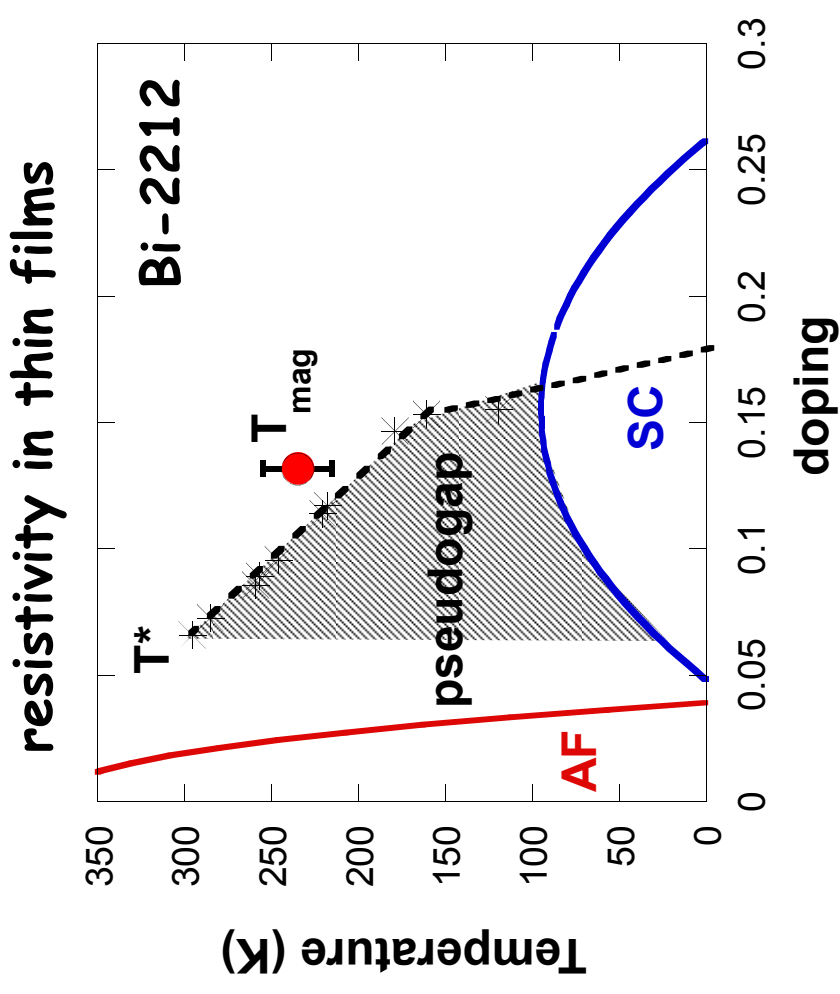
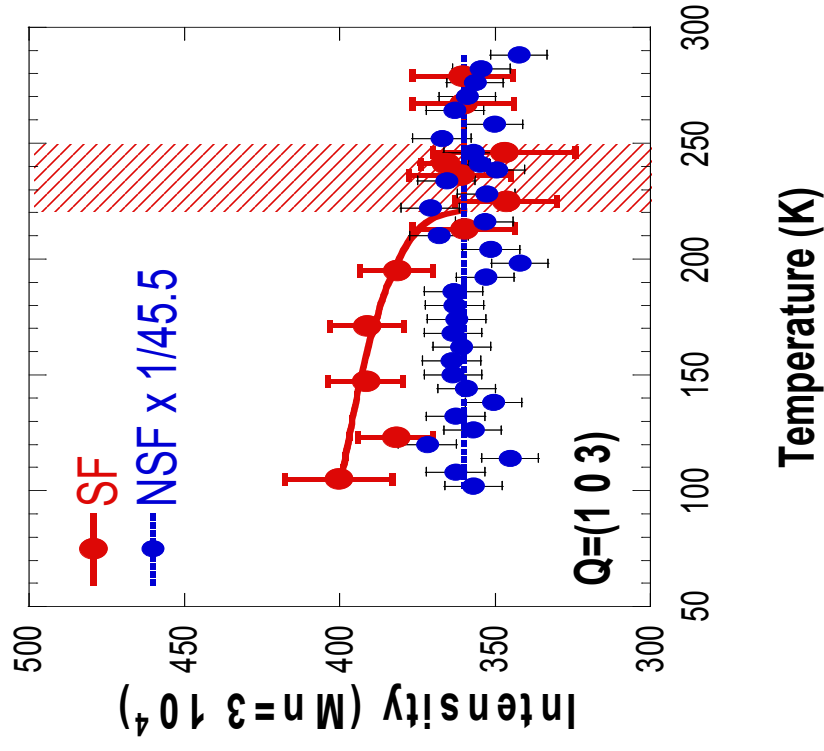
S. de Almeida
F. Giovanelli
I. Laffez
University of Tours
(France)



Annealed with $p(\text{O}_2) = 0.05$ atm at 450°C

Bragg peaks at (10L) for L odd
(Tetragonal notations)
45 deg to the inc modulation

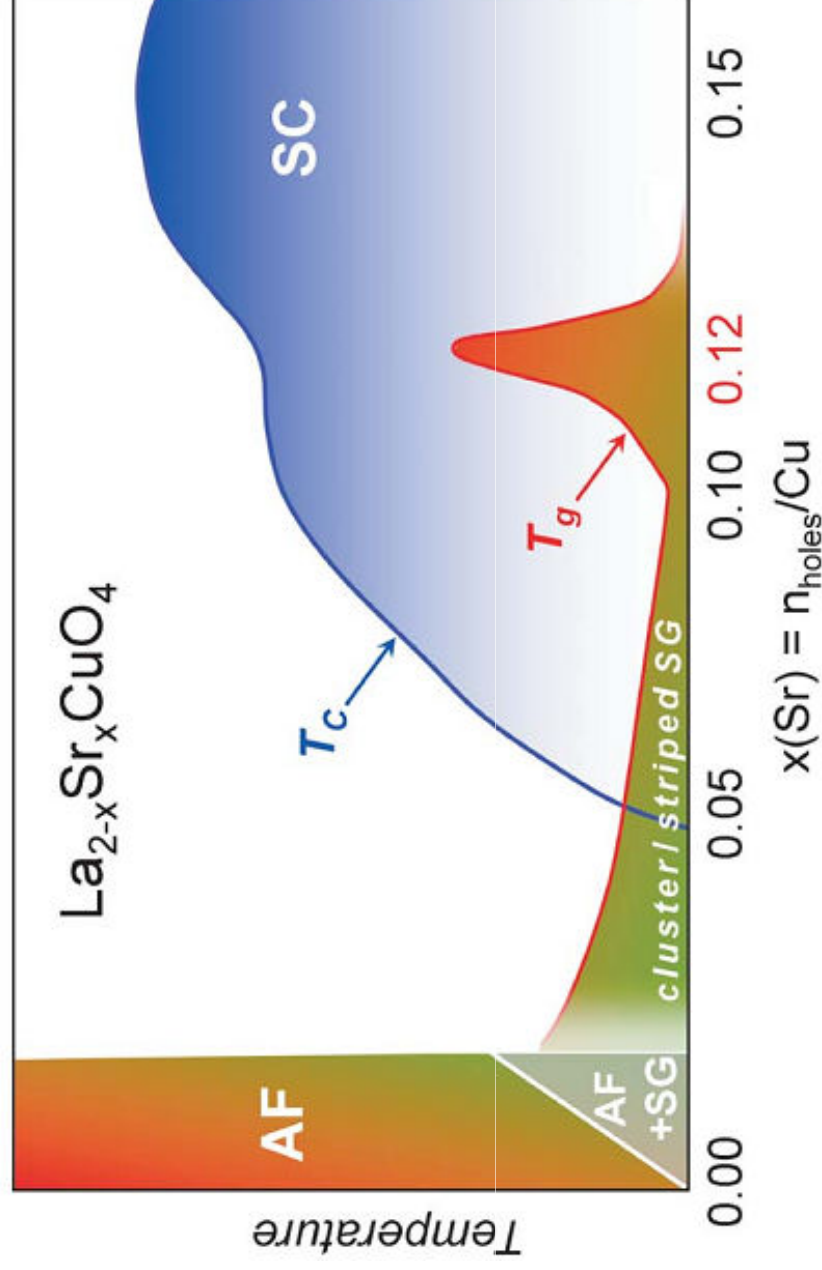
Q=0 AF Magnetic order in Bi-2212



A. Kaminski et al, PRL, 90 207003 (2003)

LSCO case: Charge inhomogeneities: Stripes

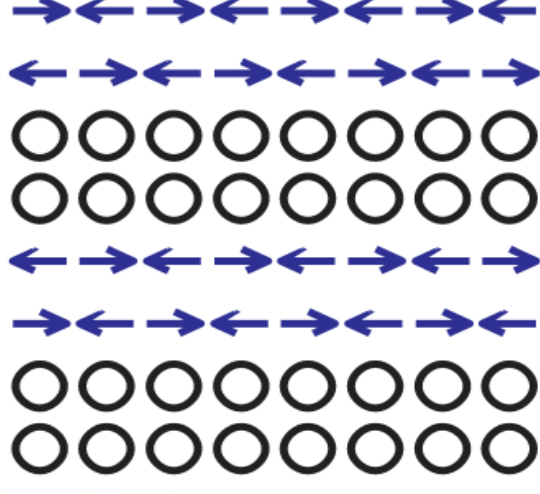
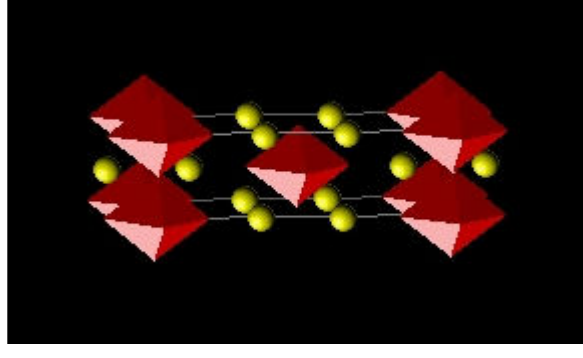
Mitrovic et al, PRB 78 014504 (2008) μ SR, NMR, ...



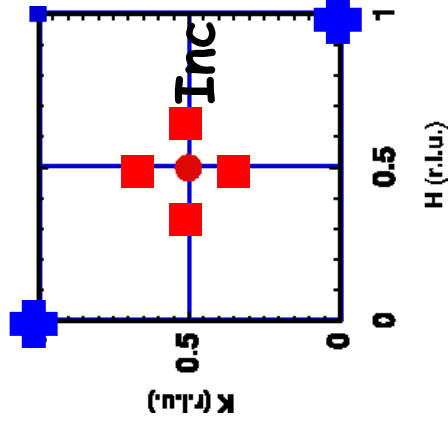
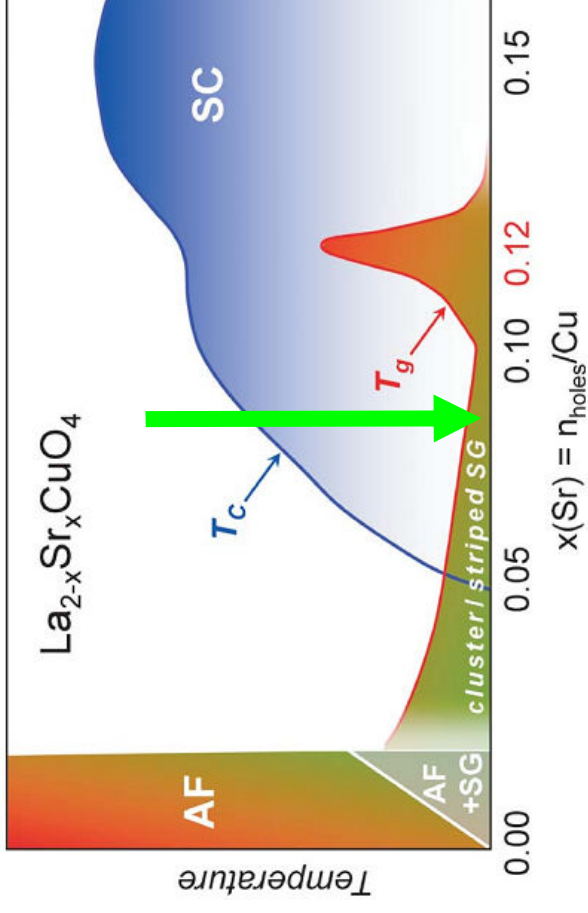
- Incommensurate AF fluctuations (Spin density wave)
- **→** Fingerprint of dynamical stripes
- Static stripe phase for Nd co-doping or with Ba

Tranquada et al, Nature 1995, 2004

Lattice invariance and fourfold symmetry broken



bond centered
charge stripes



3 Samples (7 g) (from PSI, E. Pomjakushina)
co-aligned with $x=0.085$, $T_c=22$ K.

Tetragonal notations →

Orbital-like moments at $Q=(10L)$

Search for Static Q=0 AF order in



Inverse Flipping

Ratio : R^{-1}

$(R^{-1} = \text{SF}/\text{NSF})$

$Q=(1,0,1)$

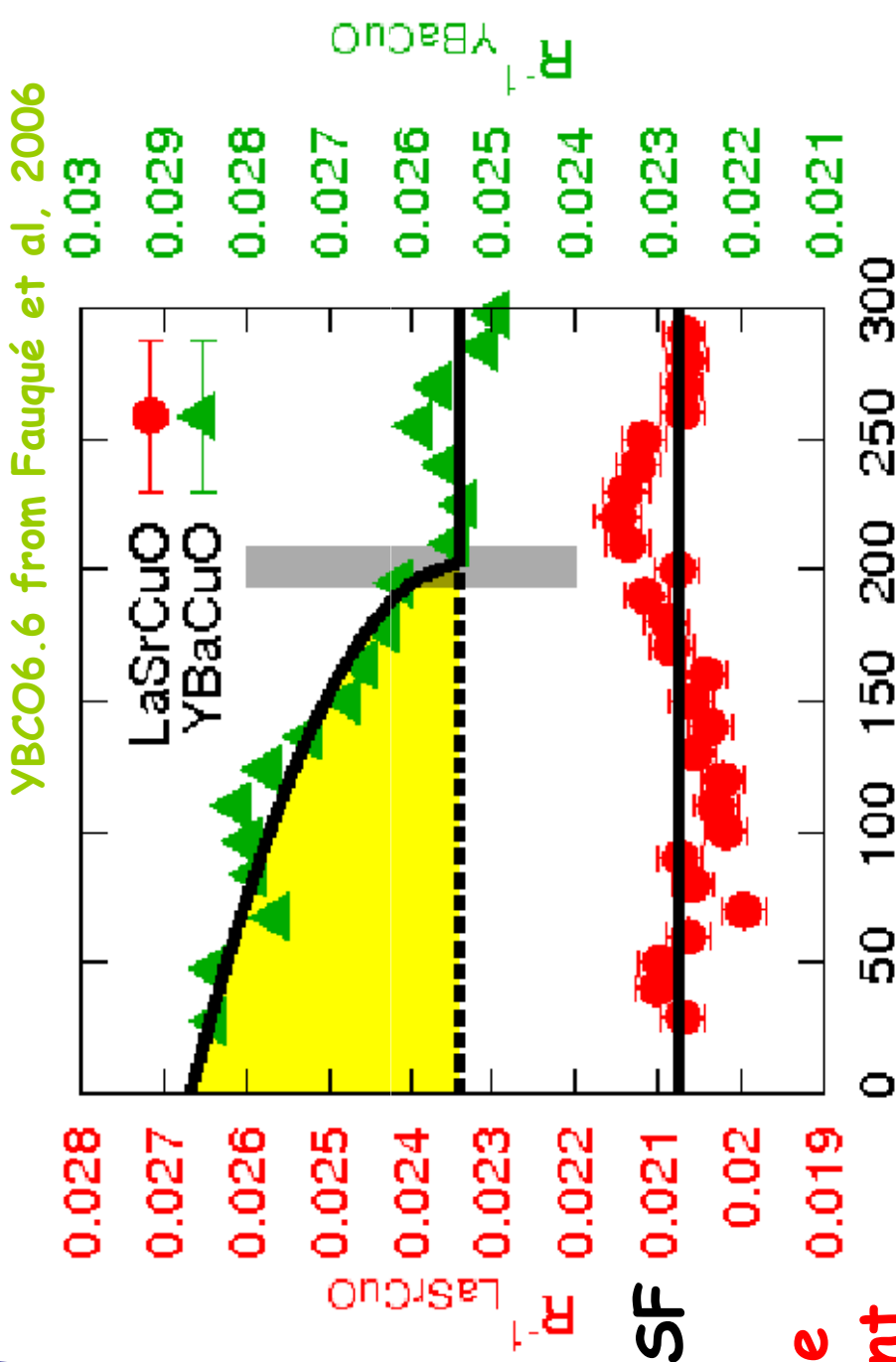
For $H//Q$

$R^{-1} = R_0^{-1} + M^2/\text{NSF}$

**No Long Range
ordered moment**

in LSCO

$(M < 0.02 \mu_B)$

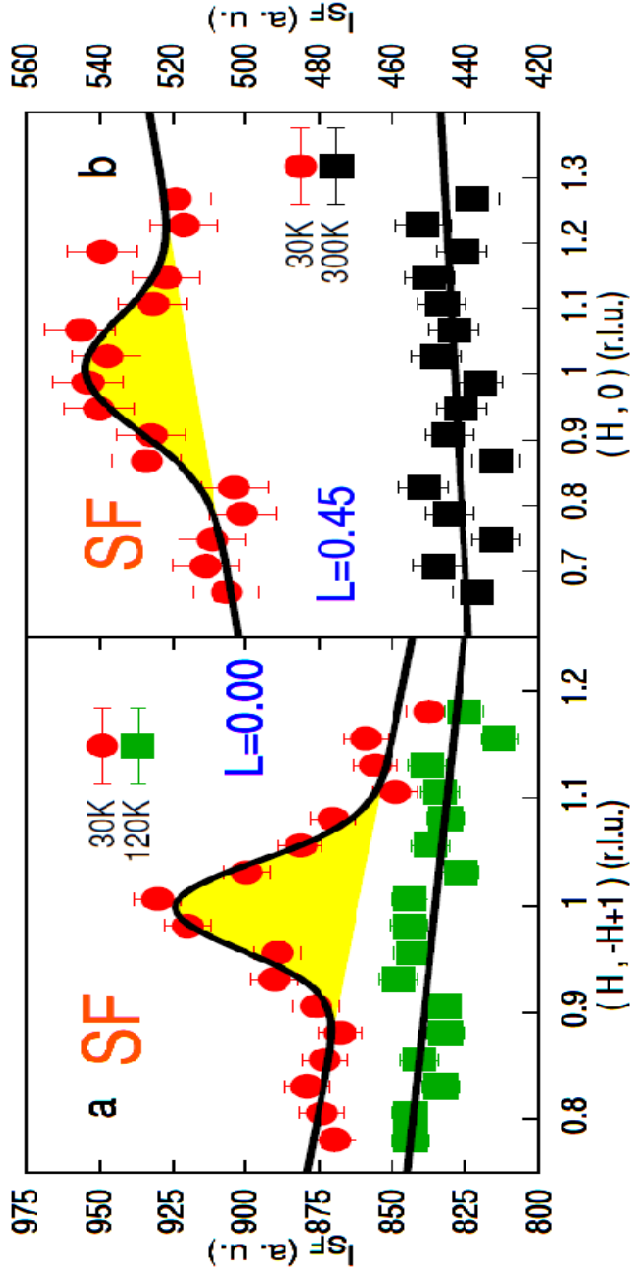


2D Short range magnetic correlations



Along 110

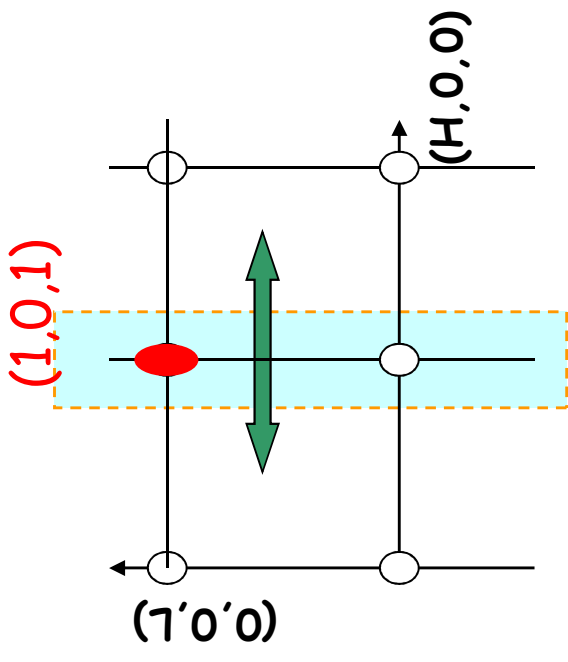
Along 100



In plane Correlations length: $\xi = 2/\Delta q$

$\xi_{110} \sim 11 \text{ \AA}$

$\xi_{100} \sim 8 \text{ \AA}$

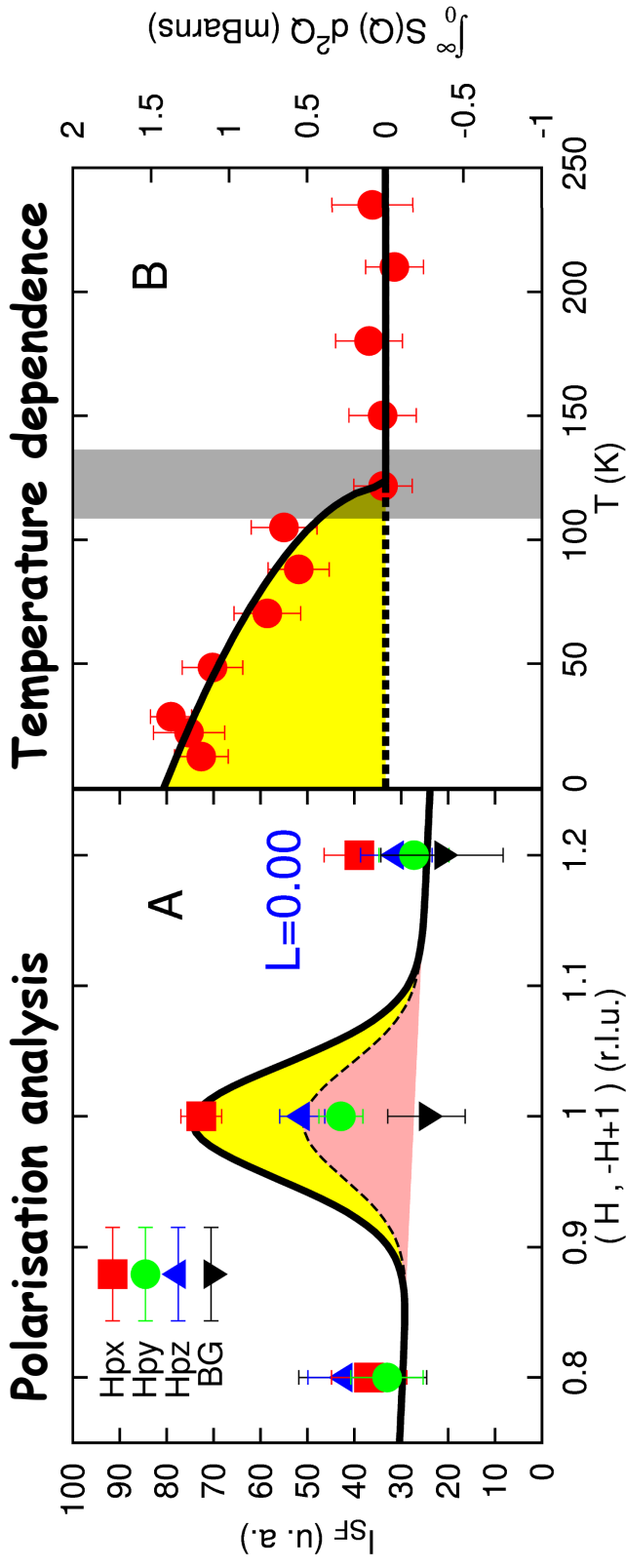


- 2D : Around $Q=(1,0,L)$ any L

- Spin-flip
- H//Q

- No signal in NSF channel

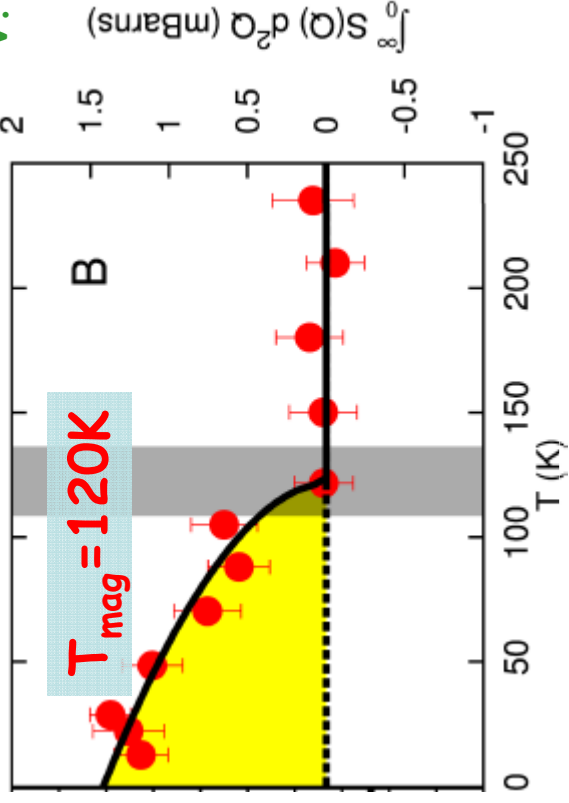
Polarization analysis and T-dependence :



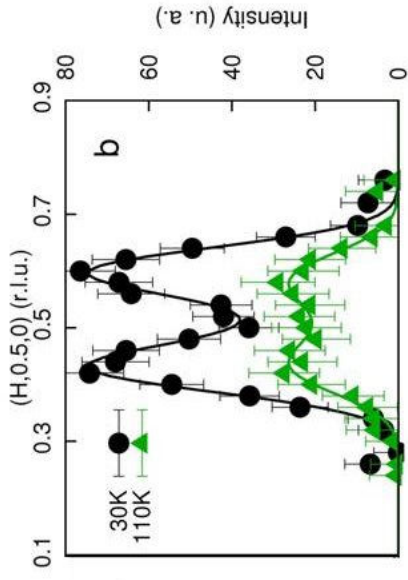
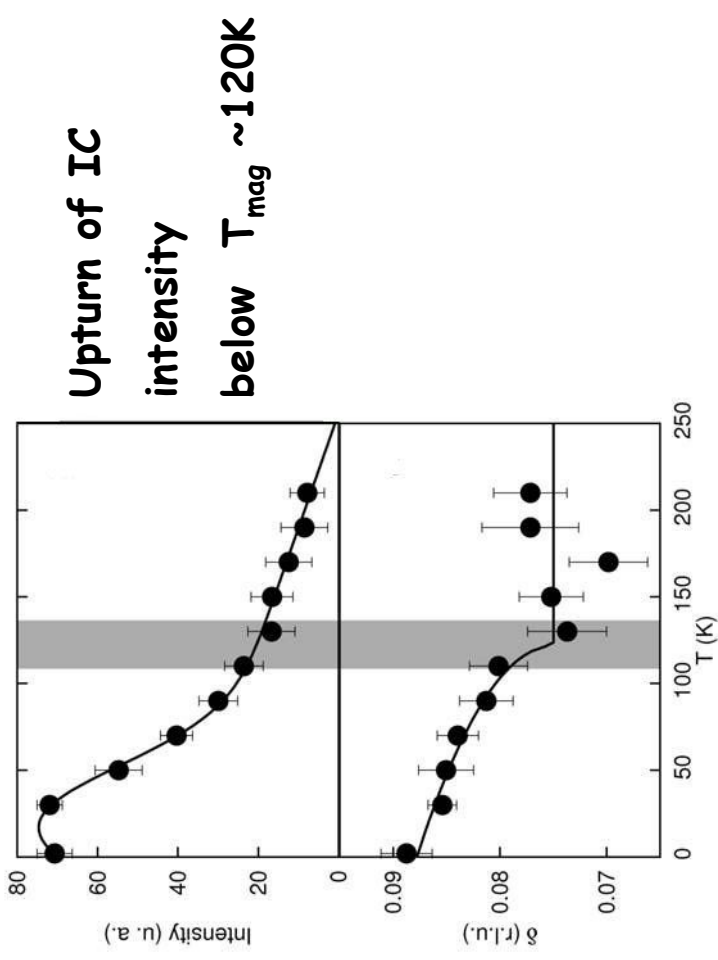
- $T_{\text{mag}} \sim 120 \text{ K}$ for 8.5% → Lower than generally expected T^*
- $I(H//Q) = I(H//z) + I(H//y)$
- $I(H//z) \sim I(H//y)$ → isotropic moments or angle (M, c^*) $\sim 45 \text{ deg}$
- Q-Integrated intensity (as Bragg scattering in $\text{YBCO}_{6.6}$)
1.2 mbarns → Local moment $\sim 0.1 \mu_B$

Temperature dependence of $Q=0$ AF order

V. Balédent et al, PRL 105, 027004 (2010)



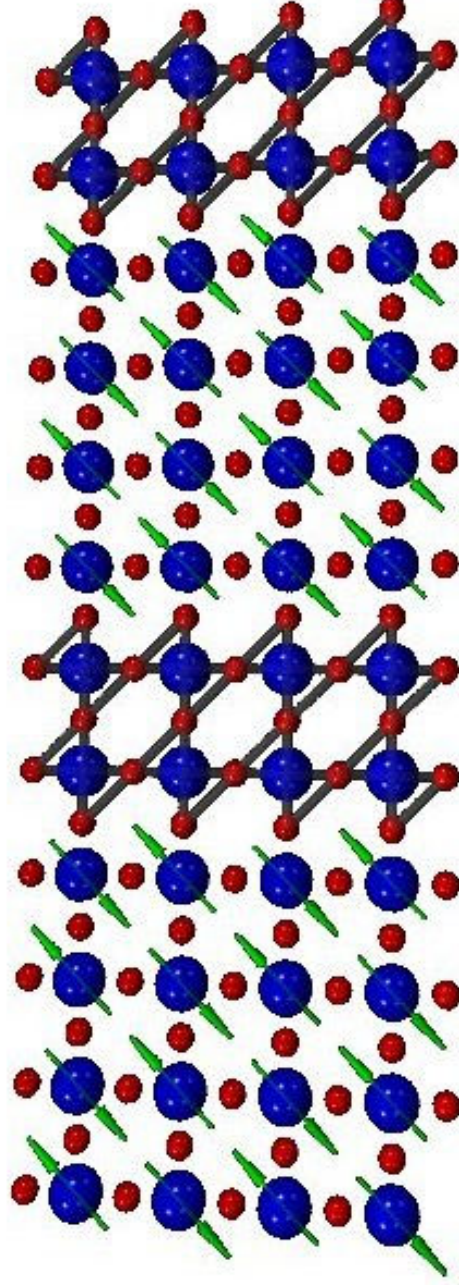
interplay between the $Q=0$
AF order and IC fluctuations



Unpolarized neutron $E = 4$ meV (2T-LLB)
Incommensurate magnetic fluctuations at $Q_{inc} = (1/2 \pm \delta, 1/2, 0)$

Mixte phase ?

For $x=0.085 \sim 1/12$



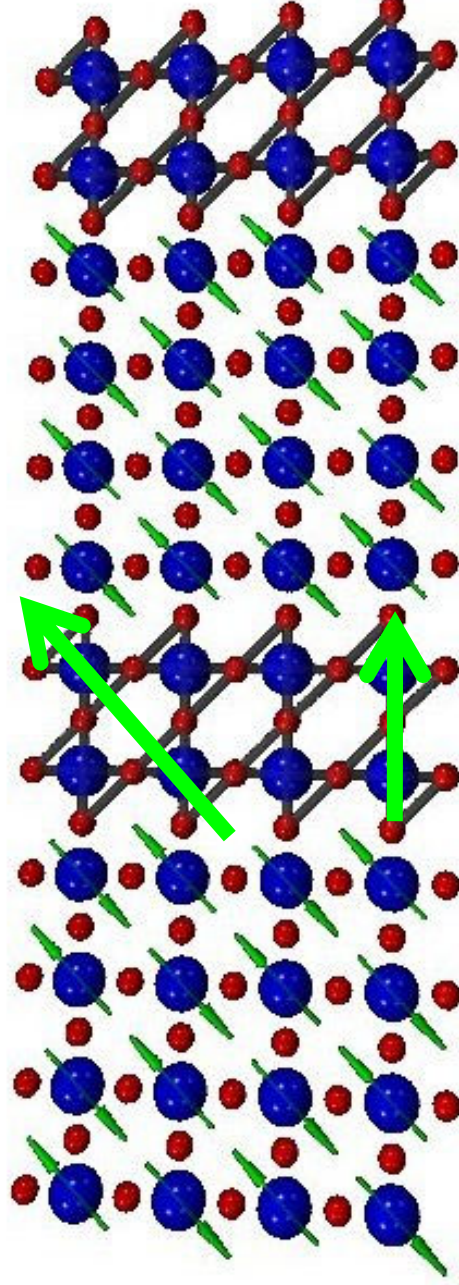
AF domains

2-leg charge
ladder

with circulating
currents

no phase coherence

$\xi_{(110)}$

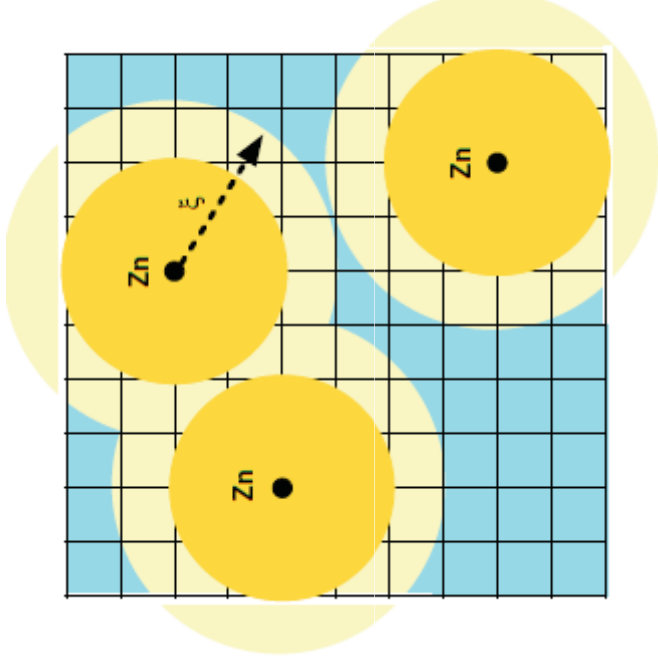
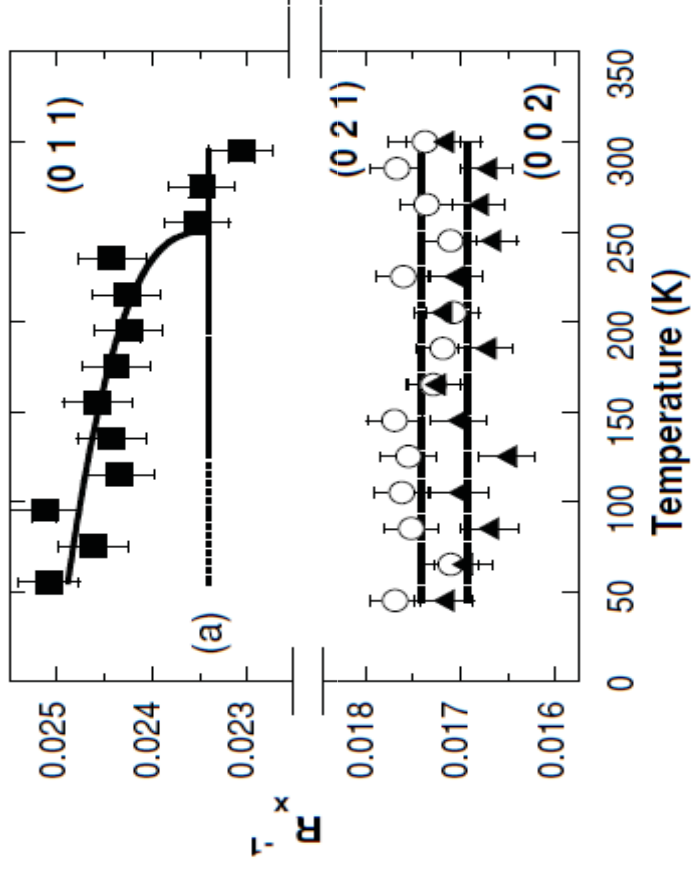


$\xi_{(100)}$

Zn(2%)-YBa₂Cu₃O_{6.6} : T_c = 30 K

Zn known not to change T* (NMR Knight shift, Alloul et al, 1993)

1/Flipping Ratio=R⁻¹=SF/NSF = R₀⁻¹ + M²/NSF



Swiss cheese model

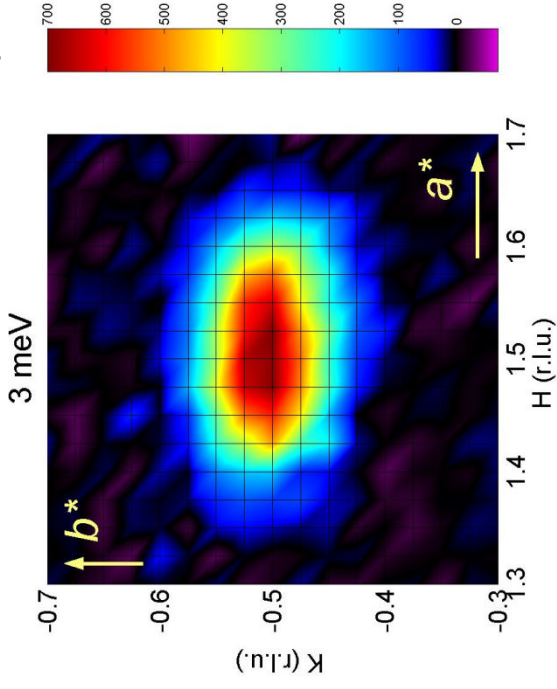
Same T_{mag} but weaker intensity

V. Balédent et al., *PRB* 83, 104504 (2011).

Zn induces low energy incommensurate spin excitations near (π, π)

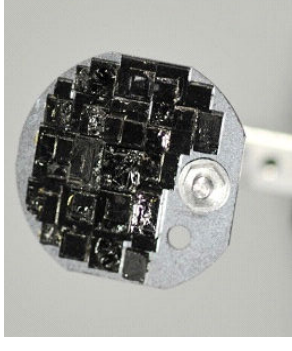
V. Suchaneck et al., *PRL*, 105, 037207 (2010).

Detwinned $\text{YBa}_2\text{Cu}_3\text{O}_{6.45}$



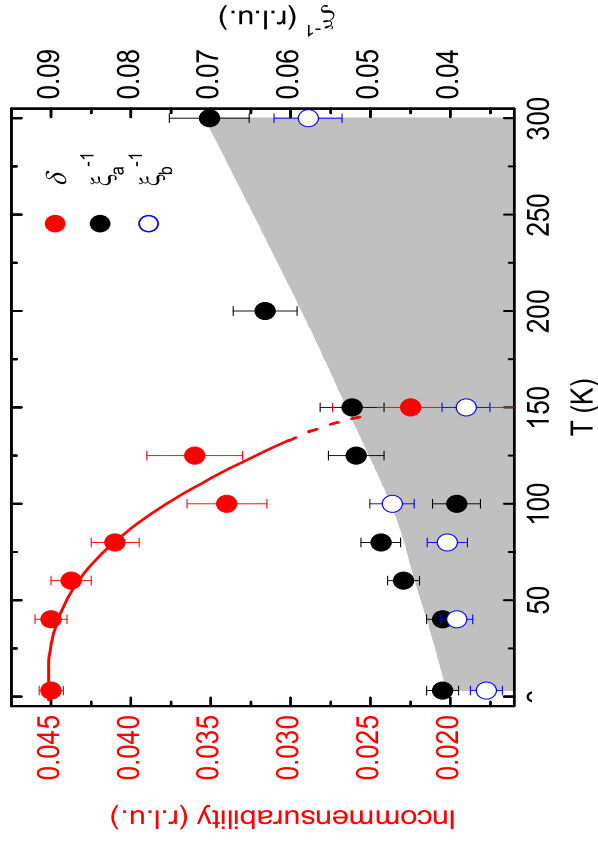
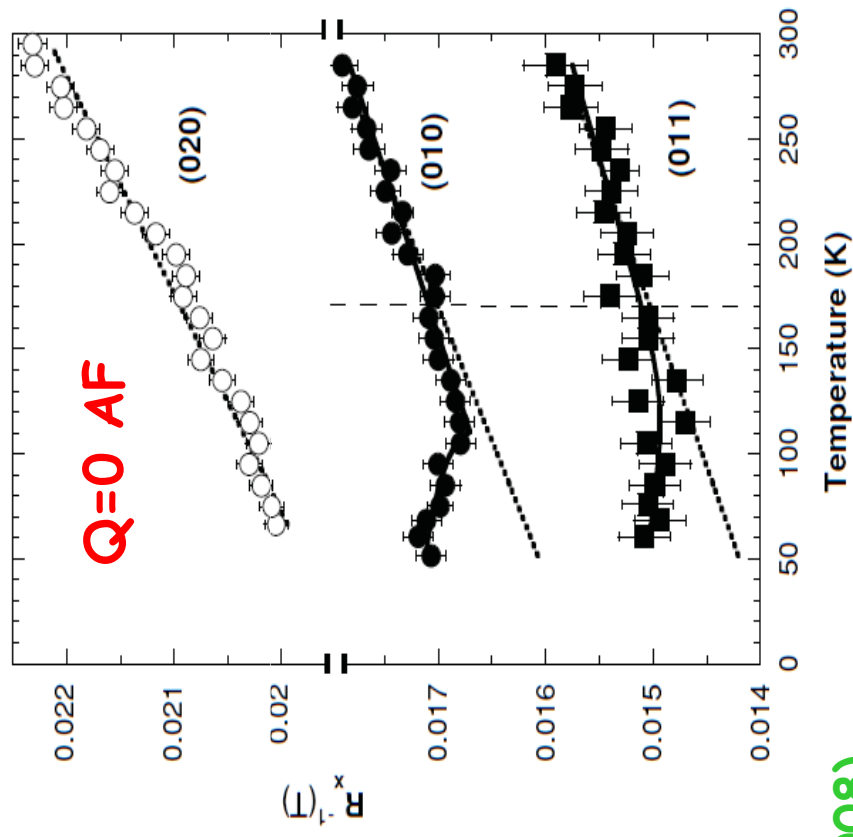
$T_c = 35\text{K}$, Hole doping = 0.085

(D. Haug, MPI Stuttgart)



$$R^{-1} = \text{SF/NSF} = R_0^{-1} + M^2/\text{NSF}$$

$Q=0$ AF



Hinkov et al., *Science* 319, 597 (2008).

\Rightarrow Electronic nematic liquid state (ELC)

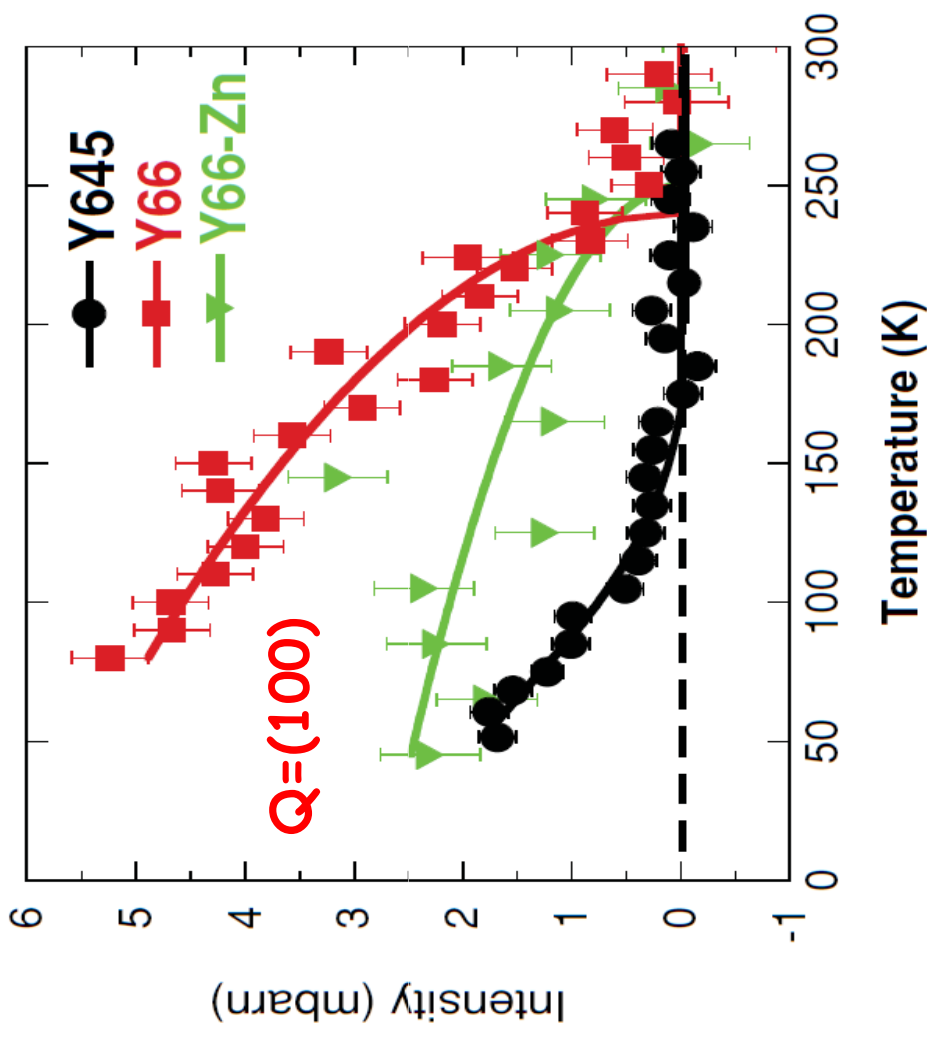
Comparison $Q=0$ AF Bragg peak intensity in YBCO

V. Baledent et al., *PRB* 83, 104504 (2011).

Near 9%
 T_{mag} decreases with doping

whereas

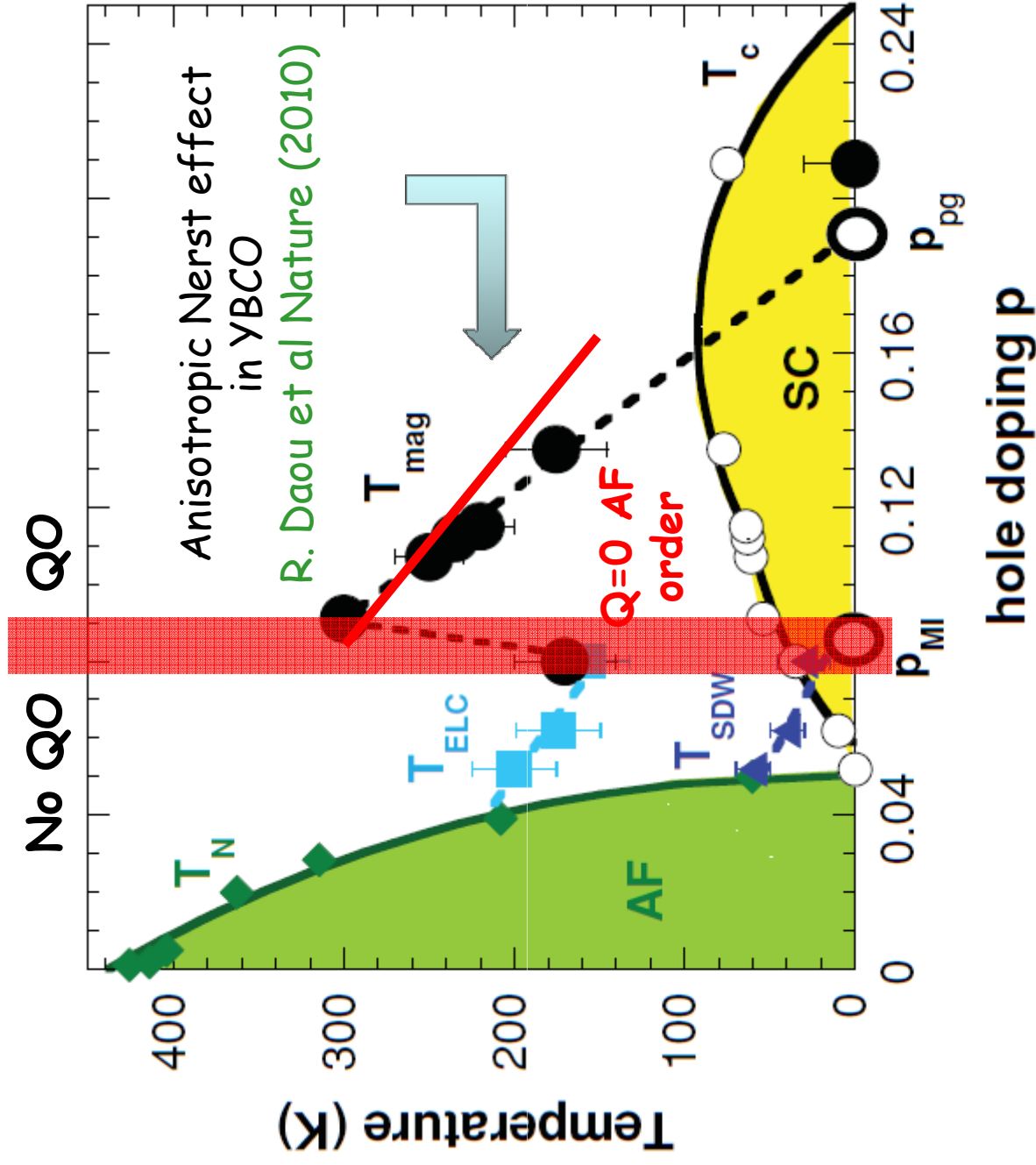
It is generally assumed
that T^* increases
(No compelling evidence that
this is correct for T^* above
Room Temperature)



IN YBCO6.45

$Q=0$ AF order develops close to $T_{\text{mag}} \sim 150\text{K} \sim T_{\text{nematic}}$

YBCO: Updated Phase diagram



V. Baledent et al., *PRB* 83, 104504 (2011).

Conclusions:

- *Technique: Polarized neutron diffraction*
 - *Result: Magnetic order in the pseudogap state of high-Tc cuprates in 4 different families: YBCO, Hg1201, LSCO, Bi2212*
 - *Implication: There is a broken symmetry below T^* which does not break the translation symmetry but breaks Time reversal symmetry*
Bourges and Sidis, arXiv: 1101.1786
- ⇒ Hidden order parameter of the pseudogap phase having the symmetry proposed for the loop current order
(talk Chandra Varma)
- ⇒ There are associated specific magnetic excitations
(next talk, Yuan Li)