

UNIVERSITY OF MINNESOTA  
ST. ANTHONY FALLS HYDRAULIC LABORATORY

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**An Evaluation of Acoustic Techniques for Measuring  
Gas Bubble Size Distributions in Cavitation Research**

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## ABSTRACT

The acoustic tone burst attenuation technique for measuring the gas bubble size distribution in water was critically reviewed. The findings reveal that the usefulness of the technique is limited to the bubble volume concentration range between 0.03 and 1.0 parts per million. Furthermore, resolution between size ranges is probably not possible with sufficient accuracy for many purposes.





AN EVALUATION OF ACOUSTIC TECHNIQUES FOR MEASURING  
GAS BUBBLE SIZE DISTRIBUTIONS IN CAVITATION RESEARCH

I. INTRODUCTION

The need to measure gas bubble size distribution in water tunnel research has been dealt with in the literature for a number of years. Many techniques and methods based on the principles of light, acoustics, gravitational separation, electronics, and photography have been investigated [1,2,3,4]\*. In recent years technological progress in optics and acoustics has provided an impetus to pursue measurement techniques based on the attenuation and scattering of sound and light for application to the problem of cavitation.

A good deal of effort has been expended at the St. Anthony Falls Hydraulic Laboratory in studying the basic phenomena of what has been variously called limited, bubbly, and transient cavitation [5]. Briefly, in this process the trajectory of an entrained nucleus carries it into a region of fluid which is superheated by virtue of the low dynamic pressures.

The basic assumption has been made in these and other tunnel investigations that these nuclei are primarily gas filled microbubbles stabilized by the turbulent processes in the tunnel. Under certain conditions, when the bubble is very small the surface tension forces prevent vaporous expansion. At some critical size these forces are overcome and the liquid rapidly changes phase at the nucleus wall, which rapidly expands the bubble.

In conjunction with this fundamental work with the physics of cavitation, methods of measuring the bubble size distribution--deemed critical to an understanding of the process by many researchers--have been sought. The effect of the entrained microbubbles upon the acoustic properties of the water seemed to offer distinct possibilities for measuring this distribution.

Bubbles have primarily two effects on the acoustic properties of the entrained liquid: they generally tend to decrease the sonic velocity and they produce an attenuation due to the absorption and scattering. Initial efforts at this laboratory utilized the change in the sonic velocity as an indication of the total free gas concentration [6]. The two main difficulties with this measurement were that the change in sound speed was very

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\* Numbers in brackets refer to references listed on page 9.

small at the free gas concentrations of interest (a few parts per million or less) and that the measurement gave only a total concentration with no indication of how the sizes were distributed. As knowledge of cavitation grew it became apparent that this total concentration of free gas bubbles was insufficient for cavitation research and that gas bubble distributions were required. The St. Anthony Falls Hydraulic Laboratory then began developing an instrumentation technique based on the attenuation of an acoustic pulse as it is transmitted through a liquid with entrained microbubbles [2] and on the fact that only bubbles whose resonant sizes are at or near the carrier frequency of the acoustic pulse contribute to the attenuation of that pulse. This is the technique to be evaluated in this report. It will be analyzed in relation to the requirements of the cavitation problem and the measurement capabilities of the existing acoustic instrumentation system.

## II. CRITERIA FOR A BUBBLE SIZE DISTRIBUTION MEASUREMENT SYSTEM

The output of any system designed to measure the size distribution would most likely take the form of a histogram with either constant or variable size ranges. In reality, of course, the size distribution is continuous, and questions arise regarding the accuracy of the bubble concentrations found in each size range and the width of size range required to adequately describe the real process.

It is first desirable to establish a set of requirements, based on a cavitation problem of interest, which a measuring instrument would have to satisfy. To accomplish this a typical cavitation example will be postulated and used to determine the accuracy of the measurement system.

The physical model to be used in the example is an axially symmetrical blunt body  $5/8$  of an inch in diameter in a six-inch water tunnel. The tunnel velocity will be 30 fps and the water temperature  $70^{\circ}\text{F}$ . The shape of the test body has been derived mathematically by placing a source disk in a uniform flow field and determining the streamline through the stagnation point. This body is shown in Fig. 1 along with its streamline and isobar patterns in the fluid. The advantage of using a body of this type is that the problem is amenable to analytical investigation.

The size distribution of the entrained microbubbles is, for the purposes of this example, described by the continuous function

$$p(r) = (r/\delta^2) \exp(-r^2/2\delta^2)$$

where  $r$  = bubble radius,  $\delta$  = most probable size, and  $p(r)$  = probability density; this distribution is shown in Fig. 2. The distribution is also assumed to be spatially homogeneous, which means that the statistical characteristics of the microbubbles entrained in any stream tube in the flow are identical. The total number of nuclei in the distribution or the total volume concentration will be varied in order to illustrate certain features and limitations of the acoustic measurement technique, but the distribution of the sizes will be considered typical and constant in the entire sample.

An analytical model which has been developed to relate the distribution of the entrained bubbles to the expected cavitation result [5] will be used. This model predicts the severity of cavitation in terms of cavitation occurrence rates based on the hydrodynamic characteristics of the test body, the pressure and velocity of the test water, and the numbers and sizes of entrained microbubbles. In this connection the cavitation occurrence rate is defined as the number of times per second that a nucleus enters the superheated region around the test body and cavitates. The number of cavitation events contributed by each stream tube is equal to the number of microbubbles flowing in the stream tube which exceed the critical size determined by the static stability theory. A good deal of experimental evidence has been accumulated concerning the cavitation occurrence rates associated with similar bodies. This experience will be used in selecting the number of bubbles in the distribution so that typical cavitation rates will occur.

It has been shown that the pressure gradient in front of the body deflects the bubbles out of the stream tube in which they were originally entrained [7]. This effect can modify the results somewhat, especially with small blunt bodies. However, the differences are expected to be small and to decrease as the body becomes more streamlined and as the relative size of the test body increases with respect to the bubble sizes. For the purposes of this example this effect will be neglected.

The distribution of sizes in this example was selected on the basis of previous experience with acoustic measurements. The number of bubbles in the entire range has been adjusted so that typical cavitation occurrence results are obtained. The analytical model will thus produce results which conform closely to observed experimental data.

The critical bubble size above which all bubbles can be assumed to cavitate in each stream tube was determined for several cavitation numbers. The distribution was integrated from the critical size to infinity to obtain the expected number of cavitation events per second in each stream tube. Summation of the contributions of all involved stream tubes at a particular cavitation number gave the total number of expected cavitation events.

As was mentioned previously, the distribution shown in Fig. 2 would normally be represented by measurements in definite size ranges, and thus the widths for these ranges must be determined. The total distribution was divided into eight, eleven, and sixteen subranges, respectively, and the results were compared. These distributions are shown in Figs. 3, 4, and 5 with subrange widths of  $2 \times 10^{-5}$ ,  $1.5 \times 10^{-5}$ , and  $10^{-5}$  ft. The cavitation occurrence characteristics as computed by the analytical model are compared with the result for the continuous distribution in Fig. 6. It is apparent that the width of the size ranges for the example presented should not exceed  $10^{-5}$  ft for an accurate prediction of the severity of cavitation.

The precision with which the probability of bubbles in any given size range can be determined is also an essential parameter used directly in the computation for the cavitation occurrence rate. The individual errors incurred in determining the probability for each size range accumulate and contribute to a total error in the occurrence rate.

### III. THE ACOUSTIC PULSE ATTENUATION TECHNIQUE

As stated previously, the measurement is based on the amplitude attenuation of an acoustic tone burst as it is propagated through the test water. The attenuation is measured as a function of the carrier frequency of the tone burst over the range of interest. The relationship between the bubble size and its resonant frequency and the fact that only bubbles at or near resonance contribute significantly to the sound attenuation are well known [8]. This effect can be measured using the tone burst technique, and the measured attenuation due to the bubbles is related to the concentration of bubbles. The ratio between the tone burst amplitudes under test conditions with bubbles present and under high-pressure conditions with bubbles absent is a measure of the attenuation due to the resonant and near resonant bubbles.

Since the acoustic attenuation is related to the volume concentration of resonance-size bubbles, and the analytical cavitation model requires the number of bubbles contained in any size range, the data reduction process must include the conversion from volume concentration to discrete numbers of bubbles.

The process of computing the bubble volume concentration in each size range from the distribution involves a somewhat complicated infinite series. However, as a first approximation the range volume concentration can be obtained by multiplying the number of bubbles flowing in the size range by the volume of the bubble size in the center of the range divided by the volume of water flowing through the test section. The accumulation of the concentrations determined in this manner differed by a factor of only approximately 6 per cent from the total concentration determined by integrating the assumed distribution function over the entire range using the relationship

$$c = \frac{4\pi N \times 10^6}{3Q\delta^2} \int_0^{\infty} r^4 e^{-r^2/2\delta^2} dr = \frac{2\sqrt{2} \pi^{3/2} N \delta^3}{Q} \times 10^6$$

where  $c$  = total volume concentration in ppm,  
 $N$  = total number of bubbles flowing through the test section per second,  
 and  $Q$  = total volume of water through the test section per second.

In the present example the total number of bubbles flowing,  $N$ , is a parameter which can be used to illustrate several points. The data of Fig. 5 are presented in Fig. 7 on the basis of volume concentration in each range for an overall volume concentration of 0.335 ppm, which in this example amounts to  $N = 10^5$  bubbles per second flowing through the test section. For other concentrations this plot is merely scaled by the value of  $N$ . It is significant that the maximum in the concentration distribution occurs at a much higher bubble size than the maximum of the nuclei number distribution.

The acoustic attenuation to be expected with the concentration and distribution of microbubbles present in a typical cavitation inception

study was calculated. The results are shown in Fig. 8 based on an assumed acoustic path length of 16 cm. The computer program developed by Brockett [9] was used to make the computations shown. The program was modified to use concentration and size as input data and produce the value of acoustic attenuation.

Based on the results of Fig. 8 a theoretical estimate can be made of the sensitivity and precision required for an attenuation measuring technique to be useful in measuring microbubble size and concentration. The attenuation ratios given in Fig. 8 would correspond to voltage amplitude ratios as measured by the pulse height tone burst method of acoustic attenuation measurement [2].

The acoustic pulse propagation method places upper and lower practical limits on the microbubble concentration which can be measured. The upper limit is imposed by an attenuation which suppresses all signals from the transmitter received at the receiving transducer. High levels of transmitted power will alter the microbubble properties. The lower limit is imposed by the change in receiving transducer output which can be identified as due to microbubbles rather than to random effects of temperature, random microbubble structure, and diffraction effects [10].

The attenuations calculated in Fig. 8 are based on a plane wave solution. Acoustic measurements using the pulse method in a water tunnel have previously required the use of spherical waves (the sonic pulse frequency and wavelength to be used are determined by bubble size, and the size of the water tunnel test section limits the size of the transducer; these two tank limitations make the generation of plane waves in the test section very difficult). Spherical waves generated by a small transducer are only approximately spherical near the transducer; consequently diffraction effects arise which may give a changing virtual path length at different frequencies and concentrations. No study has been found which gives the precise correction necessary at the low attenuations (or concentrations).

Experience in reading the pulse height gives rise to the following optimistic estimates: The test section pressure, of the order of 1/4 atmosphere, is typical. Under these conditions the resonant frequencies of interest range from approximately 25 kHz to 300 kHz, and appropriate acoustic transducers should be utilized. Ceramic disks 0.25 inches in diameter by 0.1 inches thick would be satisfactory.

For the particular size distribution in this example, the total volume concentration range within which the technique could be useful is that between 0.033 ppm and 1.0 ppm, which correspond, respectively, to  $10^5$  and  $3 \times 10^6$  nuclei per second through the test section. Similar reasonable size distributions should give results which differ very little from those reported here. Attempts to measure concentrations outside the concentration range specified above would result in attenuations either too small to be observed or so large that the signal was completely obliterated.

It is standard practice in many water tunnels to operate with a dissolved gas concentration whose relative saturation at test section pressure is unity. In most other sections of the tunnel the relative saturation based on the local pressure is therefore less than unity and the water is undersaturated. Based on atmospheric pressure rather than test section pressure, this amounts to a relative saturation of about 0.25 in the test section for the example discussed in this work.

A comparison between the results of the analytical model shown in Fig. 6 and extrapolations from experimental data taken by the authors [5] in the St. Anthony Falls six-inch water tunnel under very similar circumstances (identical flow conditions, minimum pressure coefficients, and body diameters) reveals that the free gas content of the experimental water is described by a cavitation characteristic similar to one of the lower curves of Fig. 6. In other words, the free gas content of the tunnel water in the test section is at least an order of magnitude lower than the 0.03 vppm which represents the lower limit of the acoustic attenuation method. In many practical situations, therefore, the acoustic method lacks sufficient sensitivity to be useful.

The minimum volume concentration for which the acoustic pulse attenuation technique is valid is indicated for the present example in Fig. 6 by the characteristic line representing a total concentration of 0.0335 vppm or, in terms of this example, by  $10^5$  nuclei per second passing through the test section. Under these circumstances very high cavitation occurrence rates are observed at cavitation numbers just under the absolute magnitude of the minimum pressure coefficient of the test body.

Comparison of this characteristic line with the experimental data shows that under these circumstances in the St. Anthony Falls tunnel the water must exhibit a relative saturation of 2 based upon the test section

pressure. This value substantially exceeds the standard operating conditions of many tunnels.

Certain errors are also expected within the volume concentration range in which the acoustic pulse attenuation technique is valid. In the experiments, voltages are measured which represent the acoustic pulse heights. It is estimated that this measurement can be accurate only to plus or minus 3 per cent of full scale with present techniques. When this error is applied to the expected attenuations of Fig. 8, it becomes apparent that over large portions of the total range no differences can be detected experimentally in the volume concentration measurements. This would tend to flatten the distribution shown in Fig. 7.

No difference in volume concentration could be determined in the present example over 62 per cent of the total range of the distribution at the 0.033 ppm level under the worst conditions. At the 0.165 ppm volume concentration level this proportion dropped to 31 per cent. At the 0.33 ppm level about 19 per cent of the range could not be separated under the worst conditions. As the volume concentration rose from this value to the 1.0 ppm level, the proportion rose to 31 per cent again. As previously stated, beyond this concentration the errors are so large that the technique is probably not valid.

Presented from an alternate point of view, the calculation also revealed that in the usable range the technique is sensitive to approximately one resonant bubble per cubic wavelength at the usable minimum attenuation and to approximately 35 resonant bubbles per cubic wavelength at maximum usable attenuation. This calculation is nearly independent of resonant frequency.

#### IV. CONCLUSIONS

It is concluded that in general the acoustic tone burst attenuation technique is usable in a concentration range of approximately 0.03 to 1.0 parts per million, which is about an order of magnitude higher than the expected concentration level in water tunnel test sections under most usual cavitation inception test conditions. Furthermore, within the usable range the concentration contained in a small subrange cannot be resolved or determined using present methods and techniques with enough precision to be usable in the analytical cavitation model to produce valid cavitation occurrence rate results.



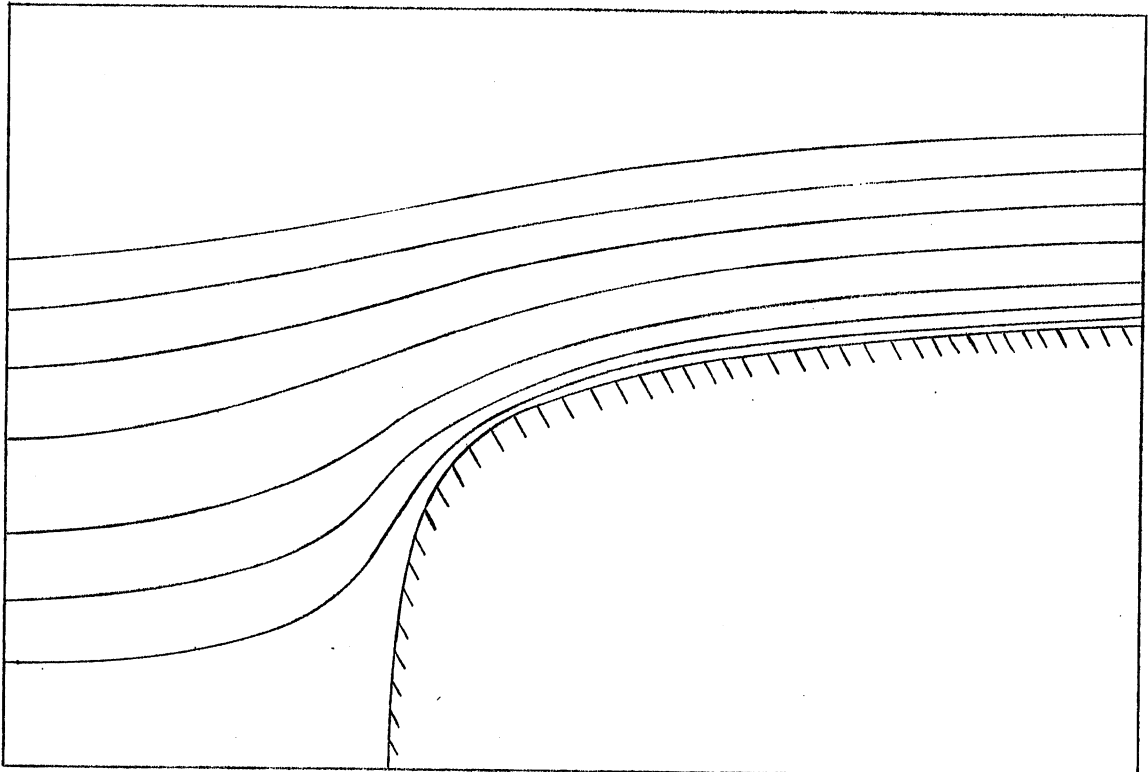
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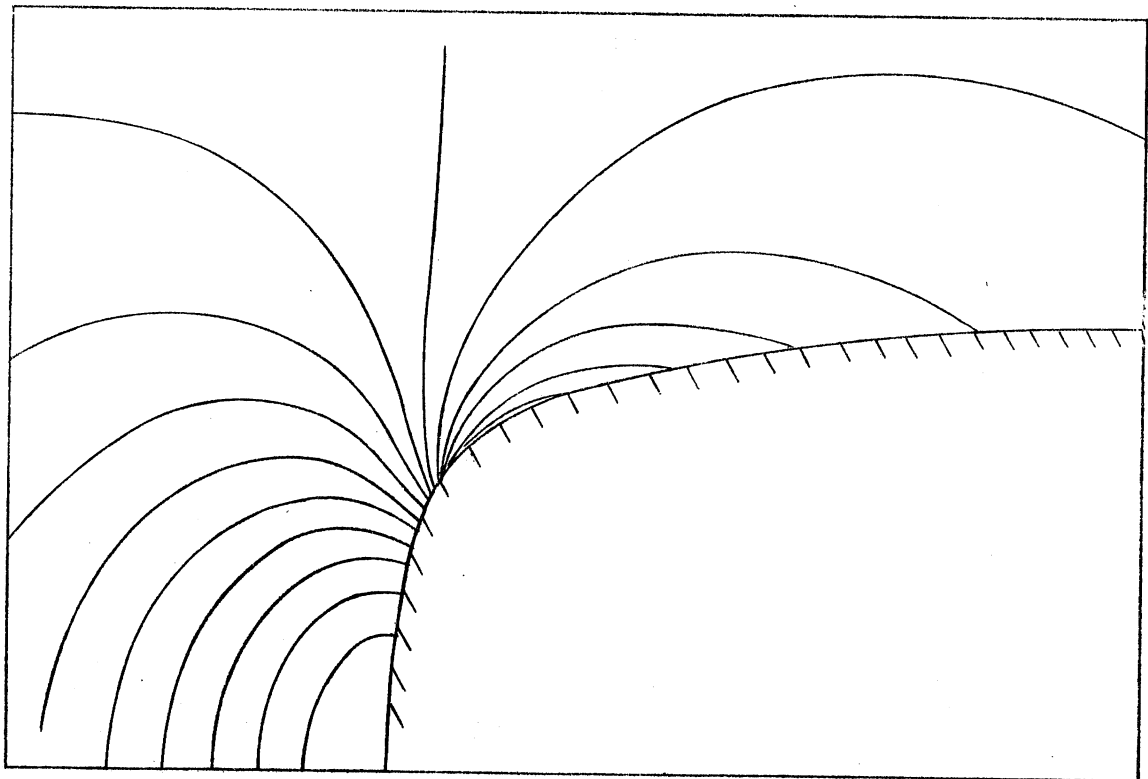


I L L U S T R A T I O N S

(Figs. 1 thru 8)



a. Body with Streamlines



b. Body with Isobars

Fig. 1 - Typical Blunt Body,  $C_{p,min} = -0.6$

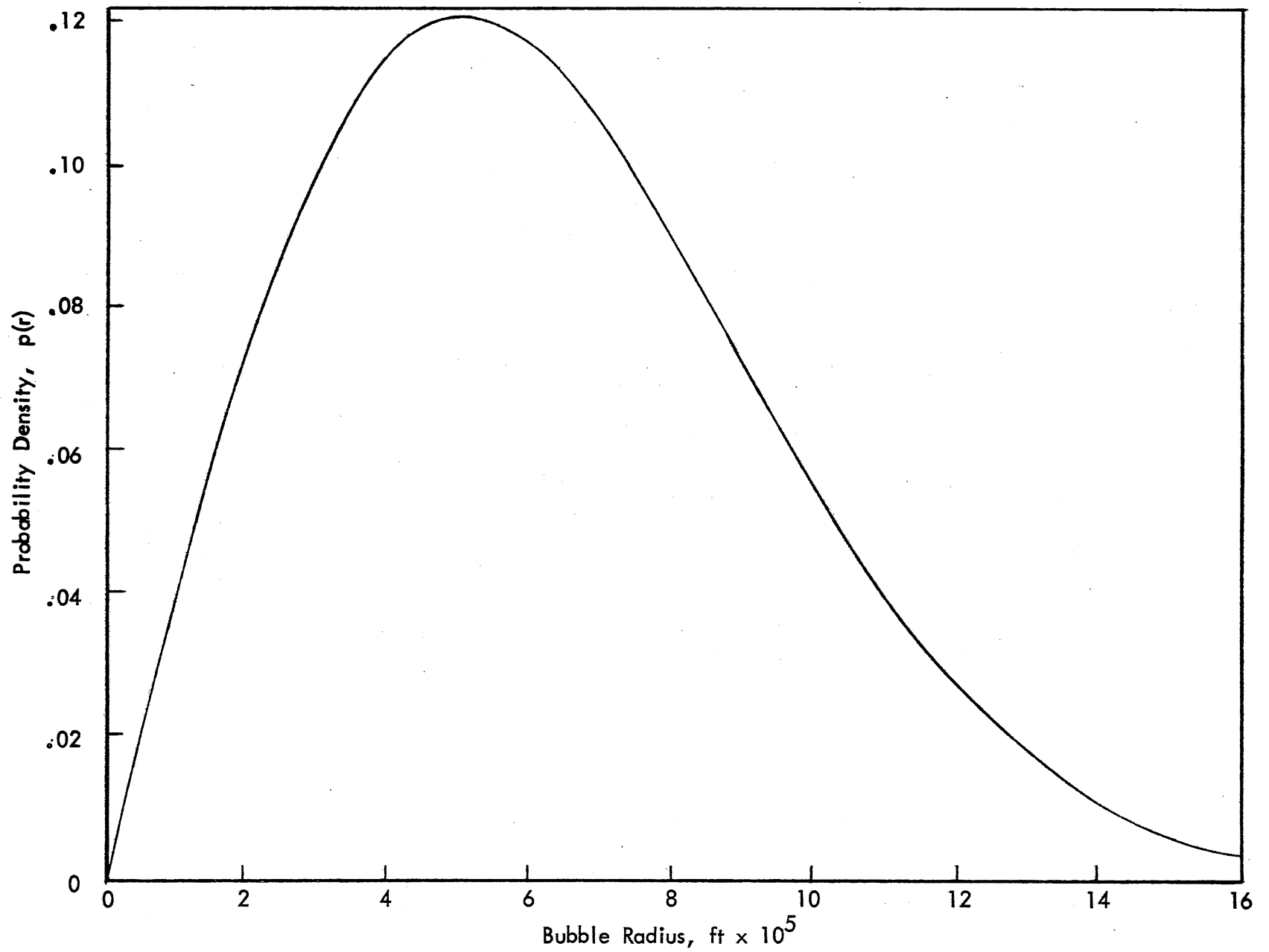


Fig. 2 - Continuous Assumed Bubble Size Distribution

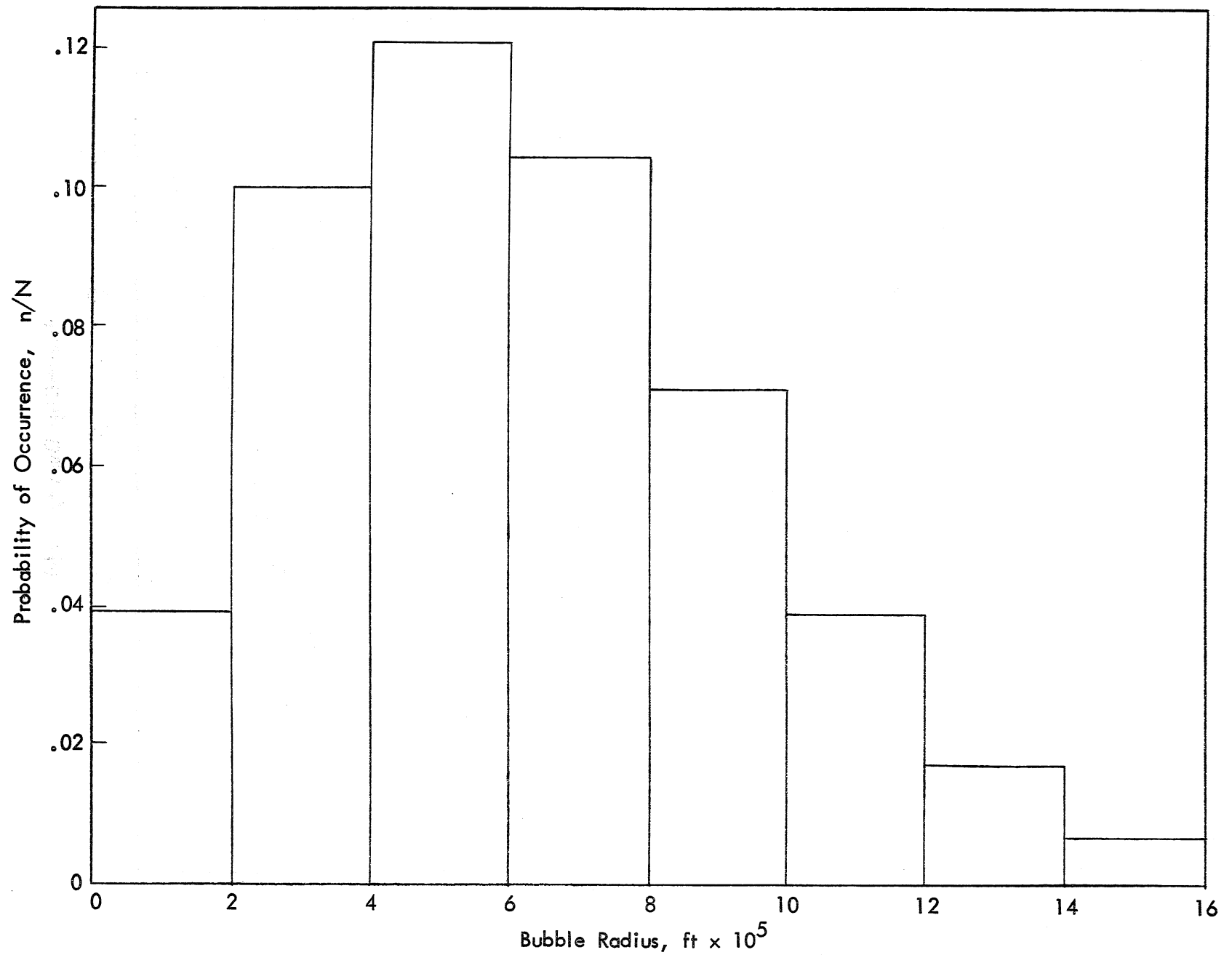


Fig. 3 - Bubble Size Distribution,  $2 \times 10^{-5}$  ft Subrange Width

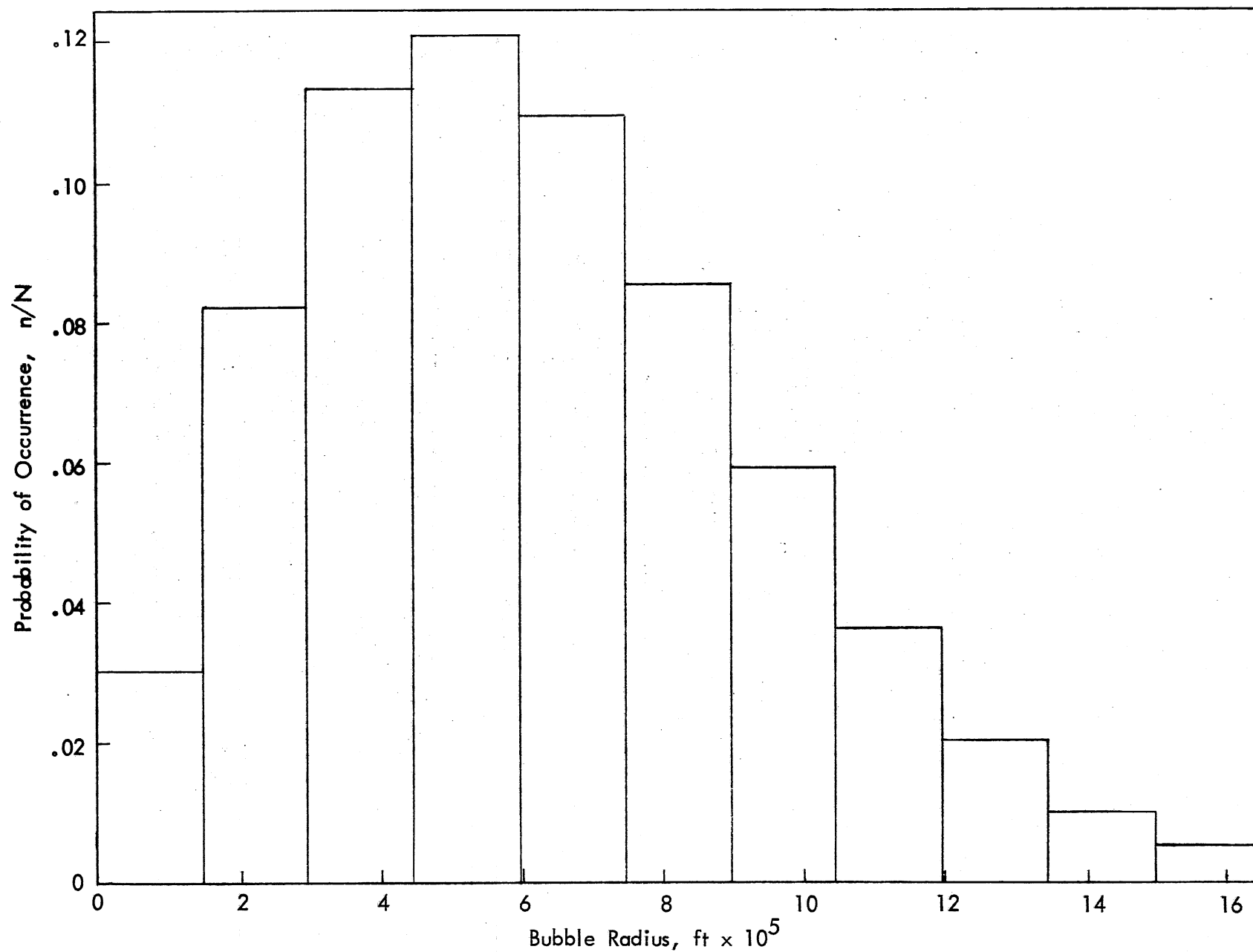


Fig. 4 - Bubble Size Distribution,  $1.5 \times 10^{-5}$  ft Subrange Width

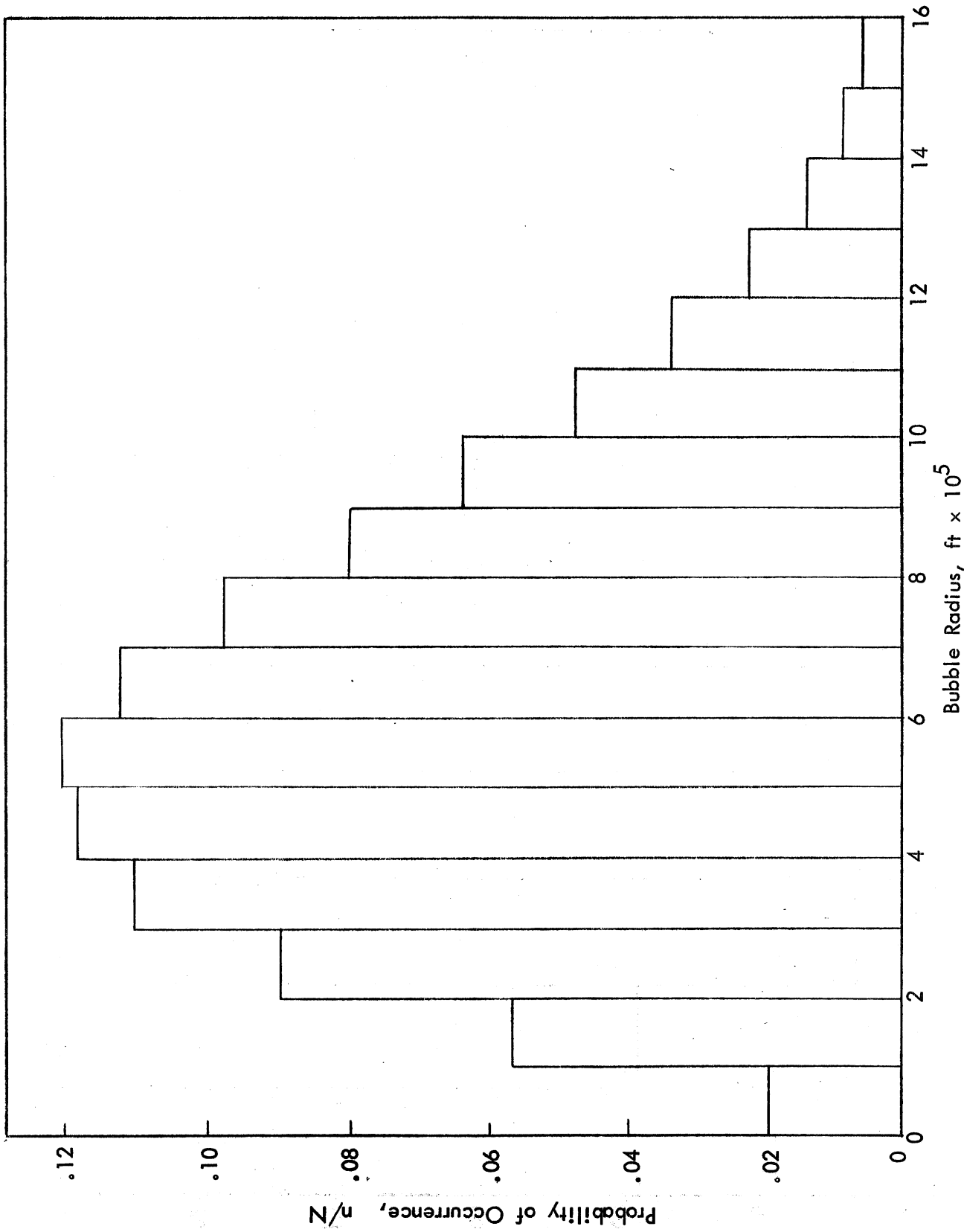


Fig. 5 - Bubble Size Distribution, 10<sup>-5</sup> ft Subrange Width



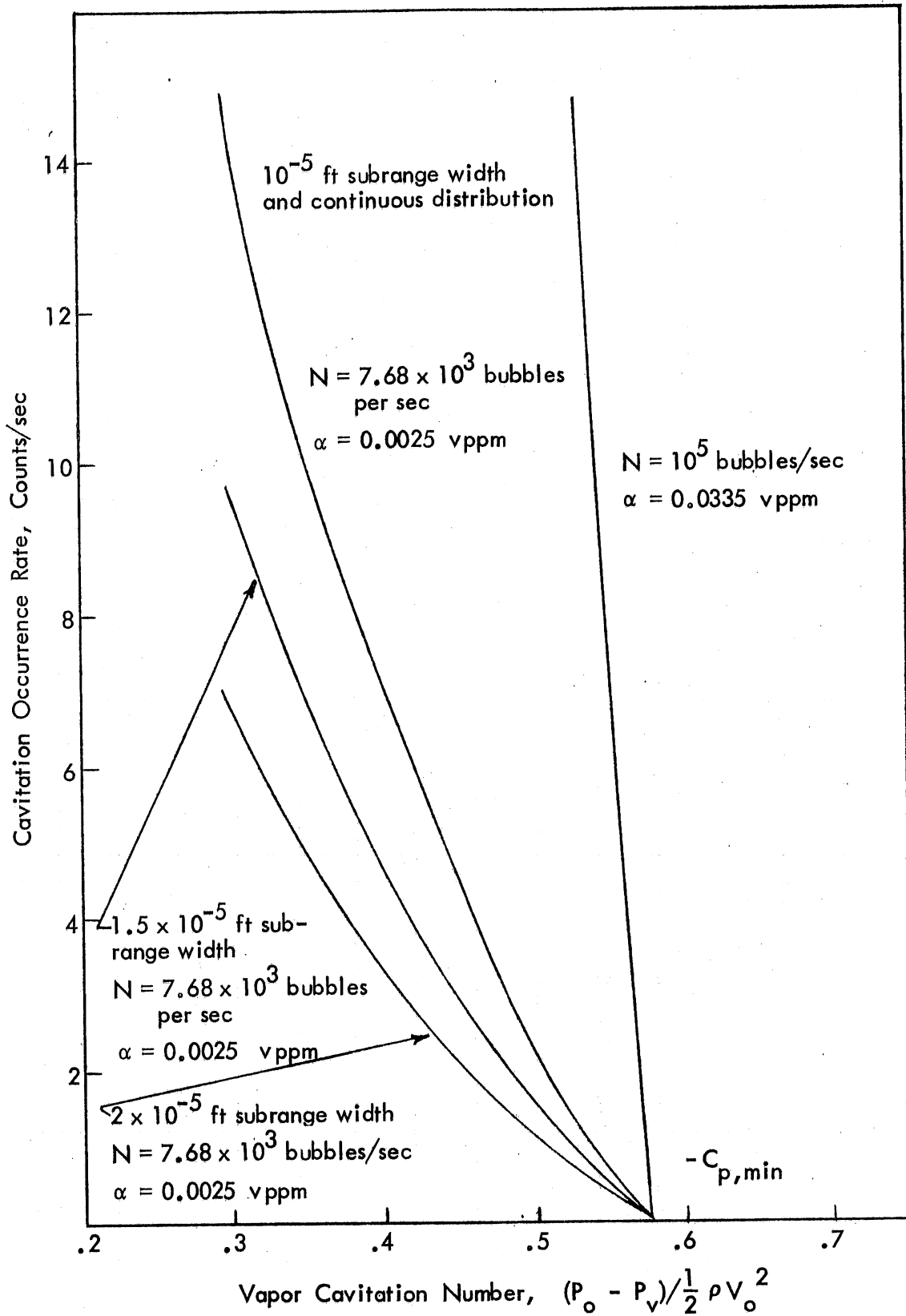


Fig. 6 - Effect of Subrange Width and Volume Concentration on the Computed Cavitation Occurrence Rate Characteristics for Source Disk Body,  $C_{p,min} = -0.6$

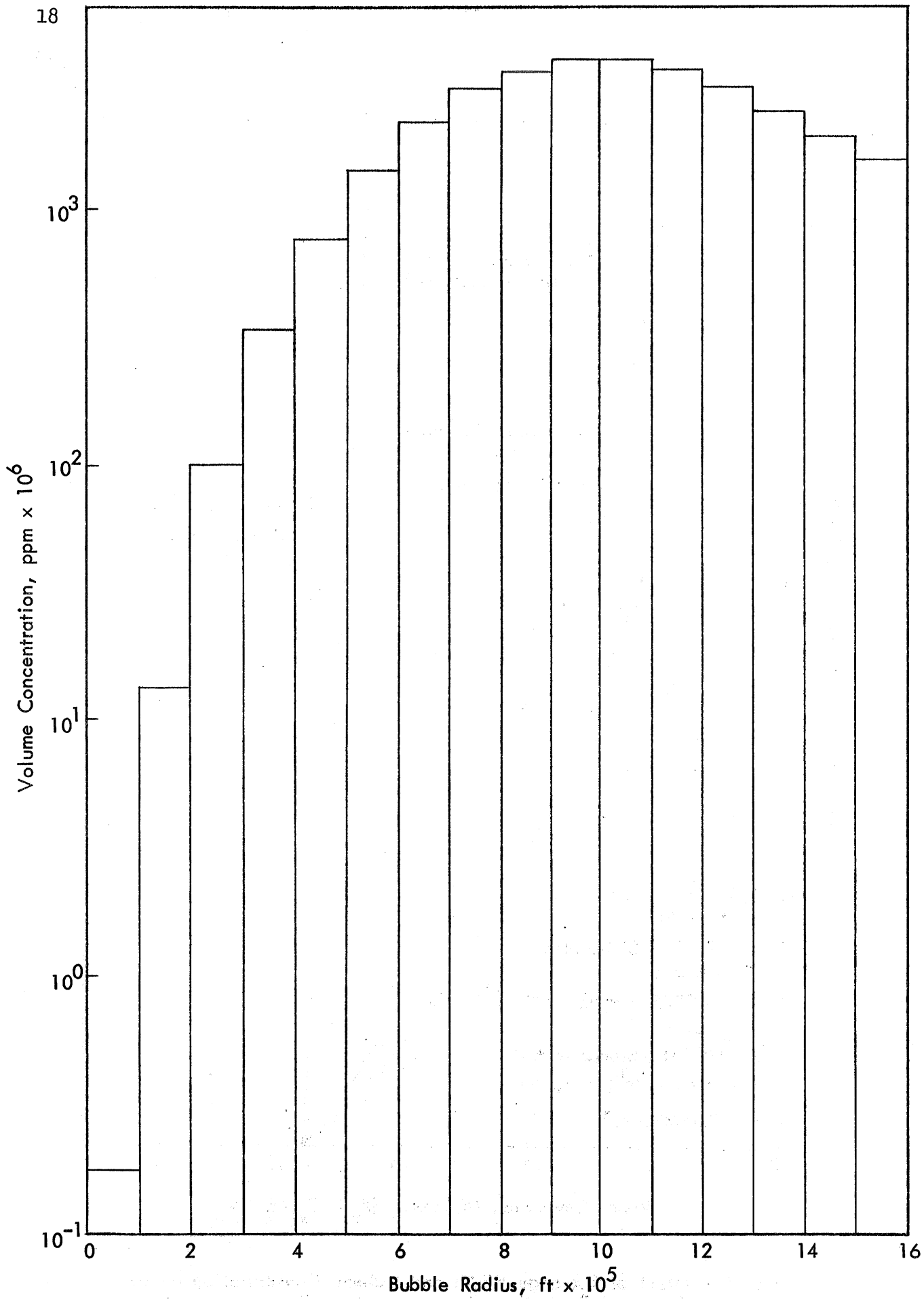


Fig. 7 - Volume Concentration Distribution,  $\alpha = 0.0335$  vppm

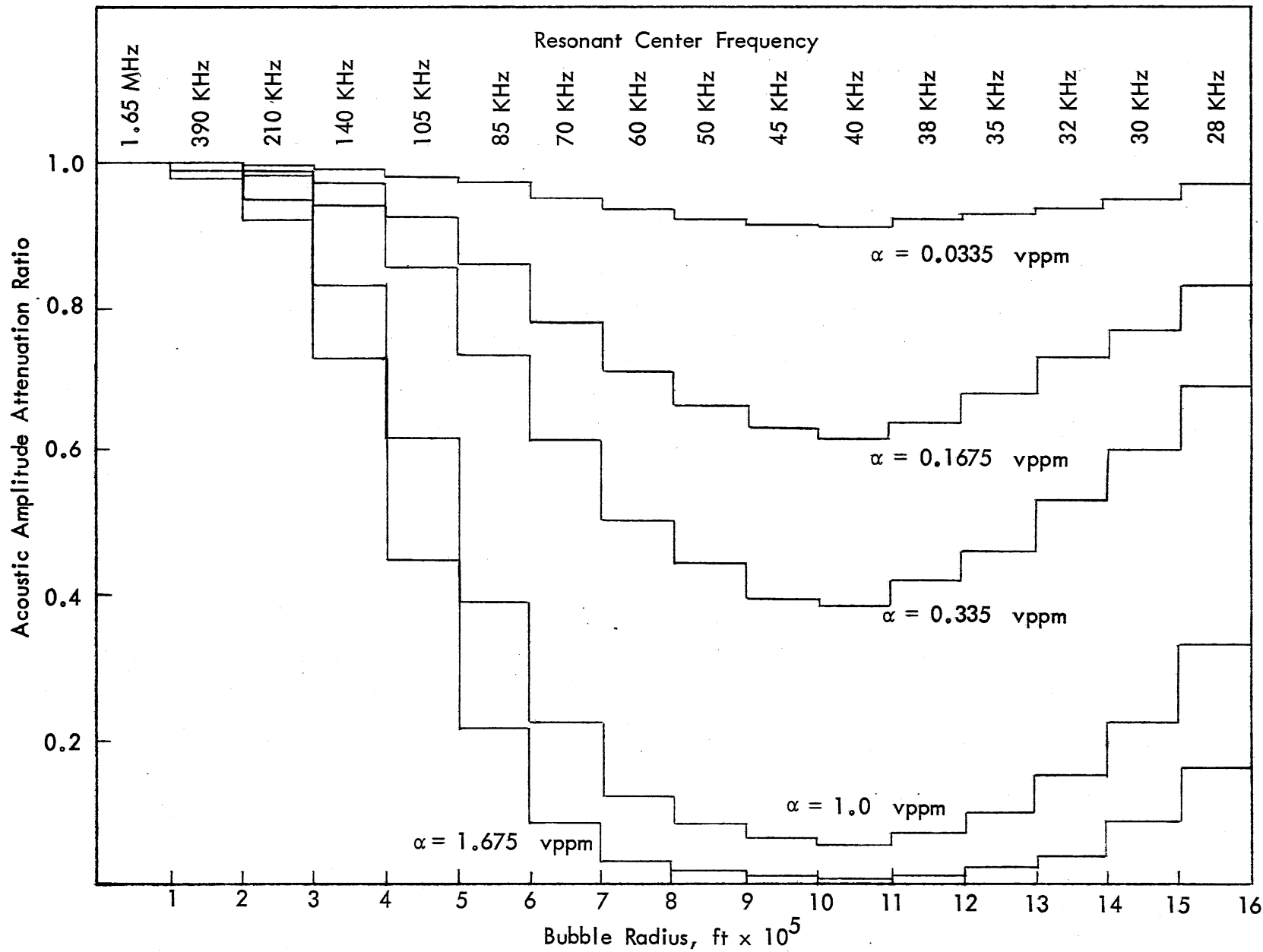


Fig. 8 - Calculated Acoustic Amplitude Attenuation Distribution for various Volume Concentration Levels



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